

Non-gaussianity in axion N-flation models

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Based on arXiv:1005.4410

by SAK, Andrew R. Liddle and David Seery (Sussex),
and earlier papers by SAK and Liddle.

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Assisted inflation

- Assisted inflation (Liddle-Mazumdar-Schnook 1998) is the observation that multiple scalar fields can cooperate to drive inflation even if each individually is unable to.
- Each field feels the acceleration from its own potential, but the collective Hubble friction from all fields.

N-flation

- N-flation (Dimopoulos et al 2008) is a realization of assisted inflation using string axions.

Motivations and assumptions

- One motivation for this idea is that sufficient inflation can be obtained with all fields maintaining **sub-Planckian** values.

Another is that it may be possible to relate assisted inflation to proper fundamental physics models.

- Focused on **adiabatic perturbations**.
- **Random** initial conditions for fields.
- Assumptions we made;
Horizon crossing and Slow-roll approximations.

N-flation phenomenology

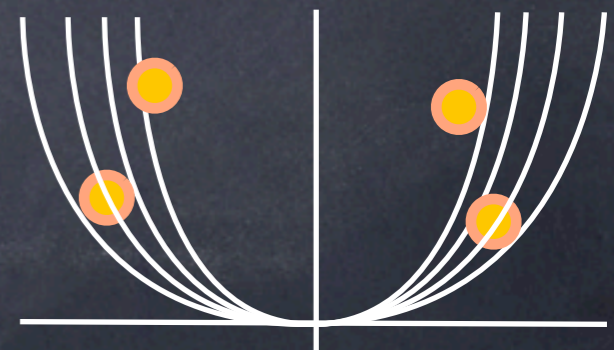
- The full string axion potential is

$$V_i = \Lambda_i^4 \left(1 - \cos \frac{2\pi\phi_i}{f_i} \right)$$

, where there are N_f fields with constants Λ_i and f_i . Throughout will ignore possible couplings btw the fields.

- This has been extensively explored in the quadratic approximation where all fields are close to their minima, in which case they are simply a set of massive fields with

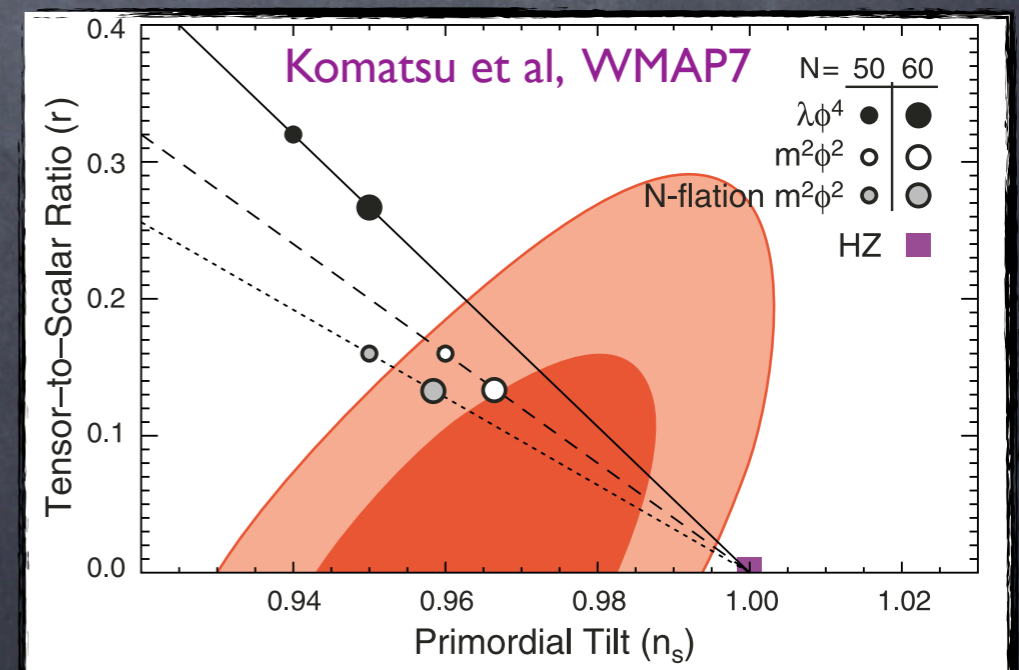
$$m_i \equiv \frac{2\pi\Lambda_i^2}{f_i}$$



