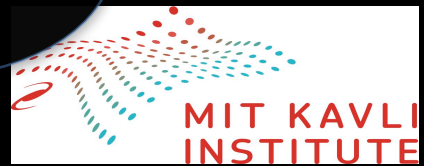


LIGO

O

Salvatore Vitale
MIT- LIGO Laboratory

Takayama, Aug 26th 2016



Characterization of compact objects with gravitational waves

Salvatore Vitale
MIT- LIGO Laboratory

Takayama, Aug 26th 2016

The dark side of astrophysics

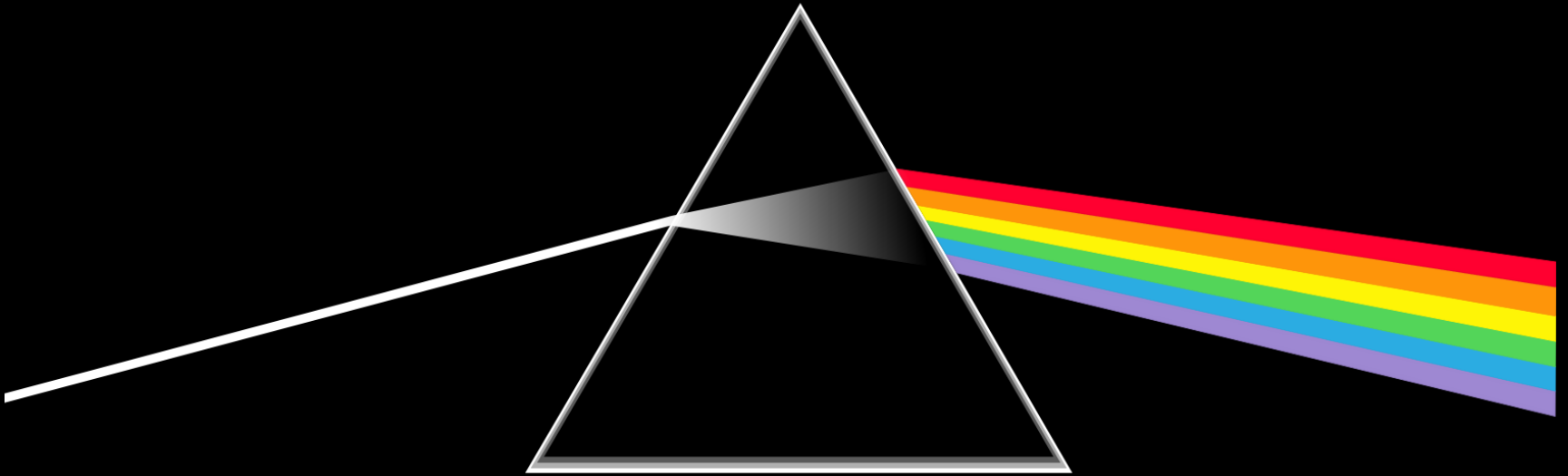
- Over the last few thousand years humanity has used electromagnetic waves to study the universe
- Light has been pivotal to our understanding of the cosmos, stars, planets
 - However light can be easily deflected, absorbed, obscured

The dark side of astrophysics

- Over the last few thousand years humanity has used electromagnetic waves to study the universe
- Light has been pivotal to our understanding of the cosmos, stars, planets
 - However light can be easily deflected, absorbed, obscured
- Now, it's the time for the dark side

LIGO

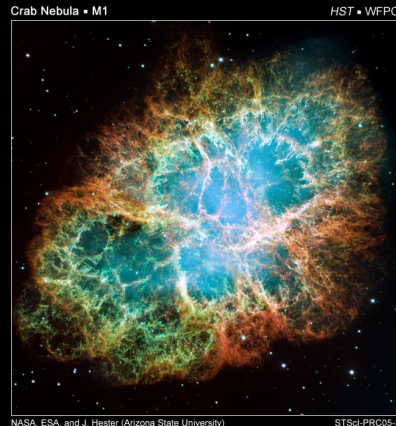
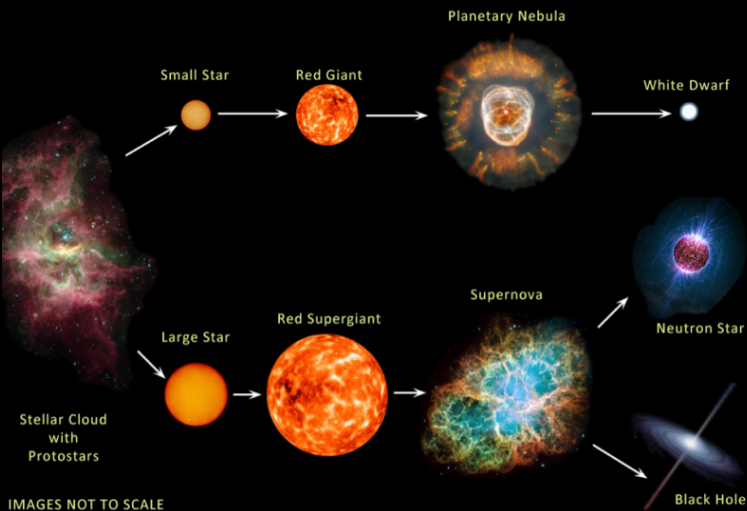
The dark side of astrophysics



Compact objects

- Compact objects such as neutron stars (NS) and black holes (BH) host some of the most extreme conditions in the universe

EVOLUTION OF STARS



Black holes

- Leftovers of massive stars
- Produce extreme gravitational fields
- Does general relativity still hold true near a BH?
- How fast can they spin?
- How big can they get?
- When did the first BHs form?

Neutron stars

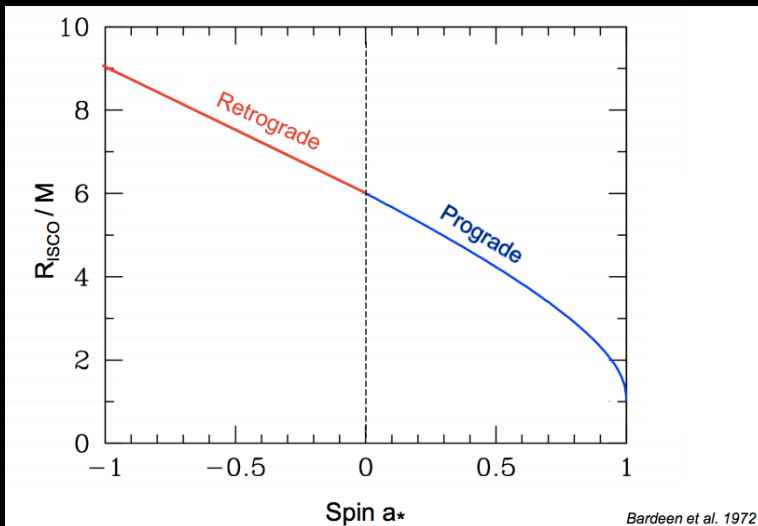
- The most dense objects we can observe
 - A mass of $1.4 M_{\odot}$ contained in a sphere with radius of 10 Km
- How does matter behave in these extreme conditions?
- Are neutron stars related to GRBs? And metal production?
- What is the maximum mass of a neutron star?

BH spins (with EM)

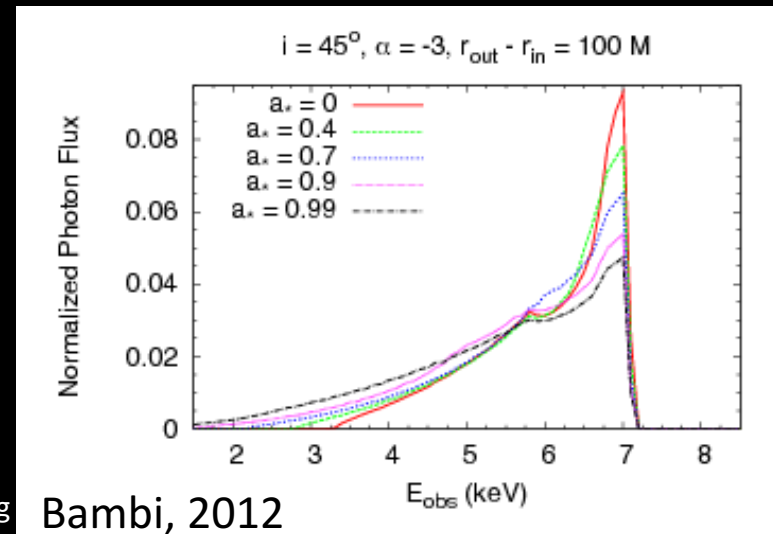
- Traditionally, the spin of black holes has been estimated through its effects on a surrounding disk
- Need an accreting black hole (e.g. in a X-ray binary)

BH spins (with EM)

- If a BH is spinning, the radius of the innermost stable circular orbit will get closer (**Continuum fitting**)
- If the debris in the disk reflect light, the spectral lines will be distorted by GR effects which depend on the spin (**FE-line**)



ale, Aug



BH spins (with EM)

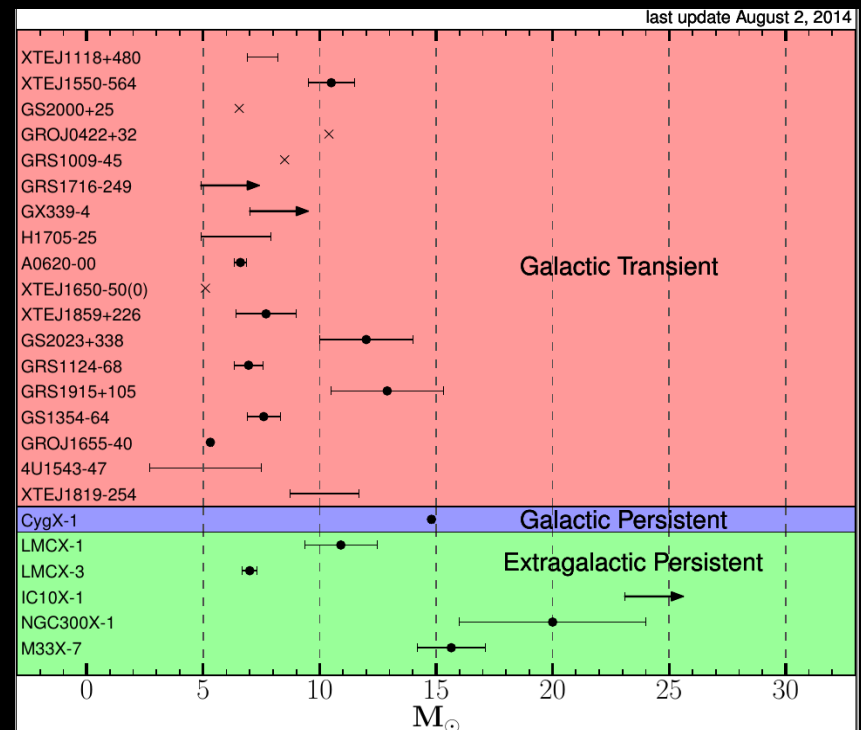
- Both methods rely on a good understanding of the disk physics and are *indirect* measurement of spin
- Sometime in tension with each other

| System | a_* (CF) | a_* (Fe line) | No. obs. | References | |
|---------------|-----------------|-----------------|----------|----------------------------------|---|
| Cygnus X-1 | > 0.983 | 0.97 ± 0.02 | 9 / 1 | Gou+ 2011, 2014 Fabian+ 2012 | ✓ |
| LMC X-1 | 0.92 ± 0.06 | $0.72 - 0.99$ | 19 / 1 | Gou+ 2009 Steiner+ 2012 | ✓ |
| GRS 1915+105 | > 0.95 | 0.98 ± 0.01 | 6 / 1 | McClintock +2006 Miller +2013 | ✓ |
| XTE J1550-564 | 0.34 ± 0.24 | 0.55 ± 0.20 | 60 / 2 | Steiner, Reis+ 2011 | ✓ |
| GRO J1655-40 | 0.8 ± 0.1 | > 0.9 | 33 / 2 | Shafee+ 2006 Reis+ 2009 | ✗ |
| 4U 1543-47 | 0.7 ± 0.1 | 0.3 ± 0.1 | 34 / 1 | Shafee+ 2006 Miller+ 2009 | ✗ |

BH mass (with EM)

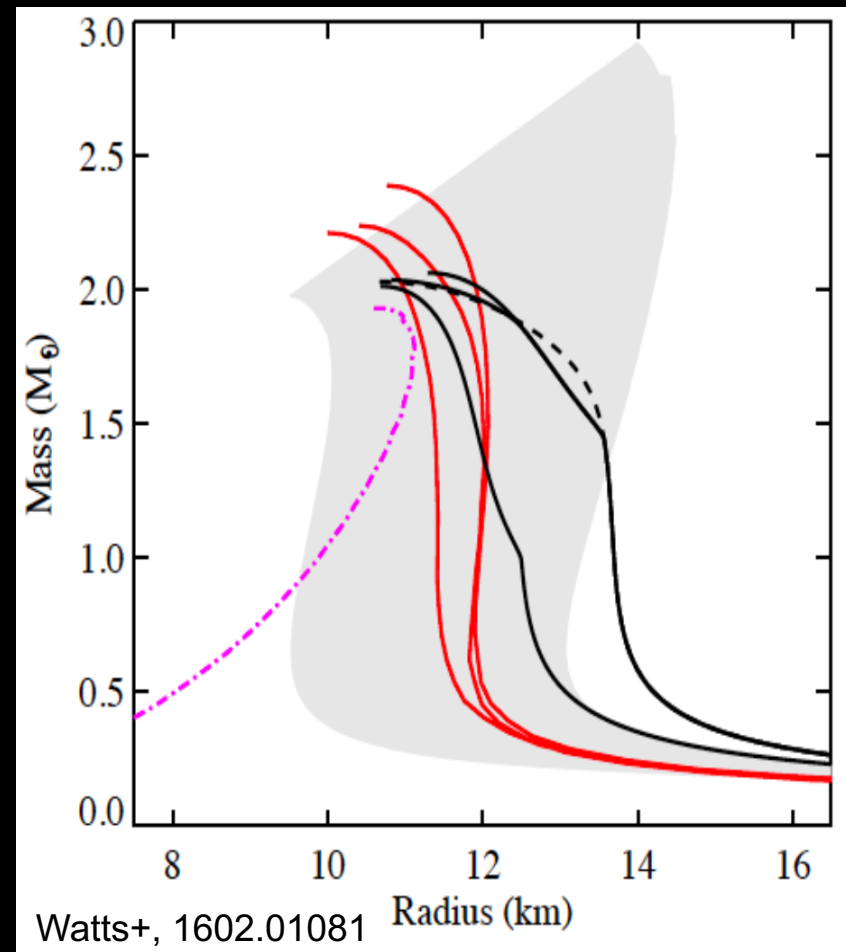
- Also rely on having a luminous companion
- Requires period, radial velocity, inclination, companion mass
- Indirect measurement

D. Lorimer



NS equation of state (with EM)

- The equation of state summarizes how matter behaves in NS
- Boils down to simultaneous measurement of mass and radius
- Possible with EM, but challenging
 - Mass estimates not always reliable
 - Radius estimates non often available and not always reliable
 - NICER to launch 2016, 5% precision



Gravitational waves

- When two compact objects orbit around each other, they emit gravitational waves (GW) that encode all of the system's properties
- Compact binary systems can thus be used to study BH and NS without the need for light just measuring the GW they emit.

What are they?

- Gravitational waves are ripples in the space-time continuum, emitted by any system with a non-constant quadrupole moment
- Yes, you can produce gravitational waves too!

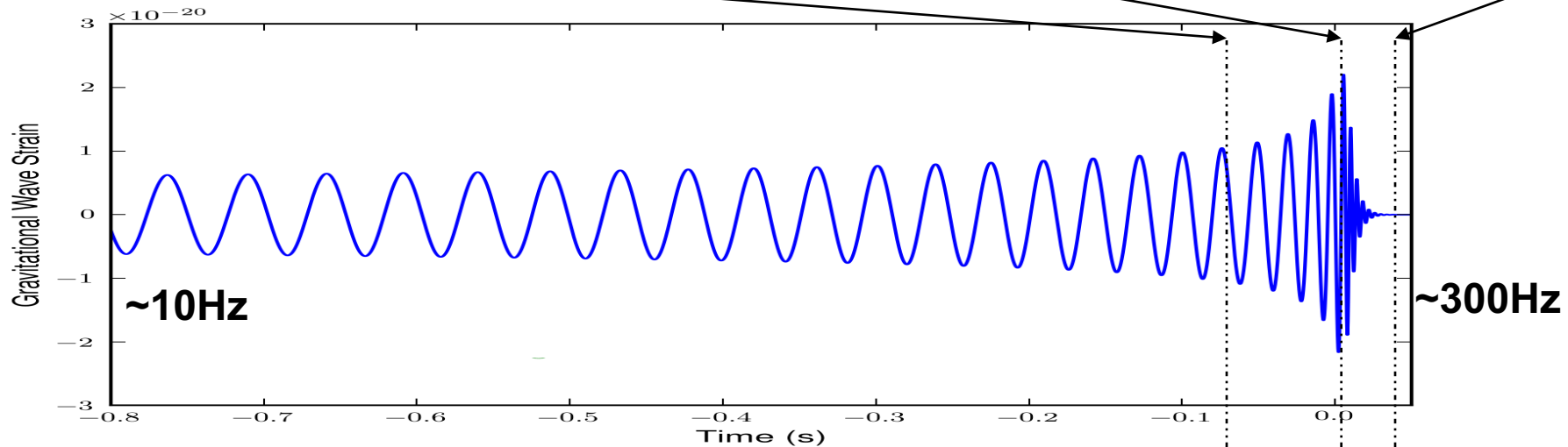
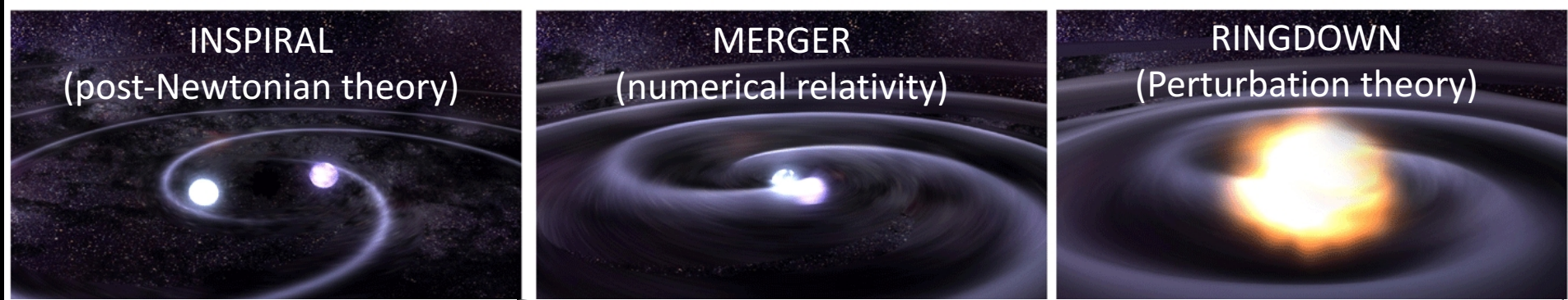
- Go through quadrupole formula and the simplest CBC signal to give an idea of

