

PAMELA's Cosmic Positron from Decaying LSP in SO(10) SUSY GUT

**arXiv:0909.3139 [ph] (PLB)
arXiv:0902.3578 [ph] (JHEP)
arXiv:0902.0071 [ph] (JCAP)**

Bumseok KYAE

(Pusan National Univ.)

“WIMP miracle”

**Weakly Interacting Massive Particle
(WIMP)
can explain naturally Dark Matter.**

Lightest **S**upersymmetric **P**article (**LSP**) in MSSM
is one of the best examples of **WIMP**.

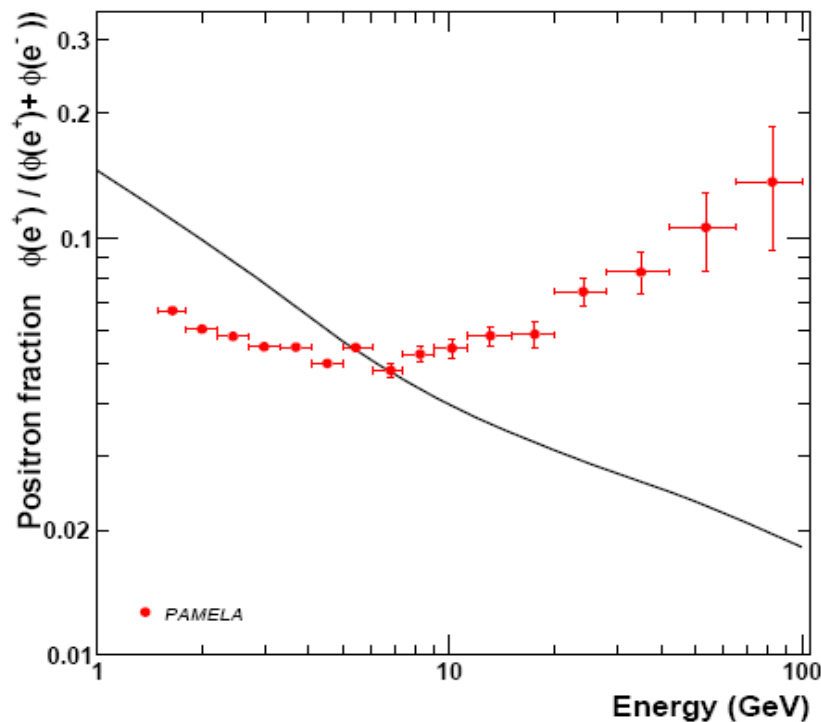
Recently **PAMELA/Fermi**
reported very challenging
observational results.

PRL102,051101(2009); Nature 458, 607 (2009)
arXiv:0905.0025(astro-ph HE)

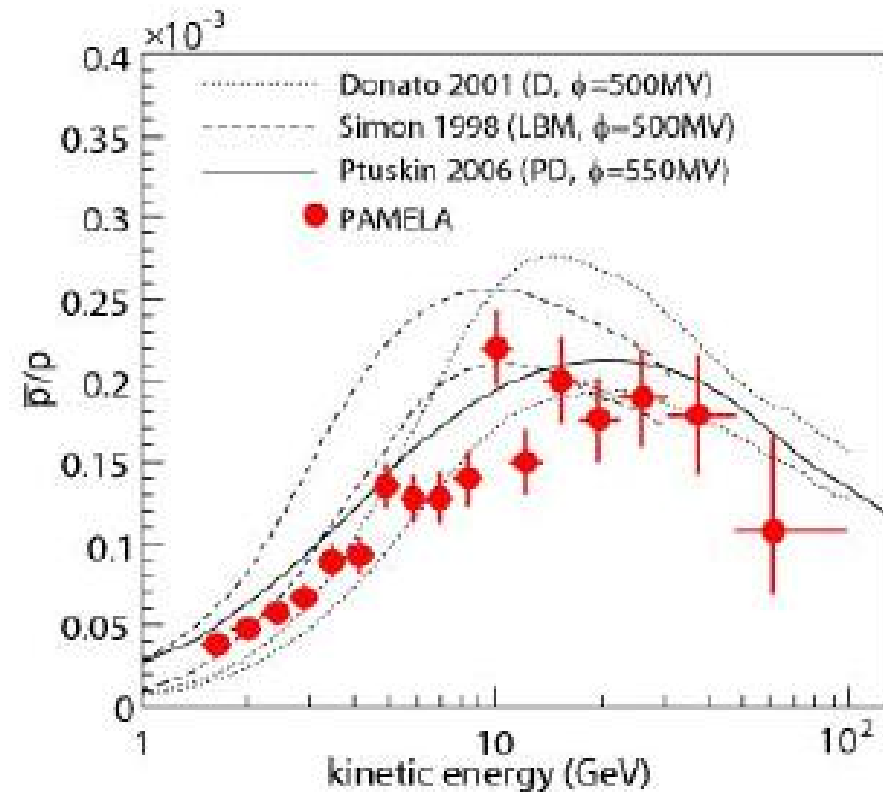
- **PAMELA** (Payload for **A**nti **M**atter **E**xploration and **L**ight nuclei **A**strophysics)
[exp. by a **SATELLITE**] measures **particles & nuclei fluxes** in cosmic ray.
- **Fermi** [exp. by a **SATELLITE**] released data on **electrons & positrons fluxes** in cosmic ray.

What are surprising?

PAMELA [arXiv.0810.4994,4995]



PAMELA positron fraction
v.s. theoretical models
(by Moskalenko & Strong '98)

