

Probing T_R at the LHC with long-lived staus.

Koichi Hamaguchi

based on, M.Endo, KH, K.Nakaji, [arXiv:1008.2307]

+ M.Endo, KH, K.Nakaji, in preparation

+ S.Asai, KH, S.Shirai, [arXiv: 0902.3754] PRL,103,141803

COSMO/CosPA 2010 at Tokyo U., September, 2010

Reheating Temperature T_R

~ highest temperature in the rad. dom. universe.

= one of the most important parameters of cosmology.

time	temperature	
??		inflation
??	T_R	Reheating ⋮ (baryogenesis, dark matter?)
1 sec	1 MeV	Big Bang Nucleosynthesis ⋮ (recombination, structure formation,)
10^{10} yr	3 K	today

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Possible probes of T_R

- Gravitational Wave Nakayama, Saito, Suwa, Yokoyama,'08
- CMB Martin, Ringeval,'10 <-- talk by C.Ringeval on Monday
- BBN with long-lived charged massive particle (for low T_R). Takayama,'07
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Q: Any (indirect) hint from LHC ??

Main message of this talk:

The T_R can be determined at the LHC.

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If **SUSY + gravitino DM + stau NLSP**
is realized in nature,

The **T_R** can be determined at the **LHC**.

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At **7 TeV 1fb^{-1}** (\approx within 2 years),
 $T_R > \text{a few } 10^8 \text{ GeV}$ can be tested
in most of the parameter space!

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If **SUSY + gravitino DM + stau NLSP** is realized in nature,

The model dependent ?

... Yes, but the model itself is **testable** at the LHC.

signal: **metastable massive charged particle**

(cf. L.Roszkowski's talk on Monday)

At $\sqrt{s} = 7$

If observed, "gravitino DM + stau NLSP" is (one of) the best candidates.

$T_R >$

in mo

(The underlying supergravity may also be tested.

Buchmuller, KH, Ratz, Yanagida, '04)

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outline

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If **SUSY + gravitino DM + stau NLSP** is realized in nature,

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The **T_R** can be determined at the **LHC**.

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SUSY + gravitino DM + stau NLSP

.... why ?

scenario: SUSY + gravitino DM + stau NLSP

why SUSY?

- naturalness, coupling unification, DM,
- most non-SUSY scenarios for BSM → low E cut-off
→ how to discuss $T > \text{cut off} ??$: inflation/reheating/baryogenesis....

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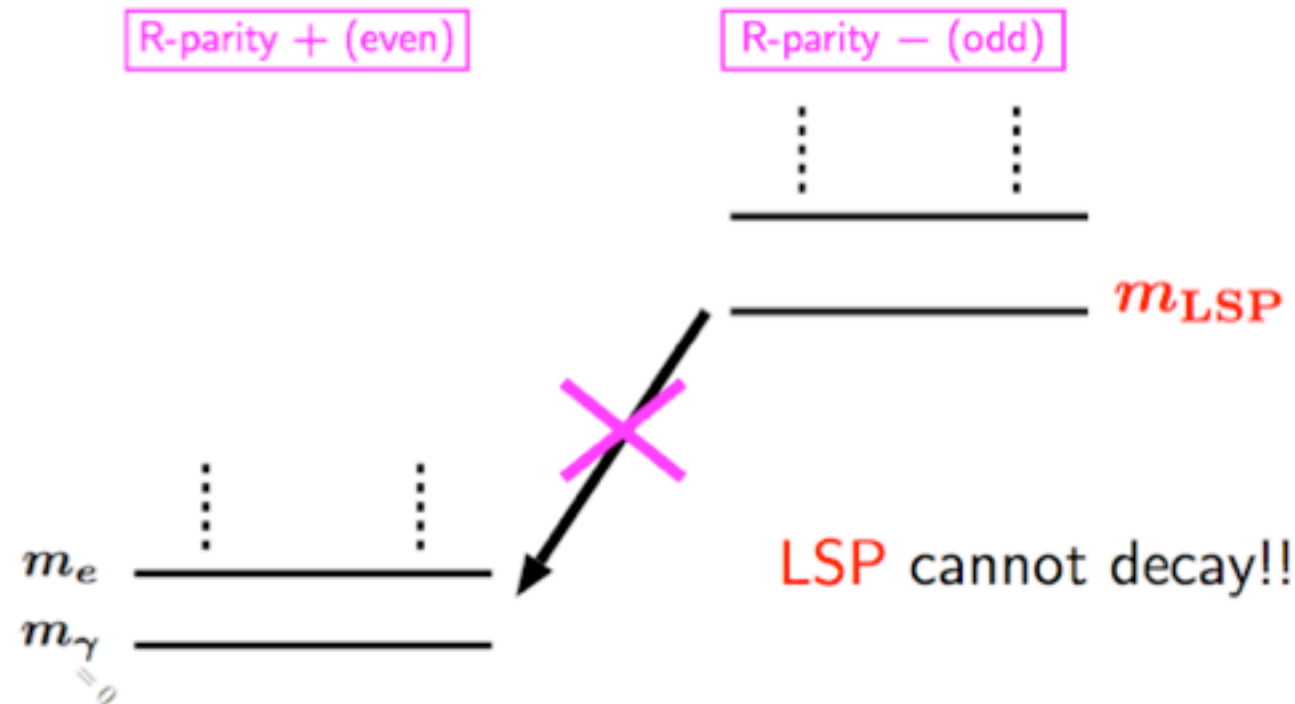
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In SUSY models + R-parity,
LSP (=Lightest SUSY Particle) is **stable**.