- $F_{GW} = 2 F_{orbit}$
- Distance dependence
- Mass dependence
- Expected strain
- Contraction and stretch in perpendicular directions

$$h_{ij} \propto \frac{4}{D} \left(\frac{\mathcal{M}G}{c^2} \right)^{5/3} \left(\frac{\pi f_{GW}}{c} \right)^{2/3}$$

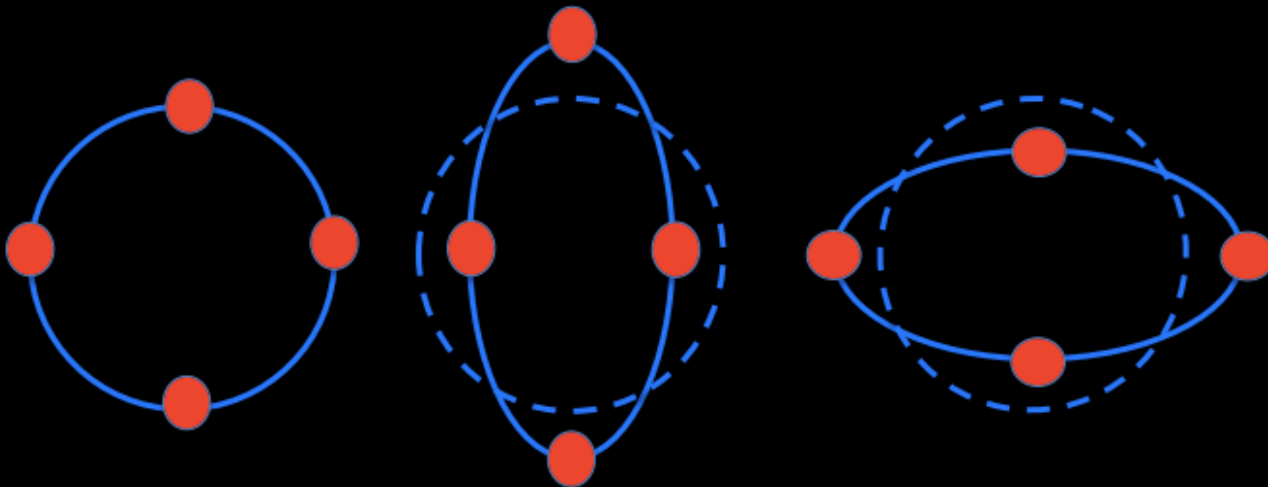
Compact Binaries Coalescences

CBC Gravitational Wave signals are very well understood



Effect of GW

- While passing through space, GW vary the distance between free floating observers
 - Distances stretch in one direction and squeeze in the perpendicular direction



Order of magnitude estimate

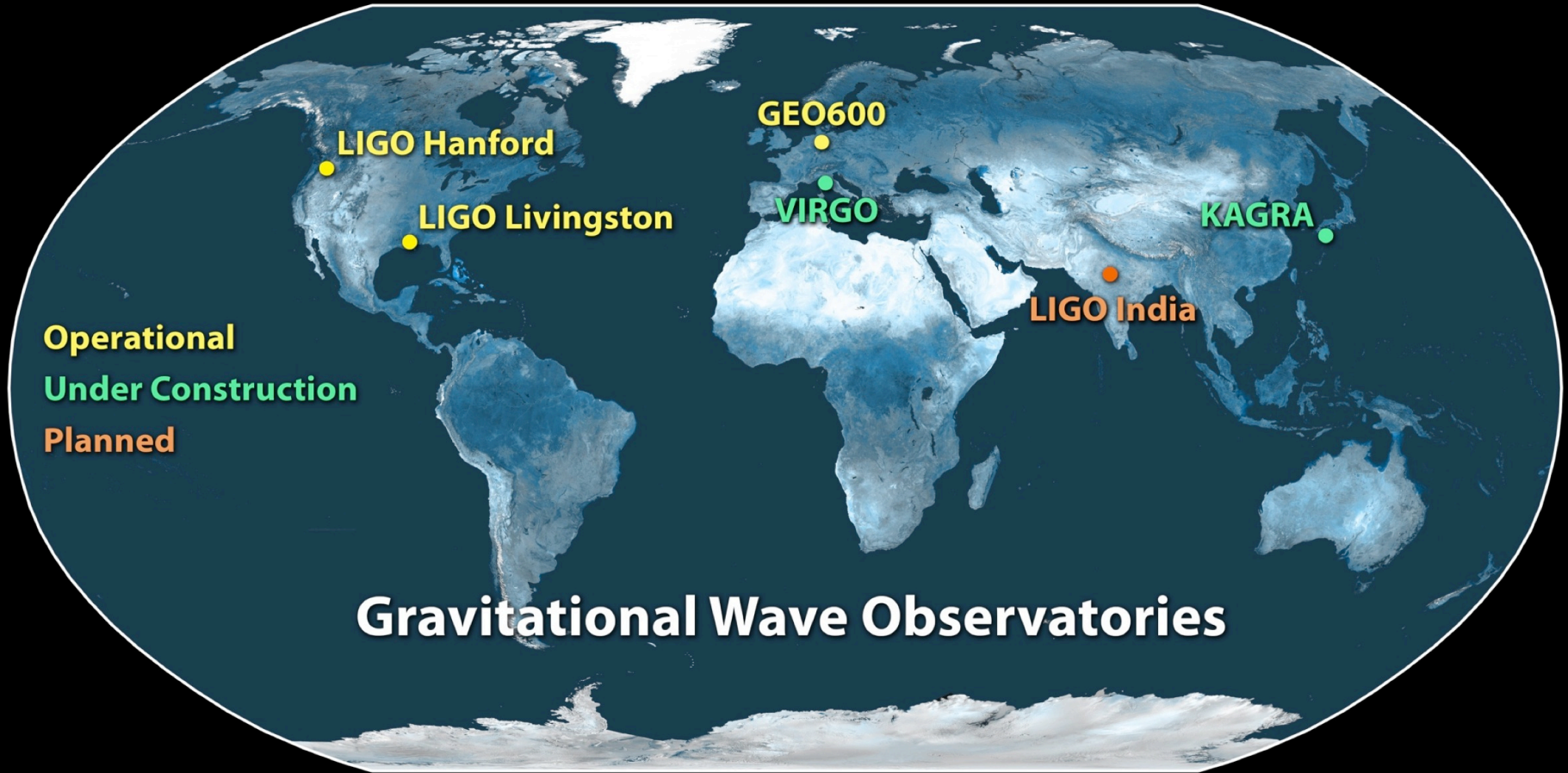
```
In [25]: Ms=1.989e30 # Kg
c=299792458 # m/s
m1=30*Ms
m2=30*Ms
G=6.67408e-11 # SI
mu=m1*m2/(m1+m2)
M=m1+m2
Mchirp=mu**(3/5.)*M**(2/5.)
freq=100. # Hz
Mpc=3.0857e24/100. # 1 Mpc in m
Distance=500.*Mpc

In [27]: 4./Distance*(Mchirp*G/c/c)**(5/3.)*(pi*freq/c)**(2/3.)

Out[27]: 1.1779380967355013e-21
```

Two $30M_{\odot}$ BHs at 500 Mpc would produce a strain (i.e. relative length variation) at Earth of roughly 1 part in 10^{21}

The network of Gravitational Wave detectors




Detections!!

- Advanced LIGO and Virgo detected 2 binary black hole coalescences in their first science run

PRL 116, 061102 (2016) week ending
12 FEBRUARY 2016

Selected for a **Viewpoint** in *Physics*
PHYSICAL REVIEW LETTERS




Observation of Gravitational Waves from a Binary Black Hole Merger

B. P. Abbott *et al.*^{*}
(LIGO Scientific Collaboration and Virgo Collaboration)
(Received 21 January 2016; published 11 February 2016)

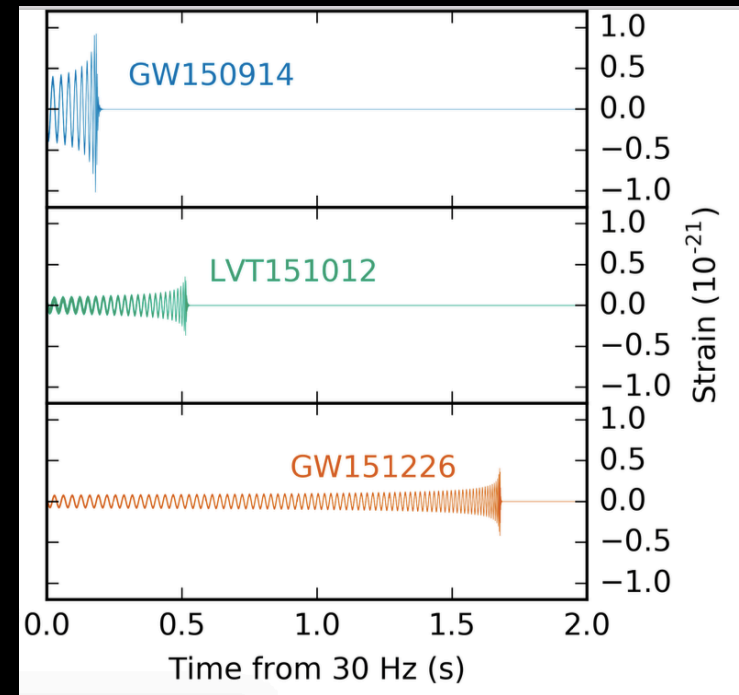
PRL 116, 241103 (2016) week ending
17 JUNE 2016

PHYSICAL REVIEW LETTERS



GW151226: Observation of Gravitational Waves from a 22-Solar-Mass Binary Black Hole Coalescence

B. P. Abbott *et al.*^{*}
(LIGO Scientific Collaboration and Virgo Collaboration)
(Received 31 May 2016; published 15 June 2016)






Parameter estimation

- The (unknown) parameters of a CBC source can be estimated using Bayesian methods
 - Explore a high dimensionality parameter space using stochastic samplings (MCMC, nested sampling)

$$p(\theta|d) \propto p(d|\theta)p(\theta)$$

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Mass estimation (with GW)

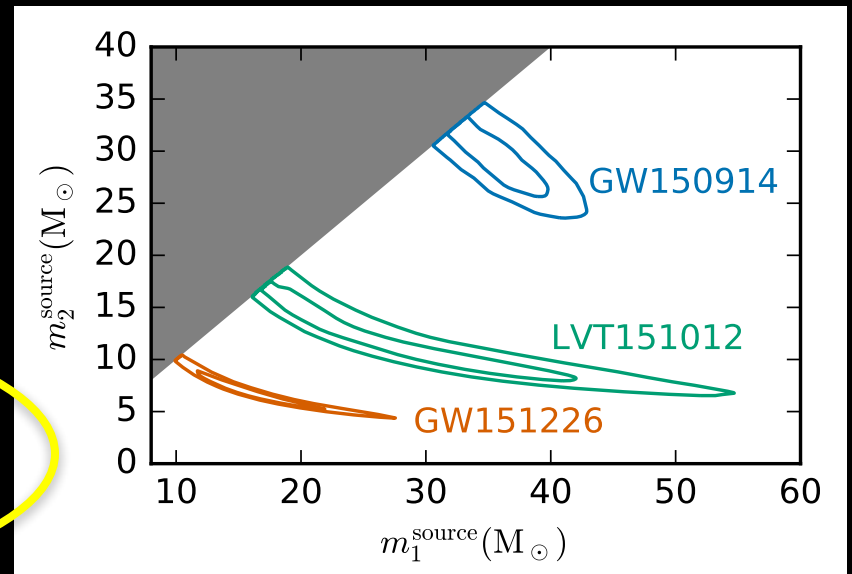
- The masses of the two objects directly affect the phasing evolution of a GW signal
 - Very good at estimating “chirp” mass
 - Worse for component masses
- This is a *direct* measurement, the masses directly affect the amount and frequency of GW emitted

Mass estimation (with GW)

- Typically, longer signals (i.e. lower masses) will lead to better estimation of masses, since we can “follow” the signal for more cycles

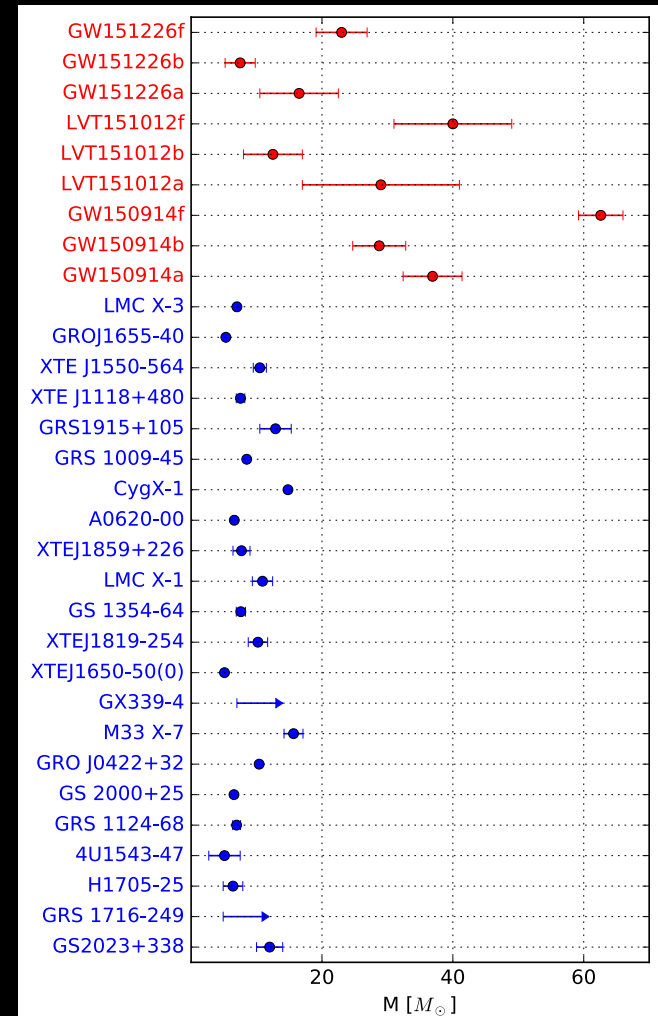
LVC, 1606.04856

| Event | GW150914 | GW151226 | LVT151012 |
|--|------------------------|------------------------|----------------------|
| Signal-to-noise ratio ρ | 23.7 | 13.0 | 9.7 |
| False alarm rate FAR/yr ⁻¹ | $< 6.0 \times 10^{-7}$ | $< 6.0 \times 10^{-7}$ | 0.37 |
| p-value | 7.5×10^{-8} | 7.5×10^{-8} | 0.045 |
| Significance | $> 5.3\sigma$ | $> 5.3\sigma$ | 1.7σ |
| Primary mass $m_1^{\text{source}}/M_\odot$ | $36.2^{+5.2}_{-3.8}$ | $14.2^{+8.3}_{-3.7}$ | 23^{+18}_{-6} |
| Secondary mass $m_2^{\text{source}}/M_\odot$ | $29.1^{+3.7}_{-4.4}$ | $7.5^{+2.3}_{-2.3}$ | 13^{+4}_{-5} |
| Chirp mass $\mathcal{M}^{\text{source}}/M_\odot$ | $28.1^{+1.8}_{-1.5}$ | $8.9^{+0.3}_{-0.3}$ | $15.1^{+1.4}_{-1.1}$ |

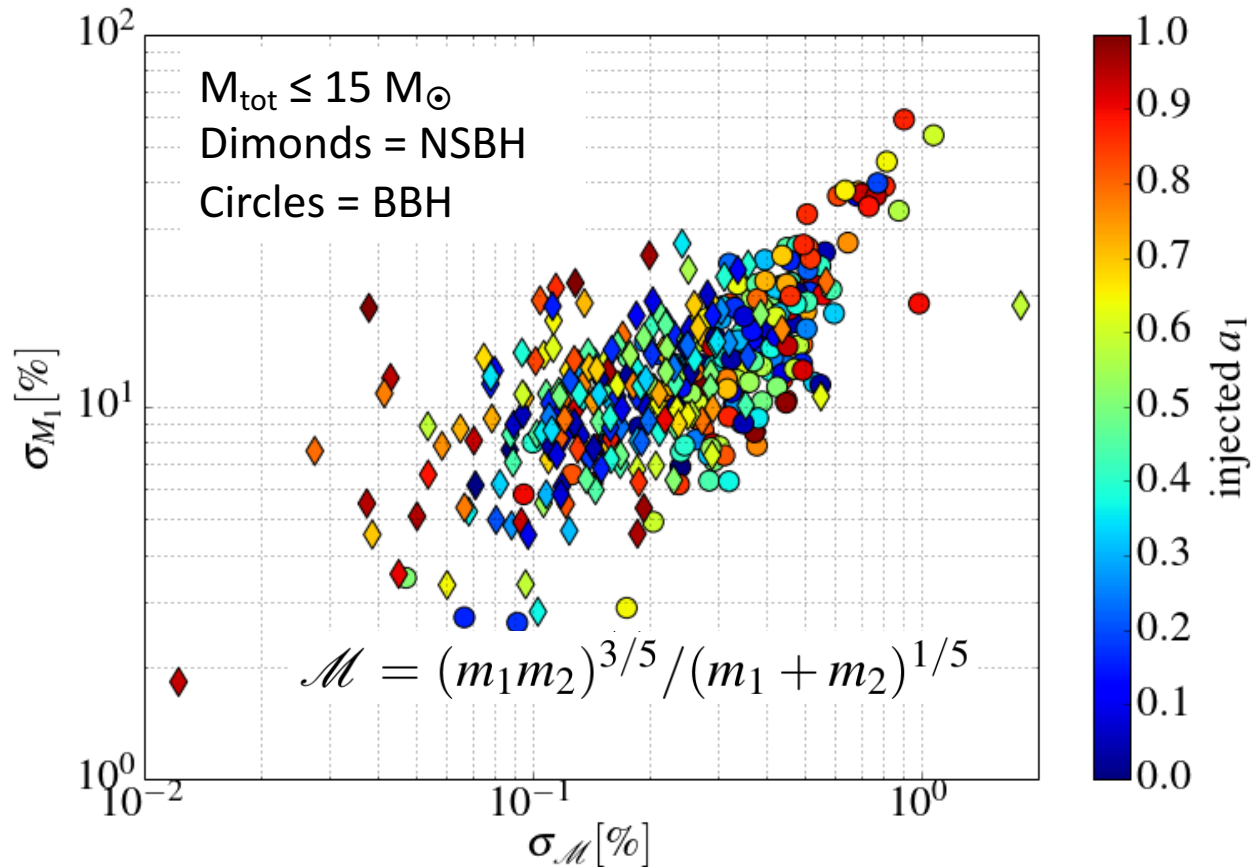


Comparison with EM

- Some of the BHs we discovered had masses significantly larger than what known from the EM
- High masses tell something about metallicity and winds of progenitor stars (LVC, ApjL 818 L22)



