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“Constraints on general second-order scalar-tensor models from
gravitational Cherenkov radiation”

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CONSTRAINTS ON GENERAL SECOND-ORDER SCALAR-TENSOR MODELS FROM GRAVITATIONAL CHERENKOV RADIATION

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BASED ON
RK AND KAZUHIRO YAMAMOTO, JCAP 07 (2012) 050

INTRODUCTION

📌 Accelerating universe

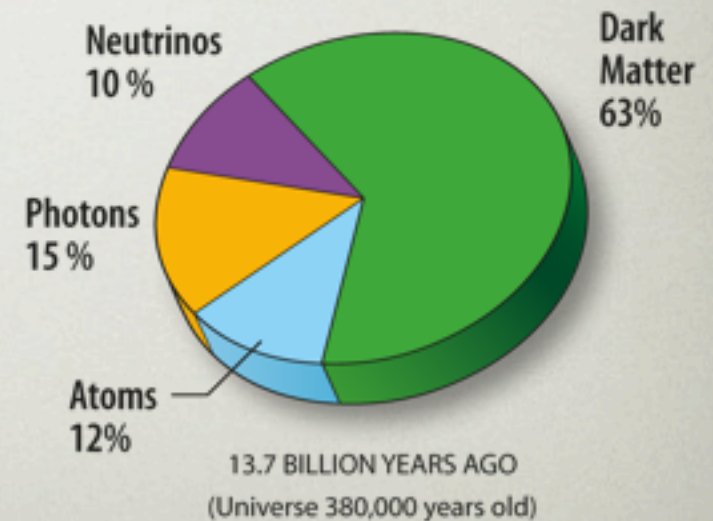
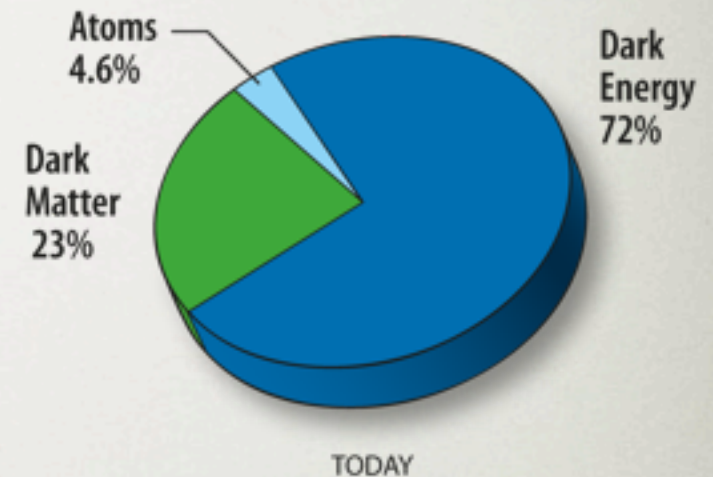
- Implication of cosmological constant?

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

- Observationally, fine !!
- **Cosmological constant problem**

121 orders of magnitude differences

Can **modification of gravity** solve this puzzle ???



(FROM WMAP WEBSITE)

GALILEON THEORY

✓ Galileon (Nicolis et al. '09)

$$\square = \nabla_{\mu} \nabla^{\mu}$$

$$\mathcal{L} \supset (\partial\varphi)^2 \underline{\square\varphi}$$

second derivative with respect to space-time

✓ Galileon term contains the second derivative term, but ...

coupling between scalar and curvature

$$\text{EOM} \supset (\square\varphi)^2 - (\nabla_{\mu} \nabla_{\nu} \varphi)^2 - R_{\mu\nu} \nabla^{\mu} \varphi \nabla^{\nu} \varphi$$

No higher-order derivative terms in EOM !!