\rightarrow If neutral, **Dark Matter** candidate!

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$$\text{squarks : } \begin{pmatrix} \widetilde{u}_L \\ \widetilde{d}_L \end{pmatrix}_i \quad \widetilde{u}_{Ri} \quad \widetilde{d}_{Ri} \quad \text{sleptons : } \begin{pmatrix} \widetilde{\nu}_L \\ \widetilde{e}_L \end{pmatrix}_i \quad \widetilde{e}_{Ri}$$

$$\text{gauginos and higgsinos : } \widetilde{\chi}_i^0, \quad \widetilde{\chi}_i^\pm, \quad \widetilde{g}$$

$$\text{gravitino : } \widetilde{G}$$

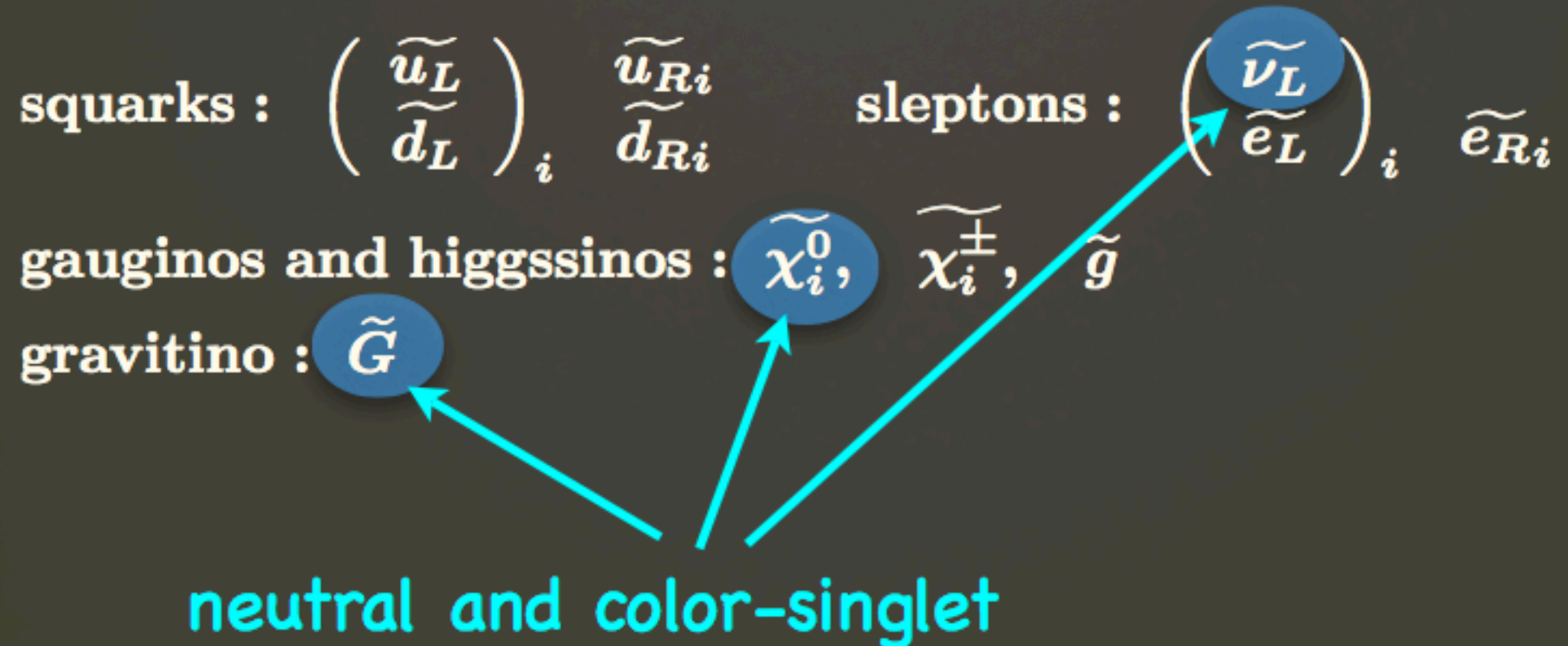
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gravitino : \widetilde{G}

~~excluded by direct
detection experiments
(cf. Falk, Olive, Srednicki, '94)~~

neutral and color-singlet

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DM candidates in minimal SUSY model

→ only **gravitino or neutralino** are allowed.

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Note:

another interesting possibility: $O(\text{eV})$ gravitino + composite DM.

Nakamura, KH, Shirai, Yanagida, '09

Nakamura, Shirai, in preparation

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why gravitino

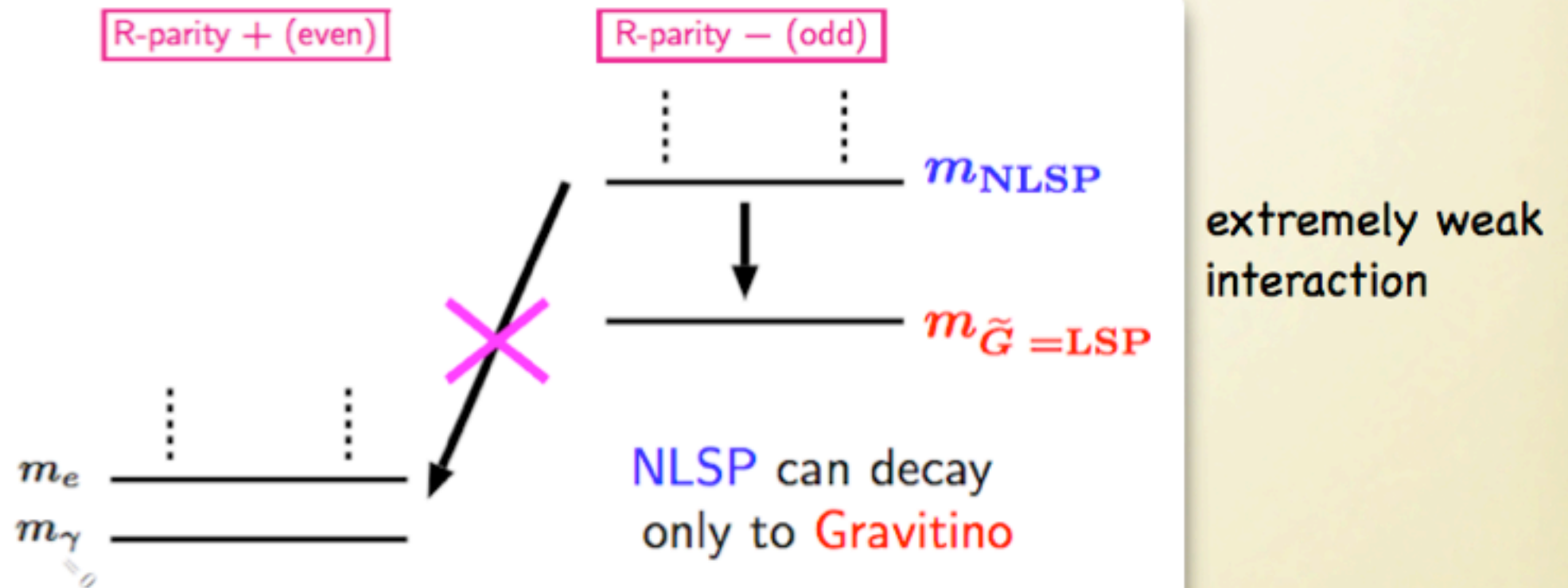
DM candidate

→ only **gr**

why stau

NLSP (Next-to-Lightest SUSY Particle)

In **Gravitino LSP** scenario, the **NLSP** is long-lived.



Lifetime e.g. for $m_{NLSP} \simeq 200 \text{ GeV}$

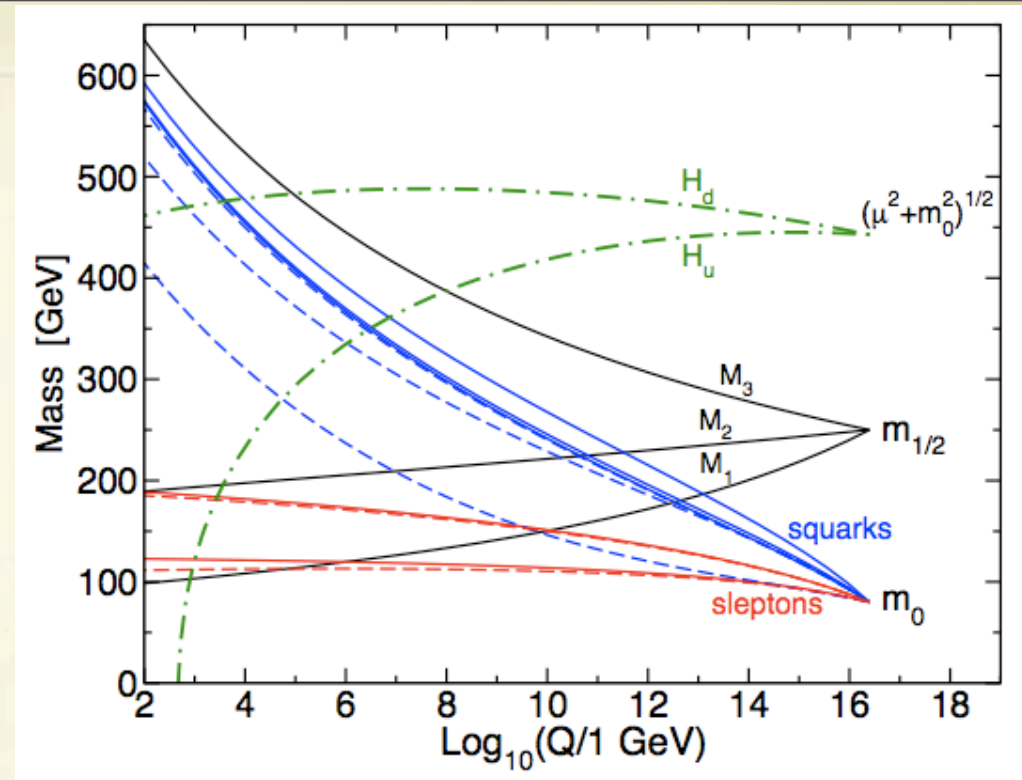
$\tau_{NLSP} \sim \mathcal{O}(\text{day})$ for $m_{\tilde{G}} \sim 10 \text{ GeV}$

$\tau_{NLSP} \sim \mathcal{O}(10 \text{ min})$ for $m_{\tilde{G}} \sim 1 \text{ GeV}$

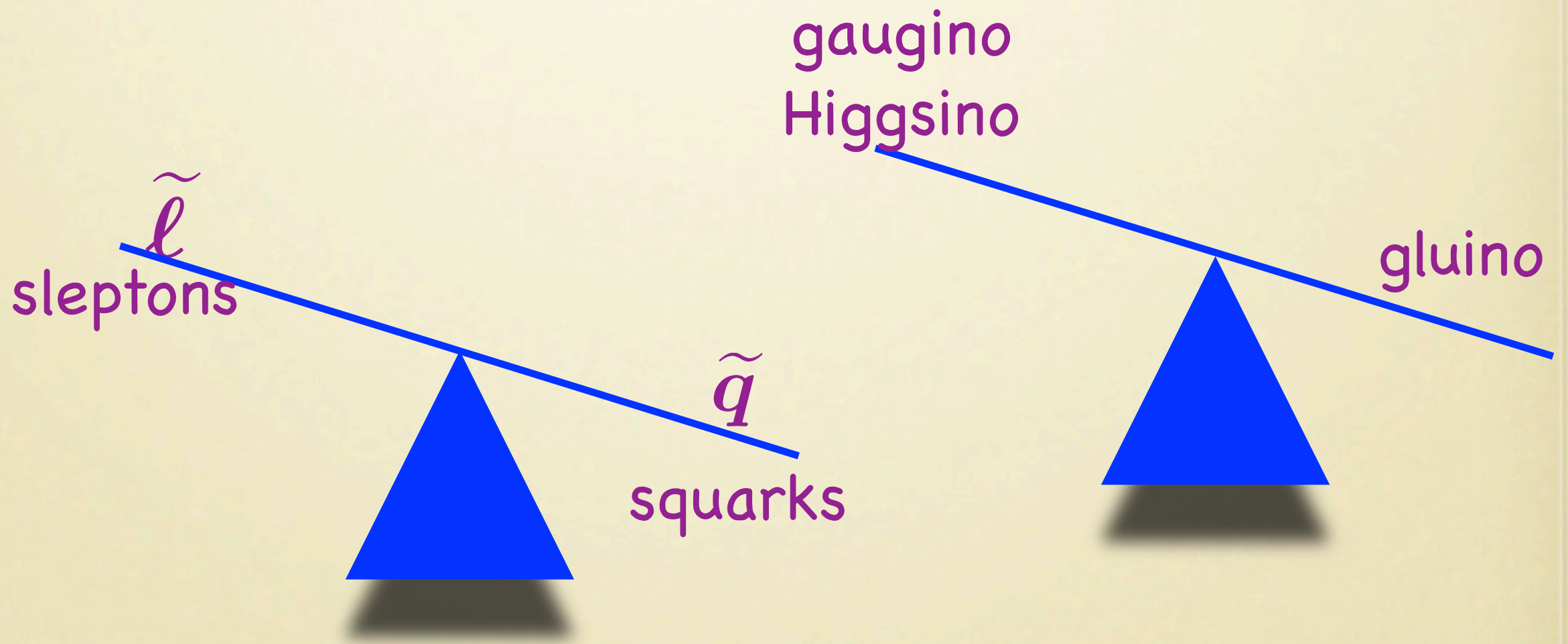
$\tau_{NLSP} \sim \mathcal{O}(10 \text{ sec})$ for $m_{\tilde{G}} \sim 0.1 \text{ GeV}$

- Why Stau NLSP ?

- In general, from RGE, tendency is
 - $M(\text{color singlet}) < M(\text{colored})$