Astrophysical distribution - Masses



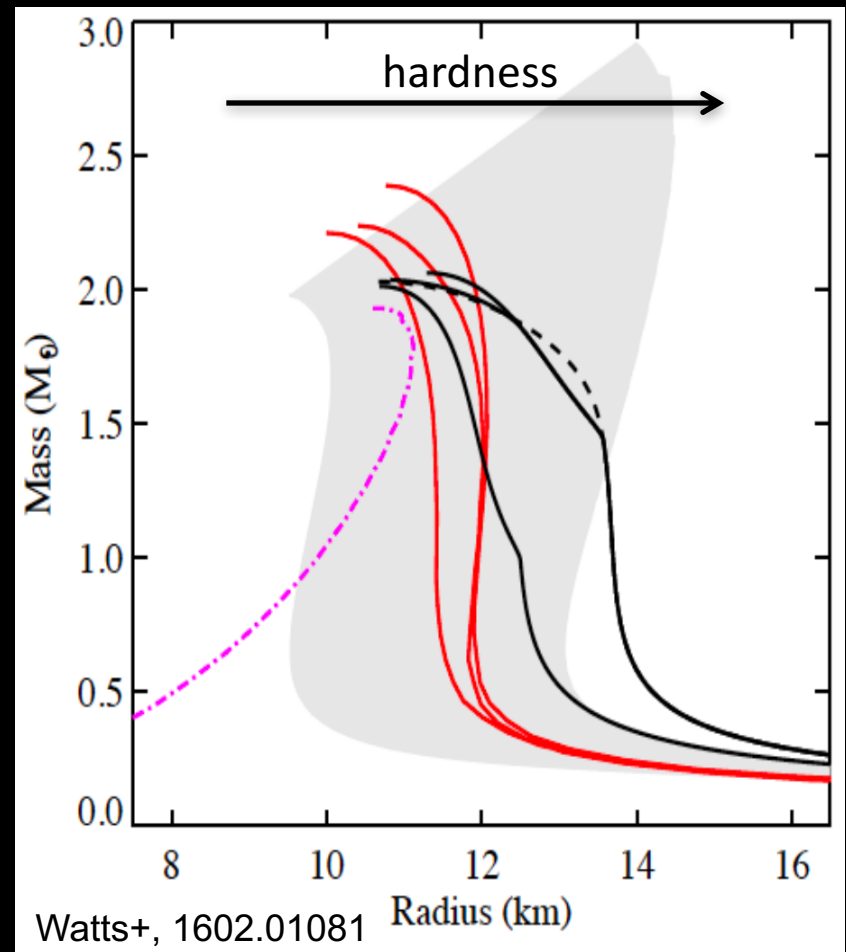
Vitale+ in prep., Vitale+ PRL 112 251101

Population inference

- As more GW will be detected, we will be able to infer the underlying mass distribution of neutron stars and black holes.
- E.g.: Are most NS masses at around $1.35 M_{\odot}$?
- Two main advantages over EM:
 - Direct measurement
 - Can potentially access many more systems

Neutron Star Equation of State

- Neutron stars host the highest densities in the (visible) universe
- Measuring their equation of state requires joint measurement of radius and mass
- Possible with EM, but challenging
 - Mass estimates not always reliable
 - Radius estimates non often available and not always reliable
 - NICER to launch 2016, 5% precision



Neutron Star Equation of State

- In a CBC, each NS will feel the tidal field of the companion, which induces a quadrupole moment

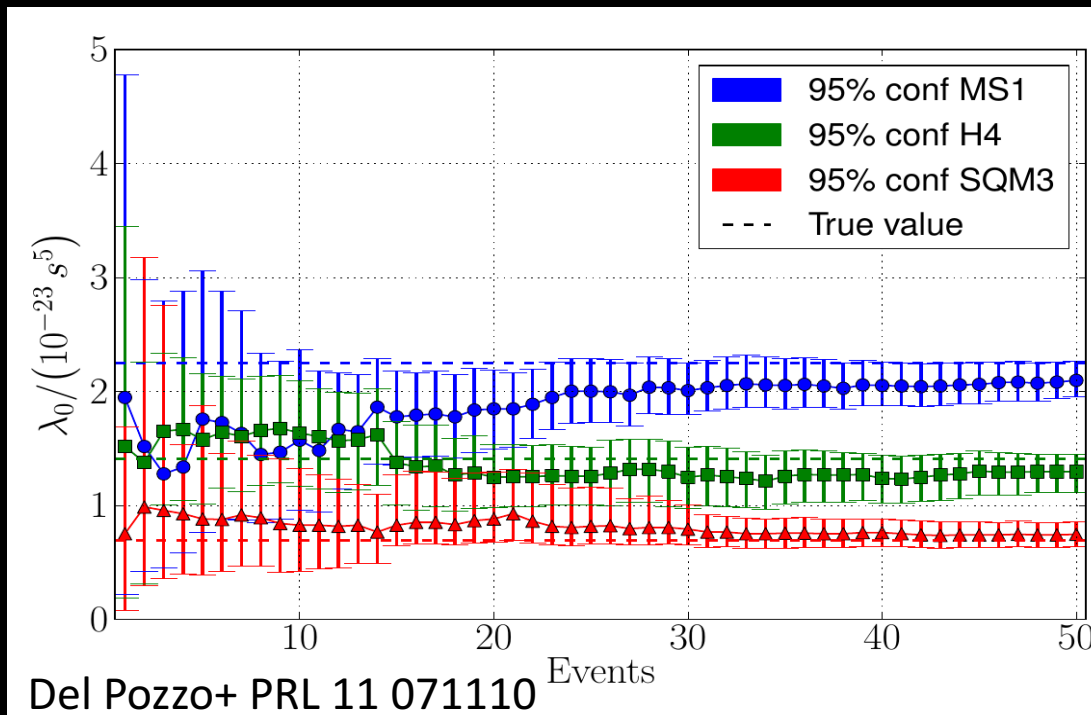
$$Q_{ij} = -\lambda(EOS, m) T_{ij}$$

- Known leading and next-to-leading effects on GW phasing
- Early studies considered single events, with contradictory findings (Read+ PRD 79 124033; Hinderer+ PRD 81 123016, many others)
- Markakis+ JPCS 189 012024 considered multiple events but used Fished matrix, unreliable at low signal-to-noise ratios (Vallisneri PRD 77 042001, Vitale+ PRD 84 104020)
- First fully Bayesian approach in 2013 (Del Pozzo+ PRL 11 071110)
- Also see Wade+, and many more.

Neutron Star Equation of State

- Two different avenues:
 - Model selection and EOS ranking
 - Parameter estimation on the tidal deformability
- Both allow to use all events
- Will focus on parameter estimation

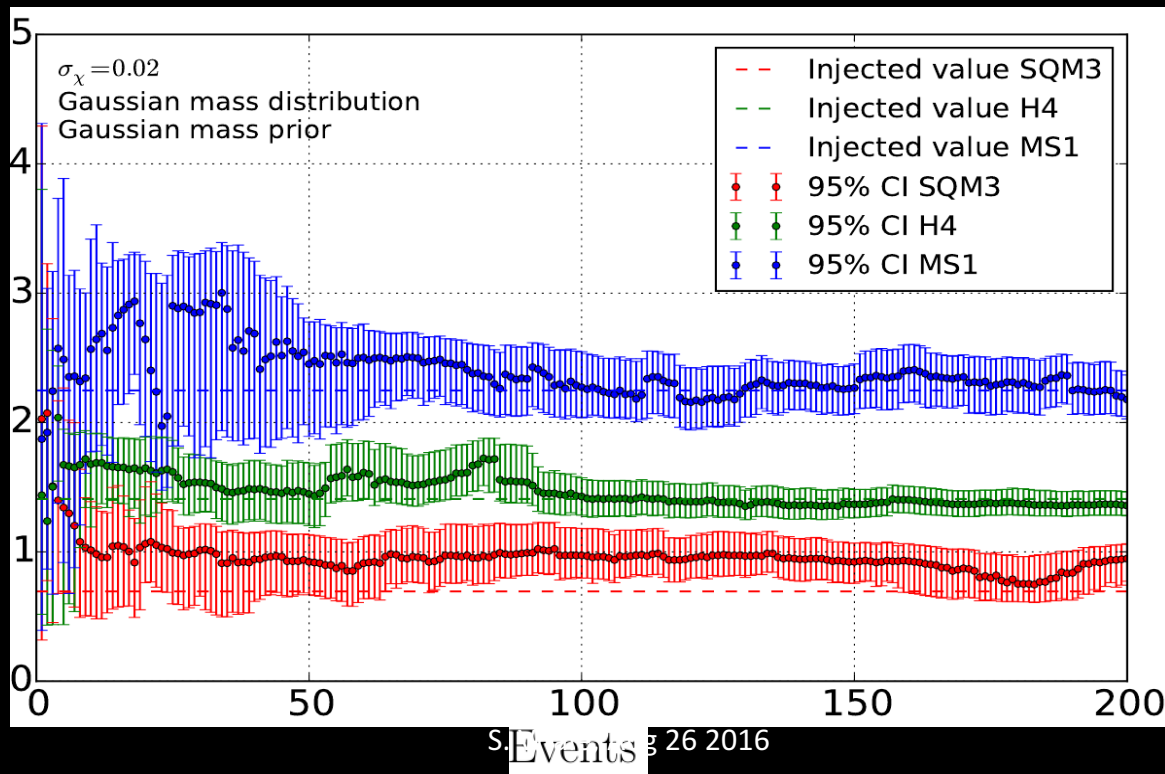
$$\lambda(m) \approx \lambda_0 + \lambda_1 (m - 1.4M)$$



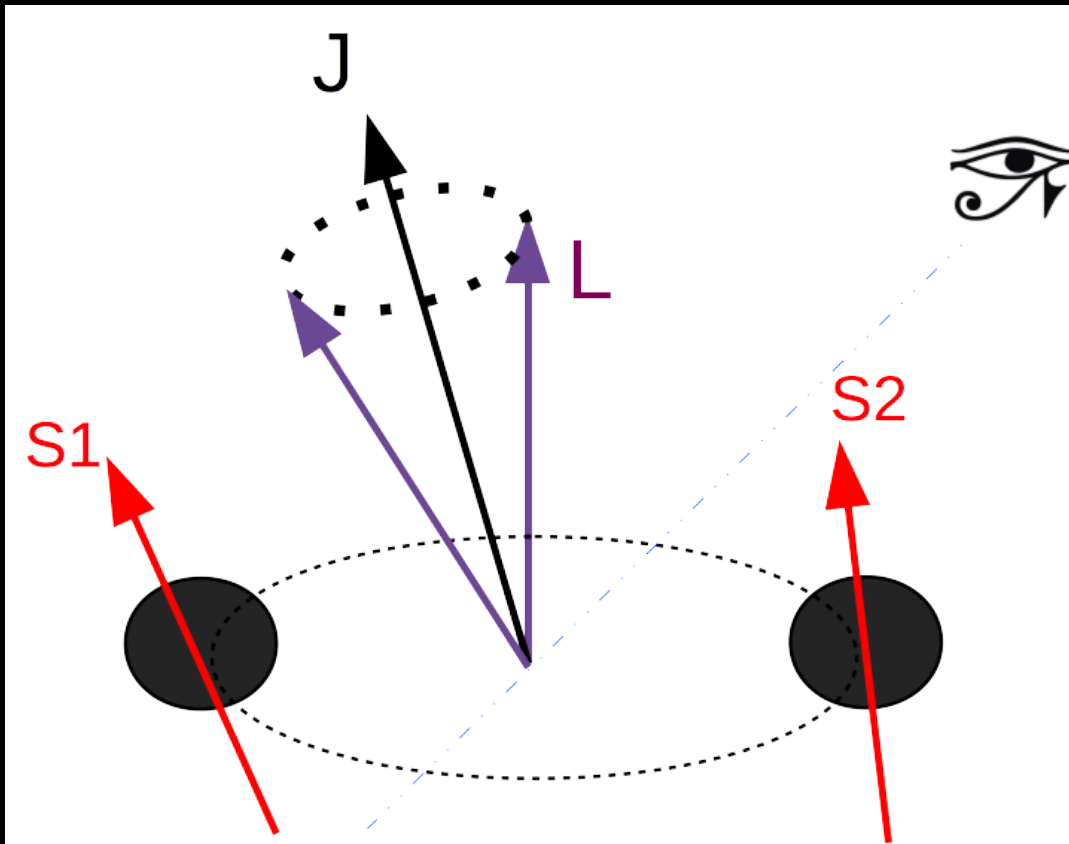
Neutron Star Equation of State

- In Agathos+ PRD 89 082001 we extended the initial study
 - Small neutron star spin
 - More physics (quadrupole-monopole term)
 - Better handling of waveform termination

$$\lambda_0 / (10^{-23} \text{ s}^5)$$

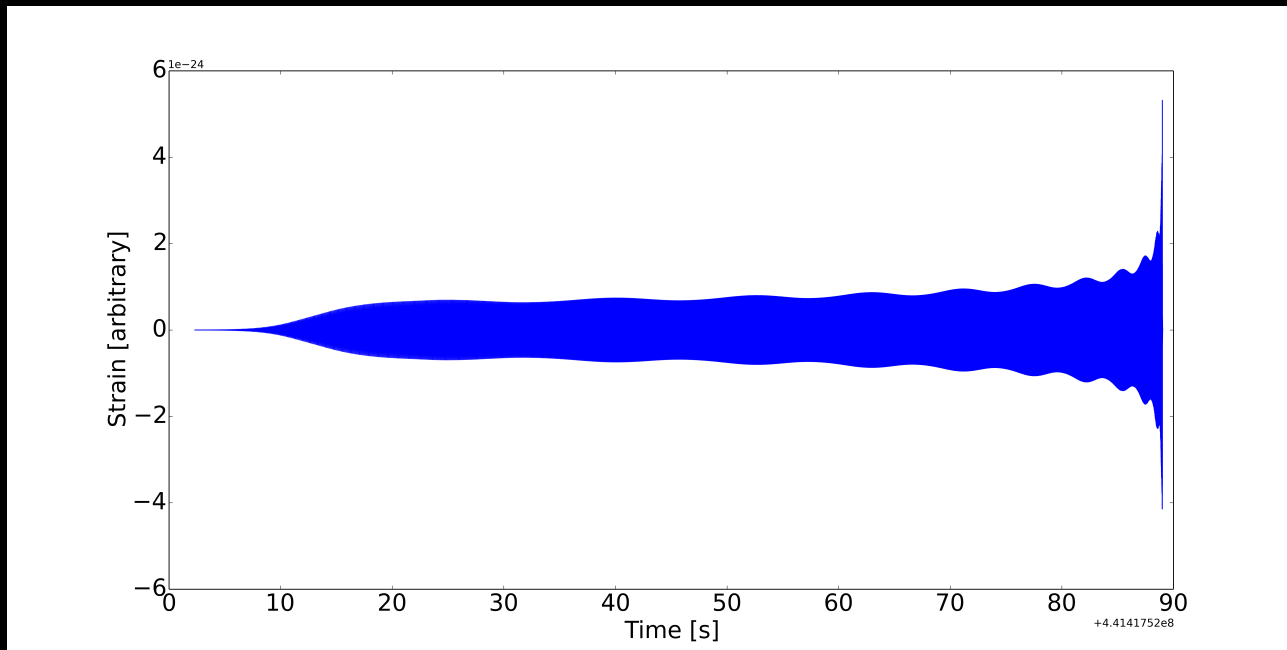


Spins



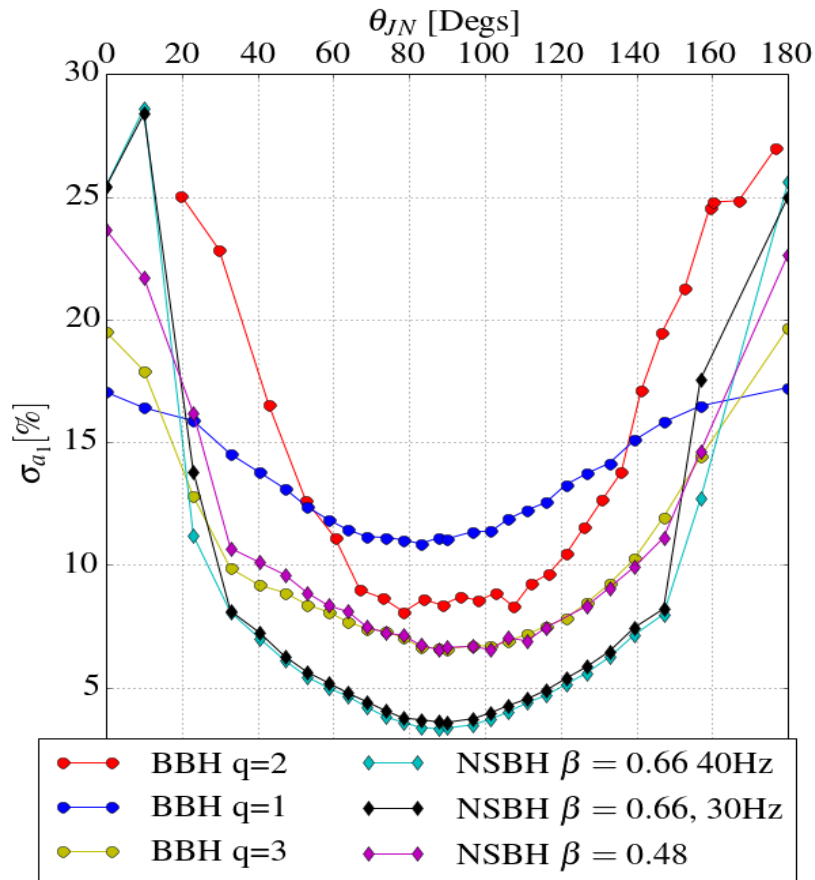
- Spin-Orbit coupling makes the orbit precess
- Waveform gets amplitude and phase modulation
- Rich physics, some degeneracies are reduced
 - Better parameter estimation

Spinning Waveforms w/ precession



- **However**, the amount of modulation visible in the *detector frame* also depends on the orientation
- Face-on CBC → Less modulation

Effects of Orbital Orientation



- Spin estimation strongly affected by orientation
- Spin errors minimum if system seen “edge-on”
- Less likely to detect than “face-on”

BH spin (with GW)

- Spins enter the waveform at higher PN orders
- They are harder to measure than mass parameters

