

# Non-Gaussianity in the CMB from nonlinear effects

Cyril Pitrou

Institute of Cosmology and Gravitation, Portsmouth

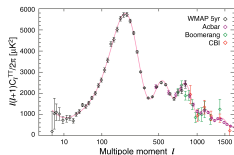
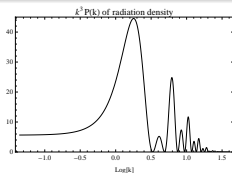
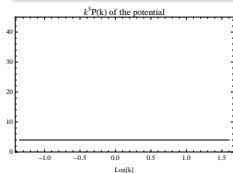
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## Standard lore of perturbation theory

- *Initial conditions*: quantization of the free theory implies Gaussian initial conditions:  $P(k)$   
 $\langle \Phi(\mathbf{k})\Phi(\mathbf{k}') \rangle = \delta(\mathbf{k} + \mathbf{k}')P(k)$
- *Evolution*: linearisation of GR.

## Transfer scheme of perturbations

- Linear equations, modes  $k$  are independent,  
 $\implies$  Gaussianity conserved.
- $P(k) \rightarrow \Theta(k, \eta) \rightarrow a_{\ell m} \rightarrow C_\ell$



# Non-Gaussianity

## non-Gaussianity (NG)

- Initial conditions non-Gaussian?  
We want to test the models of inflation with other moments of the statistics
- Non-linear dynamics is intrinsic to GR

## Statistics of the primordial gravitational potential $\Phi = \Phi^{(1)} + \frac{1}{2}\Phi^{(2)}$

- Gaussian part  $\Phi^{(1)}$  and non-Gaussian part  $\Phi^{(2)}$ :
- $\langle \Phi(\mathbf{k})\Phi(\mathbf{k}') \rangle = \delta(\mathbf{k} + \mathbf{k}')P(k)$
- $\langle \Phi(\mathbf{k}_1)\Phi(\mathbf{k}_2)\Phi(\mathbf{k}_3) \rangle = \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3)f_{\text{NL}}F(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3)$
- $F(\dots)$  = type of non-Gaussianity
- $f_{\text{NL}}$  = its amplitude.

## The transfer to temperature fluctuations $\Theta_{\ell m}$

- In general  $\Theta \equiv \mathcal{T}(\Phi)$
- Order 1  $\Theta_{\ell m}^{(1)} \equiv \mathcal{T}_L^{\ell m}(\Phi^{(1)})$
- Order 2  $\Theta_{\ell m}^{(2)} \equiv \mathcal{T}_L^{\ell m}(\Phi^{(2)}) + \mathcal{T}_{NL}^{\ell m}(\Phi^{(1)}\Phi^{(1)})$

## In Fourier space

- $\Theta_{\ell m}^{(1)}(\mathbf{k}) = \mathcal{T}_L^{\ell m}(k)\Phi_{\mathbf{k}}^{(1)}$
- $\Theta_{\ell m}^{(2)}(\mathbf{k}) = \mathcal{T}_L^{\ell m}(k)\Phi_{\mathbf{k}}^{(2)} + \int d^3\mathbf{k}_1 d^3\mathbf{k}_2 \delta^3(\mathbf{k} - \mathbf{k}_1 - \mathbf{k}_2) \mathcal{T}_{NL}^{\ell m}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}) \Phi_{\mathbf{k}_1}^{(1)} \Phi_{\mathbf{k}_2}^{(1)}$

## $f_{NL}$ or $\mathcal{T}_{NL}$ ?

$\langle \Theta_{\ell_1 m_1} \Theta_{\ell_2 m_2} \Theta_{\ell_3 m_3} \rangle \neq 0$  because of  $\langle \Theta^{(1)} \Theta^{(1)} \Theta^{(2)} \rangle$ .

# Evolution of the distribution function of radiation

## Boltzmann equation

$$L[f] = C[f]$$

- Liouville operator: Free-fall

$$L[f] = \frac{df}{ds} = p^c \nabla_c f(x, p^a) + \frac{\partial f(x, p^a)}{\partial p^c} \frac{dp^c}{ds}$$

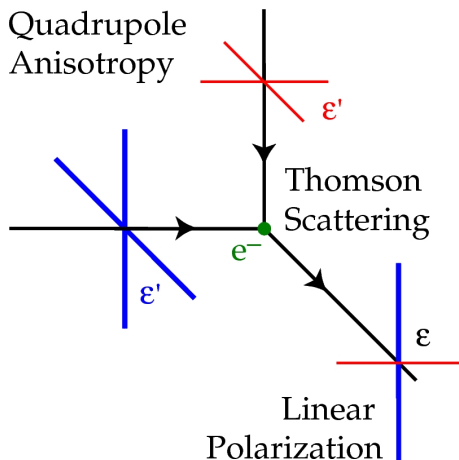
## Geodesic equation

$$p^b \nabla_b p_a = \frac{dp_a}{ds} + \omega_{bac} p^c p^b = 0$$

- Collision operator: Compton scattering on free electrons.

# Why do we also need to describe polarization?

Because if radiation has a quadrupole,  
Compton scattering generates polarisation.



# Description of polarisation by the Stokes parameters

## Tensorial distribution function

$$\text{If } n^i = (0, 0, 1): f_{ab} = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & I + Q & U + iV & 0 \\ 0 & U - iV & I - Q & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

## Covariant expression

$$f_{\mu\nu}(x, p^a) \equiv \frac{1}{2} I(x, p^a) S_{\mu\nu} + P_{\mu\nu}(x, p^a) + \frac{i}{2} V(x, p^a) e_{0\sigma}^{\rho} \epsilon_{\rho\mu\nu\sigma} n^{\sigma}$$

# Multipolar expansion

Multipoles for scalar functions ( $I$  and  $V$ )

$$I(x, p^a, n^a) = \sum_{\ell=0}^{\infty} I_{a_\ell}(x, p^a) n^{a_\ell}$$

And for polarisation, E and B modes...

$$P_{ab}(x, p^a) = \sum_{\ell=2}^{\infty} \left[ E_{abc_{\ell-2}}(x, p^a) n^{c_{\ell-2}} - n_c \epsilon^{cd} ({}_a B_b)_{dc_{\ell-2}}(x, p^a) n^{c_{\ell-2}} \right]^{TT}$$



# Steps to follow

- 1 Perturb the metric  $g_{\mu\nu} = \bar{g}_{\mu\nu} + g_{\mu\nu}^{(1)} + \frac{1}{2}g_{\mu\nu}^{(2)}$
- 2 Perturb the tetrad  $e_a^\mu = \bar{e}_a^\mu + e_a^{(1)\mu} + \frac{1}{2}e_a^{(2)\mu}$
- 3 Perturb the connections  $\omega_{abc} = \bar{\omega}_{abc} + \omega_{abc}^{(1)} + \frac{1}{2}\omega_{abc}^{(2)}$
- 4 Find the perturbed geodesic equations
- 5 Compute the perturbed Liouville operator
- 6 Compute the Thomson scattering for each electron
- 7 Sum over the electrons distribution to obtain the Collision tensor in full generalities
- 8 Expand it in perturbations
- 9 Take the multipoles  $I_{\underline{a}_\ell}$ ,  $E_{\underline{a}_\ell}$  and  $B_{\underline{a}_\ell}$  of the Boltzmann equation
- 10 Solve it or integrate it numerically

## Evolution of brightness $\mathcal{I} = T^4$ along geodesics

$$\frac{d [e^{-\bar{\tau}} \mathcal{I} E^{-4}]}{d\eta} = \bar{g}(\eta) E^{-4} [e^{\Phi} \mathcal{C}[\mathcal{I}] + \bar{\tau}' \mathcal{I}]$$

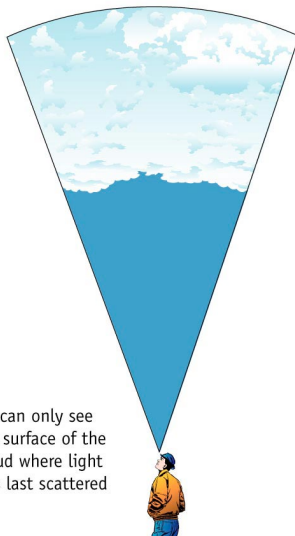
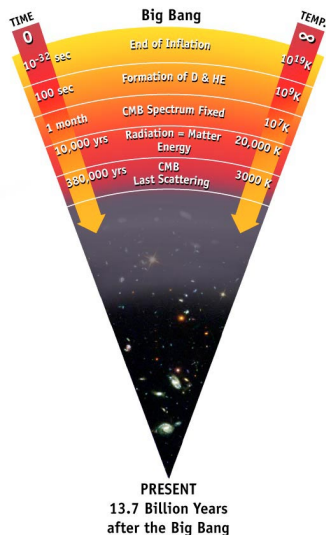
$$\frac{d}{d\eta} = \frac{\partial}{\partial \eta} + \frac{dx^l}{d\eta} \frac{\partial}{\partial x^l} + \frac{dn^i}{d\eta} \frac{\partial}{\partial n^i},$$

$$\frac{d \ln E}{d\eta} \simeq -\frac{d\Phi}{d\eta} + \Phi' + \Psi'$$

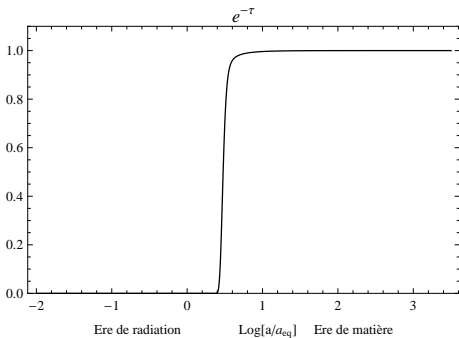
## Classification of effects:

- 1)  $dE/d\eta$ : Evolution of the energy of photons:  
*Einstein effect* (potential  $\Phi$ ) and *integrated effects*
- 2)  $\bar{g}(\eta) \dots$  Collisions on the last scattering surface (LSS):  
*Intrinsic temperature*, and *Doppler effect*.
- 3) *Lensing effect*  $\delta \left( \frac{dn^i}{d\eta} \right)$  **Order 2**
- 4) *Time-delay*  $\delta \left( \frac{dx^l}{d\eta} \right)$  **Order 2**

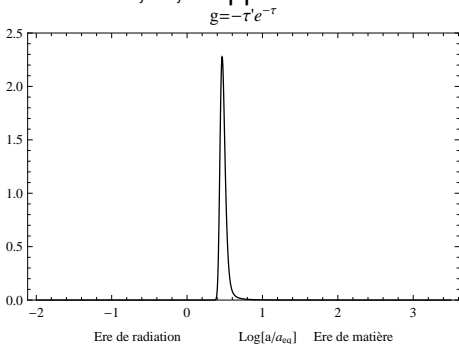
# Geometry of the background space-time



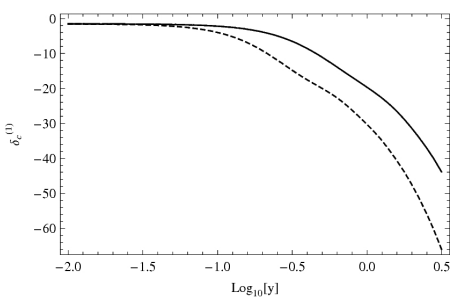
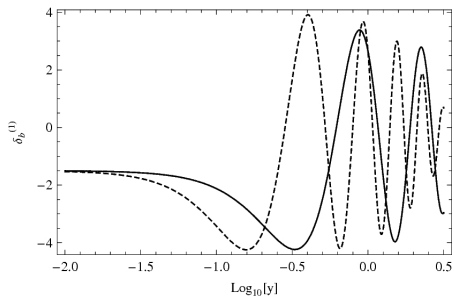
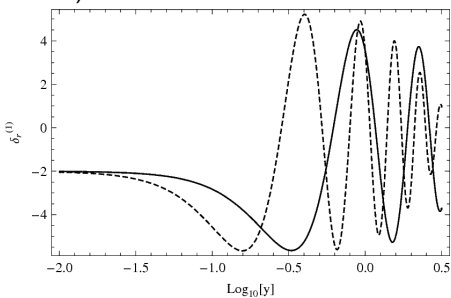
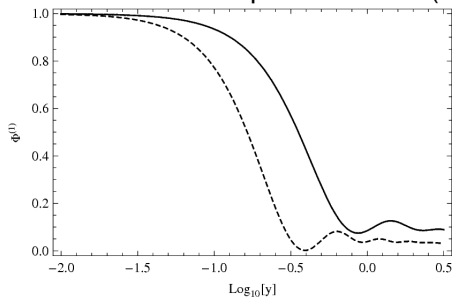
Integrated effects efficient since recombination



In the width of the LSS  
Intrinsic  $\Theta$ ,  $\Phi$ , Doppler

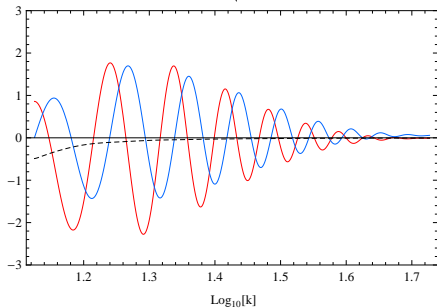


# Linear evolution of perturbations (order 1)

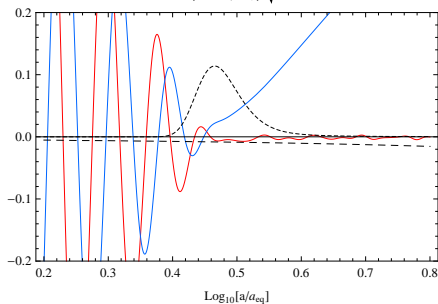


## Perturbations on the LSS

$\delta/4+\Phi, v_b/\sqrt{3}$  at LSS



$\delta_r/4+\Phi, v_b/\sqrt{3}$



- **Red:** Intrinsic  $\Theta$  and Einstein effect
- **Blue:** Doppler effect

## Linear response of radiation

$$\left[ (1 + R) \frac{\delta_r^{(1)'}}{4} \right]' + \text{visc} + \frac{k^2}{3} \frac{\delta_r^{(1)}}{4} \simeq -\frac{k^2}{3} (1 + R) \Phi^{(1)} + \left[ (1 + R) \Phi^{(1)'} \right]'$$

with  $R = 3\bar{\rho}_b / (4\bar{\rho}_r)$ .

Because of viscosity, at small scale:

$$\Theta^{(1)} \simeq \frac{\delta_r^{(1)}}{4} + \Phi^{(1)} \rightarrow -R\Phi^{(1)}$$

# Dynamics of primary effects

## Non-linear evolution (order 2)

$$\left[ (1 + R) \frac{\delta_r^{(2)'}}{4} \right]' + \text{visc} + \frac{k^2}{3} \frac{\delta_r^{(2)}}{4} \simeq -\frac{k^2}{3} (1 + R) \Phi^{(2)} + \left[ (1 + R) \Phi^{(2)'} \right]' + \text{quadr}$$

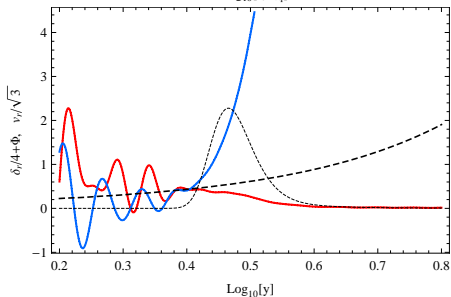
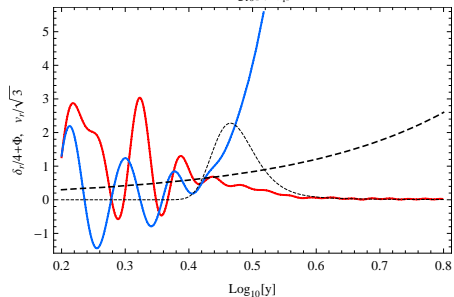
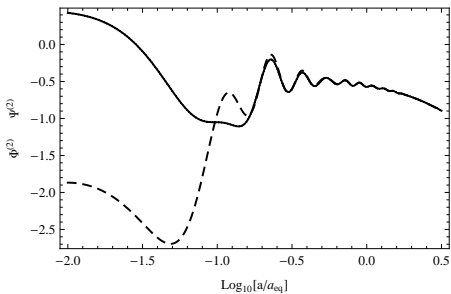
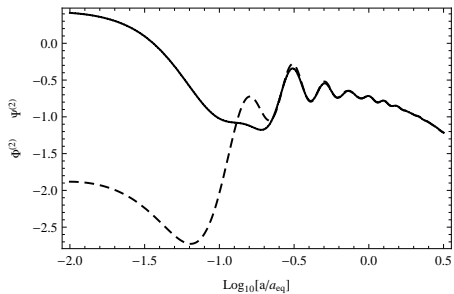
## Behaviour on small scales (viscosity)

$$\Theta^{(2)\text{intr}} \simeq \frac{\delta_r^{(2)}}{4} + \Phi^{(2)} \rightarrow -R\Phi^{(2)}$$

## Potential created by the collapse of cold dark matter

$$\frac{1}{2} \Phi^{(2)}(\mathbf{k}, \eta) \simeq -\frac{1}{6} K(\mathbf{k}_1, \mathbf{k}_2) \left( \frac{k_1 k_2 \eta}{k} \right)^2 \Phi(k_1) \Phi(k_2), \quad (1)$$





## Estimator of primordial non-Gaussianity

- We build an estimator for  $f_{\text{NL}}$  using all  $\langle \Theta_{\ell_1 m_1} \Theta_{\ell_2 m_2} \Theta_{\ell_3 m_3} \rangle$  up to  $\ell_{\text{max}}$  without taking into account the non-linear dynamics.
- If  $f_{\text{NL}} = 0$ , what measures this estimator? Answer :  $f_{\text{NL}}^{\text{eq}}$

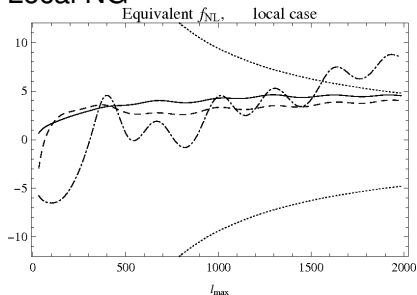
## Observational constraints on $f_{\text{NL}}$ (WMAP-5)

- Local NG:  $-9 < f_{\text{NL}} < 111$
- Equilateral NG:  $-150 < f_{\text{NL}} < 253$

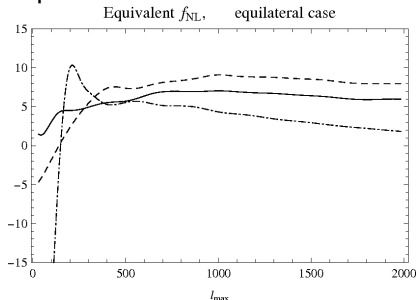
*Planck* is going to increase  $\ell_{\text{max}}$ , thus reducing the cosmic variance limitation, thus increasing the precision.

# Bispectrum generated by primary effects

## Local NG



## Equilateral NG



## Conclusion

- Secondary effect add up (same sign)
- Non-linear evolution has to be taken into account in future constraints on non-Gaussianity.

Thanks a lot for your attention.