MOST GENERAL SECOND-ORDER SCALAR-TENSOR THEORY (MGST)

✓ Horndeski found the most general Lagrangian whose EOM is second-order differential equation for ϕ and $g_{\mu\nu}$ (also known as **Generalized galileon**)

Horndeski, Int. J. Theor. Phys. 10,363 (1974), Deffayet, Gao, Steer (2011)

$$\mathcal{L}_2 = K(\phi, X) \longrightarrow \text{K-essence term} \quad \mathcal{L}_2 \supset (\partial\phi)^2, V(\phi)$$

$$\mathcal{L}_3 = -G_3(\phi, X)\square\phi \longrightarrow \text{Cubic galileon term} \quad \mathcal{L}_3 \supset (\partial\phi)^2\square\phi$$

$$\mathcal{L}_4 = G_4(\phi, X)R + G_{4,X}[(\square\phi)^2 - (\nabla_\mu\nabla_\nu\phi)(\nabla^\mu\nabla^\nu\phi)] \longrightarrow \text{Einstein-Hilbert term} \quad \mathcal{L}_4 \supset (M_{\text{Pl}}^2/2)R$$

$$\mathcal{L}_5 = G_5(\phi, X)G_{\mu\nu}(\nabla^\mu\nabla^\nu\phi) \longrightarrow \text{Non-minimal derivative coupling} \quad \mathcal{L}_5 \supset G^{\mu\nu}\nabla_\mu\phi\nabla_\nu\phi$$

$$- \frac{1}{6}G_{5,X} \left[(\square\phi)^3 - 3(\square\phi)(\nabla_\mu\nabla_\nu\phi)(\nabla^\mu\nabla^\nu\phi) \right.$$

$$\left. + 2(\nabla^\mu\nabla_\alpha\phi)(\nabla^\alpha\nabla_\beta\phi)(\nabla^\beta\nabla_\mu\phi) \right]$$

(Germani et al. 2011;
Gubitosi, Linder 2011)

$$X = -(\partial\phi)^2/2, \quad G_{iX} = \partial G_i/\partial X$$


WHY GALILEON??

 **Self-accelerating solution**

 **Free of ghost-instabilities**

 **Vainshtein mechanism** (Vainshtein 1972)

- Scalar field is effectively weakly coupled to matter in a high density region
- Reduce general relativity at small scales

 Relation with decoupling limit in **massive gravity**

(de Rham, Gabadadze, Tolley, 2010)

COSMOLOGICAL OBSERVATIONS

RK, Kazuhiro Yamamoto, JCAP 04 (2011) 025


RK, Tsutomu Kobayashi, Kazuhiro Yamamoto, Physical Review D 85 (2012) 123503

 Standard rulers (supernovae + CMB shift parameter)

- **Not powerful tools** to constrain model parameters in modified gravity theories, but **useful tools** to determine cosmological parameters

 Galaxy distribution (SDSS LRG sample)

- **The error bar is still large** to constrain model parameters

 Cross correlation between LSS and ISW

- **Excellent tool** to constrain modified gravity
- Indicates that the effective gravitational coupling **G_{eff} has to be smaller than $\sim 1.2 G_N$** , otherwise CCF becomes **negative** which contradicts with observations

Other signatures ??

SOUND SPEED OF GRAVITON IN MGST

 Quadratic action for **a tensor mode** in the most general scalar-tensor theory

Kobayashi, Yamaguchi, Yokoyama, Prog. Theor. Phys. 126, 511 (2011),

$$S_T^{(2)} = \frac{1}{8} \int dt d^3x a^3 \left[\mathcal{G}_T \dot{h}_{ij}^2 - \frac{\mathcal{F}_T}{a^2} (\vec{\nabla} h_{ij})^2 \right]$$

 Sound speed of graviton

$$c_T^2 \equiv \frac{\mathcal{F}_T}{\mathcal{G}_T}$$

$$\mathcal{F}_T \equiv 2 \left[G_4 - X \left(\ddot{\phi} G_{5X} + G_{5\phi} \right) \right]$$