LVC, 1606.04856

| | GW150914 | | | GW151226 | | | LVT151012 | | |
|--|--|--|----------------------------|--|--|----------------------------|--|--|--|
| | EOBNR | IMRPhenom | Overall | EOBNR | IMRPhenom | Overall | EOBNR | IMRPhenom | Overall |
| Detector frame | | | | | | | | | |
| Total mass M/M_{\odot} | 71.0 ^{+4.6} _{-4.0} | 71.2 ^{+3.5} _{-3.2} | 71.1 ^{+4.1±0.7} | 23.6 ^{+8.0} _{-1.3} | 23.8 ^{+5.1} _{-1.5} | 23.7 ^{+6.5±2.2} | 45 ⁺¹⁷ ₋₄ | 44 ⁺¹² ₋₃ | 44 ^{+16±5} _{-3±0} |
| Chirp mass \mathcal{M}/M_{\odot} | 30.4 ^{+2.3} _{-1.6} | 30.7 ^{+1.5} _{-1.5} | 30.6 ^{+1.9±0.3} | 9.71 ^{+0.08} _{-0.07} | 9.72 ^{+0.06} _{-0.06} | 9.72 ^{+0.07±0.01} | 18.1 ^{+1.3} _{-0.9} | 18.1 ^{+0.8} _{-0.8} | 18.1 ^{+1.0±0.5} _{-0.8±0.1} |
| Primary mass m_1/M_{\odot} | 40.2 ^{+5.2} _{-4.8} | 38.5 ^{+5.4} _{-3.3} | 39.4 ^{+5.4±1.3} | 15.3 ^{+10.8} _{-3.8} | 15.8 ^{+7.2} _{-4.0} | 15.6 ^{+9.0±2.6} | 29 ⁺²³ ₋₈ | 27 ⁺¹⁹ ₋₆ | 28 ^{+21±5} _{-7±0} |
| Secondary mass m_2/M_{\odot} | 30.6 ^{+5.1} _{-4.2} | 32.7 ^{+3.1} _{-4.9} | 31.7 ^{+4.0±0.1} | 8.3 ^{+2.5} _{-2.9} | 8.1 ^{+2.5} _{-2.1} | 8.2 ^{+2.6±0.2} | 15 ⁺⁵ ₋₆ | 16 ⁺⁴ ₋₆ | 16 ^{+5±0} _{-6±1} |
| Final mass M_f/M_{\odot} | 67.8 ^{+4.0} _{-3.6} | 67.9 ^{+3.2} _{-2.9} | 67.8 ^{+3.7±0.6} | 22.5 ^{+8.2} _{-1.4} | 22.8 ^{+5.3} _{-1.6} | 22.6 ^{+6.7±2.2} | 43 ⁺¹⁷ ₋₄ | 42 ⁺¹³ ₋₂ | 42 ^{+16±5} _{-3±0} |
| Source frame | | | | | | | | | |
| Total mass $M^{\text{source}}/M_{\odot}$ | 65.5 ^{+4.4} _{-3.9} | 65.1 ^{+3.6} _{-3.1} | 65.3 ^{+4.1±1.0} | 21.6 ^{+7.4} _{-1.6} | 21.9 ^{+4.7} _{-1.7} | 21.8 ^{+5.9±2.0} | 38 ⁺¹⁵ ₋₅ | 37 ⁺¹¹ ₋₄ | 37 ^{+13±4} _{-4±0} |
| Chirp mass $\mathcal{M}^{\text{source}}/M_{\odot}$ | 28.1 ^{+2.1} _{-1.6} | 28.1 ^{+1.6} _{-1.4} | 28.1 ^{+1.8±0.4} | 8.87 ^{+0.35} _{-0.28} | 8.90 ^{+0.31} _{-0.27} | 8.88 ^{+0.33±0.01} | 15.2 ^{+1.5} _{-1.1} | 15.0 ^{+1.3} _{-1.0} | 15.1 ^{+1.4±0.3} _{-1.1±0.0} |
| Primary mass $m_1^{\text{source}}/M_{\odot}$ | 37.0 ^{+4.9} _{-4.4} | 35.3 ^{+5.1} _{-3.1} | 36.2 ^{+5.2±1.4} | 14.0 ^{+10.0} _{-3.5} | 14.5 ^{+6.6} _{-3.7} | 14.2 ^{+8.3±2.4} | 24 ⁺¹⁹ ₋₇ | 23 ⁺¹⁶ ₋₅ | 23 ^{+18±5} _{-6±0} |
| Secondary mass $m_2^{\text{source}}/M_{\odot}$ | 28.3 ^{+4.6} _{-3.9} | 29.9 ^{+3.0} _{-4.5} | 29.1 ^{+3.7±0.0} | 7.5 ^{+2.3} _{-2.6} | 7.4 ^{+2.3} _{-2.0} | 7.5 ^{+2.3±0.2} | 13 ⁺⁴ ₋₅ | 14 ⁺⁴ ₋₅ | 13 ^{+4±0} _{-5±0} |
| Final mass $M_f^{\text{source}}/M_{\odot}$ | 62.5 ^{+3.9} _{-3.5} | 62.1 ^{+3.3} _{-2.8} | 62.3 ^{+3.7±0.9} | 20.6 ^{+7.6} _{-1.6} | 20.9 ^{+4.8} _{-1.8} | 20.8 ^{+6.1±2.0} | 36 ⁺¹⁵ ₋₄ | 35 ⁺¹¹ ₋₃ | 35 ^{+14±4} _{-4±0} |
| Energy radiated $E_{\text{rad}}/(M_{\odot}c^2)$ | 2.98 ^{+0.55} _{-0.40} | 3.02 ^{+0.36} _{-0.36} | 3.00 ^{+0.47±0.13} | 1.02 ^{+0.09} _{-0.24} | 0.99 ^{+0.11} _{-0.17} | 1.00 ^{+0.10±0.01} | 1.48 ^{+0.39} _{-0.41} | 1.51 ^{+0.29} _{-0.44} | 1.50 ^{+0.33±0.05} _{-0.43±0.01} |
| Mass ratio q | 0.77 ^{+0.20} _{-0.18} | 0.85 ^{+0.13} _{-0.21} | 0.81 ^{+0.17±0.02} | 0.54 ^{+0.40} _{-0.33} | 0.51 ^{+0.39} _{-0.25} | 0.52 ^{+0.40±0.03} | 0.53 ^{+0.42} _{-0.34} | 0.60 ^{+0.35} _{-0.37} | 0.57 ^{+0.38±0.01} _{-0.37±0.04} |
| Spin parameters | | | | | | | | | |
| Primary spin magnitude a_1 | 0.33 ^{+0.39} _{-0.29} | 0.30 ^{+0.54} _{-0.27} | 0.32 ^{+0.47±0.10} | 0.42 ^{+0.35} _{-0.37} | 0.55 ^{+0.35} _{-0.42} | 0.49 ^{+0.37±0.11} | 0.31 ^{+0.46} _{-0.27} | 0.31 ^{+0.50} _{-0.28} | 0.31 ^{+0.48±0.03} _{-0.28±0.09} |
| Secondary spin magnitude a_2 | 0.62 ^{+0.35} _{-0.54} | 0.36 ^{+0.53} _{-0.33} | 0.48 ^{+0.47±0.03} | 0.51 ^{+0.34} _{-0.46} | 0.52 ^{+0.42} _{-0.47} | 0.52 ^{+0.43±0.01} | 0.49 ^{+0.45} _{-0.44} | 0.42 ^{+0.50} _{-0.38} | 0.45 ^{+0.48±0.02} _{-0.41±0.01} |

BH spin (with GW)

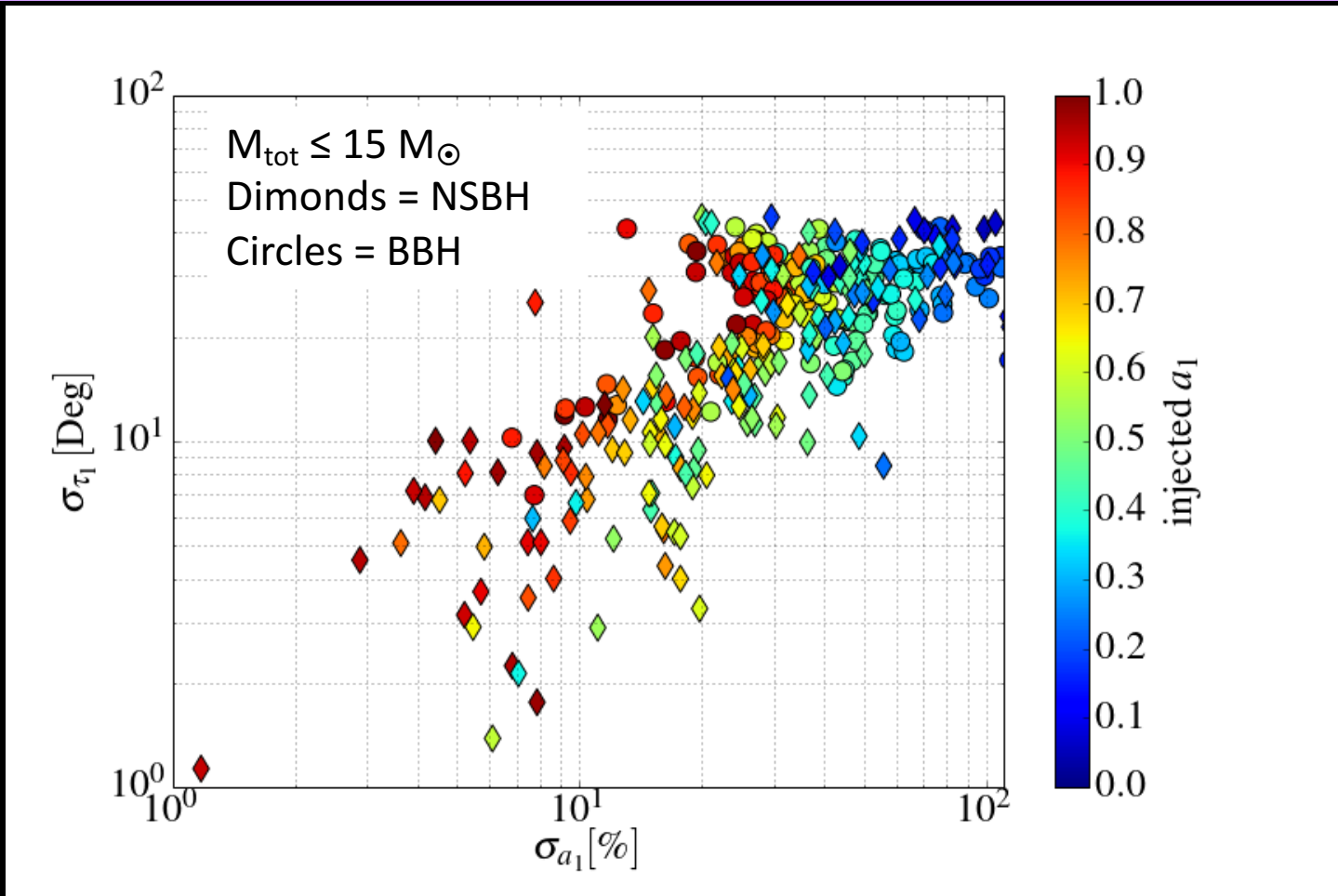
- Spins enter the waveform at higher PN orders
- They are harder to measure than mass parameters

LVC, 1606.04856

| | GW150914 | | | GW151226 | | | LVT151012 | | |
|--|------------------------|------------------------|--|------------------------|------------------------|--|------------------------|------------------------|--|
| | EOBNR | IMRPhenom | Overall | EOBNR | IMRPhenom | Overall | EOBNR | IMRPhenom | Overall |
| Detector frame | | | | | | | | | |
| Total mass M/M_{\odot} | $71.0^{+4.6}_{-4.0}$ | $71.2^{+3.5}_{-3.2}$ | $71.1^{+4.1\pm 0.7}_{-3.6\pm 0.8}$ | $23.6^{+8.0}_{-1.3}$ | $23.8^{+5.1}_{-1.5}$ | $23.7^{+6.5\pm 2.2}_{-1.4\pm 0.1}$ | 45^{+17}_{-4} | 44^{+12}_{-3} | $44^{+16\pm 5}_{-3\pm 0}$ |
| Chirp mass \mathcal{M}/M_{\odot} | $30.4^{+2.3}_{-1.6}$ | $30.7^{+1.5}_{-1.5}$ | $30.6^{+1.9\pm 0.3}_{-1.6\pm 0.4}$ | $9.71^{+0.08}_{-0.07}$ | $9.72^{+0.06}_{-0.06}$ | $9.72^{+0.07\pm 0.01}_{-0.06\pm 0.01}$ | $18.1^{+1.3}_{-0.9}$ | $18.1^{+0.8}_{-0.8}$ | $18.1^{+1.0\pm 0.5}_{-0.8\pm 0.1}$ |
| Primary mass m_1/M_{\odot} | $40.2^{+5.2}_{-4.8}$ | $38.5^{+5.4}_{-3.3}$ | $39.4^{+5.4\pm 1.3}_{-4.1\pm 0.2}$ | $15.3^{+10.8}_{-3.8}$ | $15.8^{+7.2}_{-4.0}$ | $15.6^{+9.0\pm 2.6}_{-4.0\pm 0.2}$ | 29^{+23}_{-8} | 27^{+19}_{-6} | $28^{+21\pm 5}_{-7\pm 0}$ |
| Secondary mass m_2/M_{\odot} | $3^{+5.1}_{-1.5}$ | $3^{+3.1}_{-1.1}$ | $3^{+4.0\pm 0.1}_{-1.0\pm 0.1}$ | $3^{+3.5}_{-1.2}$ | $3^{+2.5}_{-1.2}$ | $3^{+2.6\pm 0.2}_{-1.0\pm 0.1}$ | 16^{+5}_{-5} | 16^{+4}_{-6} | $16^{+5\pm 0}_{-6\pm 1}$ |
| Final mass M_f/M_{\odot} | $6^{+5.1}_{-1.5}$ | $6^{+3.1}_{-1.1}$ | $6^{+4.0\pm 0.1}_{-1.0\pm 0.1}$ | $6^{+3.5}_{-1.2}$ | $6^{+2.5}_{-1.2}$ | $6^{+2.6\pm 0.2}_{-1.0\pm 0.1}$ | 42^{+15}_{-5} | 42^{+13}_{-2} | $42^{+16\pm 5}_{-3\pm 0}$ |
| Source frame | | | | | | | | | |
| Total mass $M^{\text{source}}/M_{\odot}$ | $6^{+5.1}_{-1.5}$ | $6^{+3.1}_{-1.1}$ | $6^{+4.0\pm 0.1}_{-1.0\pm 0.1}$ | $6^{+3.5}_{-1.2}$ | $6^{+2.5}_{-1.2}$ | $6^{+2.6\pm 0.2}_{-1.0\pm 0.1}$ | 37^{+11}_{-4} | 37^{+11}_{-4} | $37^{+13\pm 4}_{-4\pm 0}$ |
| Chirp mass $\mathcal{M}^{\text{source}}/M_{\odot}$ | $2^{+1.3}_{-1.0}$ | $2^{+1.3}_{-1.0}$ | $2^{+1.3\pm 0.0}_{-1.1\pm 0.0}$ | $15.0^{+1.3}_{-1.0}$ | $15.1^{+1.3}_{-1.0}$ | $15.1^{+1.4\pm 0.3}_{-1.1\pm 0.0}$ | 23^{+16}_{-5} | 23^{+16}_{-5} | $23^{+18\pm 5}_{-6\pm 0}$ |
| Primary mass $m_1^{\text{source}}/M_{\odot}$ | 14^{+4}_{-5} | 14^{+4}_{-5} | $14^{+4\pm 0}_{-5\pm 0}$ | 35^{+11}_{-3} | 35^{+11}_{-3} | $35^{+14\pm 4}_{-4\pm 0}$ | $1.48^{+0.39}_{-0.41}$ | $1.51^{+0.29}_{-0.44}$ | $1.50^{+0.33\pm 0.05}_{-0.43\pm 0.01}$ |
| Secondary mass $m_2^{\text{source}}/M_{\odot}$ | $0.62^{+0.35}_{-0.35}$ | $0.36^{+0.53}_{-0.33}$ | $0.48^{+0.47\pm 0.03}_{-0.43\pm 0.03}$ | $0.51^{+0.42}_{-0.46}$ | $0.52^{+0.42}_{-0.47}$ | $0.52^{+0.43\pm 0.01}_{-0.47\pm 0.00}$ | $0.53^{+0.42}_{-0.34}$ | $0.60^{+0.35}_{-0.37}$ | $0.57^{+0.38\pm 0.01}_{-0.37\pm 0.04}$ |
| Final mass $M_f^{\text{source}}/M_{\odot}$ | $2.98^{+0.55}_{-0.40}$ | $3.02^{+0.36}_{-0.36}$ | $3.00^{+0.47\pm 0.13}_{-0.39\pm 0.07}$ | $1.02^{+0.09}_{-0.24}$ | $0.99^{+0.11}_{-0.17}$ | $1.00^{+0.10\pm 0.01}_{-0.20\pm 0.03}$ | $1.48^{+0.39}_{-0.41}$ | $1.51^{+0.29}_{-0.44}$ | $1.50^{+0.33\pm 0.05}_{-0.43\pm 0.01}$ |
| Energy radiated $E_{\text{rad}}/(M_{\odot}c^2)$ | $0.77^{+0.20}_{-0.18}$ | $0.85^{+0.13}_{-0.21}$ | $0.81^{+0.17\pm 0.02}_{-0.20\pm 0.04}$ | $0.54^{+0.40}_{-0.33}$ | $0.51^{+0.39}_{-0.25}$ | $0.52^{+0.40\pm 0.03}_{-0.29\pm 0.04}$ | $0.53^{+0.42}_{-0.34}$ | $0.60^{+0.35}_{-0.37}$ | $0.57^{+0.38\pm 0.01}_{-0.37\pm 0.04}$ |
| Mass ratio q | $0.02^{+0.17}_{-0.14}$ | $0.02^{+0.11}_{-0.12}$ | $0.02^{+0.14\pm 0.02}_{-0.14\pm 0.04}$ | $0.02^{+0.24}_{-0.11}$ | $0.02^{+0.15}_{-0.08}$ | $0.02^{+0.20\pm 0.07}_{-0.10\pm 0.03}$ | $0.02^{+0.31}_{-0.24}$ | $0.02^{+0.26}_{-0.17}$ | $0.02^{+0.31\pm 0.08}_{-0.20\pm 0.02}$ |
| Primary spin magnitude a_1 | $0.33^{+0.39}_{-0.29}$ | $0.30^{+0.54}_{-0.27}$ | $0.32^{+0.47\pm 0.10}_{-0.29\pm 0.01}$ | $0.42^{+0.35}_{-0.37}$ | $0.55^{+0.35}_{-0.42}$ | $0.49^{+0.37\pm 0.11}_{-0.42\pm 0.07}$ | $0.31^{+0.46}_{-0.27}$ | $0.31^{+0.46}_{-0.28}$ | $0.31^{+0.48\pm 0.03}_{-0.28\pm 0.09}$ |
| Secondary spin magnitude a_2 | $0.62^{+0.35}_{-0.35}$ | $0.36^{+0.53}_{-0.33}$ | $0.48^{+0.47\pm 0.03}_{-0.43\pm 0.03}$ | $0.51^{+0.42}_{-0.46}$ | $0.52^{+0.42}_{-0.47}$ | $0.52^{+0.43\pm 0.01}_{-0.47\pm 0.00}$ | $0.49^{+0.45}_{-0.44}$ | $0.42^{+0.50}_{-0.38}$ | $0.45^{+0.48\pm 0.02}_{-0.41\pm 0.01}$ |