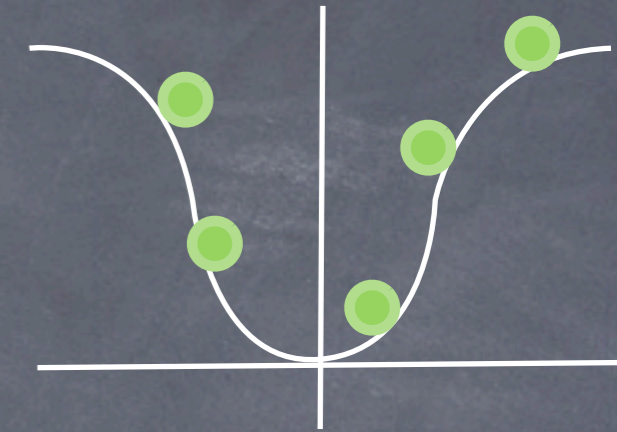
N-flation phenomenology

Regardless of these choices, the Nflation phenomenology in this approximation is remarkably simple:

- The tensor-to-scalar ratio always equals to the single-field values: $r = 8/N$ where N is the number of e-foldings.
- The scalar spectral index cannot exceed the single-field value, equalling it only in the equal-mass case:
 $n \leq 1 - 2/N$.
- The non-gaussianity f_{NL} always equals its single-field value: $f_{NL} = 2/N$ and hence is unobservably small.



The N-flation model



But..

.. in fact the quadratic approximation to the potential is unlikely to be valid. We should consider the full potential

$$V_i = \Lambda_i^4 \left(1 - \cos \frac{2\pi\phi_i}{f_i} \right)$$

Even if the potentials are all taken to be identical, assisted inflation is an attractor solution only if $d^2V/d\phi^2 > 0$ (Calcagni and Liddle 2008), which is not true near the maximum of the potential(s) where the trajectories will diverge.

[\implies From now on, called the Naxion model]

Naxion equations

- With the full potential, the observables can be calculated using the so-called δN formalism

$$\mathcal{P}_\zeta = \frac{H_*^2}{4\pi^2} \sum_i N_{,i} N_{,i} = \frac{H_*^2}{8\pi^2 M_{\text{P}}^2} \sum_i \frac{1}{\epsilon_i^*};$$

$$n - 1 = -2\epsilon_* - \frac{8\pi^2}{3H_*^2} \sum_j \frac{\Lambda_j^4}{f_j^2} \frac{1}{\epsilon_j^*} / \sum_i \frac{1}{\epsilon_i^*};$$

$$r = \frac{2}{\pi^2 \mathcal{P}_\zeta} \frac{H_*^2}{M_{\text{P}}^2} = 16 / \sum_i \frac{1}{\epsilon_i^*};$$

$$\frac{6}{5} f_{\text{NL}} \simeq \frac{\sum_{ij} N_{,i} N_{,j} N_{,ij}}{(\sum_k N_{,k} N_{,k})^2} = \frac{r^2}{128} \sum_i \frac{1}{\epsilon_i^*} \frac{1}{1 + \cos \alpha_i^*},$$

, where

$$\epsilon_i \equiv \frac{M_{\text{P}}^2}{2} \left(\frac{V'_i}{V_i} \right)^2$$

$$\epsilon \equiv -\dot{H}/H^2$$

$$\approx \sum_i (V_i/V)^2 \epsilon_i$$

- Here $\alpha_i = 2\pi\phi_i/f_i$, ϵ_i is the slow-roll parameter of each field, derivatives wrt field i and indicated by a comma, and $*$ indicates evaluation at horizon crossing

$N_{\text{axion}} : N_{\text{tot}}$

- The number of e-foldings is given by

$$N_{\text{tot}} \simeq \sum_i \left(\frac{f_i}{2\pi M_{\text{P}}} \right)^2 \ln \frac{2}{1 + \cos \alpha_i} \simeq \frac{\ln 2}{2\pi^2} \frac{f^2}{M_{\text{P}}^2} N_{\text{f}}$$

where the last expression calculates the expectation value of the sum under assumption of uniformly-distributed angles α_i . For values of f of order the Planck mass, sufficient inflation requires a large number of fields, at least hundreds.