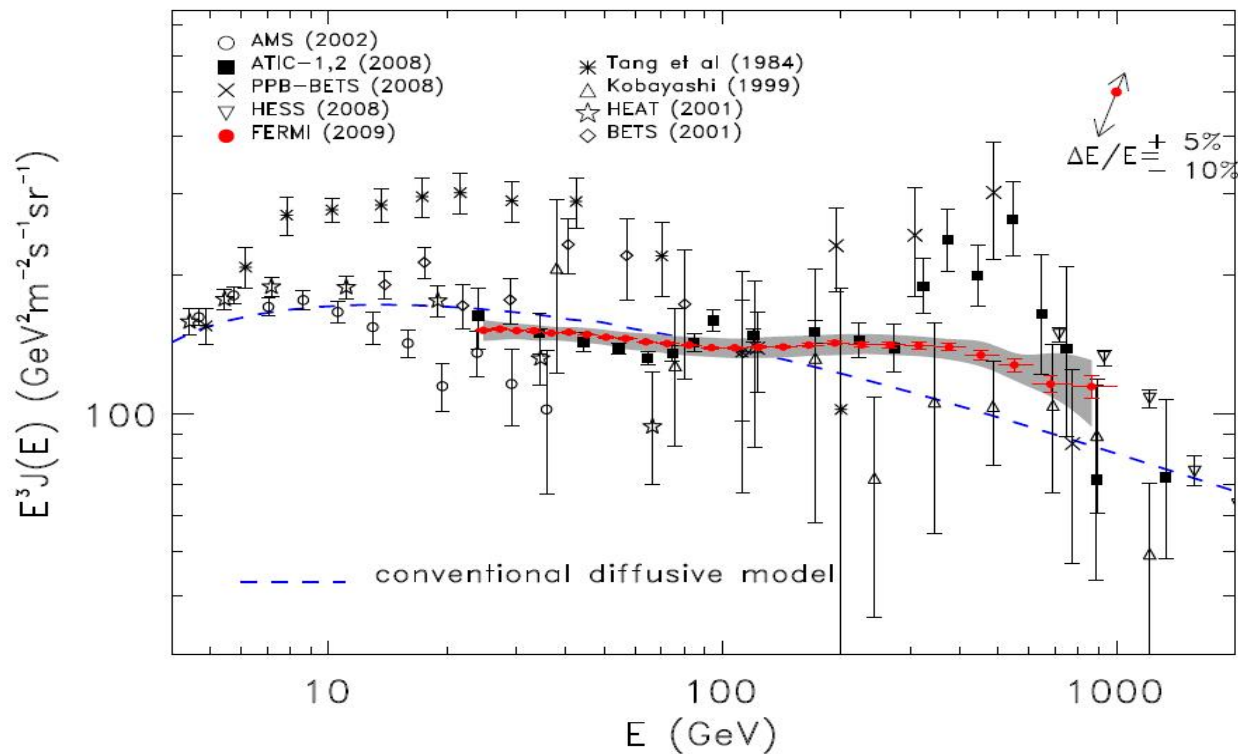


PAMELA anti-proton/proton flux ratio
v.s. theoretical calculation

What are surprising?

Fermi- LAT

[arXiv:0905.0025(astro- ph HE)]



(e⁺ + e⁻) excesses of cosmic ray are observed.

[100 GeV – 1000 GeV]

What are surprising?

(Fermi- LAT)

Positron excess keeps rising

mildly upto **1 TeV**.

As a strong possibility, it can be interpreted as a result from **TeV scale DM annihilation or decay**.

To explain the e^+ excess with annihilation,

- Should overcome “**helicity suppression,**”
to enhance **DM annihl. to e^+e^-** .
[Need a Large Boost Factor ($> 10^4$)]
- Should **suppress** the **hadronic** modes.

“**Leptophilic** annihilation !!”

Moreover,

Berstone etal. [arXiv:0811.3744]

**DM annihl. seems to be disfavored by
Gamma ray constraint,**

**if $m_{\text{DM}} \sim \text{TeV}$ (for explaining Fermi),
[$\Phi_{e^+} \propto (\rho/m_{\text{DM}})^2$] and**

if accept the galactic profile of

NFW or Einasto, because of

Bremsstrahlung at the galactic center.

DM DECAY for e^+ flux

(DM $\rightarrow e^+ e^-$, $\mu^+ \mu^-$, $\tau^+ \tau^-$ + neutral ptl.)

- We DON'T have to consider "helicity suppression."
- Gamma ray constraint is NOT serious.
 $[\Phi_{e^+} \propto (\rho/m_{DM})^2]$
- Hadronic decay should not exceed 10 %
i.e. should be "Leptophilic Decay"
- $\Gamma_{DM} \sim 10^{-26} \text{ sec}^{-1}$ for needed e^+ flux
- $m_{DM} \sim 2 \text{ TeV}$ for explaining Fermi
- Various and/or many body leptonic decays are needed for mild positron excess. [Bergtrom etal '09]

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Important Notice [comment 1]

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$$\Gamma_{\text{DM}} \sim \frac{m_{\text{DM}}^5}{192\pi^3 M_{\text{GUT}}^4} \sim 10^{-26} \text{ sec.}^{-1}$$

by **Dim. 6 operator** suppr. by M_{GUT}^2 (4 fermion int.)

for $m_{\text{DM}} \sim 2 \text{ TeV}$, $M_{\text{GUT}} \sim 10^{16} \text{ GeV}$

($m_{\text{DM}} \sim 100 \text{ GeV}$, $M_{\text{GUT}} \sim 10^{15} \text{ GeV}$)

PAMELA/Fermi's observ. might be **a signal of GUT.**

For a promising DM Decay Model

- Introduce **Leptophilic** int. between **superheavy fields** and **DM**.
- Introduce **other** (global) **symmetries** to **completely kill** the **dim. 5** operators.
- Introduce **an extra DM** component with **a TeV scale mass** for light enough **Higgs mass**.

PAMELA/Fermi anomaly
is **easily explained**
in 2-DM decay model !! [comment 2]

$$\Phi_{e^+}(E) = \left(\frac{\rho}{m_{\text{DM}}} \right) \cdot \Gamma_{\text{DM}} \times \frac{1}{4b(E)} \int_E^{m_{\text{DM}}} dE' \frac{dN_{e^+}}{dE'} I(\lambda_D),$$

In 2-DM model, $(\rho/m_{\text{DM}}) = \mathbf{n}_N$ can be **smaller**,
only if Γ_{DM} is larger,
[but $\Gamma_{\text{DM}} < 10^{-17} \text{ sec}^{-1}$, (age of univ.)⁻¹],

because the needed $\rho_{\text{DM}} \sim 10^{-6} \text{ GeV cm}^{-3}$
can be supported by χ .

2 component DM

- χ : Thermally produced, Absolutely stable,
Main comp. Relic density explained.
- N : Non-thermally produced, Meta-stable,
decay to e^+e^- , PAMELA/Fermi explained.

χ : Absolutely stable, Thermally produced.

N : Meta-stable, Non-thermally produced.

Even extremely small amount of N

[BK, Takahashi et al.]

$$[\mathcal{O}(10^{-10}) \lesssim (\Omega_N / \Omega_\chi)]$$

can produce the positron flux needed to account for PAMELA/Fermi data,

only if the decay rate is enhanced by relatively lighter M_* ,

$$[10^{12} \text{ GeV} \lesssim M_* \lesssim 10^{16} \text{ GeV}].$$

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In DM decay Models, (1st scenario)

[BK], [K.Bae, BK], [J.Huh, J.E.Kim], etc.

- Explain PAMELA/Fermi data with leptophilic YUKAWA couplings between GUT scale fields and an extra DM component \tilde{N} ($\tilde{N} \rightarrow \chi L^+ \bar{L}$).
- Even extremely small amount of \tilde{N} can explain the PAMELA/Fermi data.
- The nobleness of the MSSM (SUSY at 10^2 GeV, gauge coupling unif., χ CDM) can be maintained.
- The models can be easily embedded in Flipped SU(5).

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Disadvantages(?) in 1-DM decay Models

- **A natural explanation for $10^{-26} \text{ sec.}^{-1}$ decay rate is required. (→ Need a Model)**
- **Ad-hoc new TeV DM and new leptophilic interactions are introduced.**
- **Desired relic density is Not automatic.**
Need elaborate **Non-thermal** production of DM.
→ **Reheating temp.** should be tuned.