**RELATIVE SPIN ERRORS
CLOSE TO 100%**

Astrophysical distribution - Spin



Vitale+ in prep., Vitale+ PRL 112 251101

CBC and their formation channels

- Measuring masses and spins can help determine channel and environment in which BH and CBC are formed
- Two main formation channels
 - **Common envelope evolution**
 - Galactic fields
 - Final masses not too different
 - Aligned spins
 - **Dynamical capture**
 - Globular clusters
 - Any mass ratio (?)
 - Misaligned spins

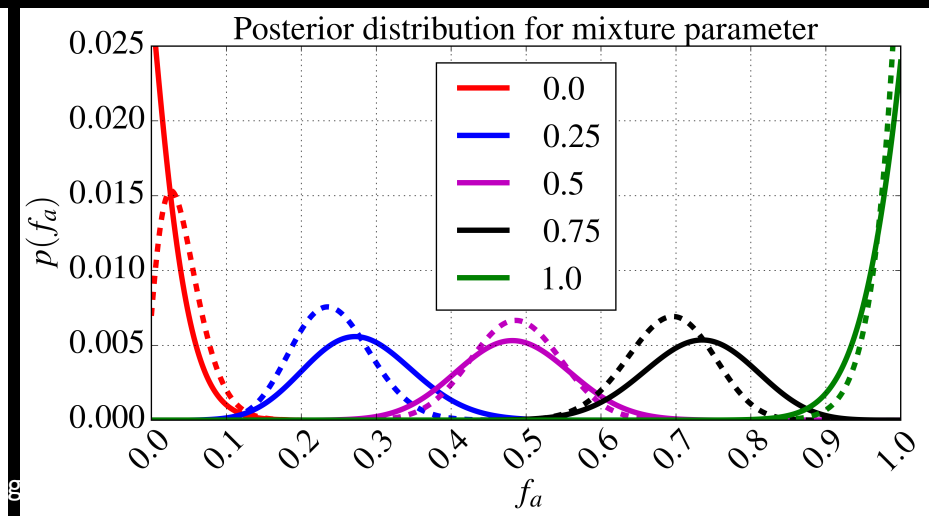
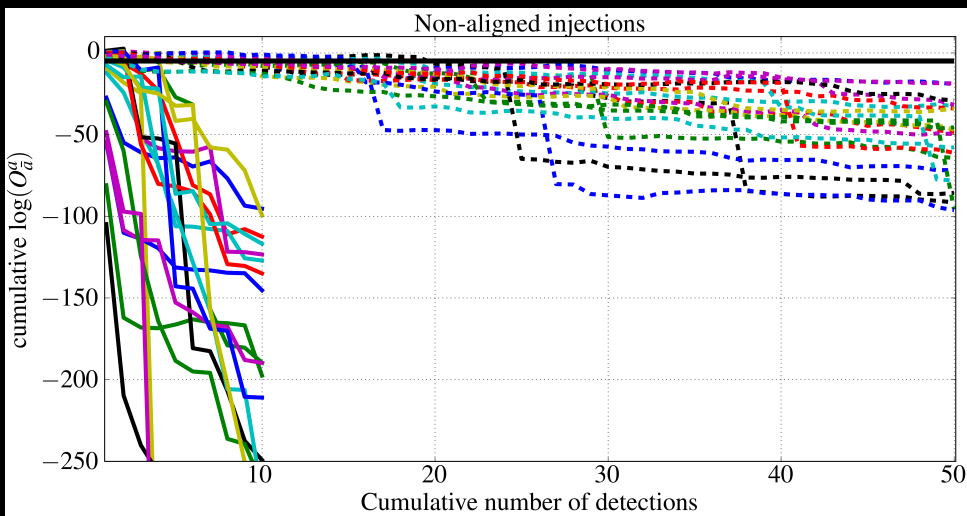
An example: spin alignment

- Most astronomers believe that CBC formed via common envelope will have aligned spins (Gerosa+, many others)
- We can use Bayesian methods to verify if and how many systems have aligned spins

-
- If time allows, show how to obtain the posterior distribution of f_a

Results

- 100 NSBH (dashed) and 200 BBH
- Astrophysical distribution
- Can measure the fraction of aligned systems with uncertainties of $\sim 15\text{-}20\%$



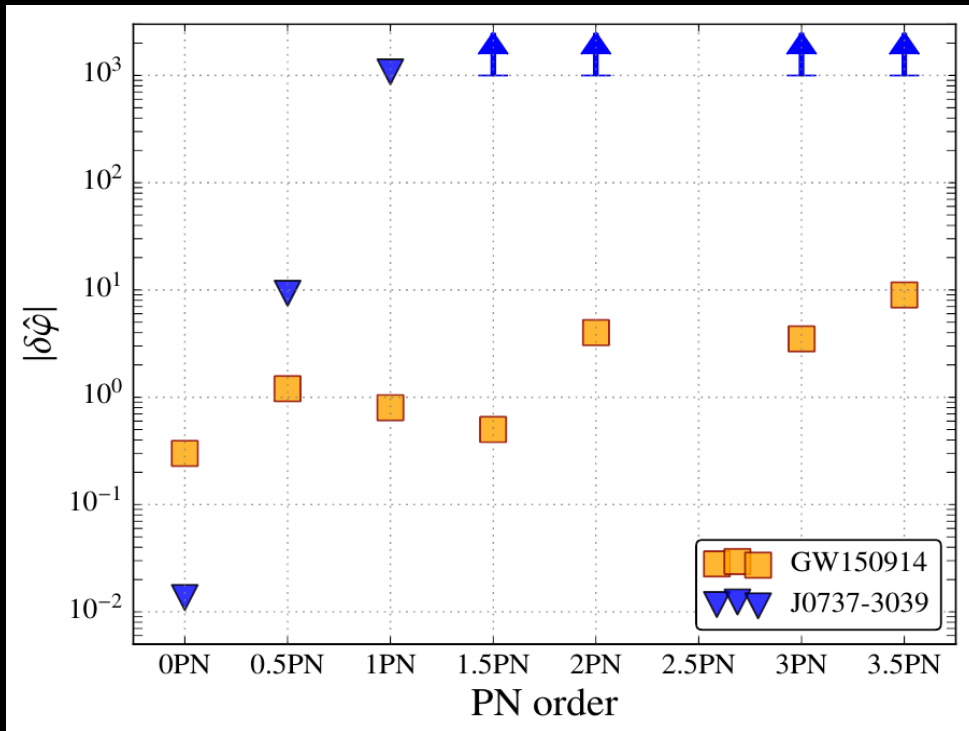
Strong field tests of General Relativity

- GWs represent our first chance to test general relativity in its strong-field dynamical regime
- Double pulsar J0737-3039 has
 - Masses $\sim M_{\odot}$
 - Speeds $\sim 1e-3 c$
 - Derivative orbital period $\sim 1e-12$
- GW150914
 - Few to tens of M_{\odot}
 - Relative speed $\sim 0.5 c$
 - Derivative orbital period ~ 1

Strong field tests of General Relativity

- There is a large number of alternative theories of gravity that we will be able to test using gravitational waves
 - Massive graviton
 - Brans-Dicke
 - Many more!
- But the real theory of gravity might not have been proposed yet. Need unmodeled tests
 - Post-Newtonian phasing tests
 - TIGER

GR tests with GW150914: phasing



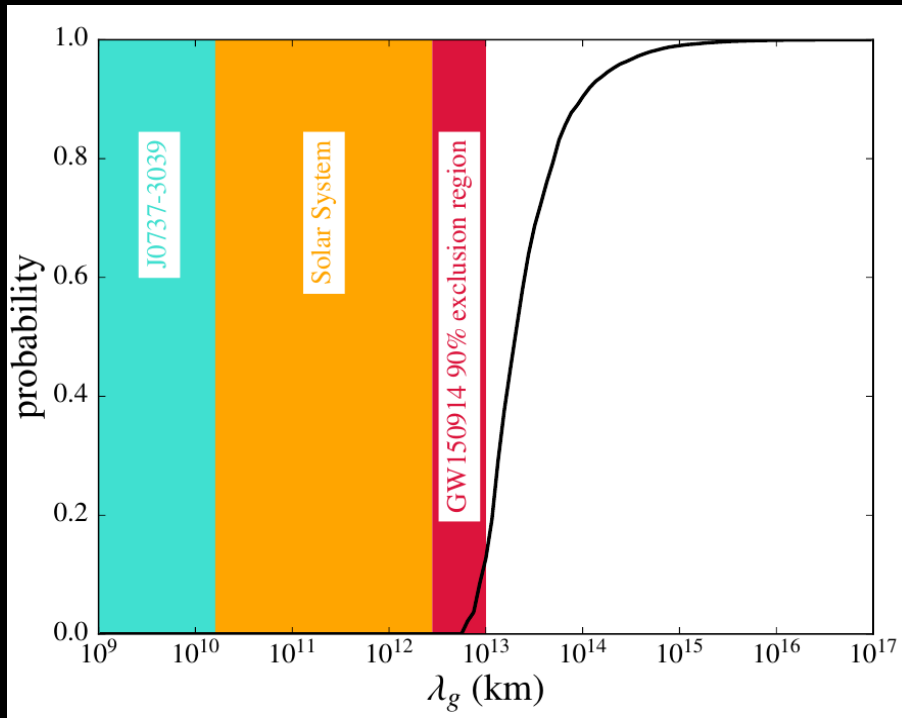
LVC PRL 116, 221101 & 1606.04856

- Post Newtonian expansion of GW phase has coefficients known within GR

$$\psi(f) \sim \sum \varphi_k (\pi M f)^{\frac{k-5}{3}}$$

- Can put upper limit on how different from GR they can be
 - Using double pulsar
 - Using GW150914
- Double pulsar already beaten at 0.5PN

GR tests with GW150914: massive graviton



- A full self-consistent theory of gravitational field mediated by massive particle is not yet available
- However, just modifying the dispersion relation one can calculate extra phasing term for GW phase (Will, PRD 57 2061)

- ✓ 3 orders of magnitude better than double pulsar
- ✓ Factor of 3 better than solar system
- ✓ Some model dependent tests do better

LVC PRL 116, 221101

$$m_g \leq 1.2 \times 10^{-22} \text{ eV}/c^2$$

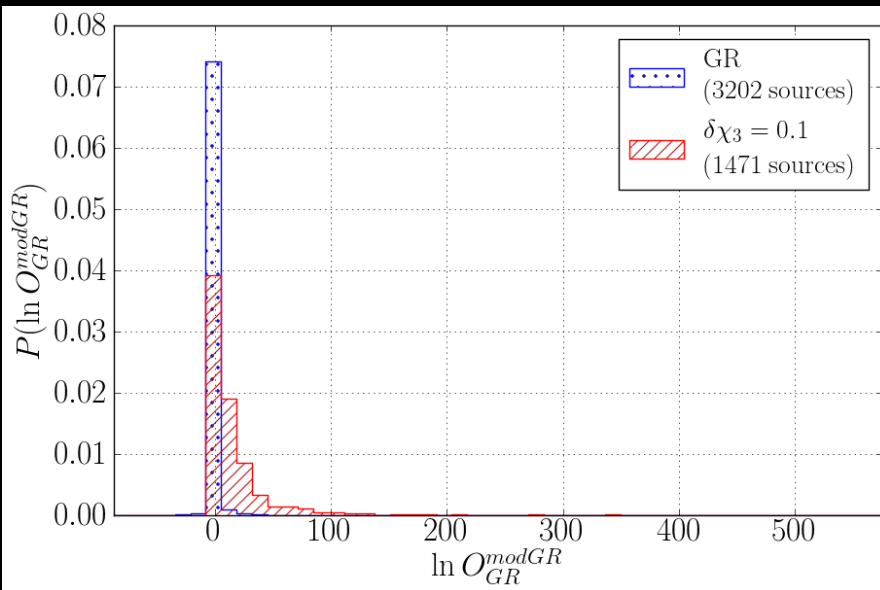
TIGER

- Test Infrastructure for General Relativity
- Look for generic deviations in an unmodeled fashion
- Based on Bayesian model selection, combine evidence from all detected events
- Does not rely on “golden events” (i.e.: high signal-to-noise ratio)
- Extensively tested on BNS (Li+ PRD 85 082003 , Agathos+ PRD 89 082001)
- In development for BBH

TIGER

- Test Infrastructure for General Relativity
- Look for generic deviations in an unmodeled fashion
- Based on Bayesian model selection, combine evidence from all detected events
- Does not rely on “golden events” (i.e.: high signal-to-noise ratio)
- Extensively tested on BNS (Li+ PRD 85 082003 , Agathos+ PRD 89 082001)
- In development for BBH
- Define two models “GR is correct” vs “GR is not correct”
 - “GR is not correct” is true if any post Newtonian or phenomenological phase parameters deviates from its GR value
- Calculate Bayesian odds (ratio of probabilities)
- Each detection contributes to odds
- Good efficiency after few tens of BNS

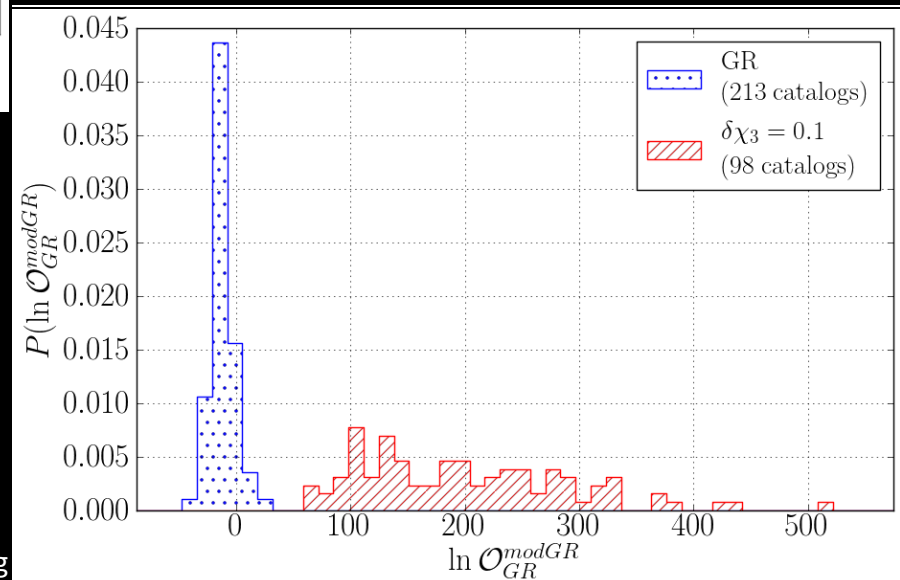
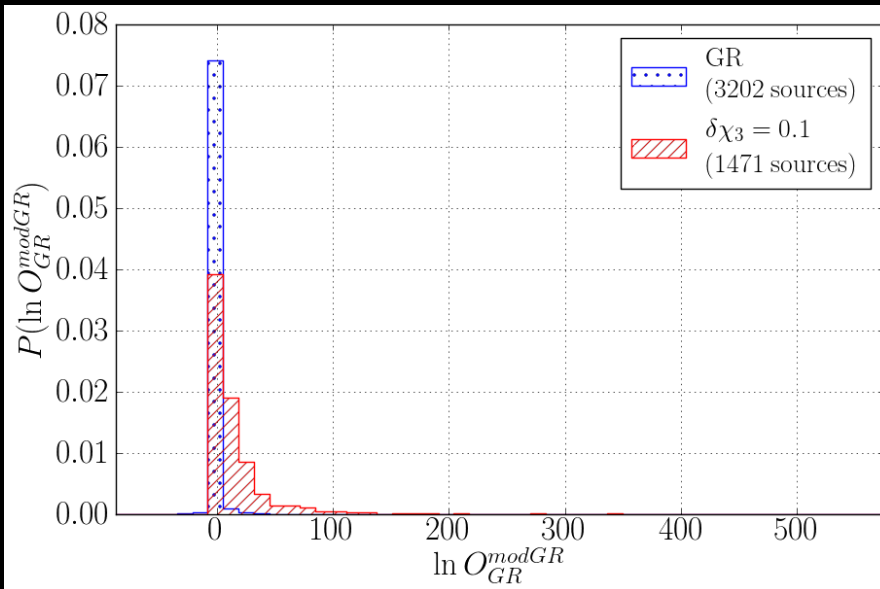
TIGER BNS example



Singe BNS source

TIGER BNS example

Singe BNS source



Catalogs of BNS sources

Cosmography with GWs

- Gravitational waves provide direct measurement of *luminosity distance*
- **If** the *redshift* of the source can be estimated in some other way one can measure cosmological parameters.

$$D_L(z) = \begin{cases} \frac{(1+z)}{\sqrt{\Omega_k}} \sinh\left[\sqrt{\Omega_k} \int_0^z \frac{dz'}{H(z')}\right] & \text{for } \Omega_k > 0 \\ (1+z) \int_0^z \frac{dz'}{H(z')} & \text{for } \Omega_k = 0 \\ \frac{(1+z)}{\sqrt{|\Omega_k|}} \sin\left[\sqrt{|\Omega_k|} \int_0^z \frac{dz'}{H(z')}\right] & \text{for } \Omega_k < 0 \end{cases}$$

$$H(z) = H_0 \sqrt{\Omega_m (1+z)^3 + \Omega_k (1+z)^2 + \Omega_\Lambda E(z, w(z))}$$

How to measure the redshift?