Naxion: n and r

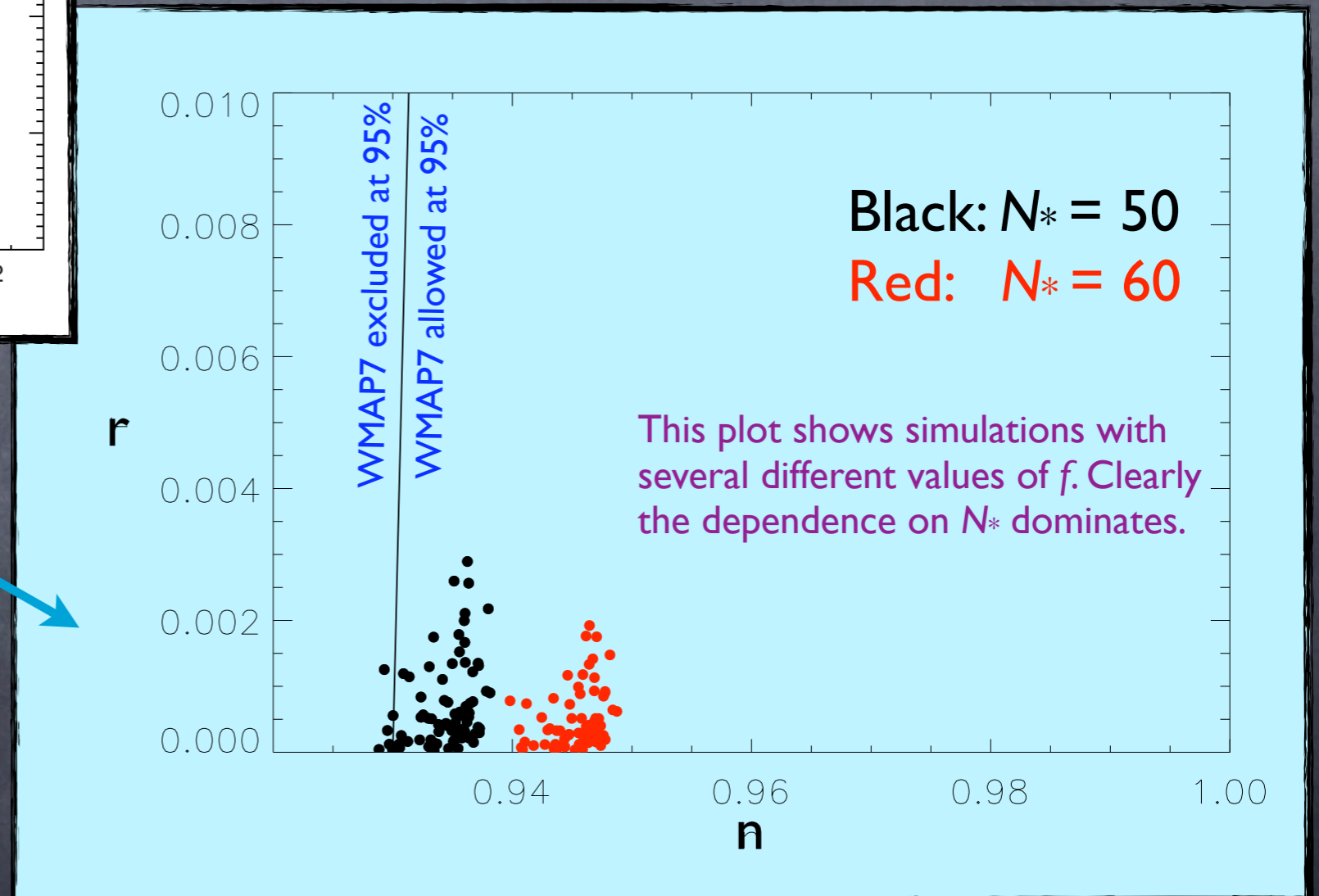
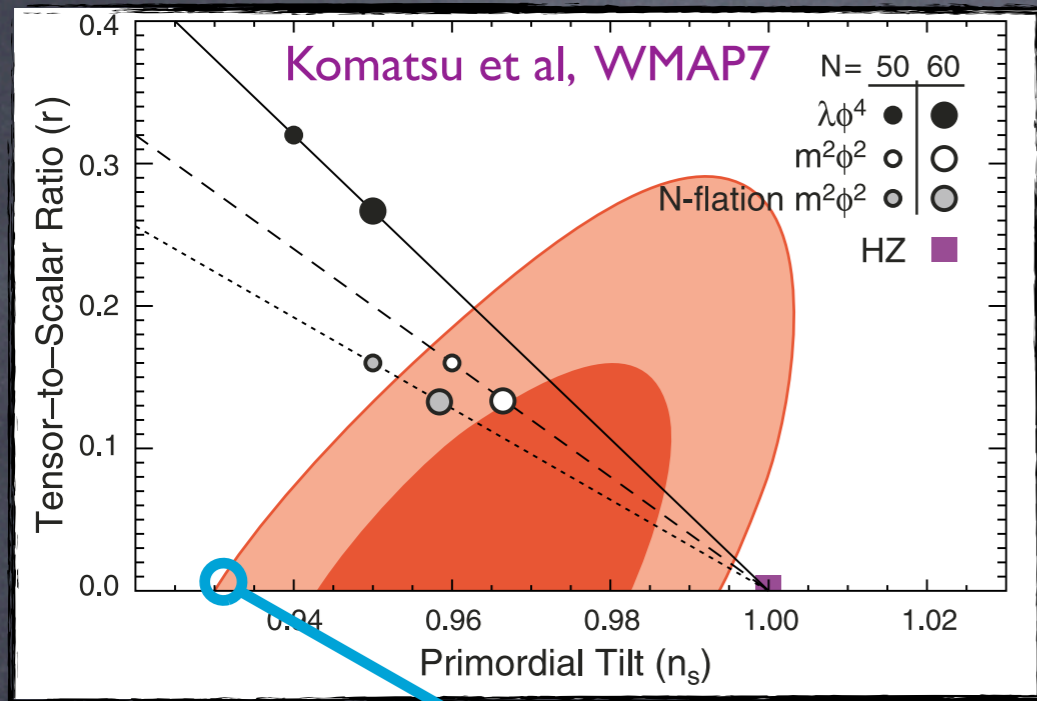
- A similar summation trick, assuming uniformly distributed fields, gives an analytic estimate of the spectral index as

$$\langle n - 1 \rangle \simeq -5 \ln 2 / N_*$$

NB: $5 \ln 2 \simeq 3.5$

- At the same time, the tensor-to-scalar ratio is highly suppressed by the small ϵ_i of fields close to the maximum.

Naxion: n and r



Naxion: non-gaussianity

- The interesting aspect of the model is the behavior of the non-gaussianity:

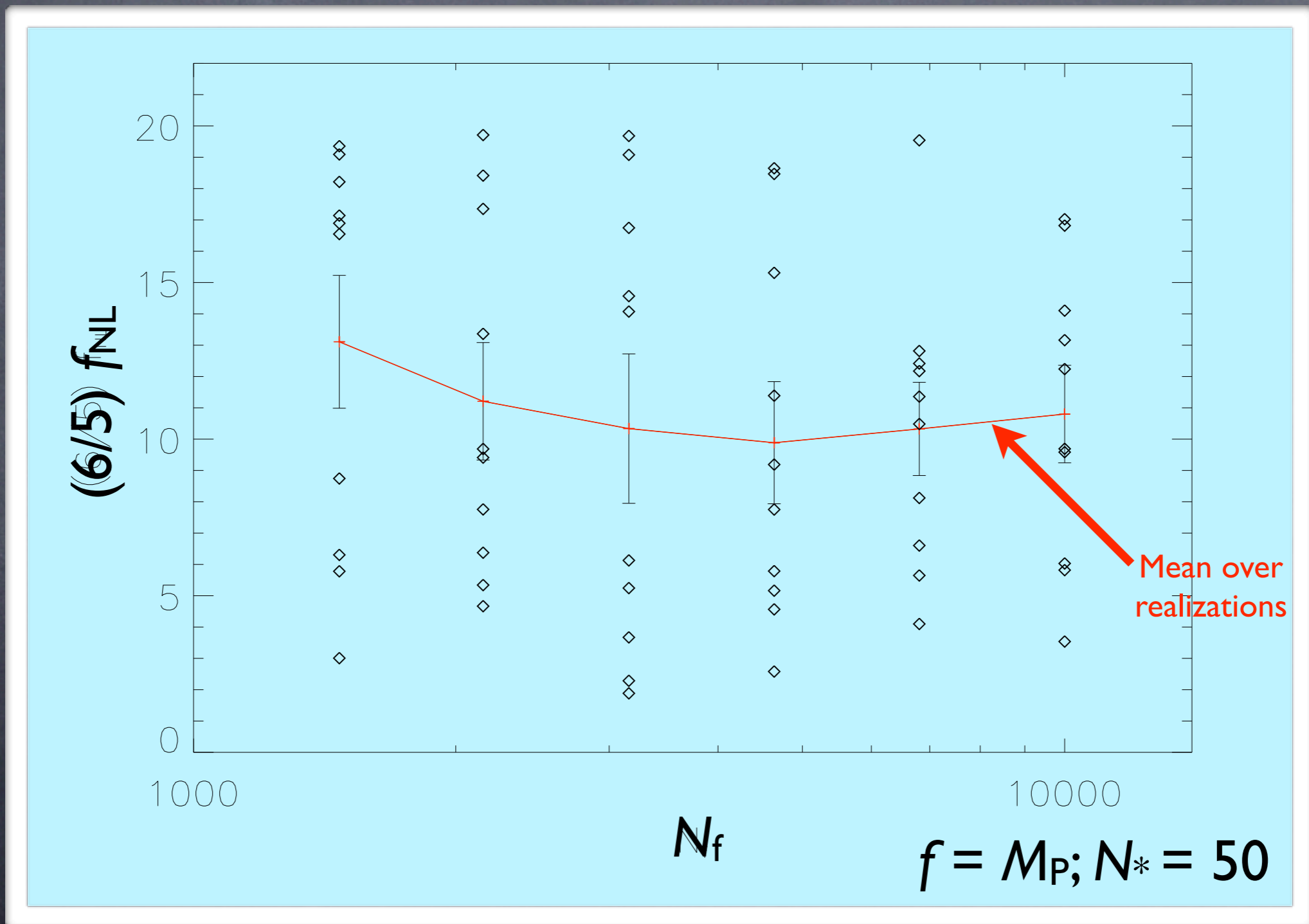
$$\frac{6}{5} f_{\text{NL}} \simeq \frac{\sum_{ij} N_{,i} N_{,j} N_{,ij}}{(\sum_k N_{,k} N_{,k})^2} = \frac{r^2}{128} \sum_i \frac{1}{\epsilon_i^*} \frac{1}{1 + \cos \alpha_i^*},$$

- The sum may be dominated by a small number of fields whose α_i is very close to π . If there are \tilde{N} fields which dominate with comparable ϵ_i , there is an approximate form

$$\frac{6}{5} f_{\text{NL}} \approx \frac{2\pi^2}{\tilde{N}} \left(\frac{M_{\text{P}}}{f} \right)^2,$$

hence for f of order M_{p} the f_{NL} may be of order tens.