I attempt to explain **PAMELA** (Fermi)
only **within the framework of** a well-
known Particle physics model, **SO(10)**
without introducing any new DM and
new special interactions. **(2nd scenario)**
[arXiv:0909.3139, BK]

From now on,

I will suppose that **DM is the bino-like LSP.**

SO(10)

$$45_G = SM + \{E, E^c\} + N \\ + \{Q', Q'^c\} + \{Q, Q^c; U, U^c\}$$

where $E = (1,1)_{-1}$, $N = (1,1)_0$,
 $Q' = (3,2)_{-5/6}$, $Q = (3,2)_{1/6}$, $U = (3,1)_{2/3}$

$$SM + \{E, E^c\} + N = LR$$

$$SM + \{Q', Q'^c\} = SU(5)$$

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$SO(10) \rightarrow SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L} = LR$
by $\langle 45_H \rangle$, $\{Q', Q'^c\}$, $\{Q, Q^c; U, U^c\}$ massive

$SO(10) \rightarrow SU(5)$ by $\langle 16_H \rangle$, $\langle 16^*_H \rangle$
 $\{E, E^c\}$, N , $\{Q, Q^c; U, U^c\}$ massive

$\langle 45_H \rangle$ is 10^{16} GeV from RG eff. of the MSSM gauge couplings, but $\langle 16_H \rangle$ is not pinned down yet.

If $\langle 45_H \rangle > \langle 16_H \rangle = \langle 16^*_H \rangle$,
masses of $\{Q', Q'^c\}, \{Q, Q^c; U, U^c\} > \{E, E^c\}, N$

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So $\{Q, Q^c; U, U^c\}$ are always heavier.

Superheavy fields in SO(10)

- **Gauge boson/Gauginos** of SO(10)/SM
- **Triplets** in $\mathbf{10}_h$ ($=\{D^c, h_d\} + \{D, h_u\}$)
e.g. by $\mathbf{10}_h \langle 45_H \rangle \mathbf{10}_h$
- **GUT breaking Higgs**
due to its VEV, they couple to MSSM fields only via non-renormalizable terms. They weakly coupled to SM

LSP decay due to sRH ν

- **If (1) R-parity is absolutely preserved, and (2) χ is the LSP, χ can not decay.**

• BUT if sRH ν develops a VEV (R viol.), or if sRH ν is lighter than χ (sRH ν LSP), χ could decay.

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LSP decay due to sRH ν

How can one obtain
e.g. **extremely small**
R-parity violating effects
NATURALLY?

LSP decay due to sRH ν

- **RH ν and sRH ν are neutral singlets under SM. But they are **charged under SO(10)**.**
- IF Yukawa couplings, e.g. $W = \ln_\mu y^c$ are somehow forbidden,
RH ν and sRH ν extremely weakly interact with SM,
since they can interact with SM only through superheavy SO(10) gauge fields / gauginos.
→ Assume one family of RH ν / sRH ν decoupled.

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Interactions of the MSSM fields and heavy gauginos

$$\tilde{e}_i^{c*} \nu_i^c \tilde{E}^c, \quad \tilde{d}_i^{c*} u_i^c \tilde{E}^c, \quad h_u^{+*} \tilde{h}_d^0 \tilde{E}^c, \quad h_u^{0*} \tilde{h}_d^- \tilde{E}^c$$

$$\tilde{\nu}_i^{c*} e_i^c \tilde{E}, \quad \tilde{u}_i^{c*} d_i^c \tilde{E}, \quad h_d^{0*} \tilde{h}_u^+ \tilde{E}, \quad h_d^{-*} \tilde{h}_u^0 \tilde{E}$$

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$$\tilde{e}_i^{c*} q_i \tilde{Q}'^c, \quad \tilde{d}_i^{c*} l_i \tilde{Q}'^c, \quad \tilde{q}_i^* u_i^c \tilde{Q}'^c$$

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For leptophilic χ decay,

- $\langle 16_H \rangle \ll \langle 45_H \rangle$, effectively **LR model**
- If sv^c is heavier than χ , a non-zero VEV $\langle sv^c \rangle$ must be assumed.
- Squarks, charged Higgs, and soft para. are much **heavier (\sim TeV)** than a slepton.
- For PAMELA, $m_\chi \sim 300$ GeV,
Fermi is explained with astrophys. source.
- One RFI γ is lighter than χ .

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**Even with 2 heavy RH ν ,
seesaw mech. is still O.K.**

$$W_\nu = y_{ij}^{(\nu)} l_i h_u \nu_j^c (j \neq 1) + \frac{1}{2} M_{i,j} \nu_i^c \nu_j^c (i, j \neq 1),$$

$$m_\nu = m_\nu^T = - \begin{pmatrix} 0 & v_{12} & v_{13} \\ 0 & v_{22} & v_{23} \\ 0 & v_{32} & v_{33} \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 \\ 0 & M_{22}^{-1} & M_{23}^{-1} \\ 0 & M_{23}^{-1} & M_{33}^{-1} \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 \\ v_{12} & v_{22} & v_{32} \\ v_{13} & v_{23} & v_{33} \end{pmatrix}$$

Still 3 LH ν can be maximally mixed.

[Frampton, Glashow, Yanagida (2002)]

If sRH ν is lighter than χ , a VEV of sRH ν is not essential. \rightarrow 4 bdy decay !!

Just for simplicity, assume a VEV of sRH ν .
(\rightarrow 3 bdy decay) e.g. by

$$W \supset \frac{1}{M_P} \langle \overline{\mathbf{16}}_H \rangle \mathbf{16}_1 S^2 + S^3$$