- If the CBC produces an EM counterpart (e.g. GRB) (Schutz, Nature 1986)
- If one knows the NS EOS (Read & Messenger PRL 108 091101, Del Pozzo+ 1506.0659)
- If the post-merger signal is observed (Messenger+ PRX 4, 041004)
- If the true distribution of NS masses is known (Taylor & Gair PRD 86, 023502)
- Even if no EM is found, but there is a reliable galaxy catalog (Del Pozzo PRD 86 043011)

Second vs third generation of GW detectors

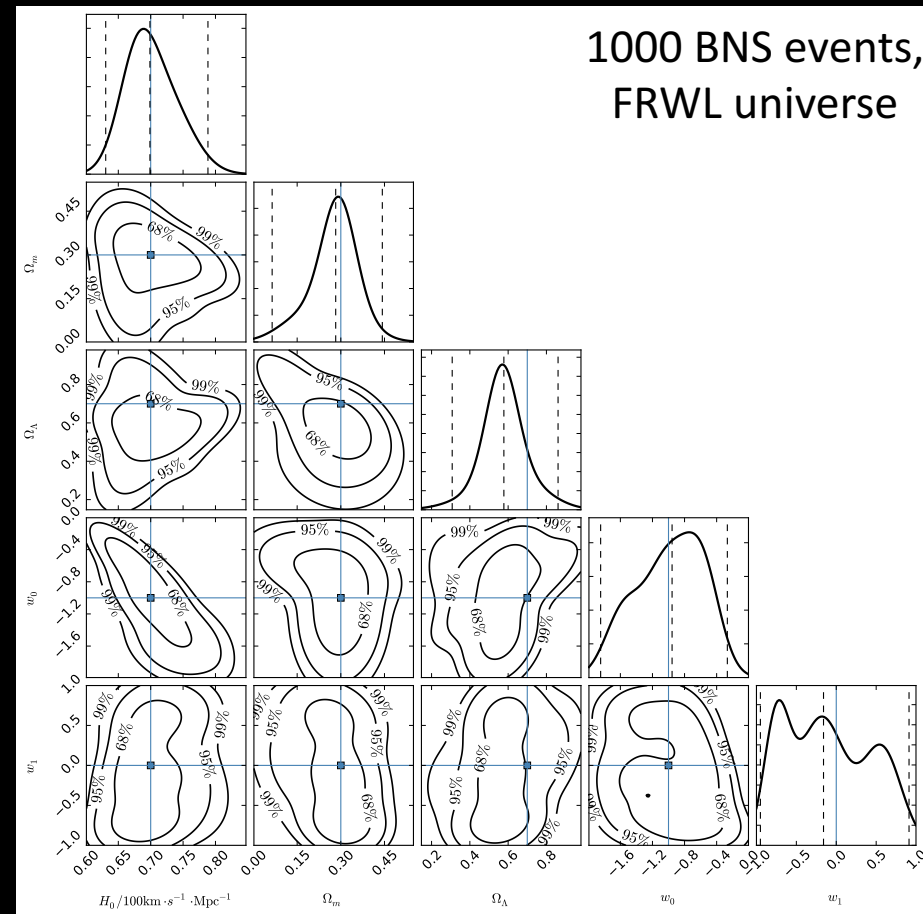
- Measuring cosmological parameters might have to wait the for the next generation of GW detectors, when thousands of detections will be made each year.
- Having thousands of events increases the chance of EM counterpart
- ... and allows for populations studies

An example: cosmology with EOS

- Del Pozzo+ (1506.0659) explores how well cosmological parameters can be estimated if the NS EOS is known.
- Thousands of BNS assuming different cosmologies
- Third generation GW detectors (Einstein telescope)

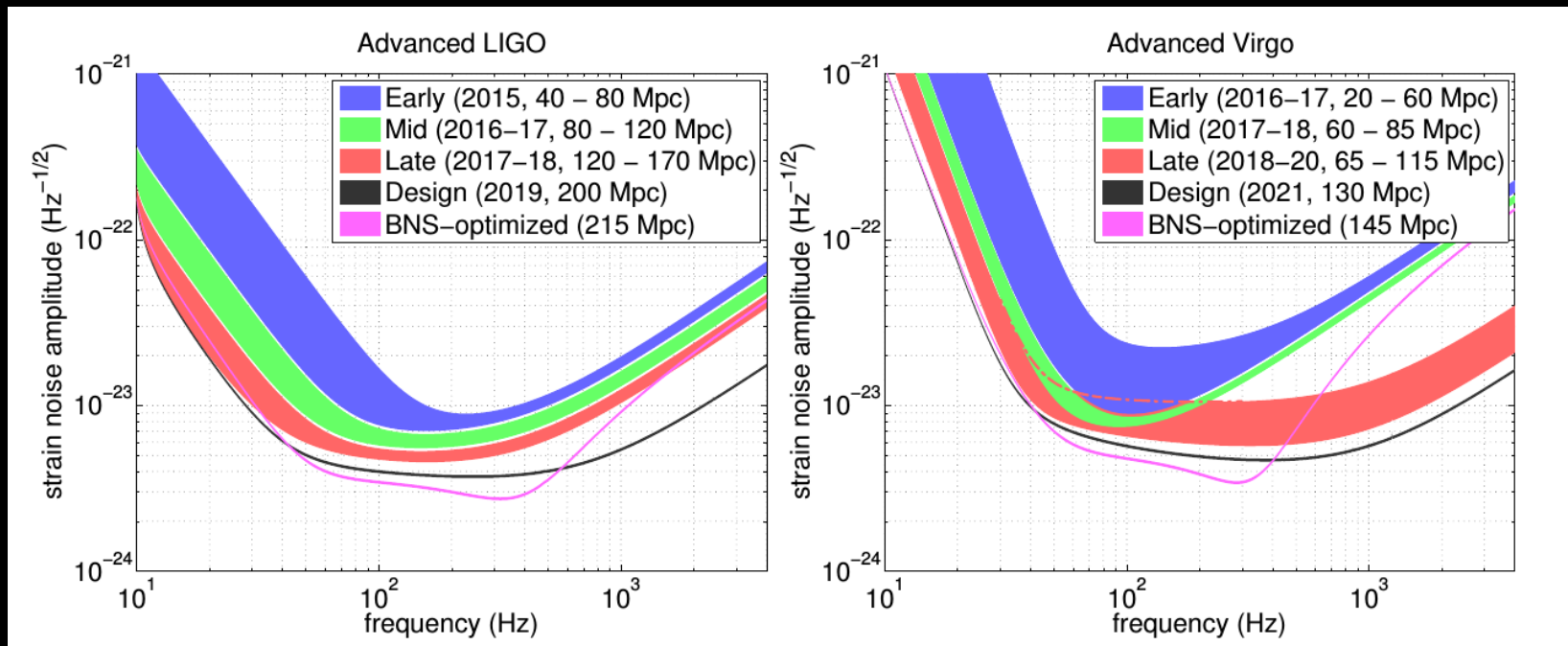
An example: cosmology with EOS

- Hubble constant measurable with precision of $\sim 10\%$ with 1000 sources
- Λ_m and Λ_k factor of >10 worse than Planck
 - These become comparable to Planck with more sources
- DE parameters (w 's) precision comparable to Planck with 1000 sources
 - Can improve EM measurement with more detections



1506.0659

Advanced LIGO not at design sensitivity yet: more detections coming soon!



- Next Science Runs:
 - 2015: 2 BBH event
 - 2016 - 2017: 5-10 significant BBH events, ~2 BNS
 - 2017-2018 : 10-30 significant BBH events, ~10 BNS
 - 2019: > 50 significant BBH events, ~20 BNS

LVC Liv. Rev. Rel 19, 1

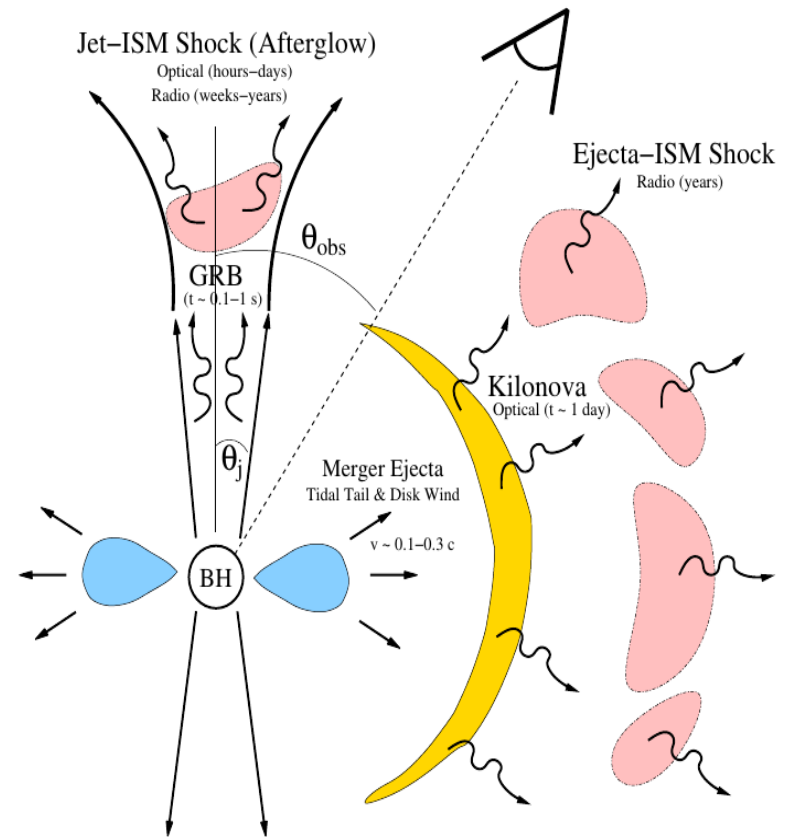


Thank you!

GW – Electromagnetic connection

- CBC containing neutron stars are expected to be bright in the EM band
 - Progenitors of short GRBs?
 - Rich variety of frequencies and timescales
- Core collapse supernovae are believed to power long GRBs
- And are BBH luminous??

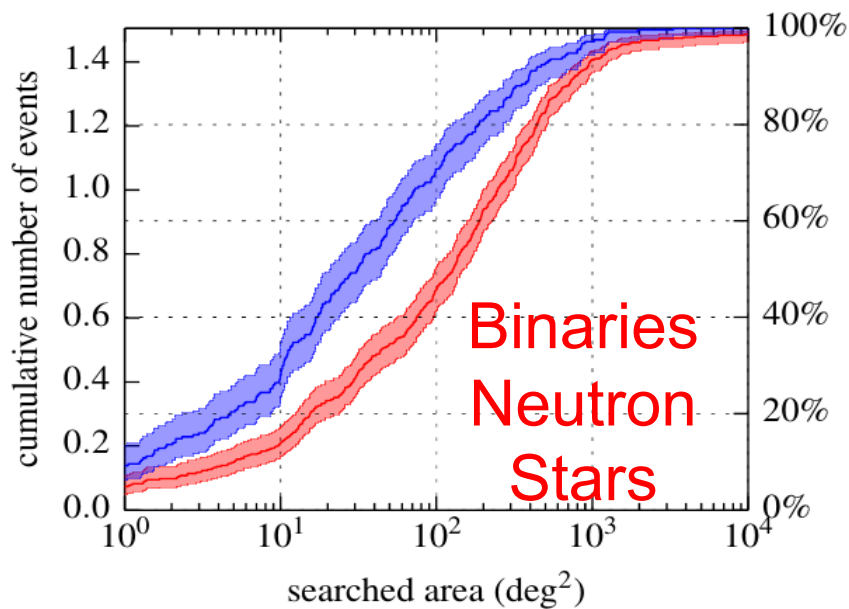
Metzger, Berger



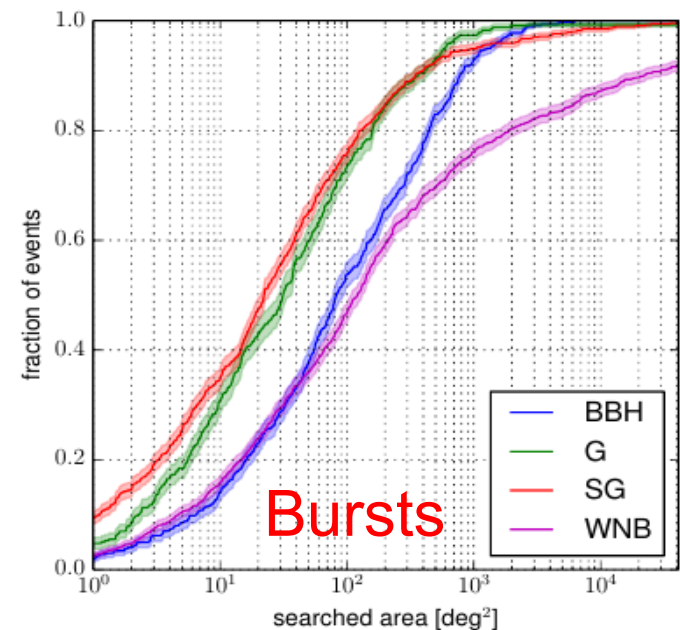
GW – Electromagnetic connection

- Challenges: searches need to be fast and precise; typical latencies few minutes
- Sky error areas from GW data are large, especially with only 2-3 detectors on-line
- Median searched areas $\leq 100 \text{ deg}^2$ in the next science run with Virgo on-line
- Not strongly dependent on model

Singer+ ApJ 795 105

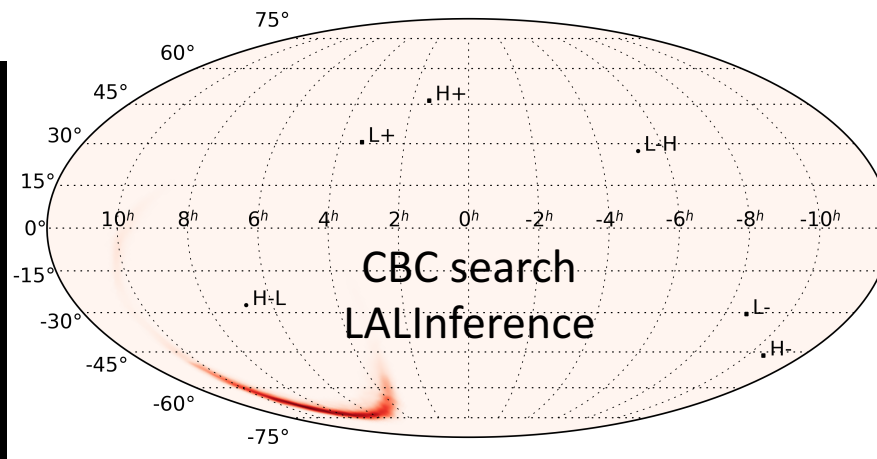
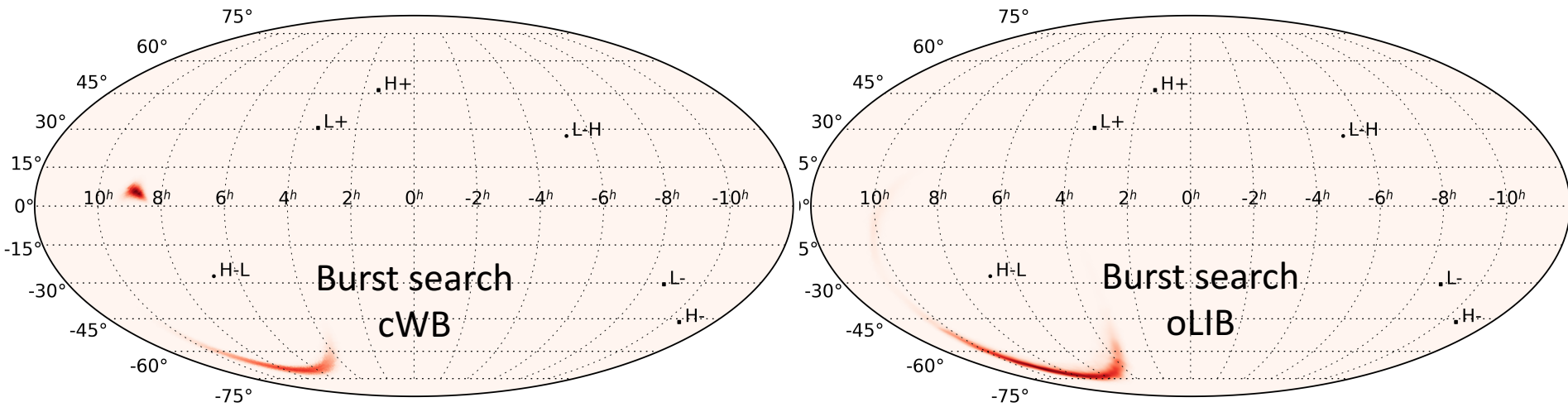


Essick+ ApJ 800 81



GW150914-Sky Localization

Sky maps generated by burst searches and shared with EM partners within 48 hours



LVC, 1602.03843

CBC targeted parameter estimation (LALInference, Veitch+ PRD 91 042003) ran with higher latency

Multi-band

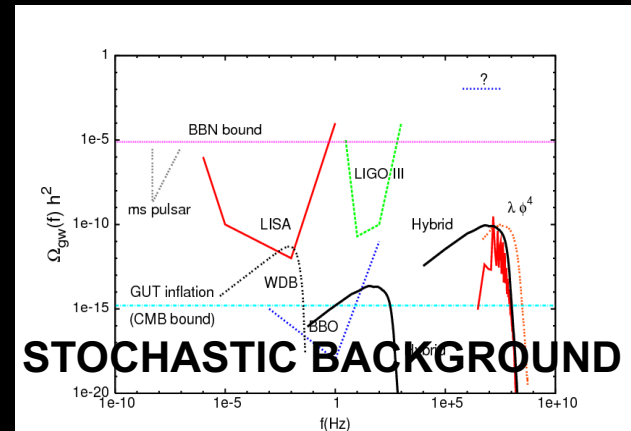
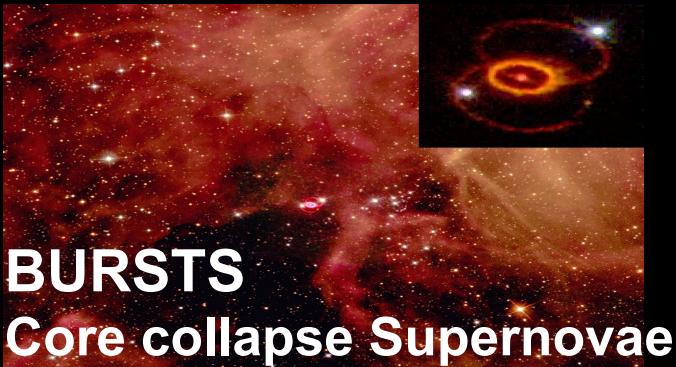
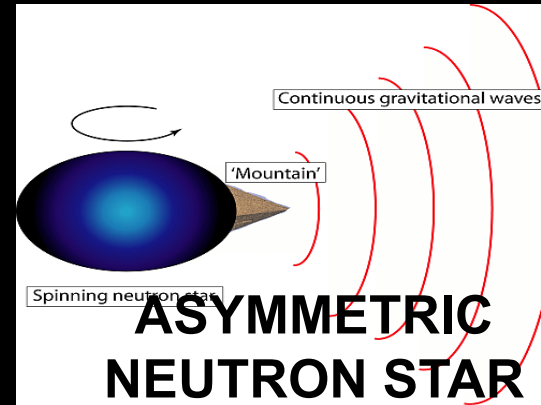
- Sesana (PRL 116, 231102) noticed that if eLISA were online, it would have detected GW150914 a few years ago
 - Pre-merger alert
 - Much smaller sky areas and telescopes ready in advance
- Benefit even if BBH are not luminous! Vitale (PRL 117, 051102) showed that BH spins can be measured a factor 2 better for those joint detections.