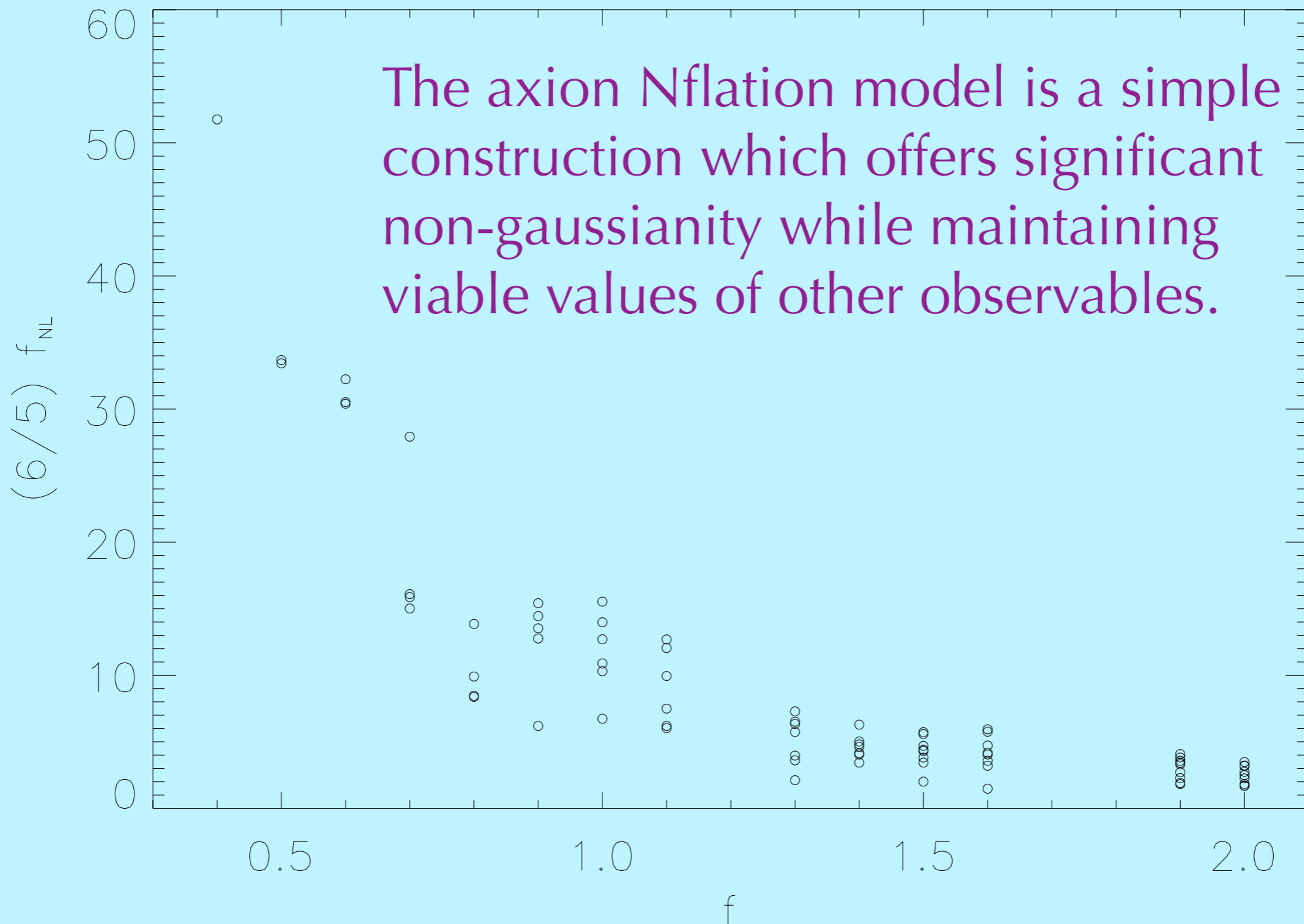
Naxion: non-gaussianity



Interpretation

- There is nothing particularly unusual about the predictions of the non-gaussianity in this model; it is in fact what one would get from a single field evolving in the axion potential.
- However a single-field model with those parameters would not be satisfactory, as it would not give sufficient inflation and the spectral index would be far from unity.
- The scenario works because the assisted inflation mechanism strongly alters the predicted spectral index, but has only a marginal effect on the non-gaussianity

Conclusion



Naxion: trispectrum

- A similar analysis yields an estimate of the trispectra

$$\tau_{\text{NL}} = (4\pi^4 / \bar{N}^2)(M_{\text{P}}^4 / f^4)$$

$$= (6/5 f_{\text{NL}})^2$$

$$(54/25)g_{\text{NL}} = (8\pi^4 / \bar{N}^2)(M_{\text{P}}^4 / f^4)$$

- As seen in the bispectrum plot, there is a large spread of predictions due to the randomness of initial conditions.

However there are predicted correlations within a realization, for instance between r and f_{NL} .