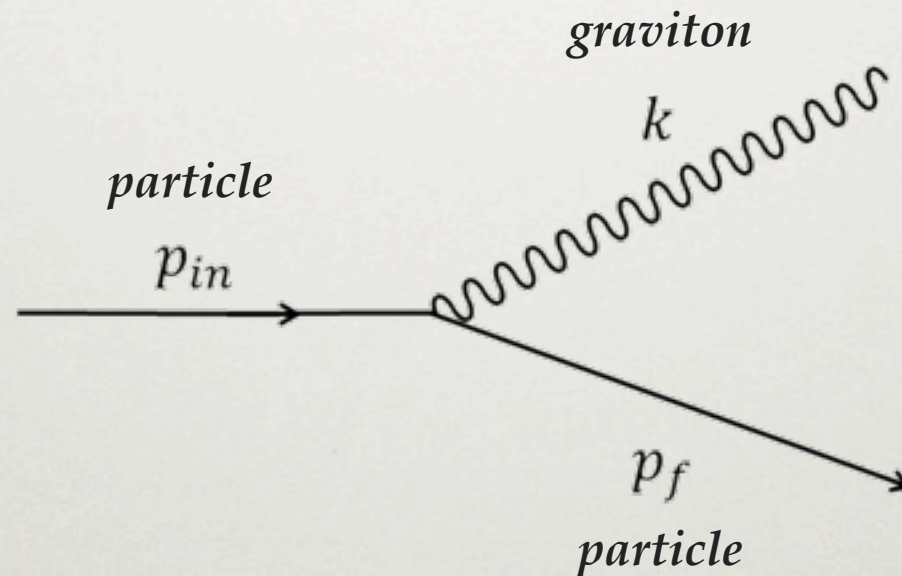
$$\mathcal{G}_T \equiv 2 \left[G_4 - 2X G_{4X} - X \left(H \dot{\phi} G_{5X} - G_{5\phi} \right) \right]$$

Sound speed of graviton could be **different from speed of light !!!**

GRAVITATIONAL CHERENKOV RADIATION

- 📌 If the sound speed of graviton is smaller than the speed of light, **particle should emit graviton** through the similar process to Cherenkov radiation

Moore and Nelson (2001)



- 📌 Highest energy cosmic ray ($p \sim 3 \times 10^{11}$ GeV) can provide us **the lower bound on the sound speed of graviton**

GRAVITATIONAL CHERENKOV RADIATION

Consider the complex scalar in a **FRW background**

$$S_m = \int d^4x \sqrt{-g} \left[-g^{\mu\nu} \partial_\mu \Psi^* \partial_\nu \Psi - m^2 \Psi^* \Psi - \xi R \Psi^* \Psi \right]$$

Quantize the complex scalar and tensor field as

$$\hat{\Psi}(\eta, \mathbf{x}) = \frac{1}{a} \int \frac{d^3p}{(2\pi)^{3/2}} \left[\hat{b}_{\mathbf{p}} \psi_p(\eta) e^{i\mathbf{p}\cdot\mathbf{x}} + \hat{c}_{\mathbf{p}}^\dagger \psi_p^*(\eta) e^{-i\mathbf{p}\cdot\mathbf{x}} \right] \quad [\hat{b}_{\mathbf{p}}, \hat{b}_{\mathbf{p}'}^\dagger] = \delta(\mathbf{p} - \mathbf{p}')$$
$$\hat{h}_{\mu\nu} = \frac{1}{a} \sqrt{\frac{2}{\mathcal{G}_T}} \sum_\lambda \int \frac{d^3k}{(2\pi)^{3/2}} \left[\varepsilon_{\mu\nu}^{(\lambda)} \hat{a}_{\mathbf{k}} h_k(\eta) e^{i\mathbf{k}\cdot\mathbf{x}} + \varepsilon_{\mu\nu}^{(\lambda)} \hat{a}_{\mathbf{k}}^\dagger h_k^*(\eta) e^{-i\mathbf{k}\cdot\mathbf{x}} \right] \quad [\hat{c}_{\mathbf{p}}, \hat{c}_{\mathbf{p}'}^\dagger] = \delta(\mathbf{p} - \mathbf{p}')$$

Mode functions satisfy

$$\left(\frac{d^2}{d\eta^2} + p^2 + m^2 a^2 \right) \psi_p(\eta) = 0$$
$$\left(\frac{d^2}{d\eta^2} + c_T^2 k^2 - \frac{a''}{a} \right) h_k(\eta) = 0$$

GRAVITATIONAL CHERENKOV RADIATION

The total radiation energy from the complex scalar field

$$E = \sum_{\lambda} \sum_{\mathbf{k}} (\omega_{\mathbf{k}}/a) \langle \hat{a}_{\mathbf{k}}^{\dagger(\lambda)} \hat{a}_{\mathbf{k}}^{(\lambda)} \rangle$$

where

$$\langle \hat{a}_{\mathbf{k}}^{\dagger(\lambda)} \hat{a}_{\mathbf{k}}^{(\lambda)} \rangle = 2\Re \int_{t_{\text{in}}}^t dt_2 \int_{t_{\text{in}}}^{t_2} dt_1 \langle H_I(t_1) \hat{a}_{\mathbf{k}}^{\dagger(\lambda)} \hat{a}_{\mathbf{k}}^{(\lambda)} H_I(t_2) \rangle$$

$$H_I = a \int d^3x h_{ij} \partial_i \Psi \partial_j \Psi^*$$

Graviton emission rate (using sub-horizon approximation)

$$\frac{dE}{dt} \simeq \frac{G_N p_{\text{in}}^4}{a^4} \frac{4(1 - c_T)^2}{3(1 + c_T)^2}$$

A particle with momentum p **cannot** possibly have been traveling for **longer than**

$$t \sim \frac{a^4}{G_N} \frac{(1 + c_T)^2}{4(1 - c_T)^2} \frac{1}{p^3}$$

GRAVITATIONAL CHERENKOV RADIATION

📌 Observations from cosmic rays tells

Time scale that cosmic ray turn into radiation energy of graviton

>

Time scale that cosmic ray travels from origin to us

📌 The highest energy cosmic ray

- Energy
- Distance

$$E_{\text{highest}} \sim 10^{11} \text{ GeV}$$

$$ct \sim 10 \text{ kpc}$$

📌 **Constraint on the sound speed of graviton**

$$1 - c_T \lesssim 2 \times 10^{-15}$$

TOY MODEL 1

📌 Gubitosi and Linder model (Gubitosi and Linder 2011)

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{Pl}}^2}{2} R + X + \frac{\lambda}{M_{\text{Pl}}^2} G^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi + \mathcal{L}_m[g_{\mu\nu}, \psi] \right]$$

$$K = X$$

$$G_3 = 0$$

$$G_4 = M_{\text{Pl}}^2/2$$

$$G_5 = -\lambda\phi/M_{\text{Pl}}^2$$

📌 Condition for existence of self-accelerating solution and avoiding the ghost-instability

$$-\frac{1}{18} \frac{M_{\text{Pl}}^2}{H_0^2} < \lambda < -\frac{1}{30} \frac{M_{\text{Pl}}^2}{H_0^2}$$

λ is always negative

📌 Sound speed of graviton

$$c_T^2 = \frac{M_{\text{Pl}}^2 + 2\lambda\dot{\phi}^2/M_{\text{Pl}}^2}{M_{\text{Pl}}^2 - 2\lambda\dot{\phi}^2/M_{\text{Pl}}^2} < 1$$

Inconsistent with the constraint from the gravitational Cherenkov radiation...

TOY MODEL 2

$$c_T > 1 - \epsilon$$

$$\epsilon = 2 \times 10^{-15}$$

 Extended galileon model (De Felice and Tsujikawa 2011)

$$K = -c_2 M_2^{4(1-p_2)} X^{p_2},$$

$$G_3 = c_3 M_3^{1-4p_3} X^{p_3},$$

$$G_4 = \frac{1}{2} M_{\text{pl}}^2 - c_4 M_4^{2-4p_4} X^{p_4},$$

$$G_5 = 3c_5 M_5^{-(1+4p_5)} X^{p_5},$$

$$p_2 = p$$

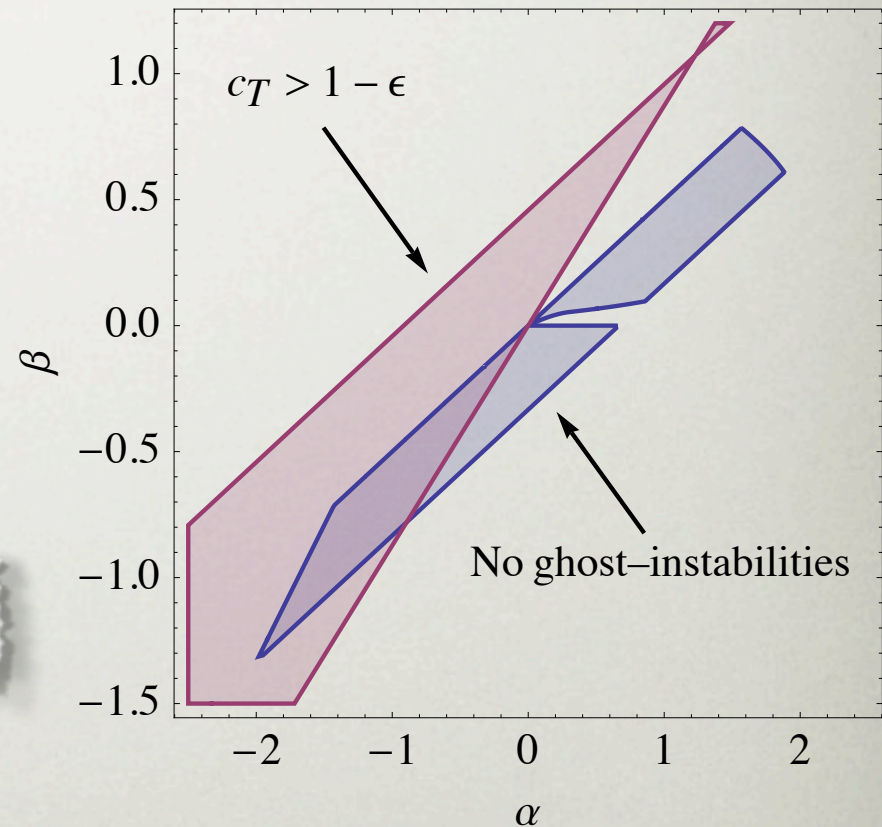
$$p_3 = p + (2q - 1/2)$$

4 model parameters $p, q, \alpha(c_3), \beta(c_4)$

$$p_4 = p + 2q$$

$$p_5 = p + (6q - 1)/2$$

$p = 1$ and $q = 1/2$



Strong constraints for the model parameters α and β

SUMMARY


- In the most general scalar-tensor theory, sound speed of graviton could be different from speed of light
- The constraints from **gravitational Cherenkov radiation** would be a powerful probe.
- **Gravitational Cherenkov radiation** could be a criteria for the construction of modification of gravity

GRAVITATIONAL CHERENKOV RADIATION FOR MASSIVE GRAVITON

 Quadratic action for **a tensor mode** for massive graviton

Gumrukcuoglu, Kuroyanagi, Lin, Mukohyama, Tanahashi (2012)

$$I_{tensor}^{(2)} = \frac{M_{Pl}^2}{8} \int dt d^3x N a^3 \sqrt{\Omega} \left[\frac{1}{N^2} \dot{\gamma}^{ij} \dot{\gamma}_{ij} + \frac{c_g^2(t)}{a^2} \gamma^{ij} (\Delta - 2K) \gamma_{ij} - M_{GW}^2(t) \gamma^{ij} \gamma_{ij} \right]$$

 Dispersion relation

$$\omega_k^2 = c_g^2 k^2 + a^2 M_{GW}^2$$

For $c_g=c$, there is **no gravitational Cherenkov radiation** even if $m \neq 0$

Currently, checking the case $c_g \neq c$ and $m \neq 0$...