Gravitational Wave Sources

TRANSIENT



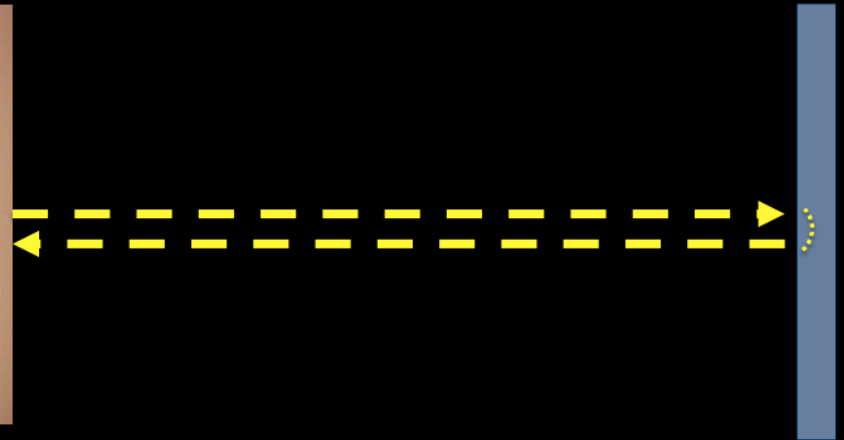
PERSISTENT



MATCH FILTER UNMODELED

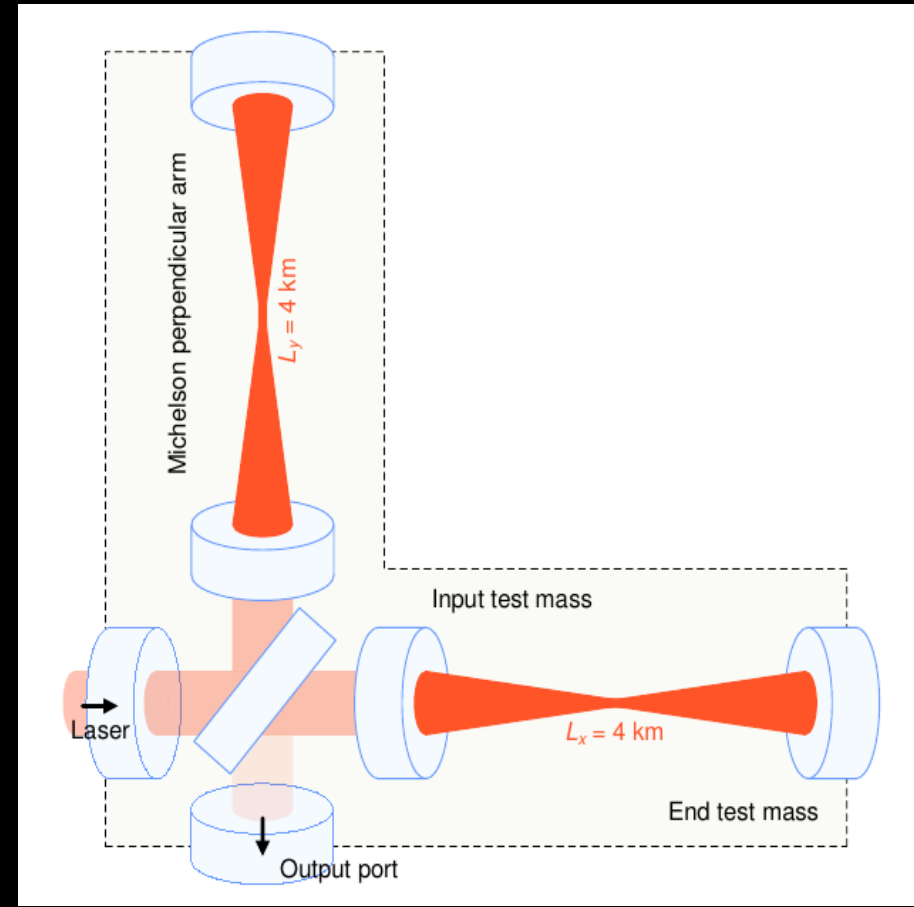
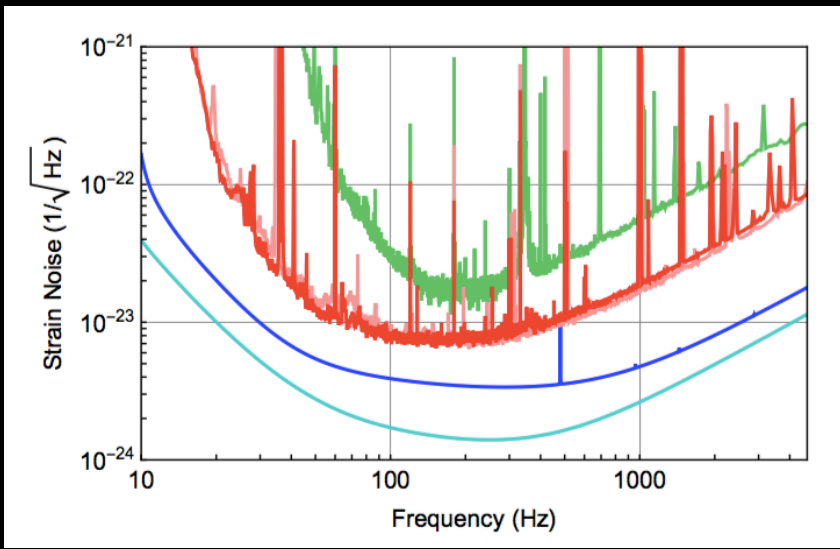
How to detect them?

- We basically need to be able to measure the distance between objects *very* well
- I mean really well:
If the masses are 4km apart, you must be able to monitor their distance to 10^{-18} meters
- Use lasers!



Ground based GW detectors

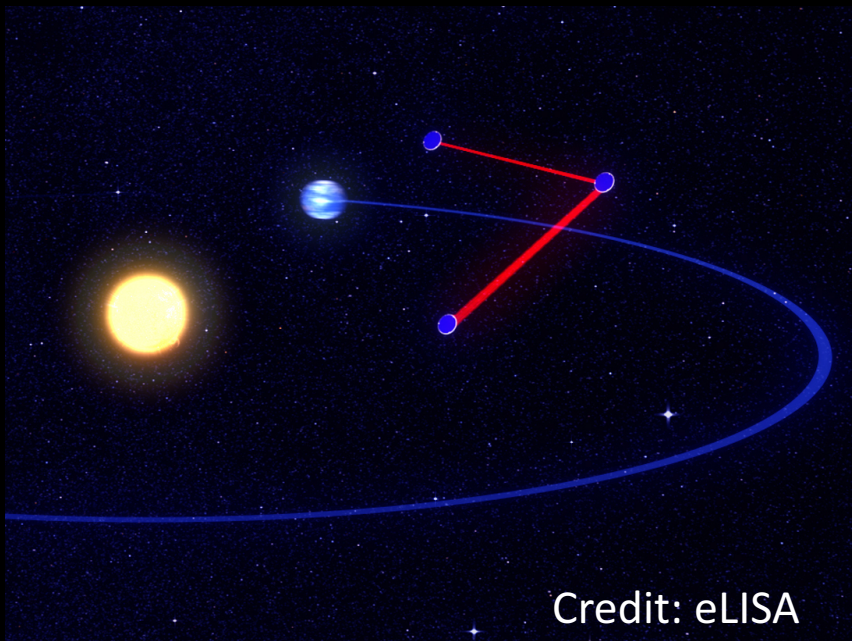
- Use laser interferometry to measure small differential length variations



LVC, 1602.03838

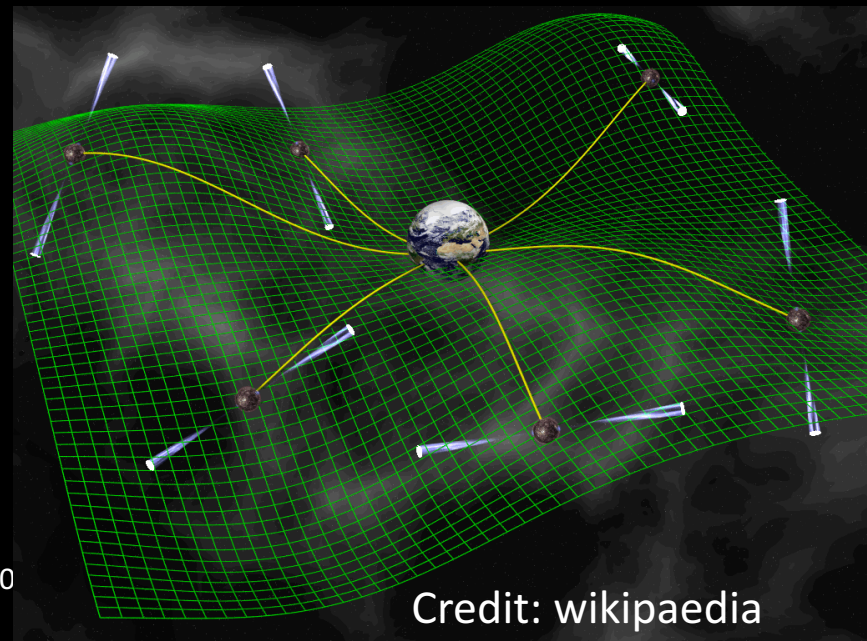
Space based GW detectors

- eLISA
 - Interferometry in space
 - Targets more massive objects than LIGO
 - Launch > 2030
- Pulsar timing array
 - Use pulses time-travel as a clock
 - Targets supermassive BH



Credit: eLISA

ale, Aug 26 20



Credit: wikipedia

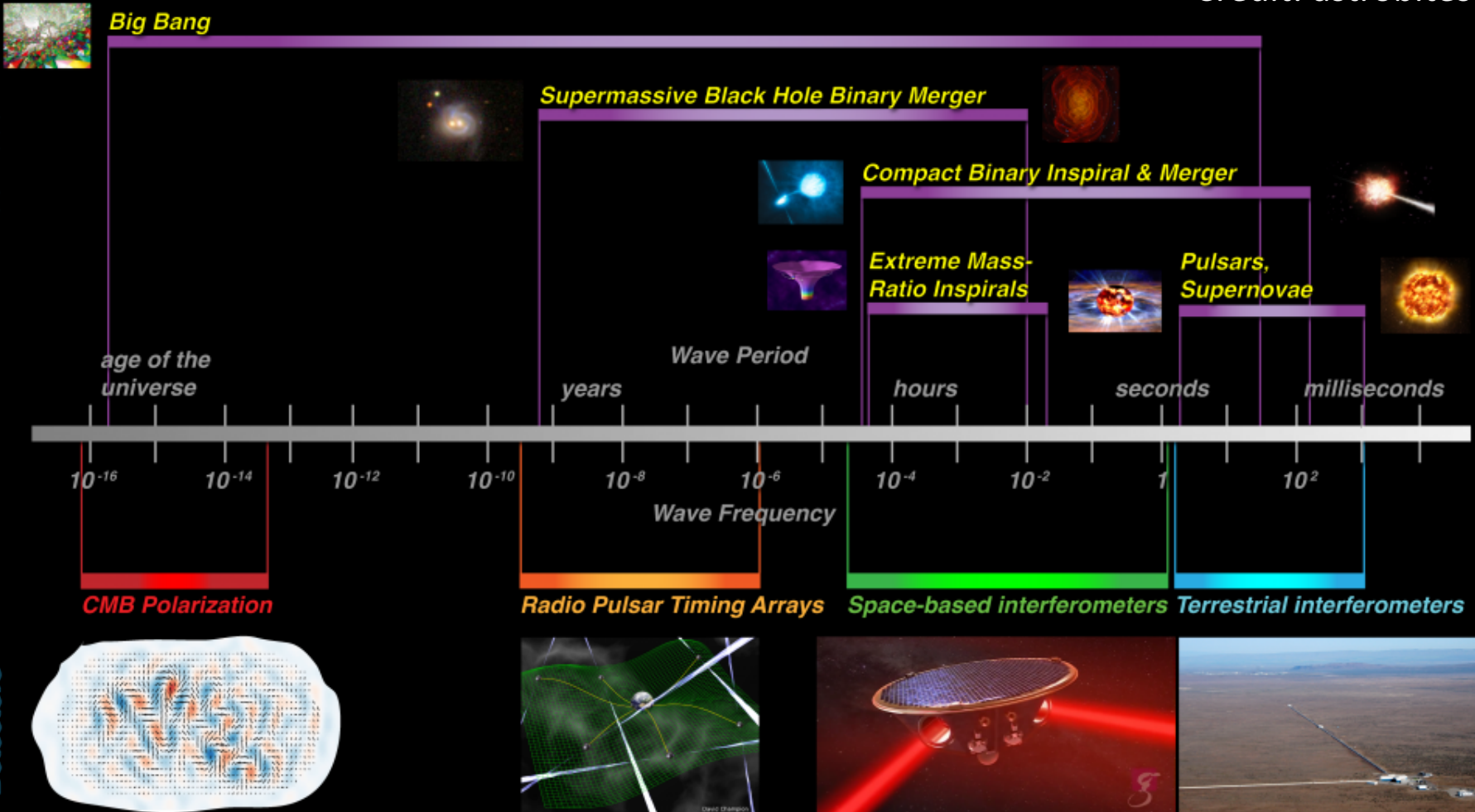
The GW spectrum

Credit: astrobit.es

The Gravitational Wave Spectrum

Sources

Detectors

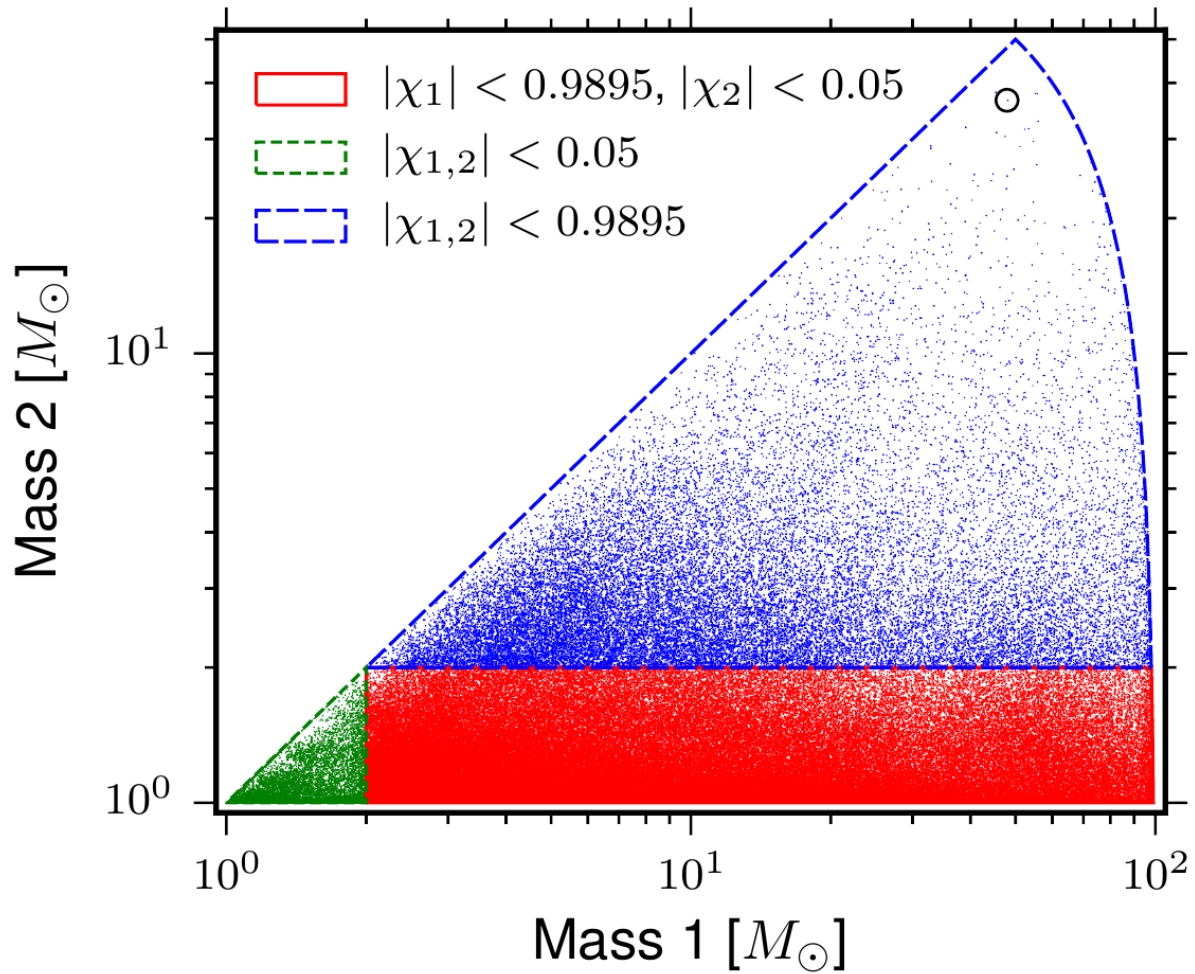


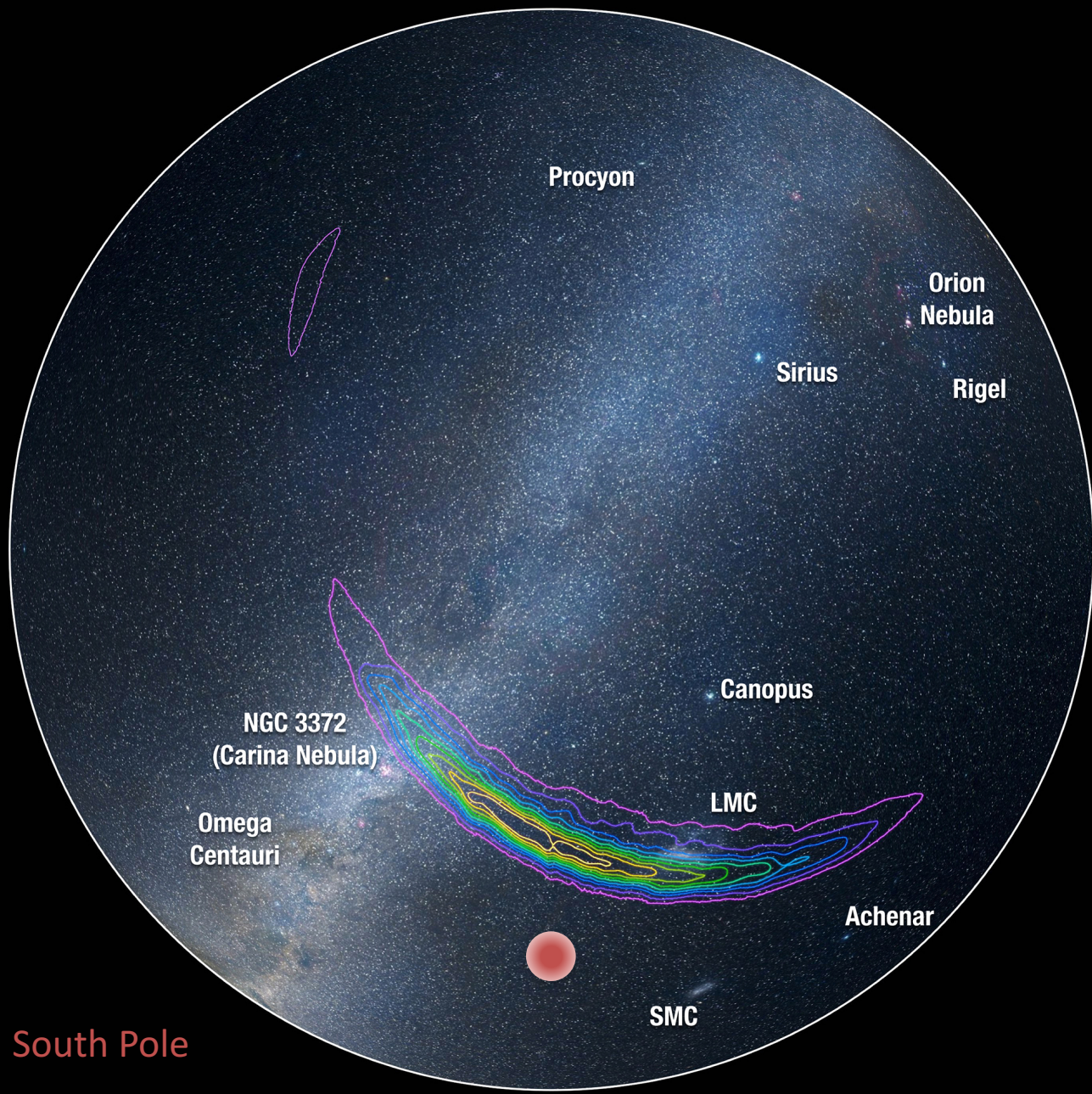
In the next five years...

