

Kazunori Nakayama, JGRG 22(2012)111209

“Probing dark radiation with inflationary gravitational waves”

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**RESCEU SYMPOSIUM ON  
GENERAL RELATIVITY AND GRAVITATION**

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# Probing dark radiation with inflationary gravitational waves

Kazunori Nakayama (The University of Tokyo)

R.Jinno, T.Moroi, KN, arXiv:1208.0184

JGRG22 @ University of Tokyo (2012/11/12)

# Contents

- Observational evidence of dark radiation
- Effects of dark radiation on inflationary gravitational waves

# Dark radiation

- Radiation energy density

$$\rho_{\text{rad}} = \left[ 1 + N_{\text{eff}} \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} \right] \rho_{\gamma}$$

$N_{\text{eff}} = 3.04$  in the standard model

- Helium abundance

$$N_{\text{eff}} = 3.68^{+0.80}_{-0.70} (2\sigma)$$

Izotov, Thuan, 1001.4440

- WMAP+ACT+BAO

$$N_{\text{eff}} = 4.56 \pm 0.75 (68\%)$$

Dunkley et al., 1009.0866

- WMAP+SPT+BAO

$$N_{\text{eff}} = 3.86 \pm 0.42 (68\%)$$

Keiser et al., 1105.3182

- WMAP+ACT+SPT+BAO

$$N_{\text{eff}} = 4.08^{+0.71}_{-0.68} (95\%)$$

Archidiacono, Calabrese, Melchiorri, 1109.2767

# Dark radiation

$$\Delta N_{\text{eff}} \simeq 1 \longrightarrow \text{Dark radiation ?}$$

- Dark radiation (X) should satisfy :
  - X interaction is negligibly small
  - X is relativistic at the CMB epoch
- Many models are proposed so far...

Ichikawa, Kawasaki, KN, Senami, Takahashi (2007), KN, Takahashi, Yanagida (2010),  
Fischler, Meyers (2011), Kawasaki, Kitajima, KN (2011), Hasenkamp (2011)  
Menestrina, Scherrer (2011), Jeong, Takahashi (2012), K.Choi et al (2012) and many others

What is unique signature of dark radiation ?

# Inflationary GWs

- Inflation generates primordial GWs as quantum tensor fluctuations in de-Sitter spacetime

$$ds^2 = a^2(t) [-d\tau^2 + (\delta_{ij} + h_{ij}) dx^i dx^j]$$

$$h_{ij} = \frac{1}{M_P} \sum_{\lambda=+,-} \int \frac{d^3 k}{(2\pi)^{3/2}} h_k^\lambda(t) e^{i\mathbf{k}\cdot\mathbf{x}} e_{ij}^\lambda$$

Quantization

$$\langle h_k^\lambda h_{k'}^{\lambda'} \rangle = \frac{H_{\text{inf}}^2}{2k^3} \delta^3(k - k') \delta^{\lambda\lambda'}$$

- Dimensionless power spectrum almost scale invariant

$$\Delta_h^2(k) = \left( \frac{H_{\text{inf}}}{2\pi M_P} \right)^2$$

# Evolution of GW

- Eq.of.m of GW (without dark radiation)

$$\ddot{h}_\lambda + 3H\dot{h}_\lambda + (k/a)^2 h_\lambda = 0 \quad \rightarrow \quad \begin{aligned} h_\lambda &\sim \text{const} \quad \text{for } k \ll aH \\ h_\lambda &\propto a(t)^{-1} \quad \text{for } k \gg aH \end{aligned}$$

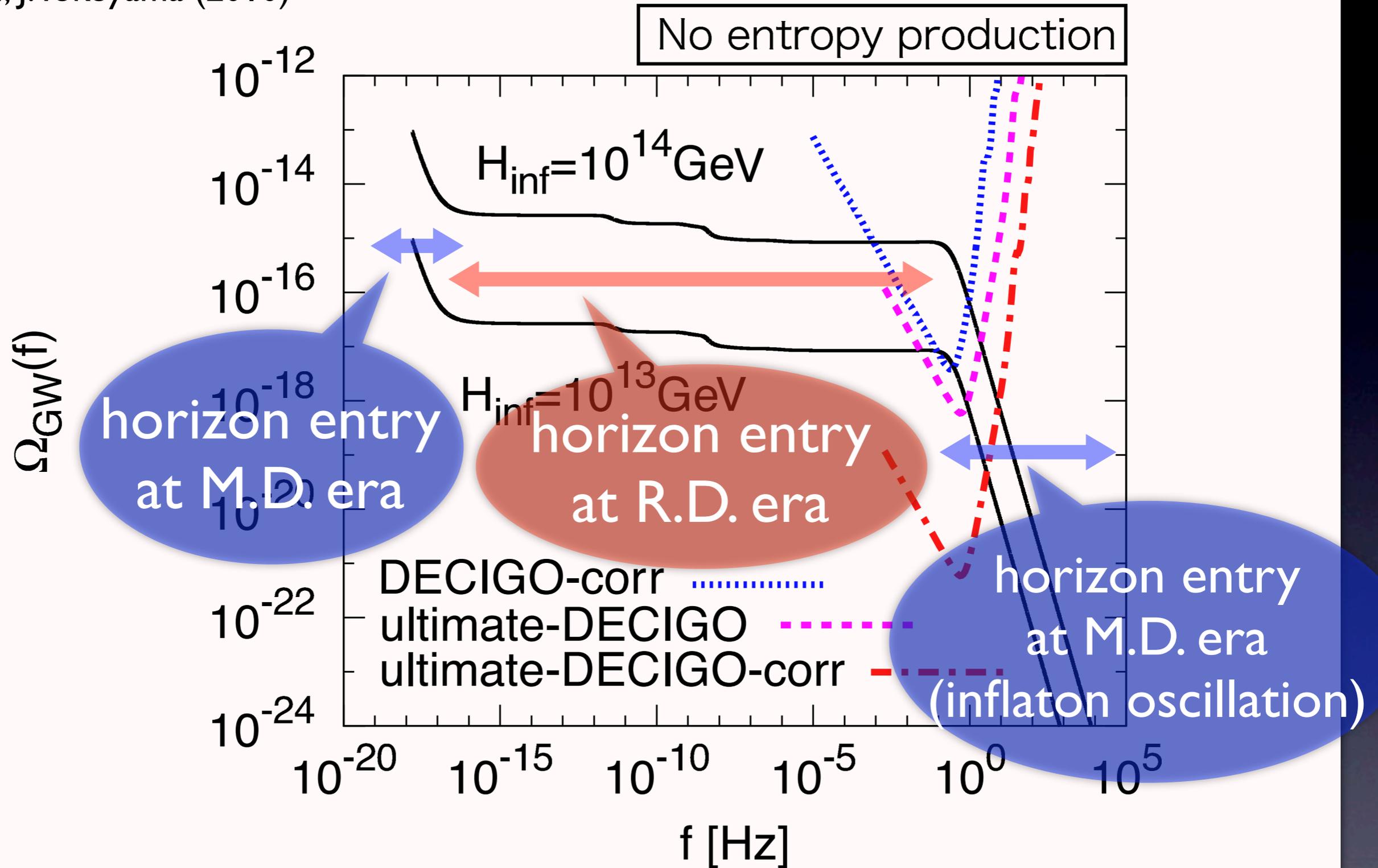
- GW energy density at horizon entry

$$\rho_{\text{GW}}(k) \sim M_P^2 \Delta_h^2(k) (k/a)^2 \sim M_P^2 H_{\text{in}}(k)^2 \Delta_h^2(k)$$

$$\rho_{\text{tot}} \sim M_P^2 H_{\text{in}}(k)^2$$

$$\rightarrow \quad \Omega_{\text{GW}}(k) = \frac{\rho_{\text{GW}}(k)}{\rho_{\text{tot}}} \sim \Delta_h^2(k) \sim \text{const at horizon entry}$$

$$\rightarrow \quad \Omega_{\text{GW}}^0(k) \sim \Omega_{\text{rad}}^0 \Delta_h^2(k) \quad \text{at present for } k \gg k_{\text{eq}}$$



GW spectrum traces thermal history of the Universe !

N.Seto, J.Yokoyama (2003), Boyle, Steinhardt (2005), KN, Saito, Suwa, Yokoyama (2008)

# Dark radiation and GW

- Dark radiation affects GW spectrum in **two** ways

$$\ddot{h}_{ij} + 3H\dot{h}_{ij} + (k/a)^2 h_{ij} = 16\pi G\Pi_{ij}$$

Modified expansion rate

Anisotropic stress of X

cf ) For standard neutinos, see

S.Weinberg (2003), Y.Watanabe, E.Komatsu (2005)

- **Modified expansion rate** by parent field of X
  - Modification on GW spectrum at **high** frequency
- **Anisotropic stress** is turned on after X production
  - Modification on GW spectrum at **low** frequency

# A model

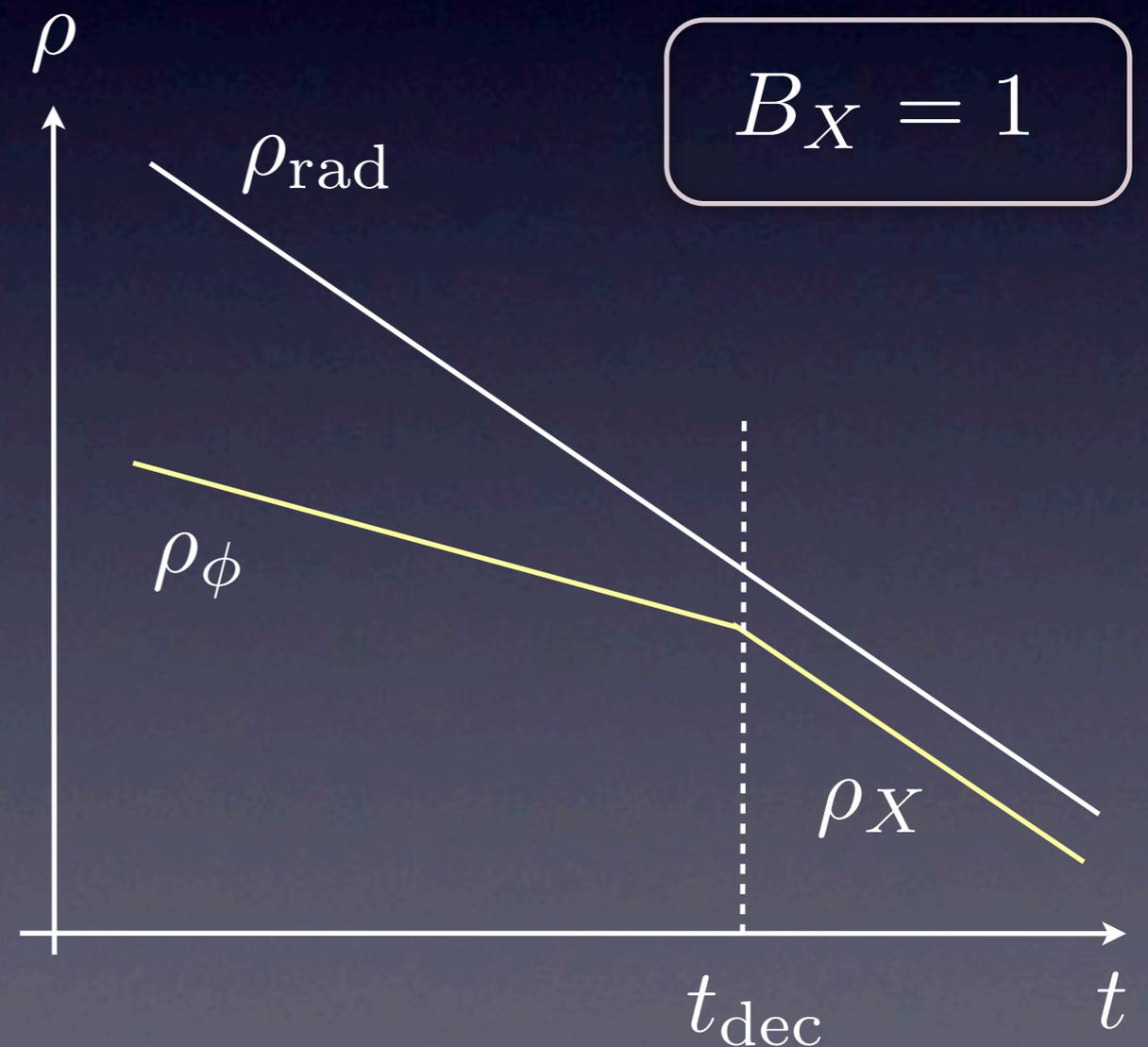
- A scalar field  $\phi$  decays into  $X$  at  $H \sim \Gamma_\phi$  with branching ratio  $B_X$

- Background evolution :

$$\begin{aligned}\dot{\rho}_\phi + 3H\rho_\phi &= -\Gamma_\phi\rho_\phi, \\ \dot{\rho}_{\text{rad}} + 4H\rho_{\text{rad}} &= \Gamma_\phi(1 - B_X)\rho_\phi, \\ \dot{\rho}_X + 4H\rho_X &= \Gamma_\phi B_X\rho_\phi,\end{aligned}$$

- $\phi$  nearly dominate at decay for  $\Delta N_{\text{eff}} \simeq 1$

- Example)  $\phi$  : saxion  
 $X$  : axion



# A model

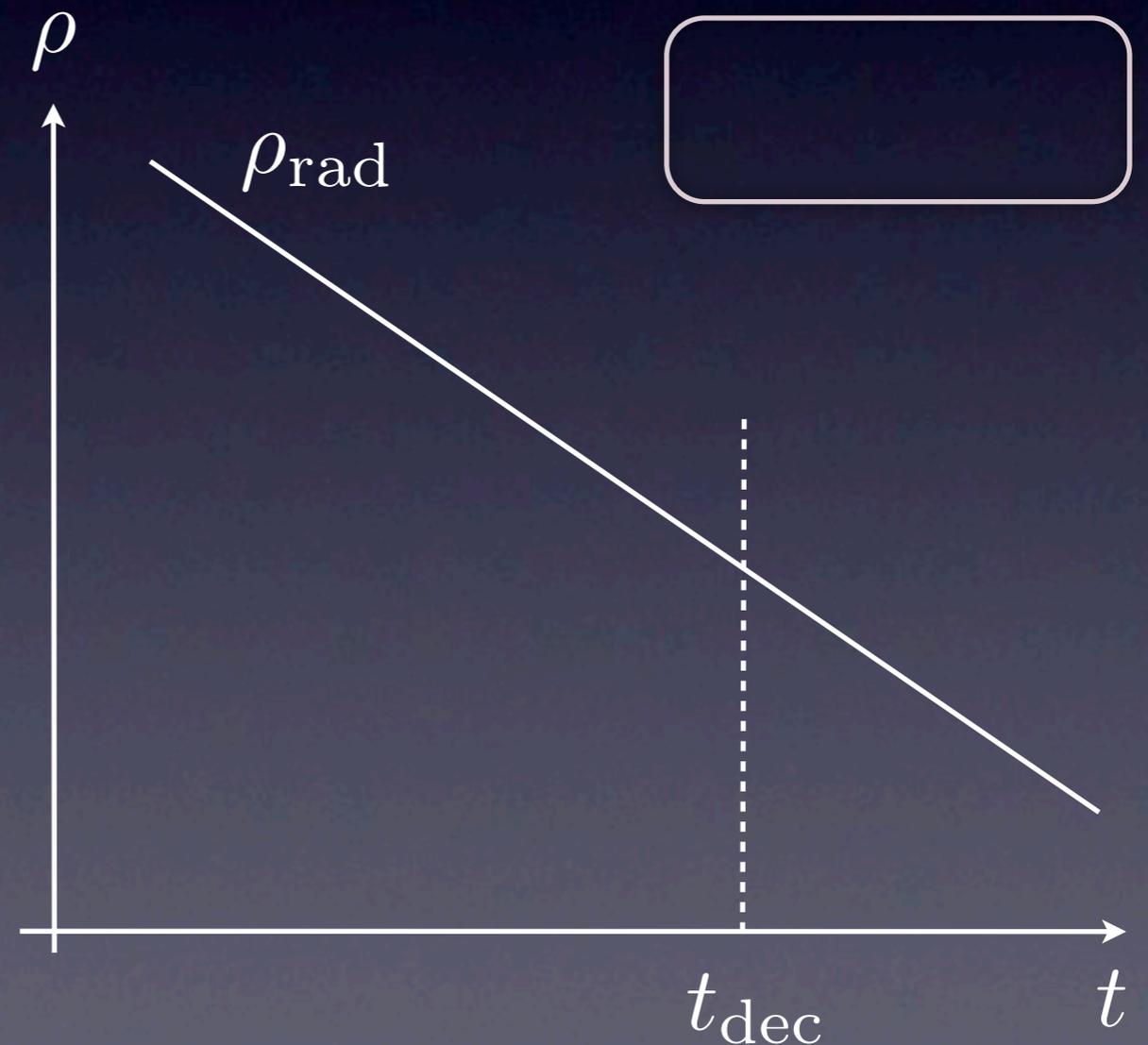
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# A model

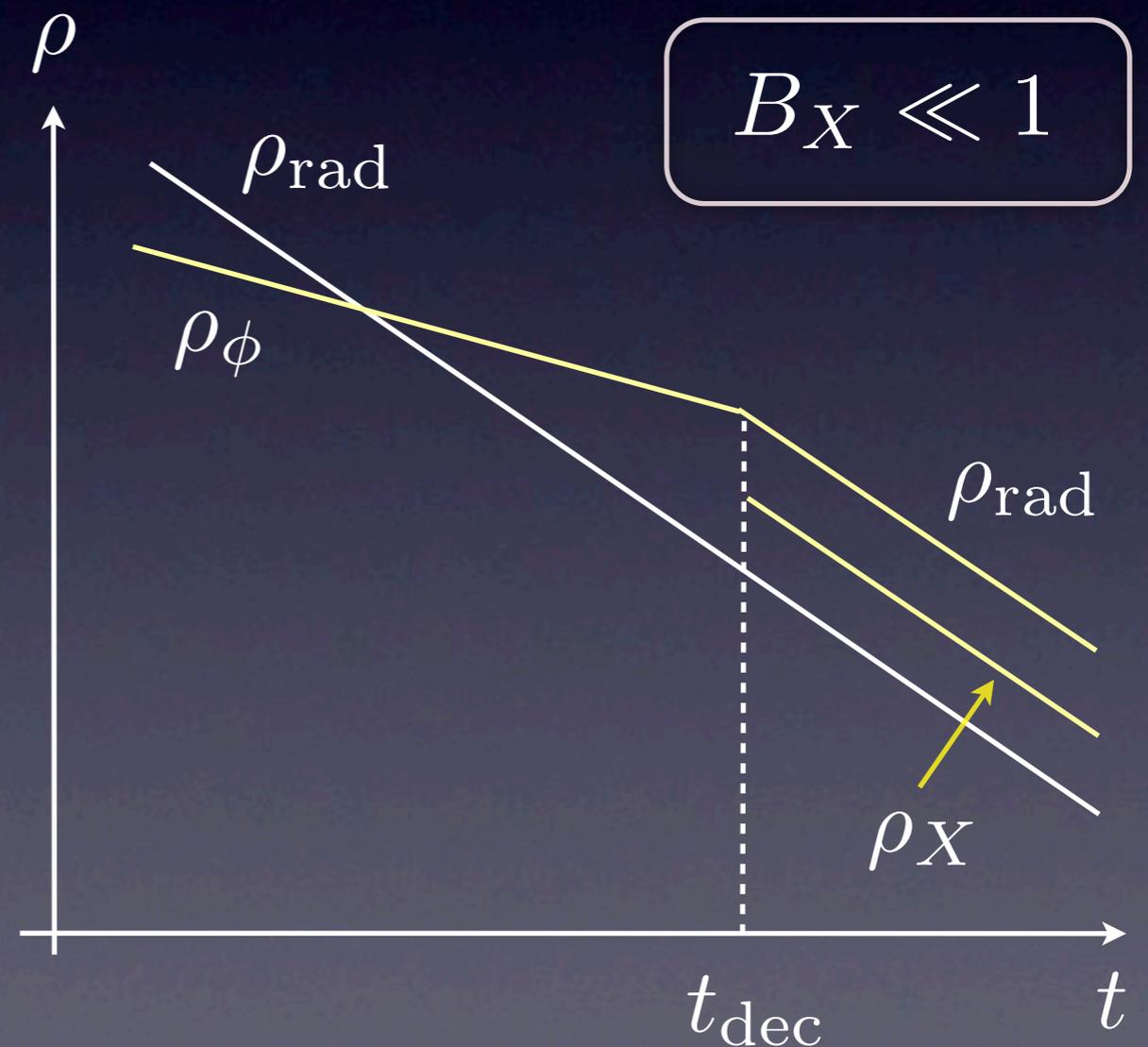
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- Background evolution :

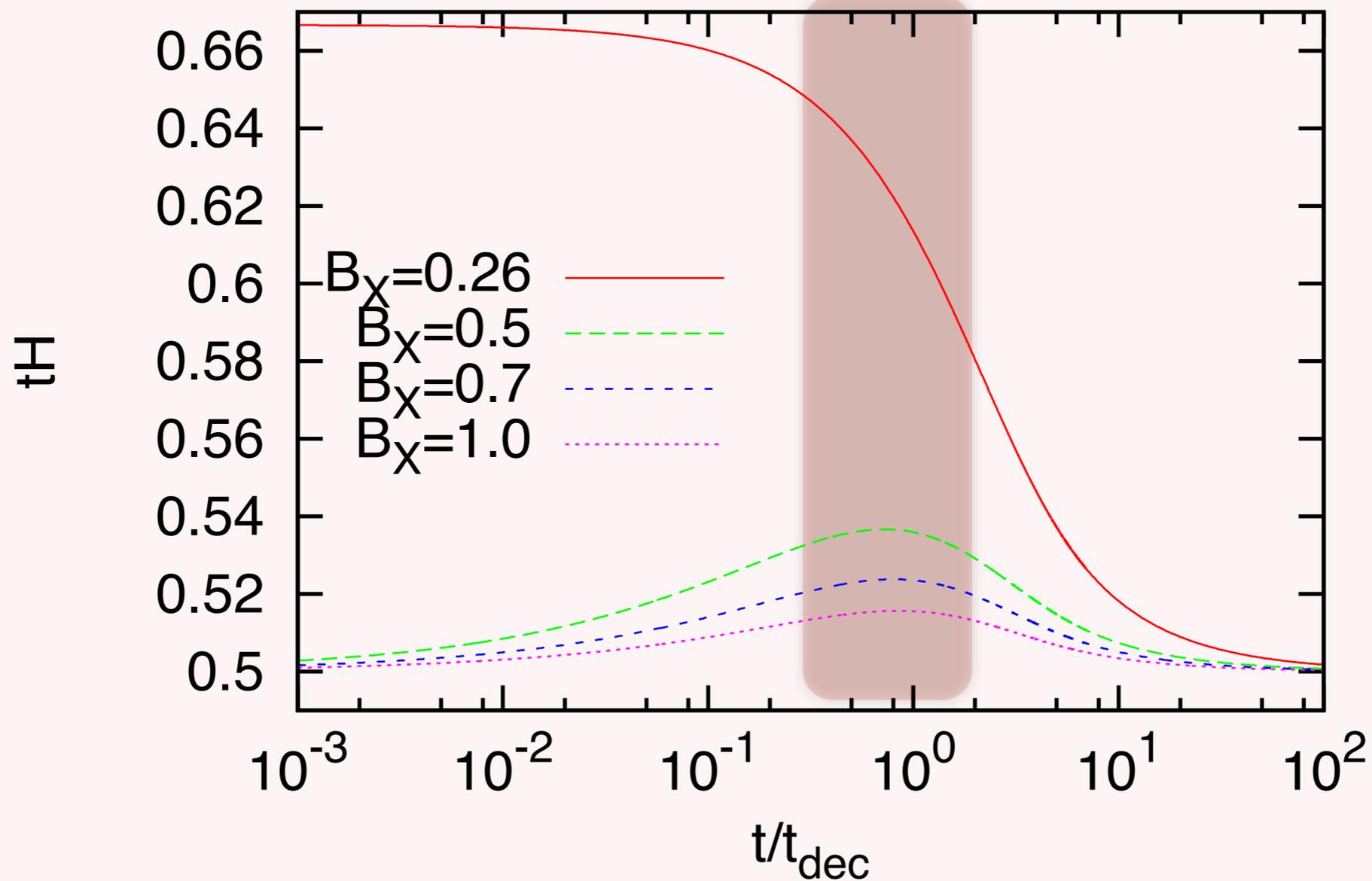
$$\begin{aligned}\dot{\rho}_\phi + 3H\rho_\phi &= -\Gamma_\phi\rho_\phi, \\ \dot{\rho}_{\text{rad}} + 4H\rho_{\text{rad}} &= \Gamma_\phi(1 - B_X)\rho_\phi, \\ \dot{\rho}_X + 4H\rho_X &= \Gamma_\phi B_X\rho_\phi,\end{aligned}$$

- $\phi$  nearly dominate at decay for  $\Delta N_{\text{eff}} \simeq 1$

- Example)  $\phi$  : saxion  
 $X$  : axion



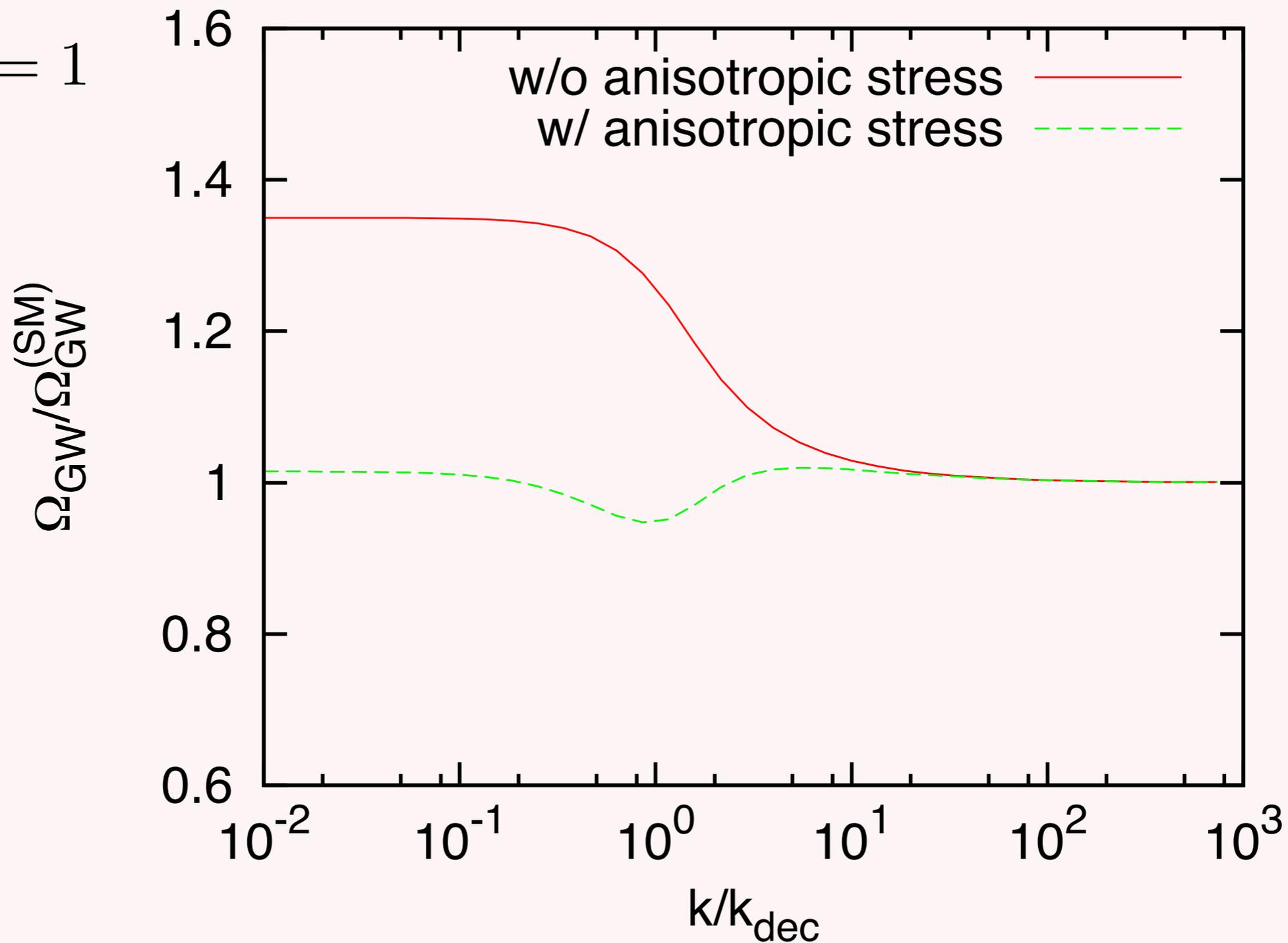
# A model



Deviation from R.D.,  $tH=0.5$ , around  $\Phi$  decay

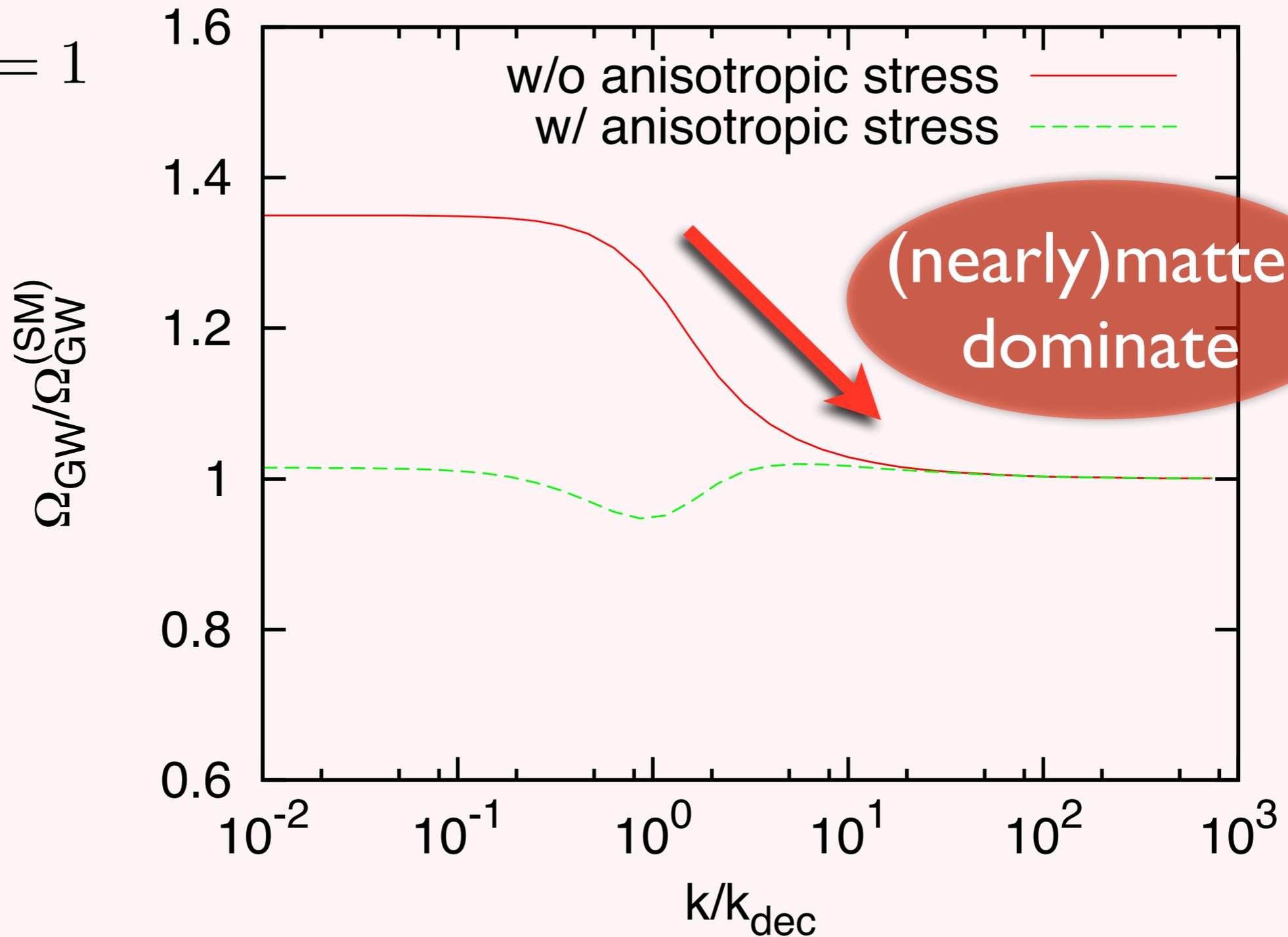
# Numerical result

$$B_X = 1$$



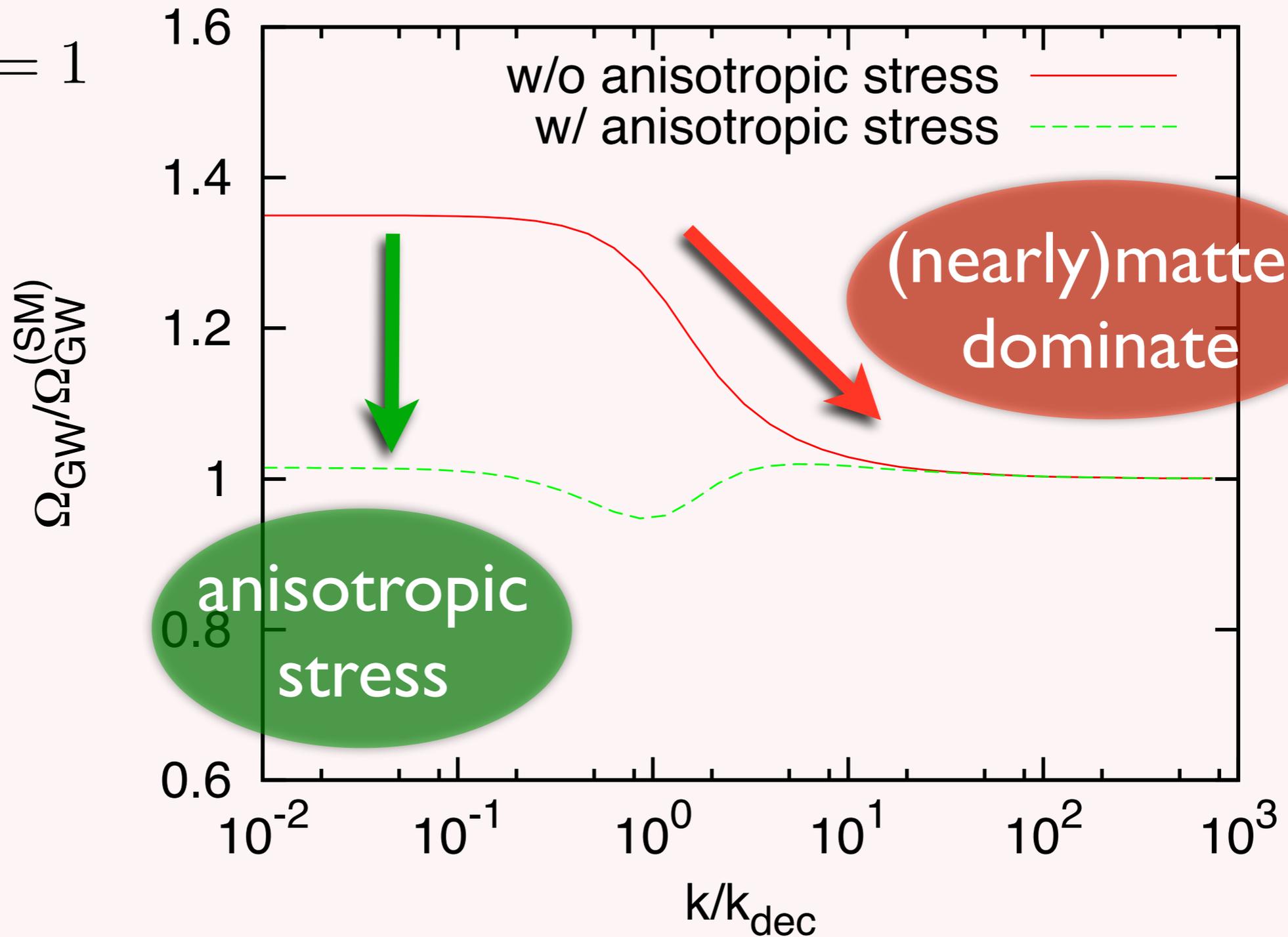
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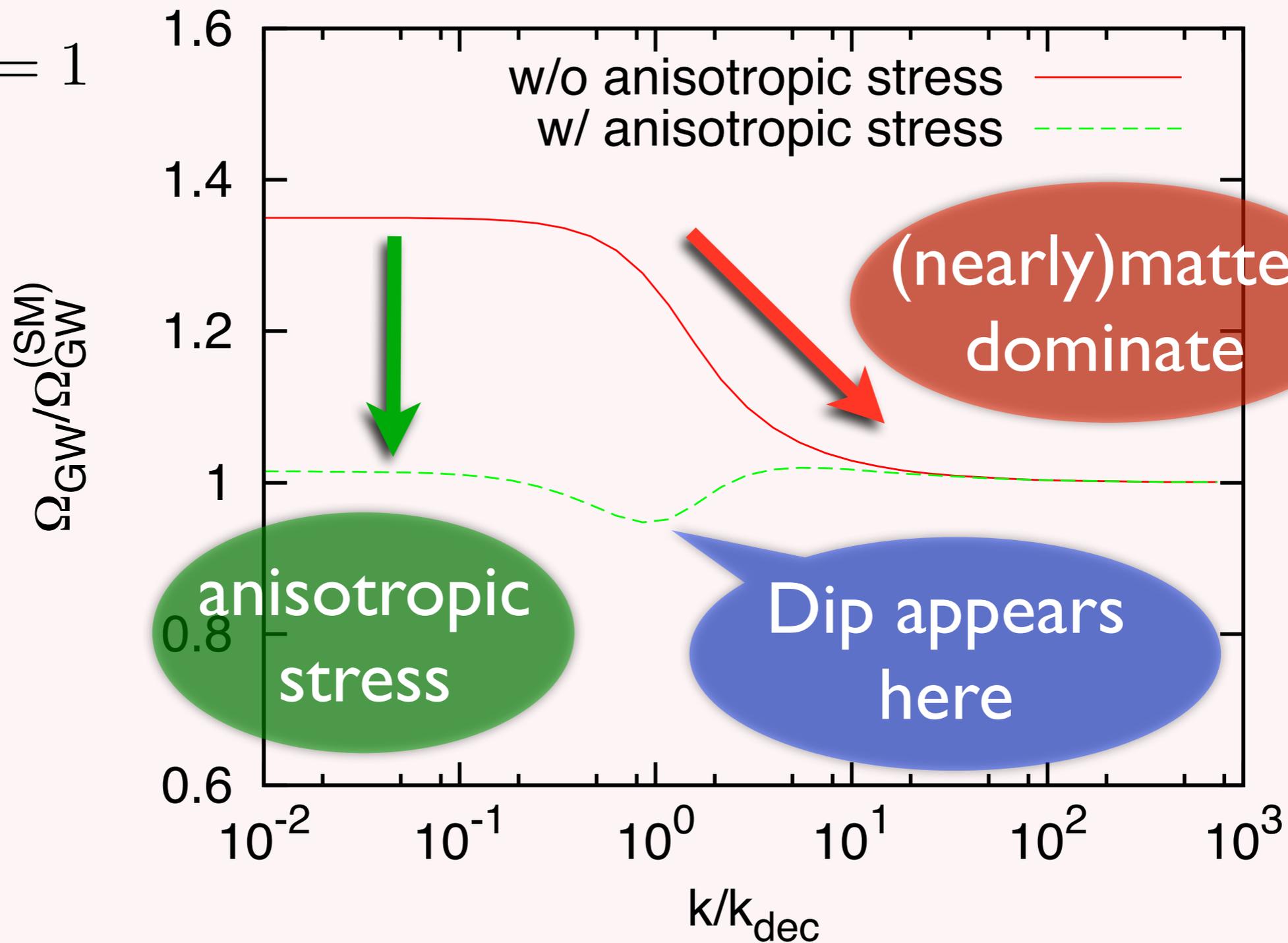
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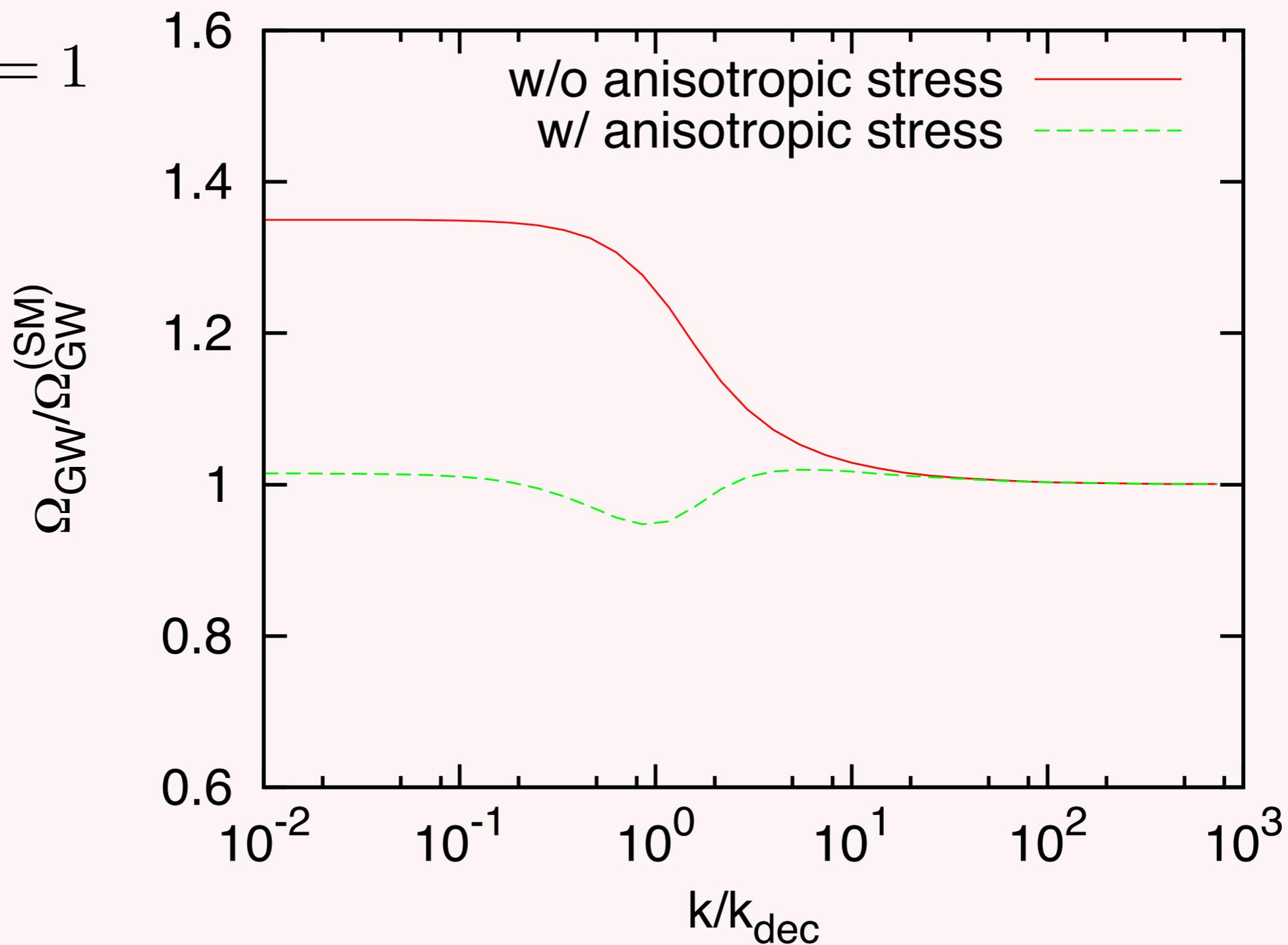
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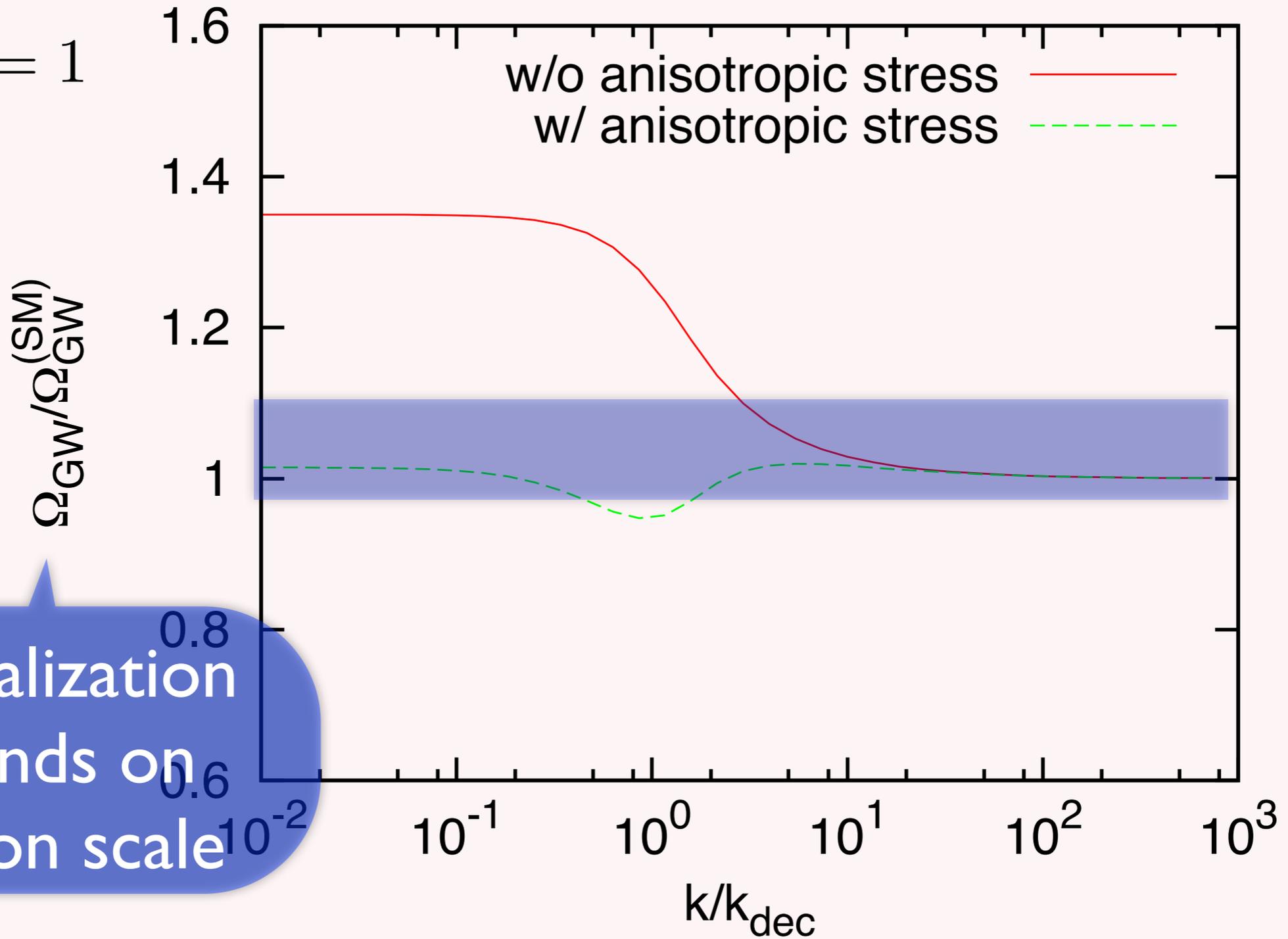
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# Numerical result

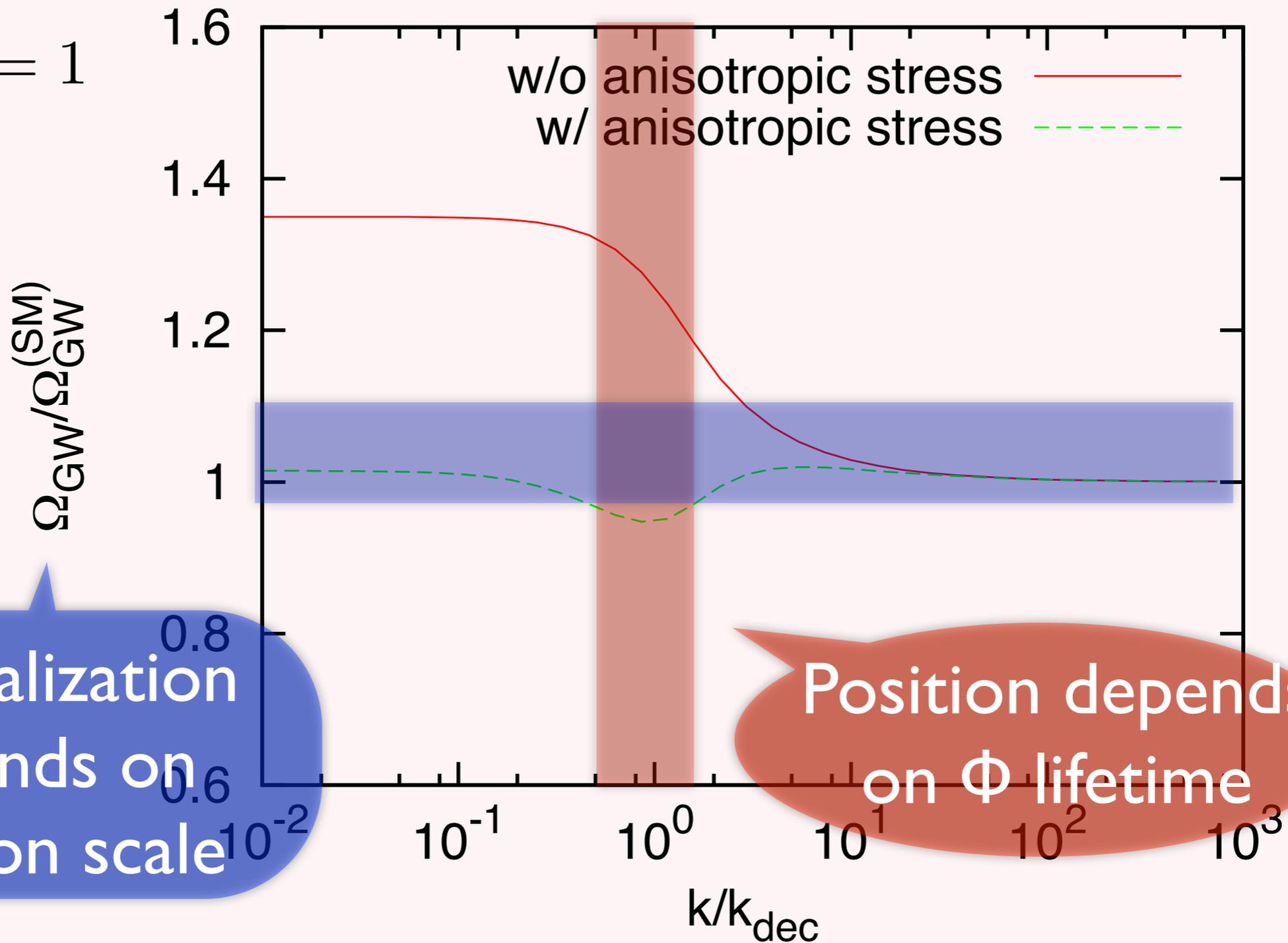
$$B_X = 1$$



Normalization depends on inflation scale

# Numerical result

$$B_X = 1$$

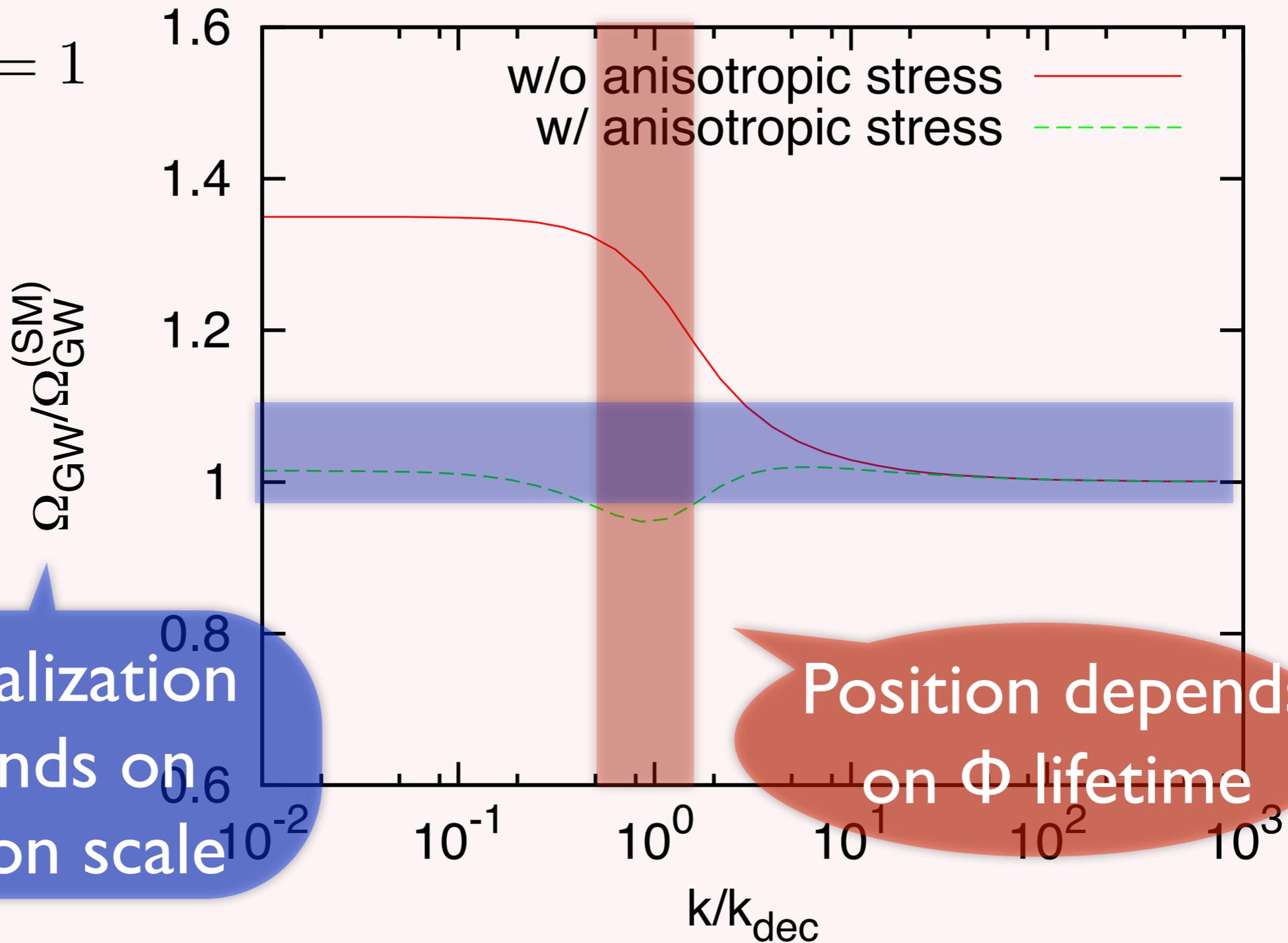


Normalization depends on inflation scale

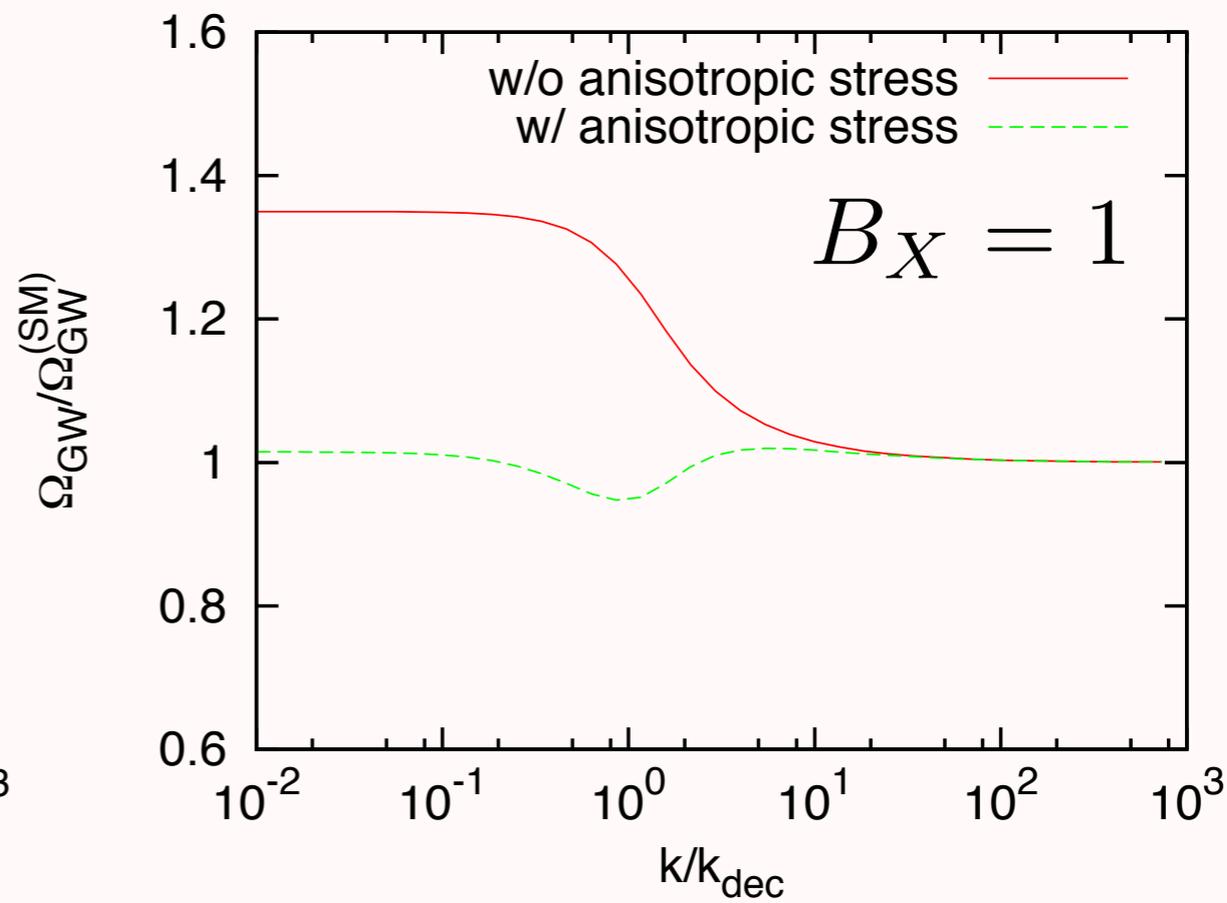
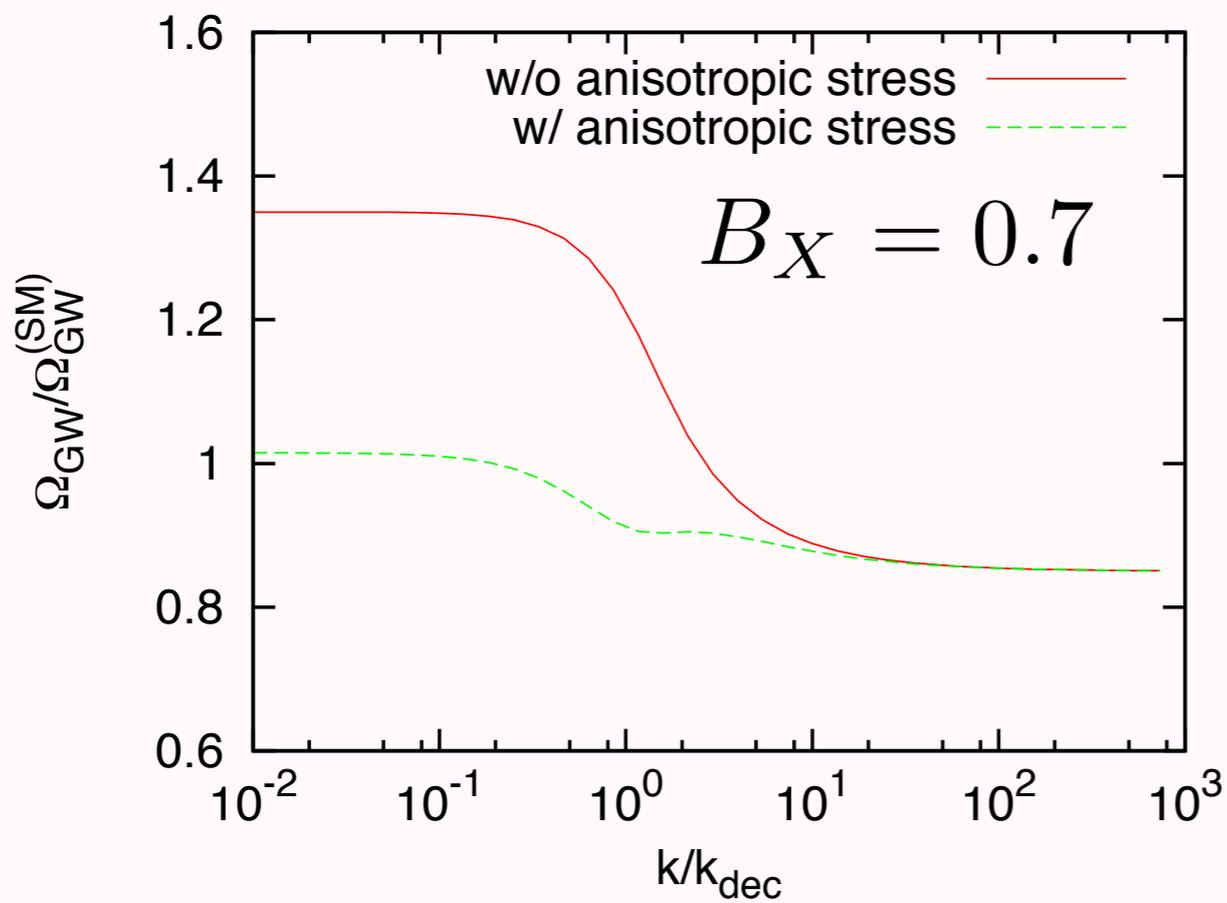
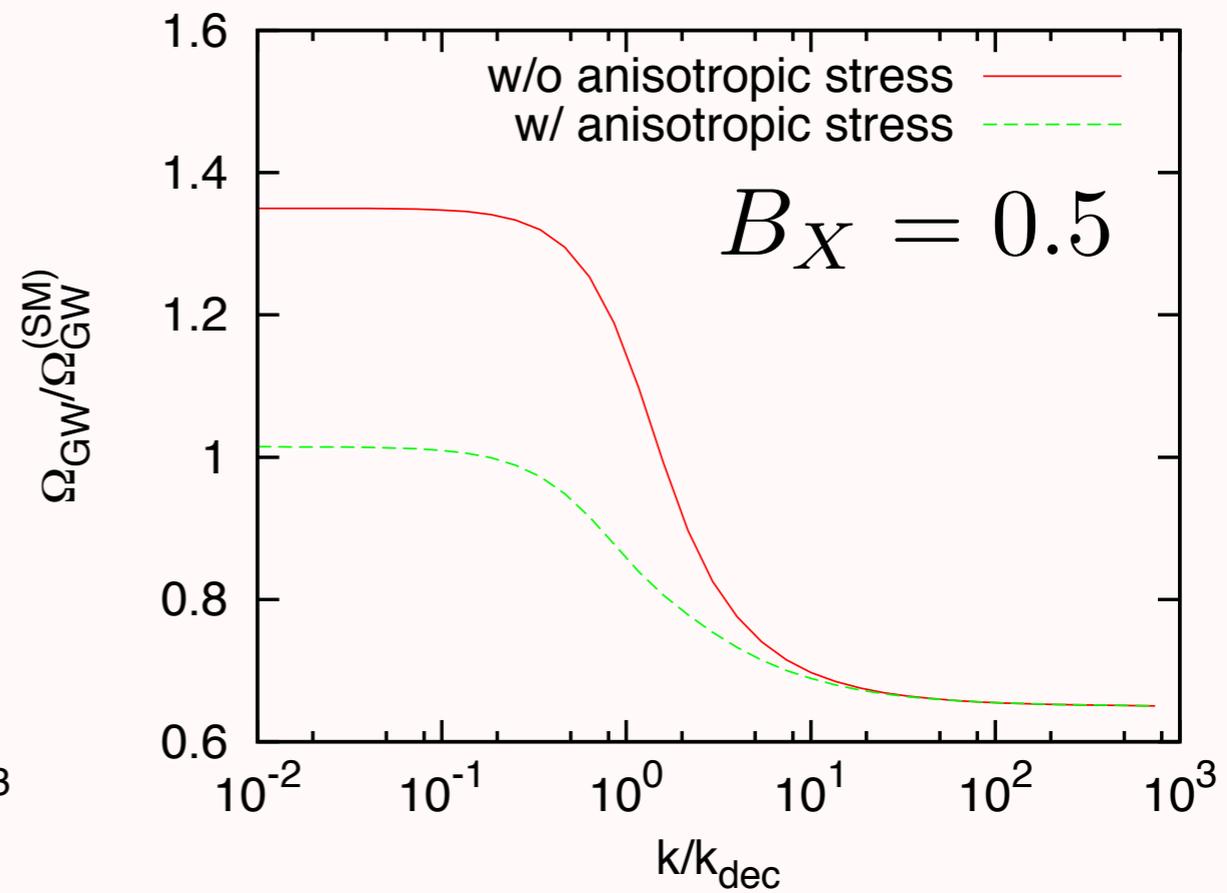
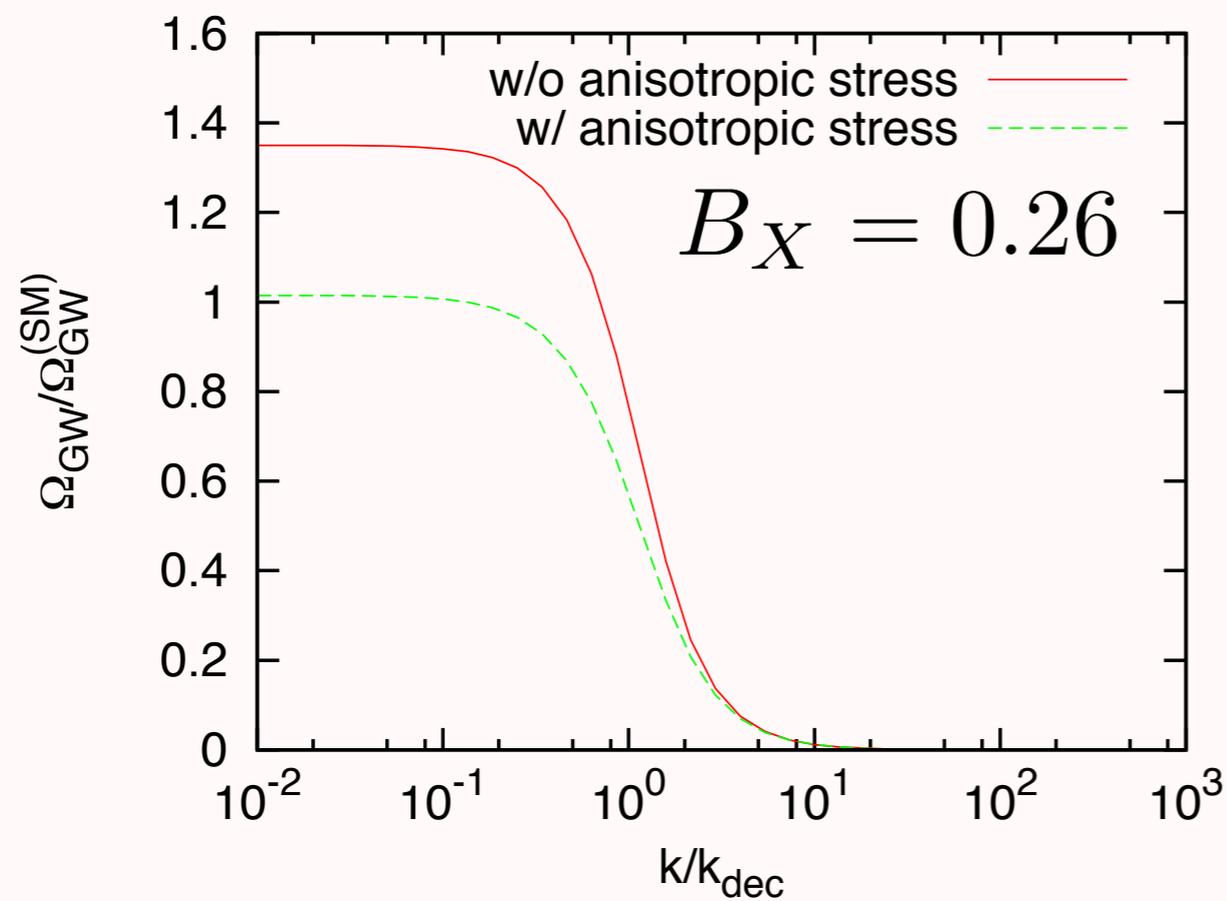
Position depends on  $\Phi$  lifetime

# Numerical result

$$B_X = 1$$



Detectable at DECIGO for  $r \gtrsim 10^{-3}$  and  $T_\phi \sim 10^7 \text{ GeV}$



# Summary

- Recent observation suggest extra light species : **dark radiation**
- **Dark radiation** leaves characteristic signature in primordial GW spectrum
- It also contains information on the production mechanism of dark radiation.