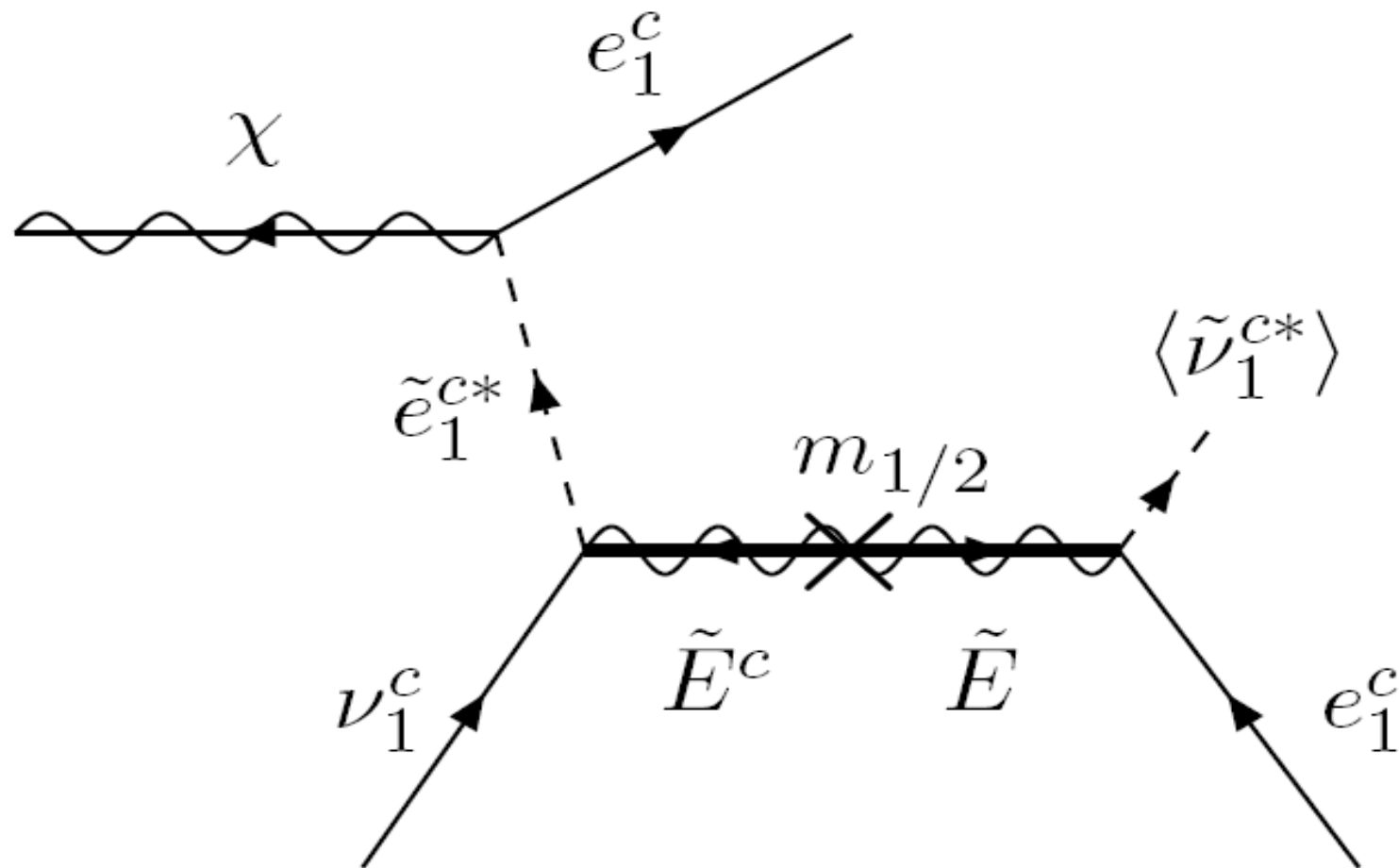
$$\mathbf{R}(\mathbf{16}_1) = \mathbf{R}(S) = 2/3$$

$$\mathbf{R}(\mathbf{16}^*_H) = 0$$

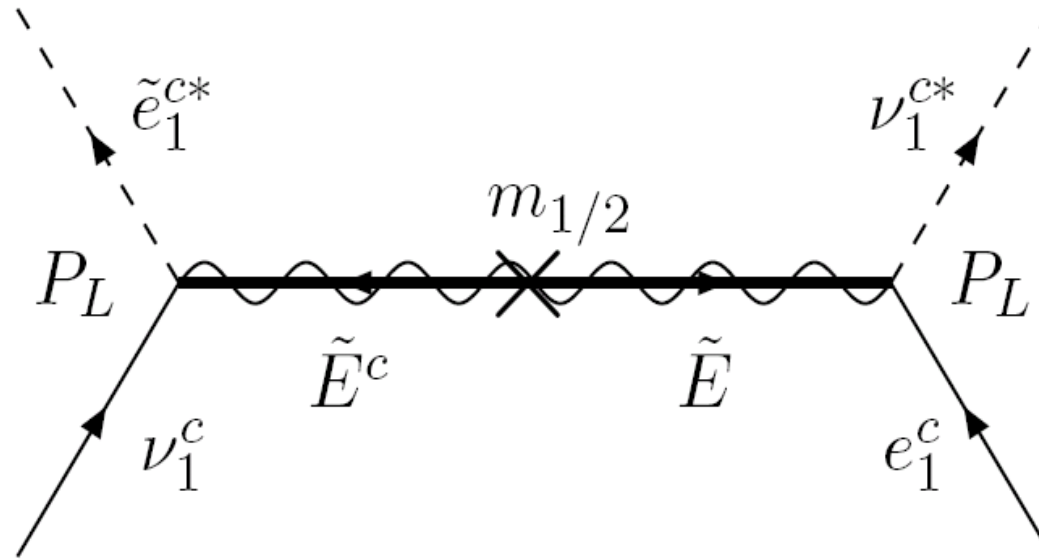
including soft terms in V,

$$\langle \tilde{\nu}_1^c \rangle \sim m_{3/2} \times \frac{M_E}{M_P}$$

LSP decay diagram



Charged gaugino mediation



Dirac mass M_E by
Gauge sym. breaking

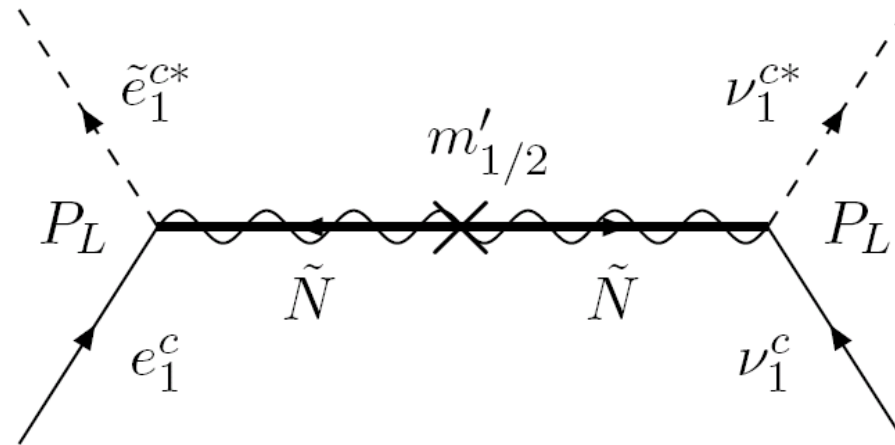
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Majorana mass $m_{1/2}$
by SUSY breaking

This diagram is suppressed by

$$m_{1/2} / M_E^2$$

Neutral gaugino mediation



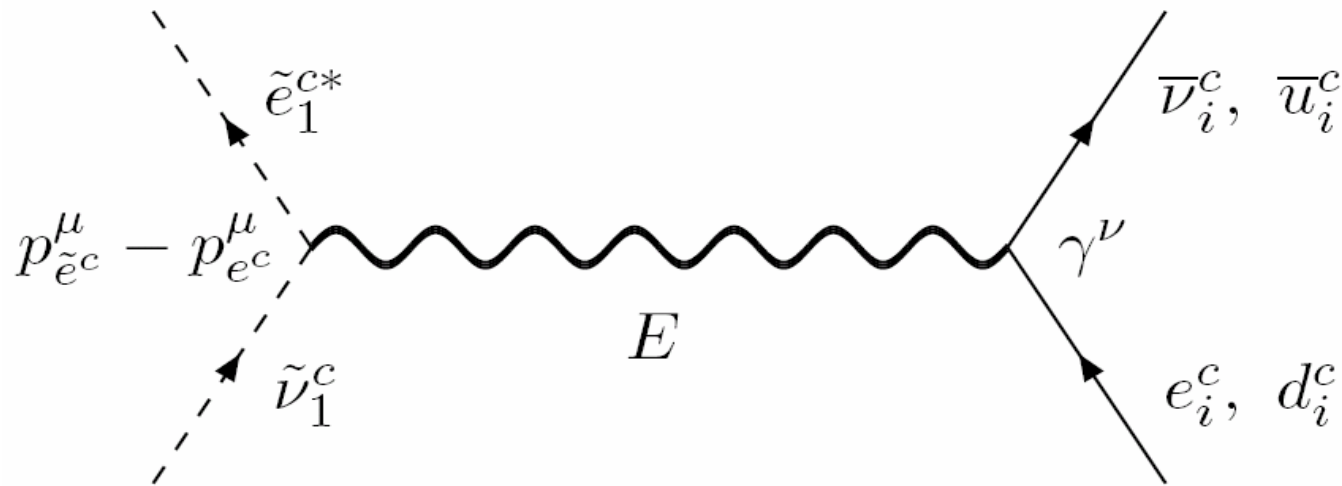
$$g_{LR} = (2/3)^{1/2} g_{B-L} = g_{10}$$

$$M_N = M_E \times (5/2)^{1/2}$$

Eff. coupling is $\frac{1}{4}$ of the C.C. case.

Suppressed by $2/5 \times 1/4 = 1/10$
Compared to the C.C. case

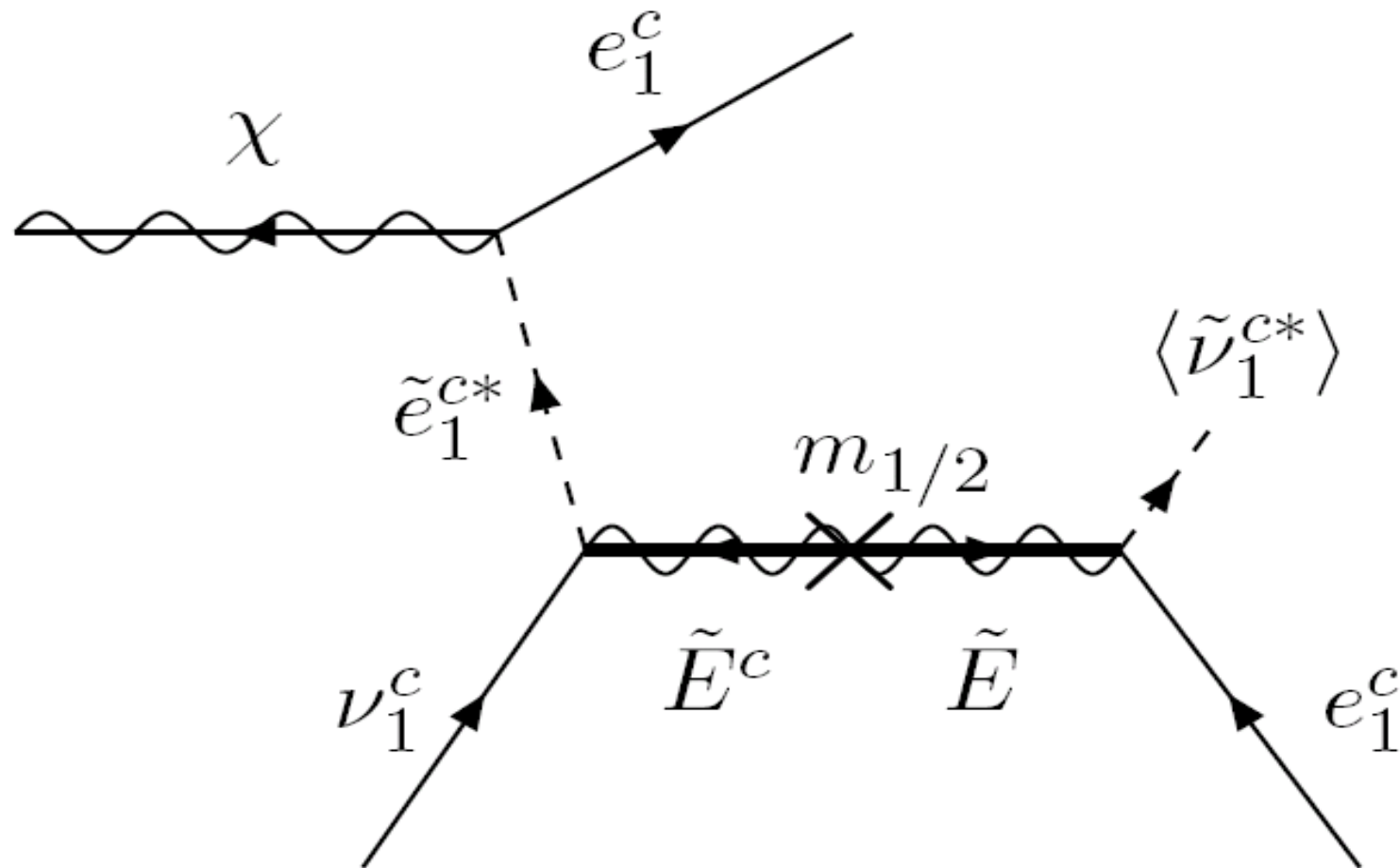
Charged gauge field mediation



A derivative coupling is involved.

Since $m_{1/2} \gg m_\chi$ this diagram is suppressed.

LSP decay diagram



The 1st realization of $\Gamma_\chi \sim 1 / (M_{\text{GUT}})^4$ from the gauge interaction.

The **decay rate of χ** is

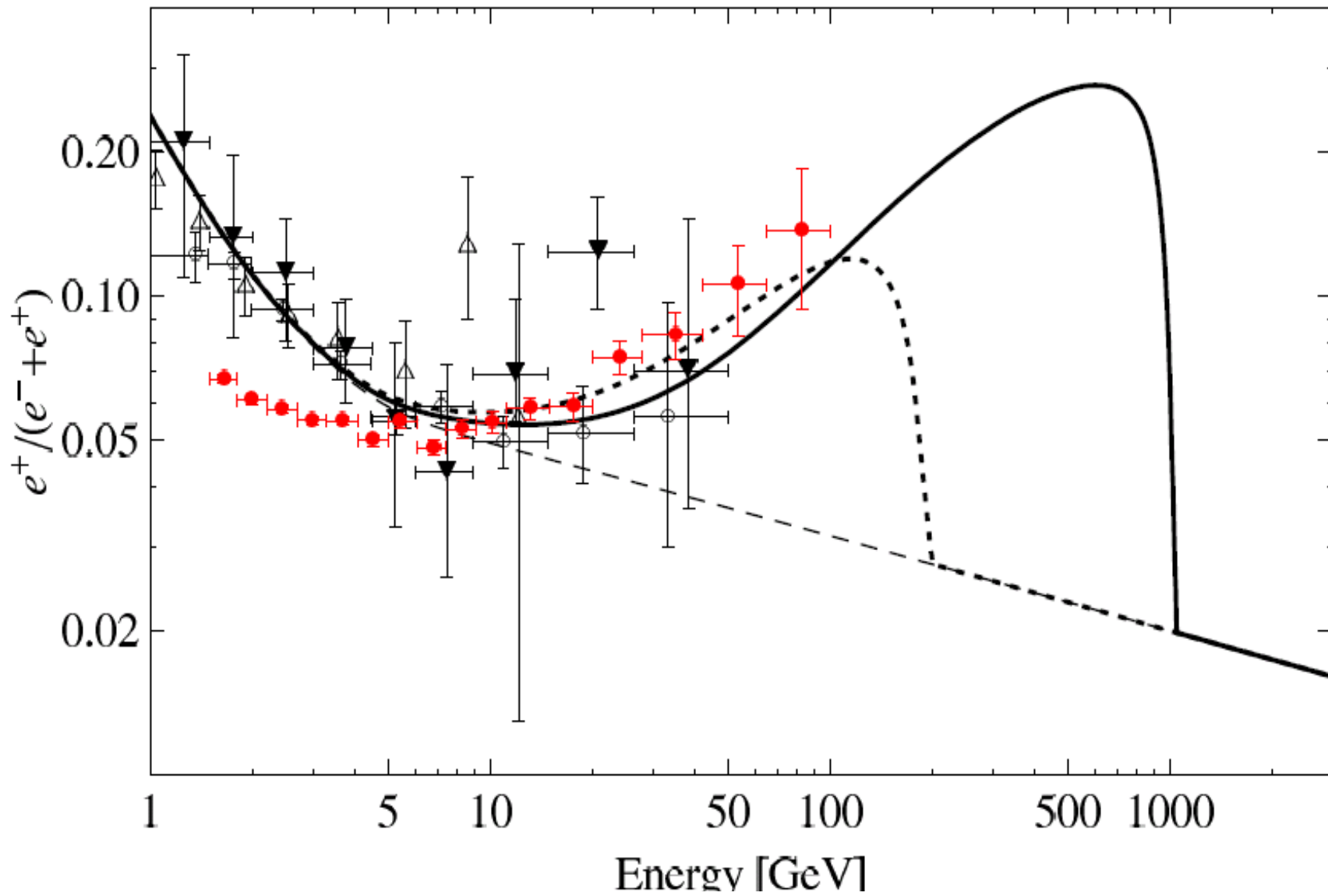
$$\Gamma_\chi = \frac{\alpha_{10}^2 \alpha_Y m_\chi^5}{96 M_E^4} \left(\frac{m_{3/2} \langle \tilde{\nu}_1^c \rangle}{m_{\tilde{e}_1^c}^2} \right)^2 \sim \frac{\alpha_{10}^2 \alpha_Y m_\chi^5}{96 M_E^2 M_P^2} \left(\frac{m_{3/2}}{\kappa m_{\tilde{e}_1^c}} \right)^4 \sim 10^{-26} \text{ sec.}^{-1},$$

To be consistent with the **PAMELA**'s data,

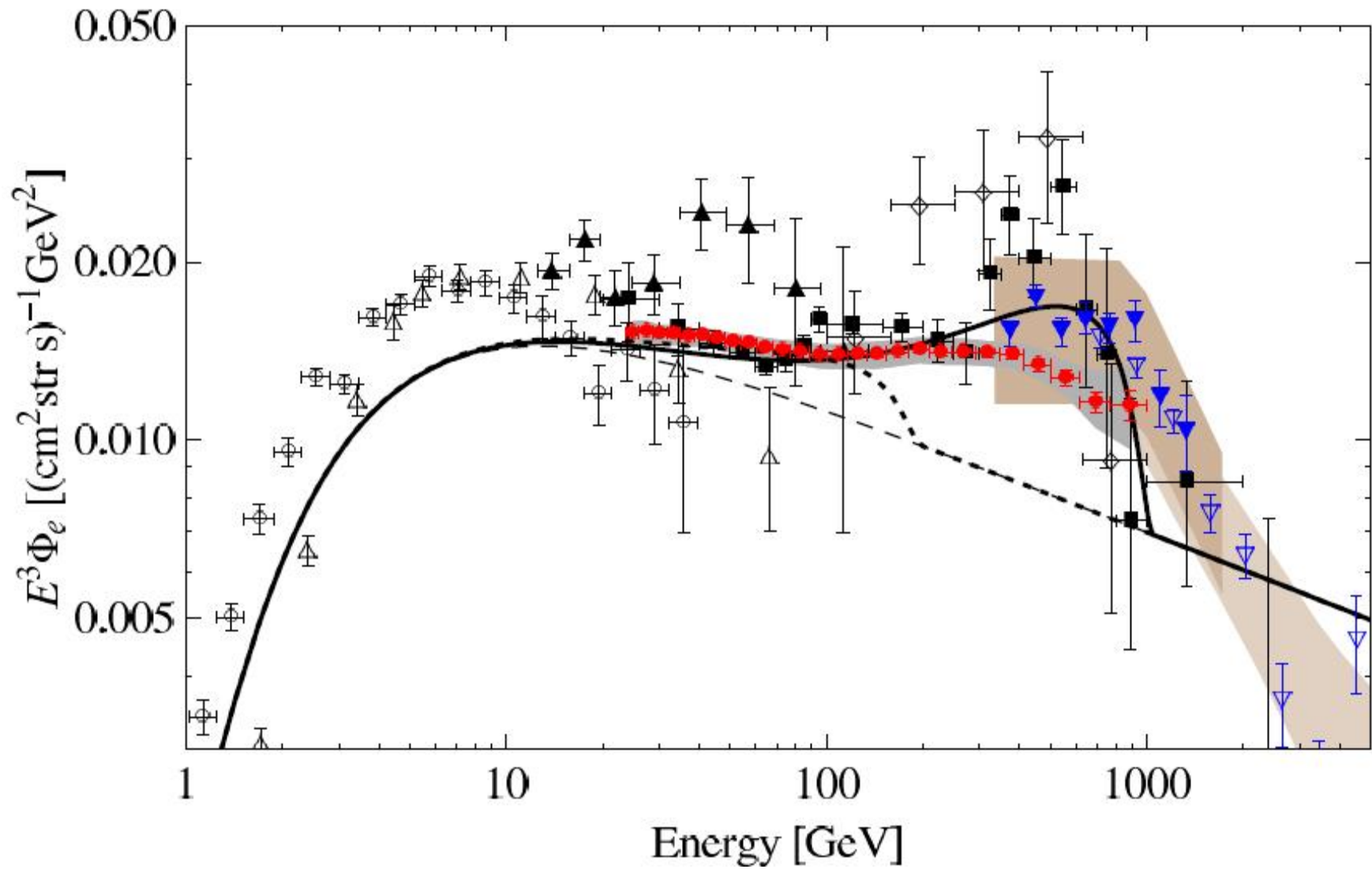
$$\mathbf{M_E \sim \langle 16_H \rangle \sim 10^{14} \text{ GeV}}$$