- More binary black holes detection will lead to:
 - Understanding of black hole mass and spin distributions, formation channels
 - Tests of General Relativity in its strong field regime
- We will probably detect binary neutron stars
 - Joint Electromagnetic/Gravitational Wave discovery: progenitors, environment, opening angle,..
 - Rank equations of state
 - Maximum mass of neutron star?
- We will never stop listening!
 - Gravitational wave bursts from supernovae, other violent events, unknown sources

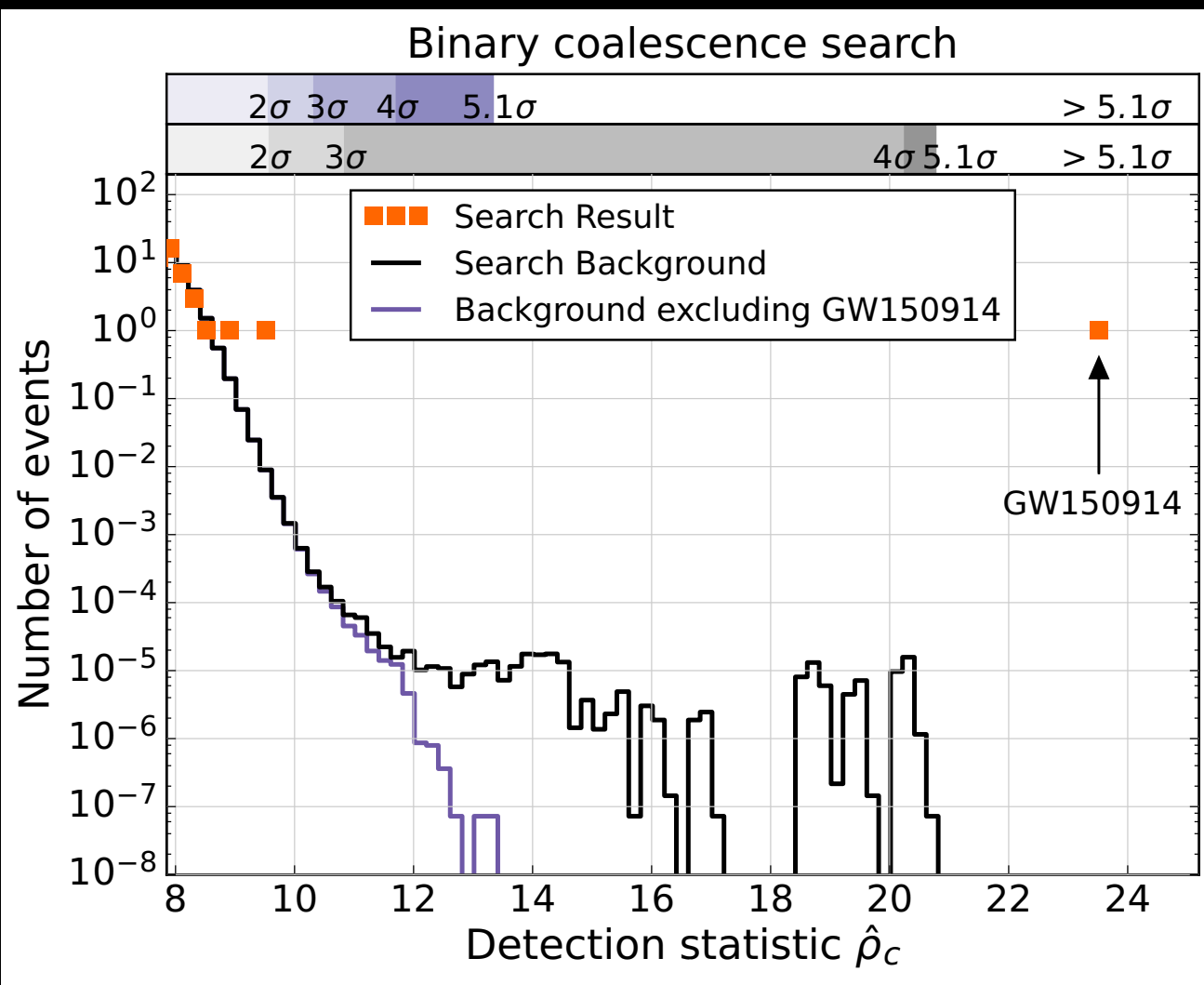
More science with early upgrades to Advanced LIGO (10 years)

- Early upgrades envisioned beyond Advanced LIGO, larger science output (Miller+ PRD 91 062005):
 - Squeezed states of light can improve both high and low frequencies, improving sky localization, measurements of neutron star tidal deformability (Lynch+ PRD 91 044032)
 - Coating and beam size can improve sweet spot sensitivity
 - more detections
 - Suspensions and heavier masses affect < 50 Hz region:
 - higher SNR and parameter estimation for BBH
- Overall up to a factor of 2 in reach, ~ 8 in volume
- Evolving science goals

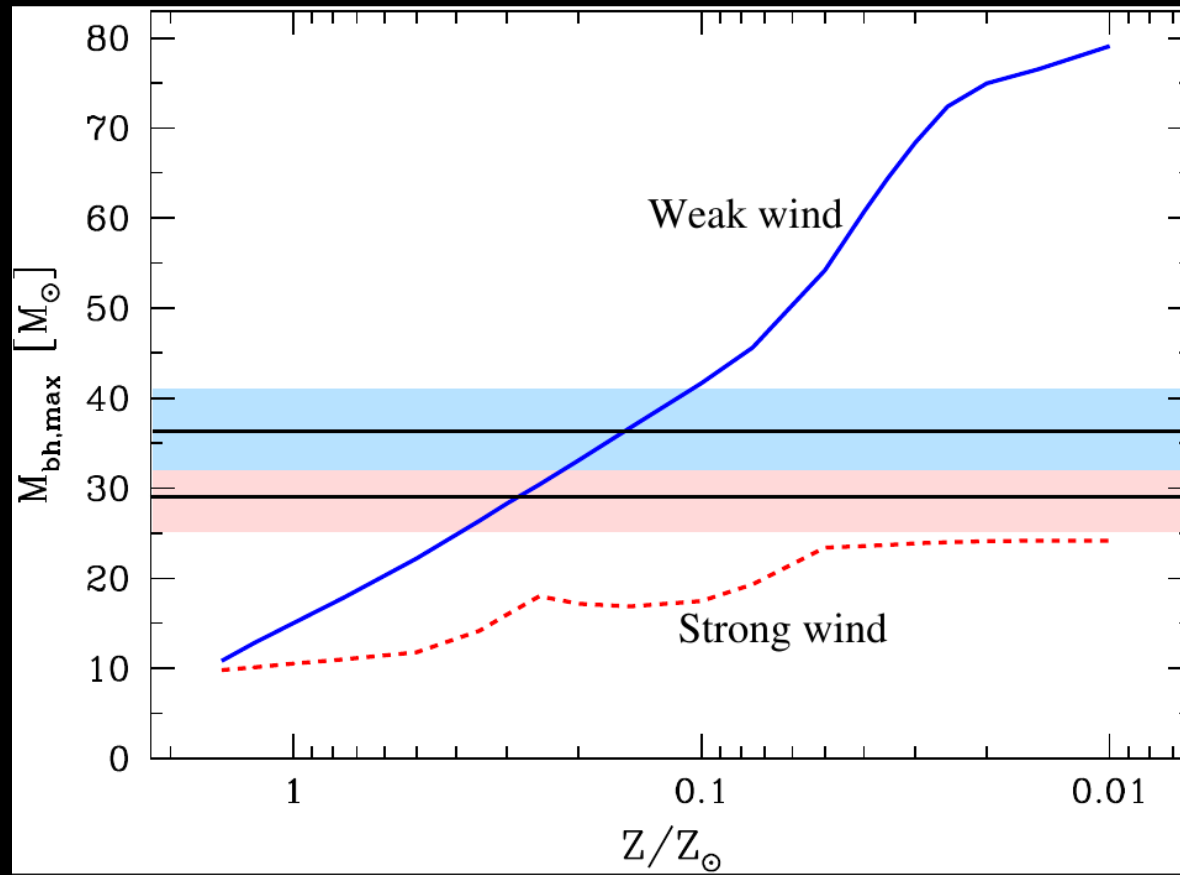




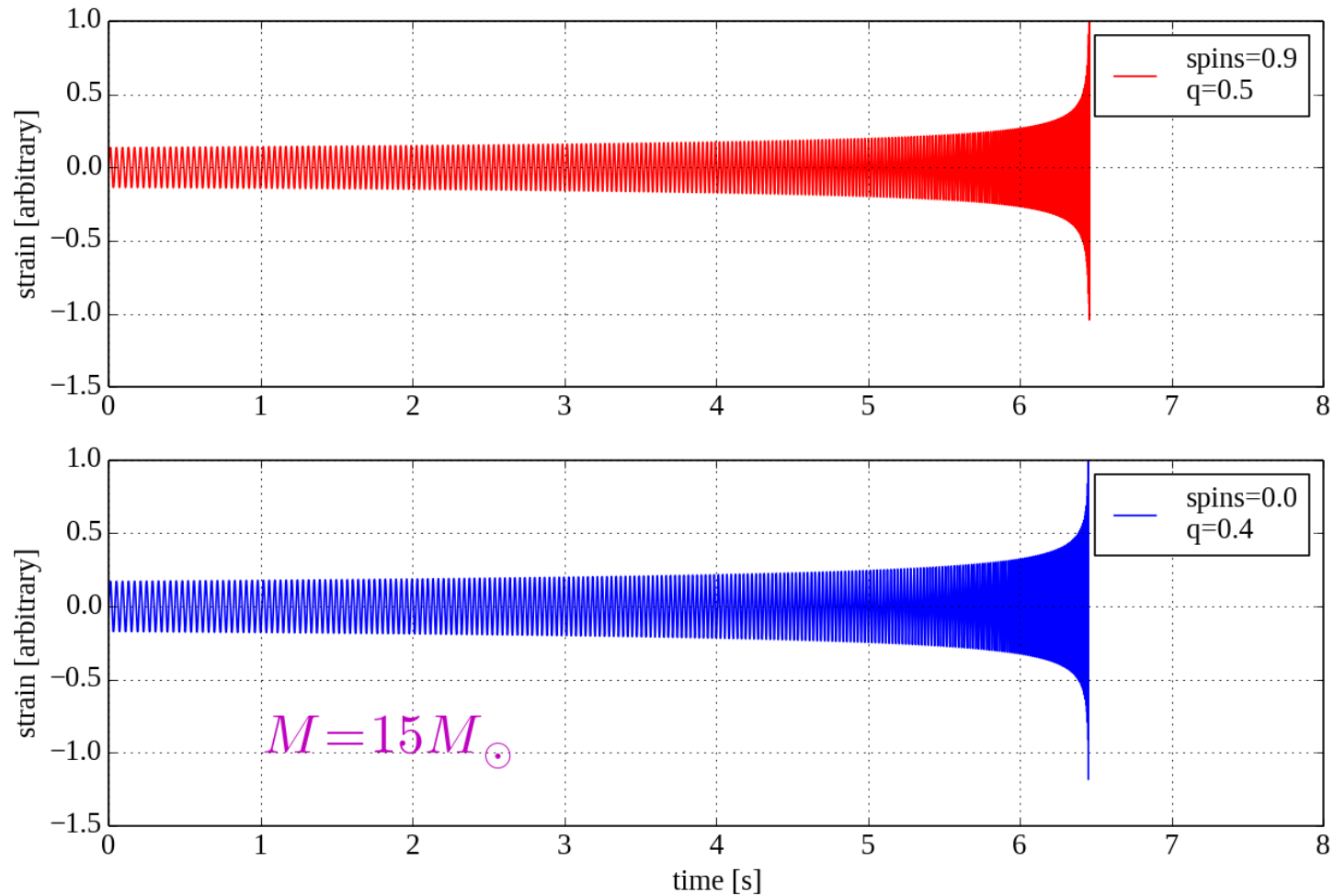
South Pole



Astrophysical implications

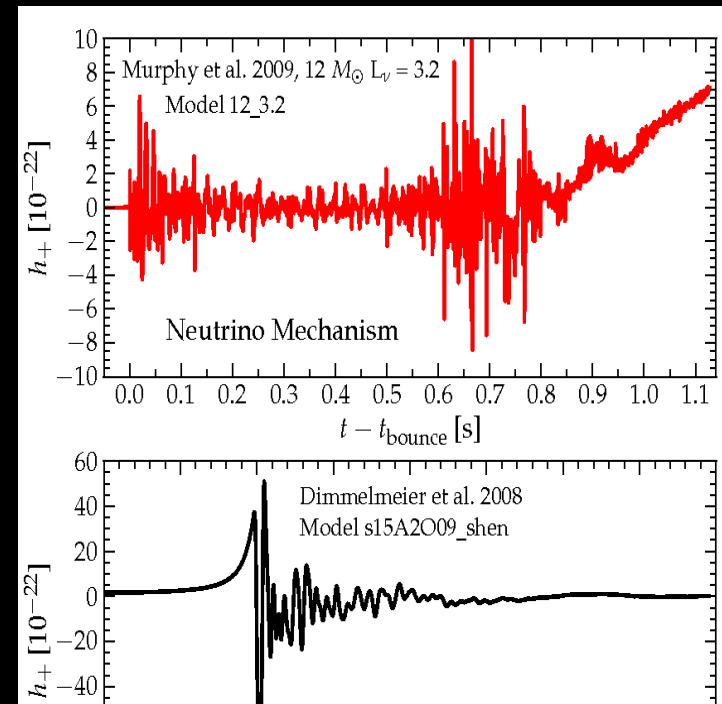


Aligned-spin mass-spin degeneracy



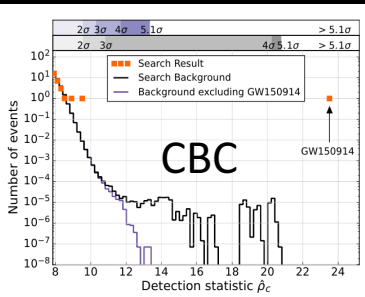
Unmodeled Sources

- Any violent astrophysical or cosmological phenomenon
 - Core collapse supernovae
 - Cosmic strings
 - Post-merger signals from hypermassive neutron stars
 - Something unexpected
- Uncertain or no knowledge of gravitational wave signal morphology
- Less advertised, but not less interesting than CBCs
 - Also, high mass binary black holes escaped CBC net

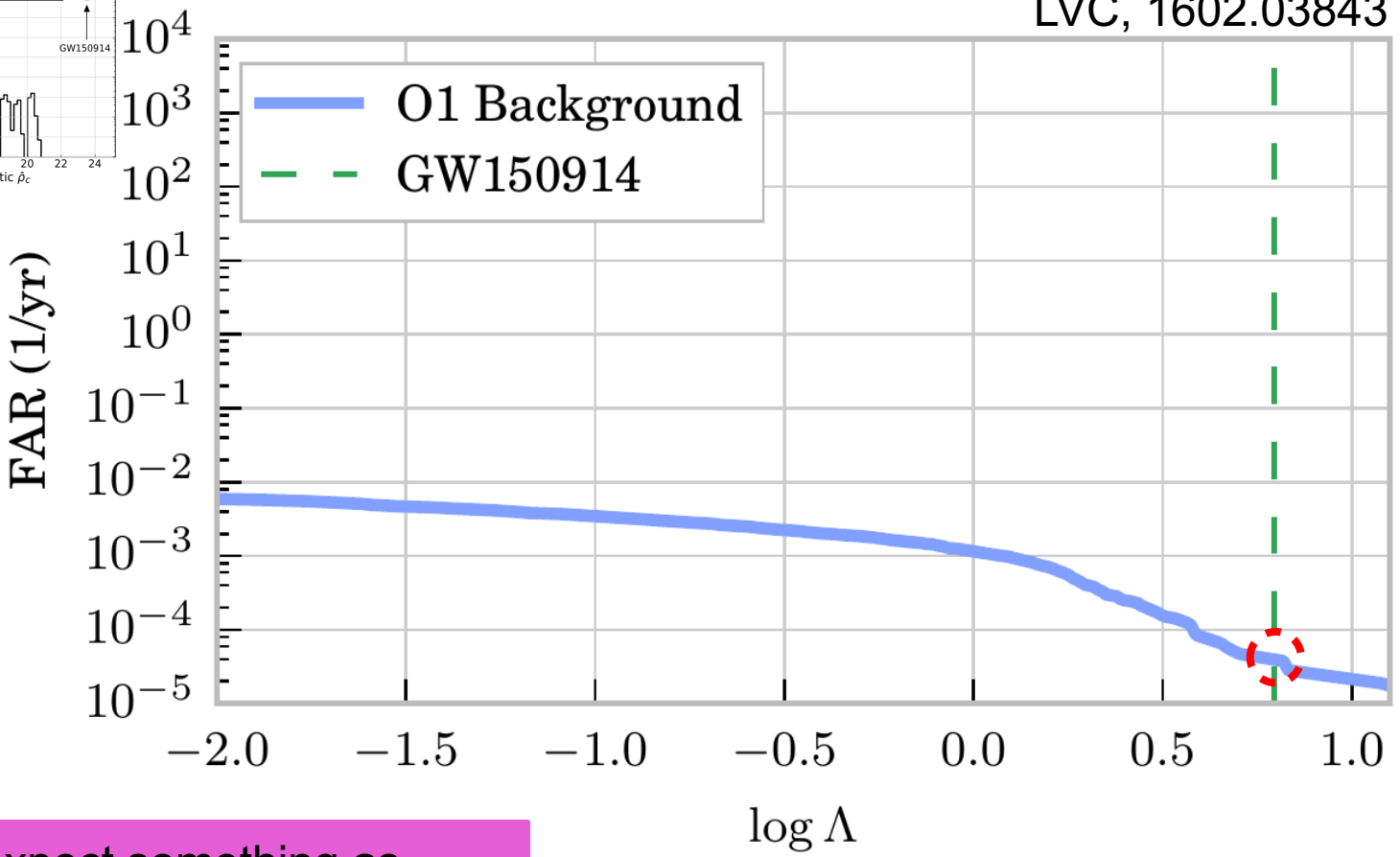


Logue+, PRD 86 044023

oLIB results for GW150914



LVC, 1602.03843



More rare

FAR (1/yr)

log Λ

More GW-like

We expect something as significant as GW150914 to only happen once every 27,000 years: