

# $U(1)_{B-L}$ Extra-Natural Inflation with Standard Model on a Brane

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# Plan

- Introduction
- $U(1)_{B-L}$  Extra-Natural Inflation
- $U(1)_{B-L}$  Extension of Standard Model
- Summary

# Motivations

## Why extra-natural inflation?

- Natural large field inflation model
- EFT approach valid

## Why $U(1)_{B-L}$ ?

- RH neutrinos necessitated by gauge anomaly cancellation
- Ubiquitous in beyond the SM scenarios
  - Appears in GUT and string theory embeddings
  - R-parity made exact in SUSY versions of SM

## The Scenario

- (4+1)-dimensional bulk spacetime
- Extra-dimension compactified on  $S^1$  of radius  $L_5$
- SM on (3+1)-dimensional brane localized in extra dimension
- $U(1)_{B-L}$  gauge theory in bulk
- Inflaton from Wilson loop of  $U(1)_{B-L}$  gauge field

# Ingredients of Extra-Natural Inflation

## The potential:

$$V(\phi) = \frac{V_0}{2} \left[ 1 - \cos\left(\frac{\phi}{f}\right) \right], \quad V_0 = \frac{c_0}{\pi^2} \frac{1}{(2\pi L_5)^4}$$

- $c_0$  “counts”  $U(1)_{B-L}$ -charged fields with  $m \lesssim L_5^{-1}$
- **Decay constant:**  $f^{-1} = g_4(2\pi L_5)$ 
  - $g_4$  4D gauge coupling

For small quantum gravity correction:  $\ell_5^3 \equiv (L_5 M_5)^3 \gg 1$

- $M_P^2 = M_5^3 (2\pi L_5)$ ,  $M_P \simeq 2.4 \times 10^{18}$  GeV
- $\ell_5 = \frac{1}{2\pi} \left( \frac{M_P}{g_4 f} \right)^{2/3}$

# Slow-roll Inflation

**Slow-roll parameters:**

$$\epsilon(\phi) \equiv \frac{M_P^2}{2} \left( \frac{V'}{V} \right)^2 = \frac{M_P^2}{2f^2} \frac{1 + \cos \frac{\phi}{f}}{1 - \cos \frac{\phi}{f}} \ll 1$$

$$\eta(\phi) \equiv M_P^2 \frac{V''}{V} = \frac{M_P^2}{f^2} \frac{\cos \frac{\phi}{f}}{1 - \cos \frac{\phi}{f}} \ll 1$$

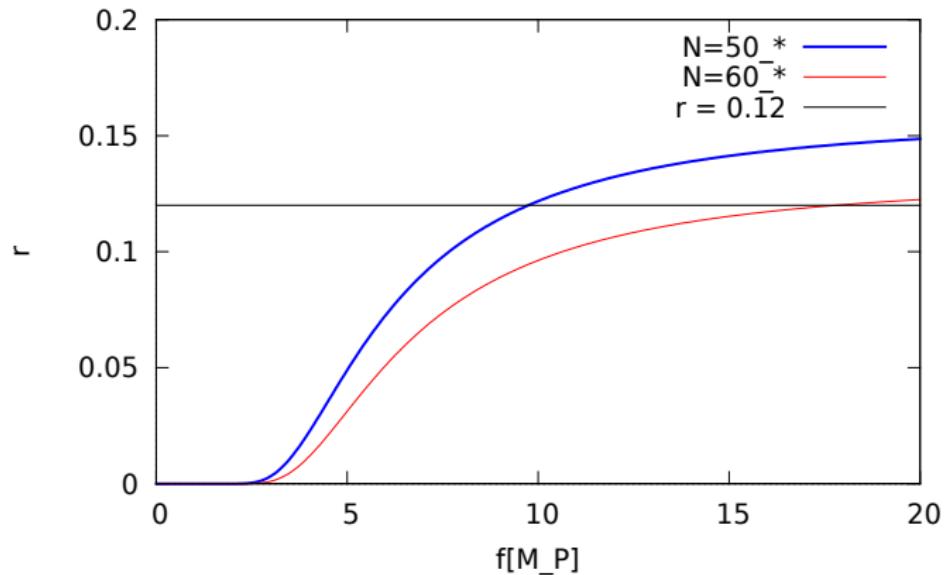
**Number of e-folds:**

$$N(\phi) = - \int_{\phi}^{\phi_e} \frac{d\phi}{M_P^2} \frac{V}{V'} = - \left( \frac{f}{M_P} \right)^2 \log \left[ \frac{1}{2} \left( 1 + \cos \frac{\phi}{f} \right) \right]_{\phi_e}^{\phi}$$