$$\mathbf{2 \text{ RH } \nu \text{ masses } \sim 10^{10} \text{ GeV}}$$

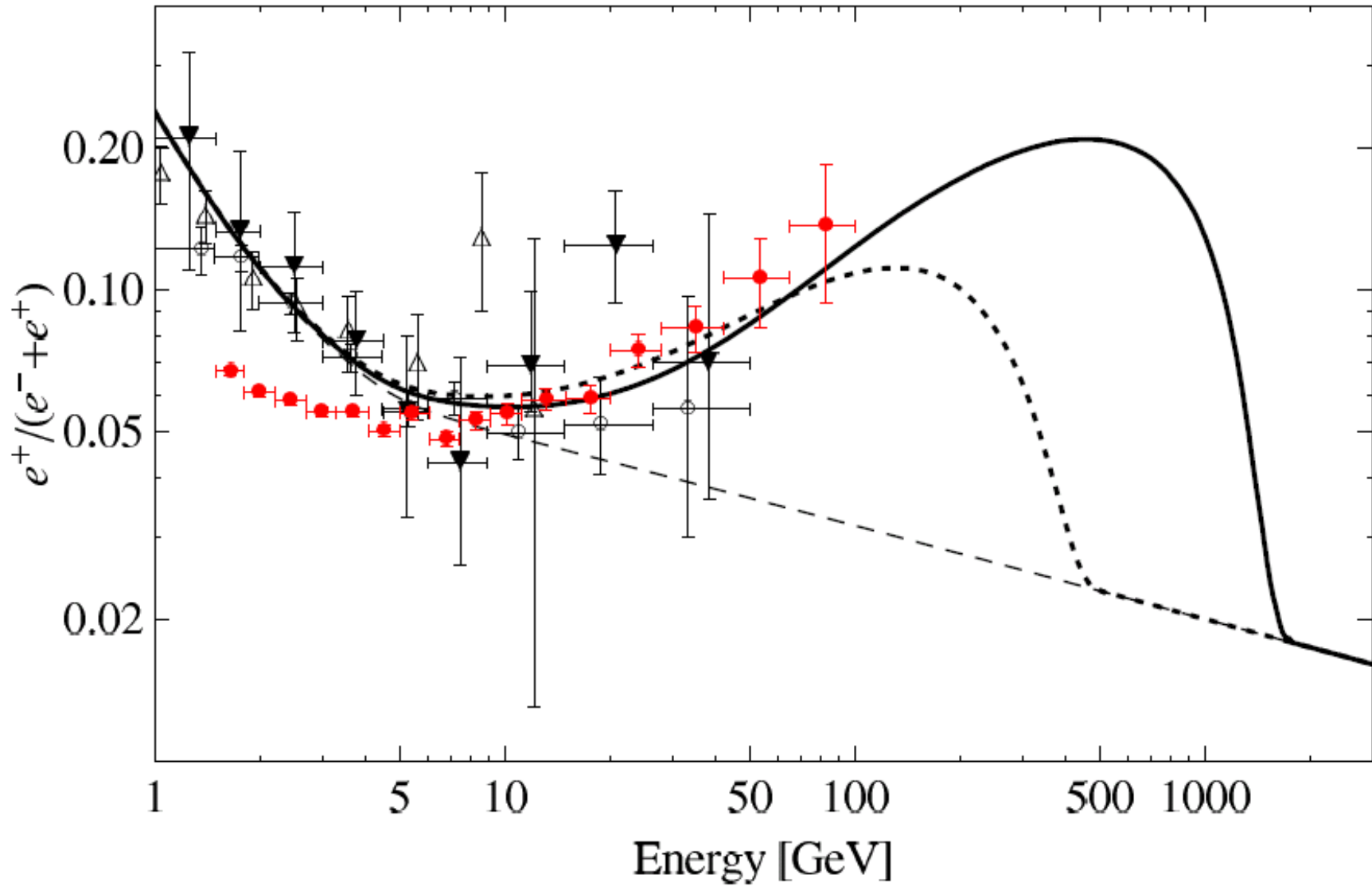
from $W \supset \frac{1}{M_P} \langle \overline{16}_H \rangle \langle \overline{16}_H \rangle 16_i 16_j (i, j \neq 1) \supset (10^{10} \text{ GeV}) \times \nu_i^c \nu_j^c (i, j \neq 1)$