# Backup Slides

# GW normalization

- Standard model
  - GW spectrum at horizon entry

$$\Omega_{\text{GW}}(k = aH) = \frac{\Delta_h^2(k)}{24} \quad \Delta_h^2(k) \equiv \frac{8}{M_P^2} \left( \frac{H_{\text{inf}}}{2\pi} \right)^2 \left( \frac{k}{k_0} \right)^{n_t}$$

- GW spectrum at present ( $k \gg k_{\text{eq}}$ )

$$\Omega_{\text{GW}}^{(\text{SM})}(k) = \gamma^{(\text{SM})} \Omega_{\text{rad}}^{(\text{SM})} \times \Omega_{\text{GW}}(k = aH),$$

Expansion history :

$$\gamma^{(\text{SM})} = \left[ \frac{g_*(T_{\text{in}}(k))}{g_{*0}^{(\text{SM})}} \right] \left[ \frac{g_{*s0}^{(\text{SM})}}{g_{*s}(T_{\text{in}}(k))} \right]^{4/3},$$

# GW normalization

- Standard model plus dark radiation

- GW spectrum at present ( $k \gg k_{\text{eq}}$ )

$$\Omega_{\text{GW}}(k) = \gamma \Omega_{\text{rad}} \times \Omega_{\text{GW}}(k = aH),$$

Expansion history  
modified by X :

$$\gamma = \frac{1 + \frac{7}{43} \left( \frac{g_{*s}(T_\phi)}{10.75} \right)^{1/3} \Delta N_{\text{eff}}}{1/\gamma^{(\text{SM})} + \frac{7}{43} \left( \frac{g_{*s}(T_\phi)}{10.75} \right)^{1/3} \Delta N_{\text{eff}}},$$

Radiation  
density :

$$\Omega_{\text{rad}} = \Omega_{\text{rad}}^{(\text{SM})} \times (g_{*0}/g_{*0}^{(\text{SM})}) \quad g_{*0} = 2 \left[ 1 + N_{\text{eff}} \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} \right]$$

Overall normalization is affected

# GW normalization

- Parameterize normalization

$$\frac{\Omega_{\text{GW}}(k)}{\Omega_{\text{GW}}^{(\text{SM})}(k)} = C_1 \times C_2$$

Modified BG by X :

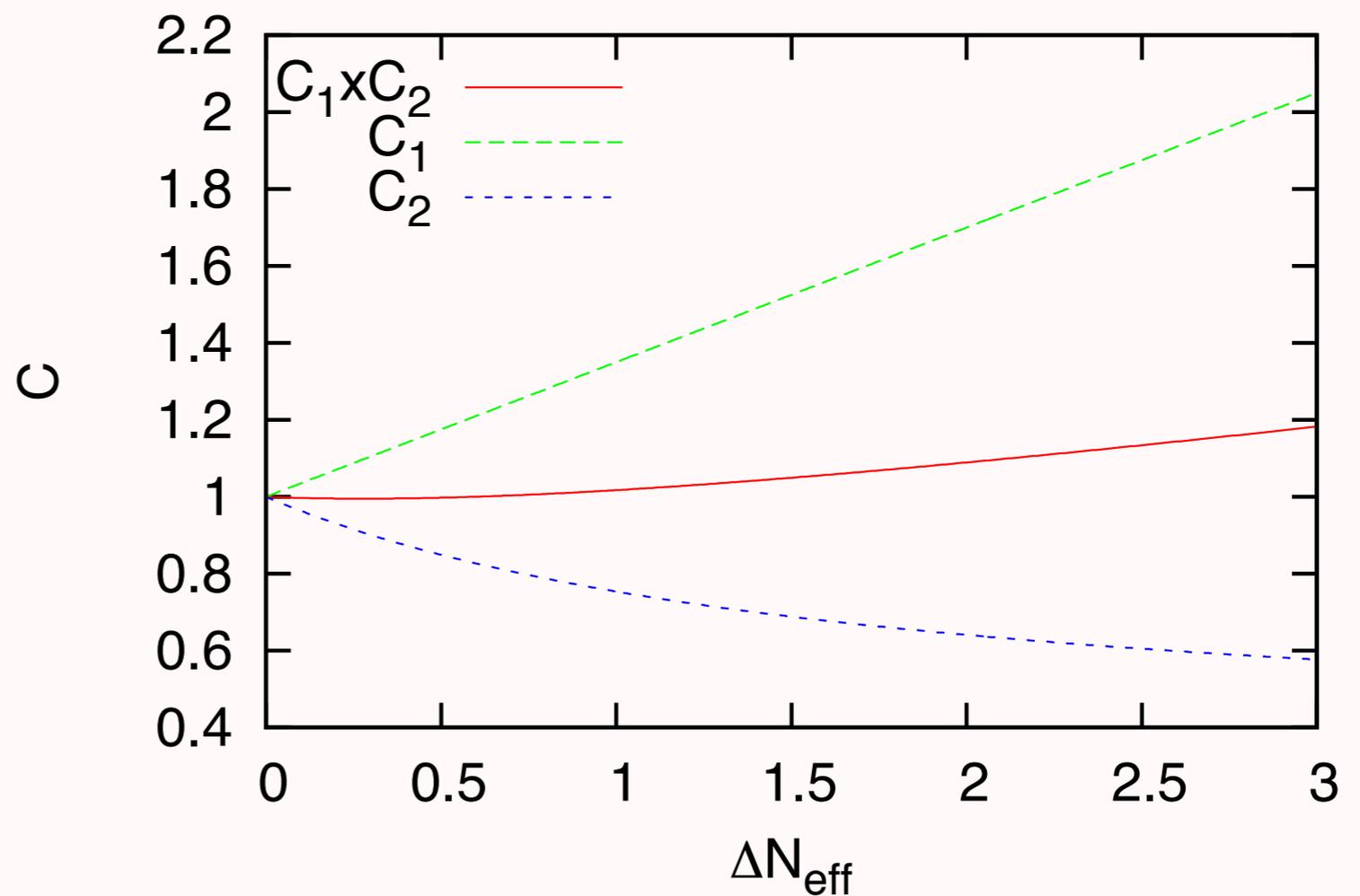
$$C_1 \equiv \frac{\gamma}{\gamma^{(\text{SM})}} \frac{g_{*0}}{g_{*0}^{(\text{SM})}}$$

Anisotropic stress X :

$C_2$

analytically  
derived in

Dicus, Repko (2004)



$C_1 \times C_2$  accidentally close to unity

# Anisotropic stress

- Boltzmann eq. for  $X$

$$\frac{dF}{dt} = \frac{B_X}{4\pi(p^0)^3} \Gamma_{\phi\rho\phi} \delta \left( p^0 - \frac{m_\phi}{2} \right)$$

$F$  : distribution function of  $X$

$$\frac{dF}{dt} = \frac{\partial F}{\partial t} + \frac{p^i}{p^0} \frac{\partial F}{\partial x^i} + \frac{1}{2} g_{ij,k} \frac{p^i p^j}{p^0} \frac{\partial F}{\partial p_k}$$

cf) Geodesic eq.

$$\frac{dp_i}{dt} = \frac{1}{2} g_{jk,i} \frac{p^j p^k}{p^0}$$

GW effect here

- Perturbed :

$$\frac{\partial(\delta F_1 + \delta F_2)}{\partial t} + \frac{\bar{p}^i}{\bar{p}^0} \frac{\partial(\delta F_1 + \delta F_2)}{\partial x^i} + \frac{1}{2} (\delta g_{jk})_{,i} \frac{\bar{p}^j \bar{p}^k}{\bar{p}^0} \frac{\partial \bar{F}}{\partial p_i} = a \frac{\partial^2 \bar{F}}{\partial p \partial t} \delta p^0$$

$$\delta F_1(t, x^i, p_i) \equiv \bar{F}(t, (g^{ij} p_i p_j)^{1/2} / a) - \bar{F}(t, p)$$

$$\delta F_2(t, x^i, p_i) \equiv F - \bar{F} - \delta F_1$$



$$\frac{\partial \delta F_2}{\partial t} + \frac{\hat{p}_i}{a} \frac{\partial \delta F_2}{\partial x^i} = \frac{1}{2} \frac{\partial h_{ij}}{\partial t} \frac{\partial \bar{F}}{\partial p} p \hat{p}_i \hat{p}_j$$

Contributes to  
anisotropic stress

# Anisotropic stress

- EM tensor of X

$$\delta T_{ij}^{(X)} = \frac{1}{a^3} \int d^3p \left[ (\delta F_1 + \delta F_2) \frac{p_i p_j}{\bar{p}^0} + \bar{F} p_i p_j \delta \left( \frac{1}{p^0} \right) \right] = \frac{1}{a^2} \int d^3p \delta F_2 p \hat{p}_i \hat{p}_j + \frac{1}{3} a^2 h_{ij} \rho_X$$

Anisotropic stress  $a^2 \Pi_{ij}$

- From Boltzmann eq :

$$\delta F_2 = \int_0^\tau d\tau' \frac{1}{2} \frac{\partial h_{ij}}{\partial \tau}(\tau') \frac{\partial \bar{F}}{\partial p}(\tau') p \hat{p}_i \hat{p}_j e^{-ik\mu(\tau-\tau')}$$

- Eq.of.m of GW (with dark radiation)

$$\ddot{h}^{(\lambda)} + 3H\dot{h}^{(\lambda)} + \frac{k^2}{a^2} h^{(\lambda)} = -24H^2 \frac{1}{a^4(t)\rho_{\text{tot}}(t)} \times \int_0^t a^4(t') \rho_X(t') K \left( k \int_{t'}^t \frac{dt''}{a(t'')} \right) \dot{h}^{(\lambda)}(t', \mathbf{k}) dt',$$

Anisotropic stress of X induced by GWs