- $\phi_e$  defined by  $\epsilon(\phi_e) = 1$
- $\phi = \phi(f, N)$

# Scalar-to-Tensor Ratio

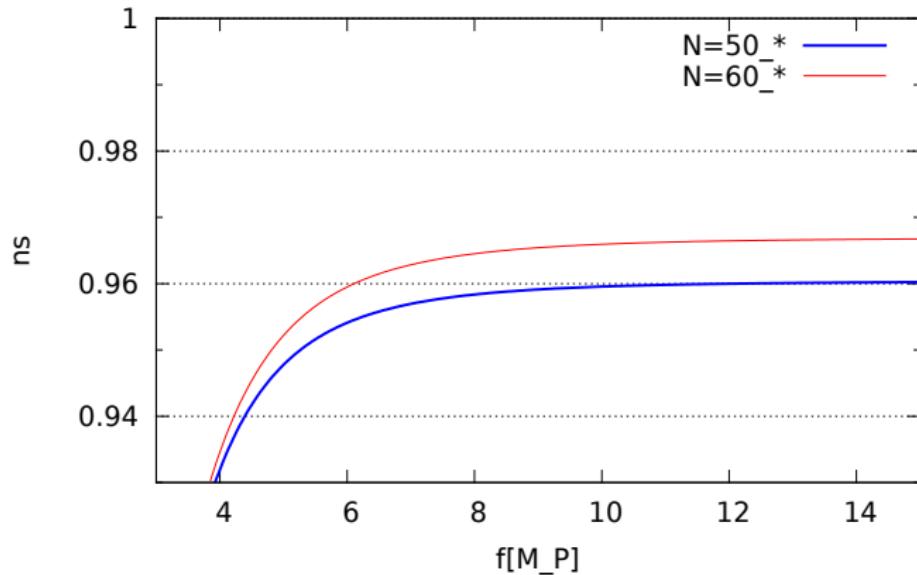
$$r = 16\epsilon$$



Pivot scale  $k_* = 0.002 Mpc^{-1}$

# Spectral Index

$$n_s = 1 - 6\epsilon + 2\eta$$



Pivot scale  $k_* = 0.002 \text{ Mpc}^{-1}$

# Slow-roll Power Spectrum

For spatially flat universe in slow-roll approximation:

- **Friedman equation:**  $3H^2 M_P^2 = \rho = V(\phi)$

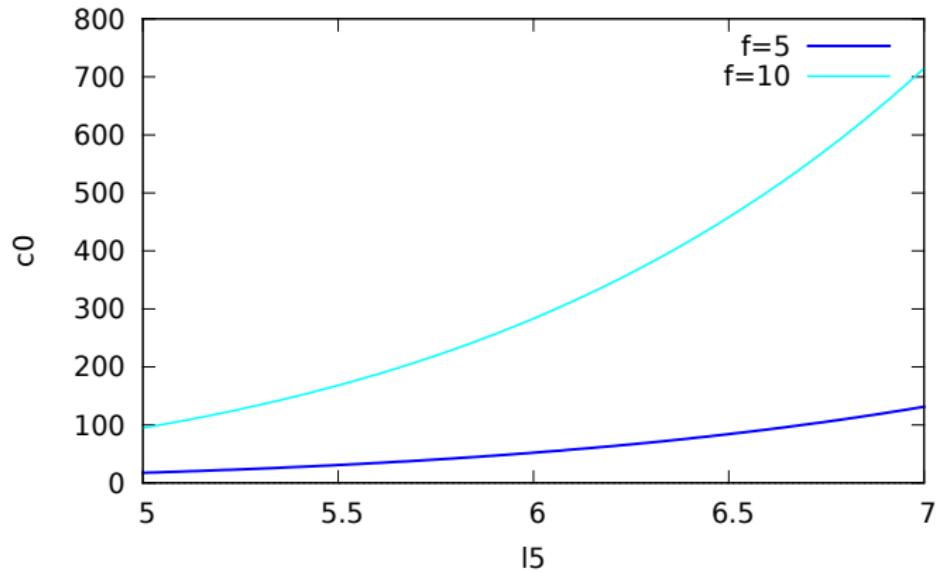
**Power spectrum:**

$$\begin{aligned} P_\zeta(k_*) &= \frac{H^2}{8\pi^2 M_P^2 \epsilon} = \frac{V(\phi_*)}{24\pi^2 M_P^4 \epsilon_*} \\ &= \frac{c_0}{2\pi^2} \frac{1}{(2\pi\ell_5)^6} \left[ 1 - \cos \frac{\phi(f, N_*)}{f} \right] \frac{1}{24\pi^2 \epsilon(f, N_*)} \end{aligned}$$

- **Observed value:**  $P_\zeta(k_*) = 2.2 \times 10^{-9}$
- $c_0$  determined once  $\ell_5$ ,  $f$ ,  $N_*$  given

## Behaviour of $c_0$

$$c_0 \sim \ell_5^6$$



$$N_* = 50$$

# Natural Parameter Values

## Naturalness:

- $N_* \simeq 50$  (from inflaton decay; see later)
- $c_0 \sim \mathcal{O}(10)$

## Observational constraints:

- $r \lesssim 0.12 \Rightarrow f \lesssim 10M_P$
- $n_s > 0.94 \Rightarrow f \gtrsim 5M_P$

**Preferred value:**  $\ell_5 \simeq 5, f \simeq 5$

$$g_4 \simeq 10^{-3}, L_5^{-1} \simeq 10^{17} \text{ GeV}$$

## $U(1)_{B-L}$ Standard Model

**Particle content and charge assignment:**

	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)_{B-L}$
$q_L^i$	<b>3</b>	<b>2</b>	$+1/6$	$+1/3$
$u_R^i$	<b>3</b>	<b>1</b>	$+2/3$	$+1/3$
$d_R^i$	<b>3</b>	<b>1</b>	$-1/3$	$+1/3$
$\ell_L^i$	<b>1</b>	<b>2</b>	$-1/2$	$-1$
$\nu_R^i$	<b>1</b>	<b>1</b>	$0$	$-1$
$e_R^i$	<b>1</b>	<b>1</b>	$-1$	$-1$
$H$	<b>1</b>	<b>2</b>	$-1/2$	$0$
$\Sigma$	<b>1</b>	<b>1</b>	$0$	$+2$

- **4D brane fields:** SM and RH neutrinos
- **5D bulk fields:**  $U(1)_{B-L}$  gauge field  $A_M$ , complex scalar  $\Sigma$

## Scalar Sector

**4D scalar potential:**

$$V(H, \Sigma) = \mu_H^2 H^\dagger H + \mu_\Sigma^2 \Sigma_0^* \Sigma_0 + \frac{\lambda_1}{2} (H^\dagger H)^2 + \frac{\lambda_2}{2} (\Sigma_0^* \Sigma_0)^2 + \lambda_3 H^\dagger H \Sigma_0^* \Sigma_0$$

- $\Sigma_0$  zero mode of  $\Sigma$  in the extra dimension

**After SSB:**  $\Sigma$  breaks  $U(1)_{B-L}$ ,  $m_W$  fixes  $v_H = 246$  GeV

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_H + h \end{pmatrix}, \quad \Sigma_0 = \frac{v_\Sigma + s}{\sqrt{2}}$$
$$\frac{v_H^2}{2} = \frac{-\mu_H^2 \lambda_2 + \mu_\Sigma^2 \lambda_3}{\lambda_1 \lambda_2 - \lambda_3^2}, \quad \frac{v_\Sigma^2}{2} = \frac{-\mu_\Sigma^2 \lambda_1 + \mu_H^2 \lambda_3}{\lambda_1 \lambda_2 - \lambda_3^2}$$

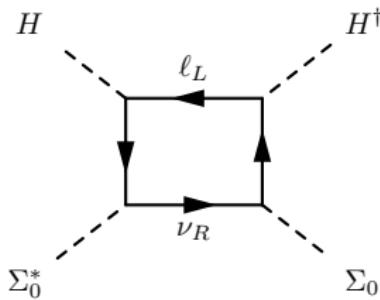
## Neutrino Sector

See-saw mechanism:

$$\mathcal{L} \supset -Y_D^{ij} \overline{\nu_R^i} H^\dagger \ell_L^j - \frac{1}{2} Y_N^{ij} \overline{\nu_R^{ic}} \nu_R^j \Sigma_0 + h.c.$$

- **Active neutrino mass:**  $m_\nu \sim Y_D^2 v_H^2 / M_N$
- **RH neutrino mass:**  $M_N = Y_N v_\Sigma / \sqrt{2}$

Scalar mixing  $H^\dagger H \Sigma_0^* \Sigma_0$  induced at one-loop:



## Constraints

**Higgs mass contribution:**

$$\delta m_H^2 \sim \frac{Y_D^2 Y_N^2}{(4\pi)^2} \frac{v_\Sigma^2}{2} \sim \frac{m_\nu M_N^3}{(4\pi)^2 v_H^2}$$

**Naturalness:**  $m_H = 126$  GeV,  $\sqrt{\delta m_H^2} \lesssim \mathcal{O}(100)$  GeV

**Bench mark:**  $m_\nu \sim 0.1$  eV

- **B-L breaking scale:**  $M_N \lesssim 10^7$  GeV,  $v_\Sigma \lesssim 10^7 / Y_N$  GeV
- $|\lambda_3| \sim \frac{m_\nu M_N}{(4\pi)^2 v_H^2} Y_N^2 \lesssim 10^{-10} Y_N^2$
- $m_{Z'} \sim g_4 v_\Sigma \lesssim \mathcal{O}(10)$  TeV
- **Collider limit:**  $m_{Z'} \geq g_4 \times (6 \text{ TeV}) \gtrsim 6$  GeV

## Inflaton Decay

$\mathbb{Z}_2$  symmetry:  $x^5 \rightarrow -x^5$ ,  $A_5 \rightarrow -A_5$

- Inflaton absolutely stable if exact

$\mathbb{Z}_2$ -breaking Chern-Simons coupling:

$$S_{CS} = \frac{k}{48\pi^3} \int \mathcal{A} \mathcal{F}^2 = g_4^2 \frac{k}{16\pi^2} \int d^4x \frac{\phi}{2\pi f} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Dominant decay width:

$$\Gamma_{\phi \rightarrow AA} \simeq \frac{g_4^4}{16\pi} \left( \frac{k}{32\pi^3} \right)^2 \frac{m_\phi^3}{f^2}$$

Inflaton mass:  $\ell_5 \simeq 5$ ,  $f \simeq 5M_P$

$$m_\phi^2 = \frac{V_0}{2f^2} = \frac{c_0}{2\pi^2} g_4^4 f^2 \simeq 10^{13} \text{ GeV}$$

# Reheating

**Reheating temperature:**

$$\begin{aligned} T_R &= \left( \frac{90}{\pi^2 g_*(T_R)} \right)^{1/4} \sqrt{\Gamma M_P} \\ &\simeq \left( \frac{90}{g_*(T_R)} \right)^{1/4} \frac{g_4^2}{4\pi} \frac{|k|}{32\pi^3} \sqrt{\frac{m_\phi^3 M_P}{f^2}} \simeq |k| \times 1 \text{ GeV} \end{aligned}$$

- $T_R \simeq 1 \sim 10 \text{ GeV}$ ,  $g_*(T_R) \simeq 60 \sim 90$

**Standard estimate:**

$$N_* \simeq 49 + \frac{2}{3} \ln \left( \frac{\rho_*^{1/4}}{10^{16} \text{ GeV}} \right) + \frac{1}{3} \ln \left( \frac{T_R}{1 \text{ GeV}} \right)$$

- $N_* \simeq 50$  natural

# Summary

- Cosmological observations constrain the  $U(1)_{B-L}$  gauge coupling to  $g_4 \lesssim 10^{-3}$ , which in turn constrains the particle physics scenario at high energy assuming naturalness.
- Allowed interaction between the inflaton and SM particles restricted by the higher-dimensional set-up, which enables a simple determination of the reheating temperature  $T_R$ .

## Future directions:

- Baryon asymmetry
- Dark matter
- String theory embedding