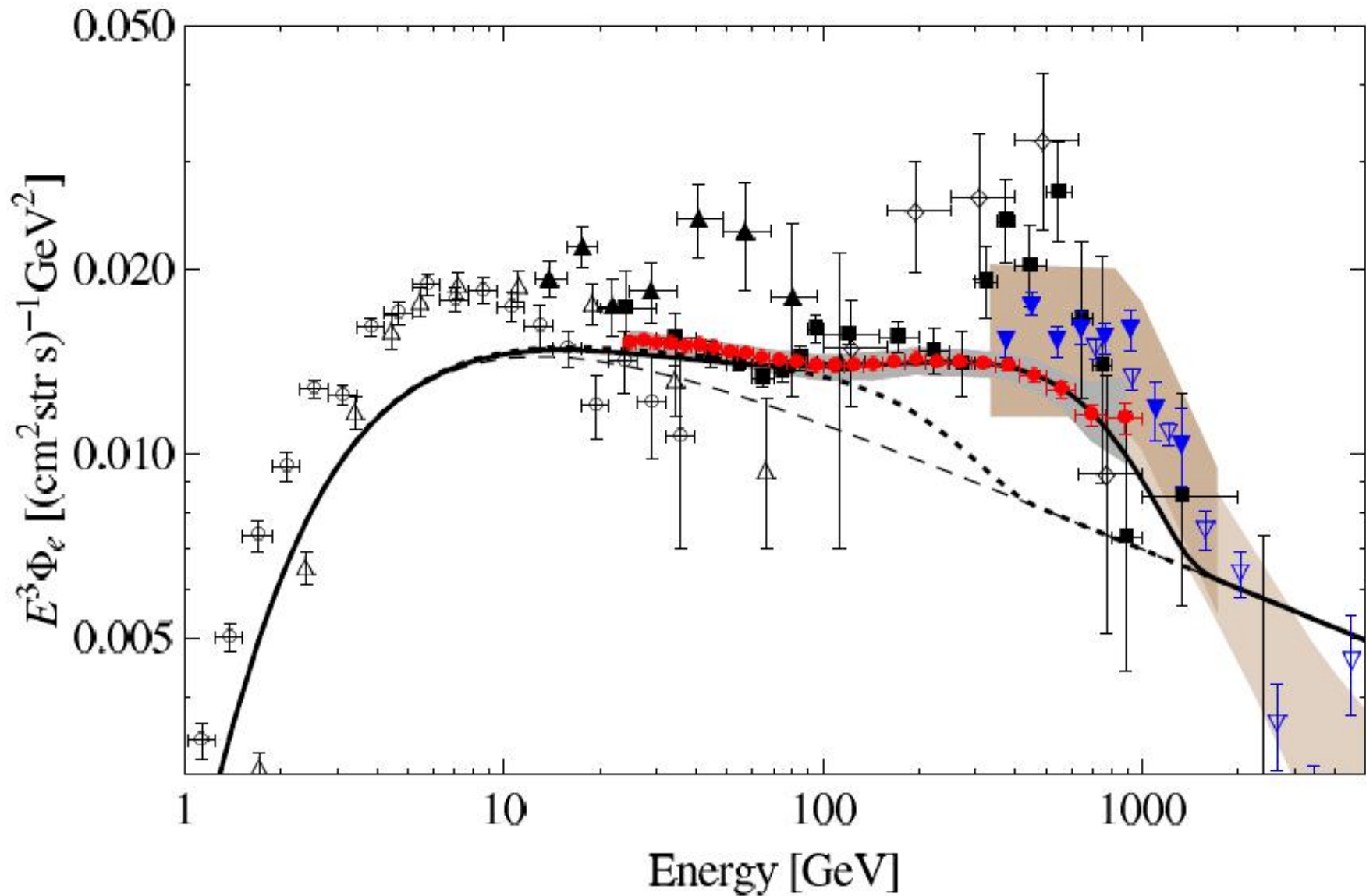
Dotted Line: $\chi \rightarrow e^+e^- \nu$, $M_\chi=400$ GeV, [Ibarra,etal '09]



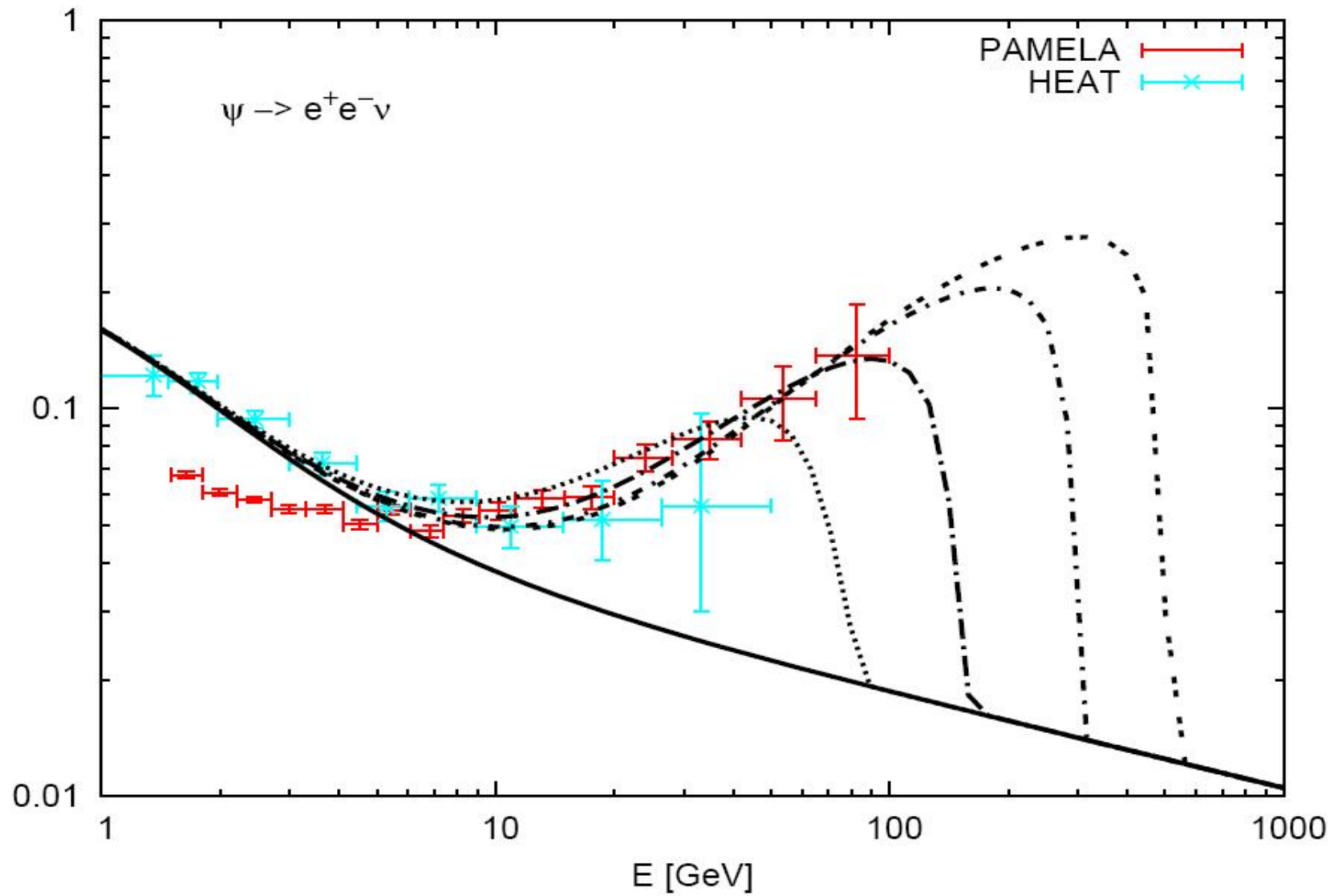
Dotted Line: $\chi \rightarrow e^+e^- \nu$, $M_\chi=400$ GeV, [Ibarra,etal '09]



Solid Line: $\chi \rightarrow \mu^+\mu^-\nu$, $M_\chi=3.5$ TeV, [Ibarra,etal '09]



Solid Line: $\chi \rightarrow \mu^+ \mu^- \nu$, $M_\chi = 3.5$ TeV, [Ibarra, etal '09]



Thick Line: $\chi \rightarrow e^+e^-\nu$, $M_\chi=300$ GeV, [Ibarra,etal '08]

Conclusions

- Still the bino-like LSP DM scenario is consistent with PAMELA,
if $s_{RH} \nu$ develops a VEV or is lighter than bino, and a RH ν is light enough.
- SO(10) provides a relatively predictable explanation.
- In the specific case, LR breaking scale is 10^{14} GeV, and the seesaw scale is 10^{10} GeV.