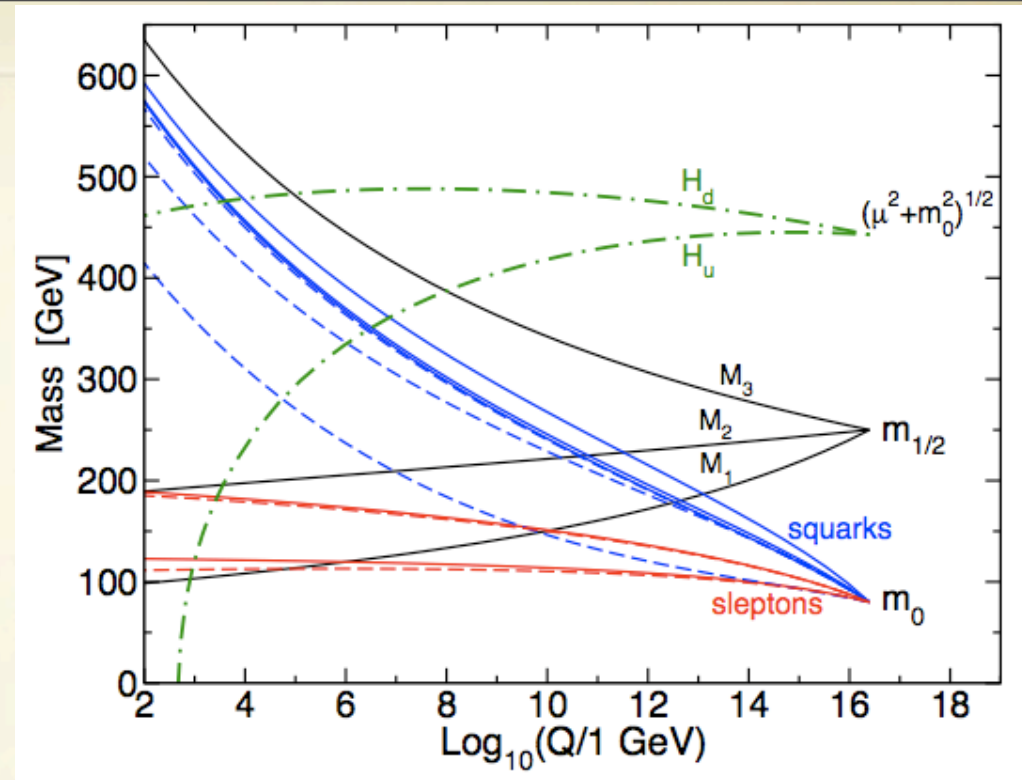


typical RG evolution (from S.P.Martin, hep-ph/9709356)

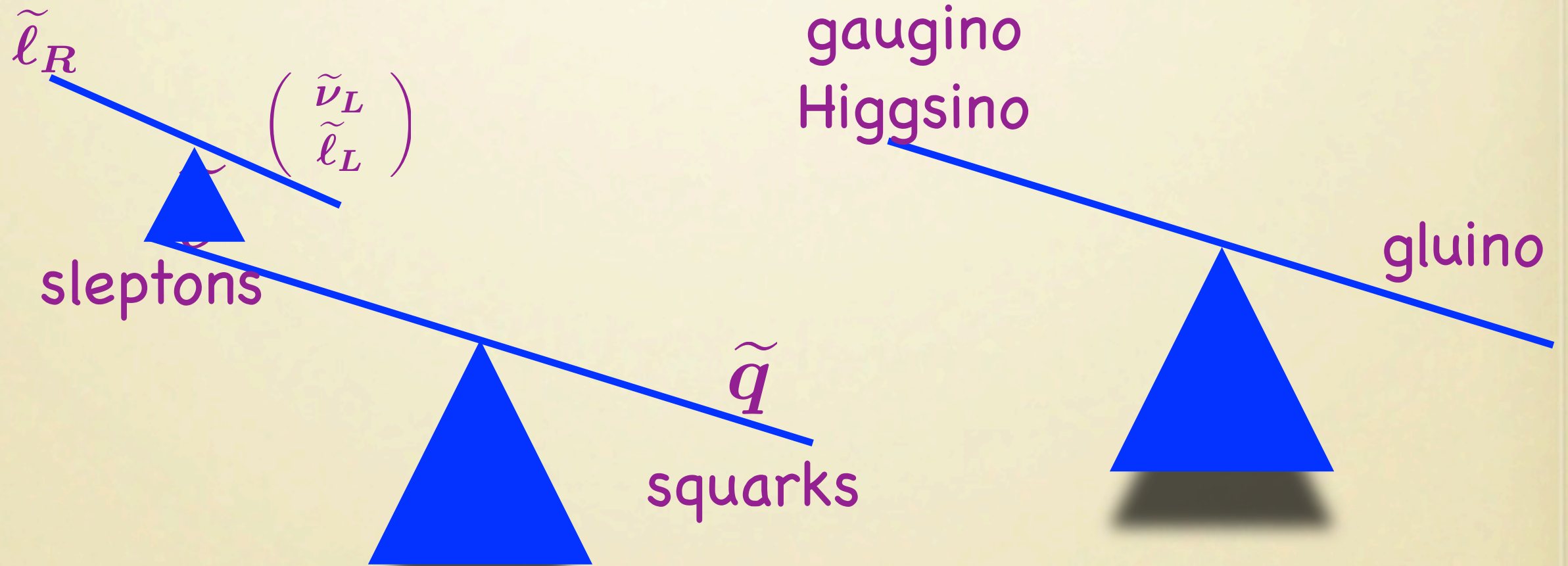


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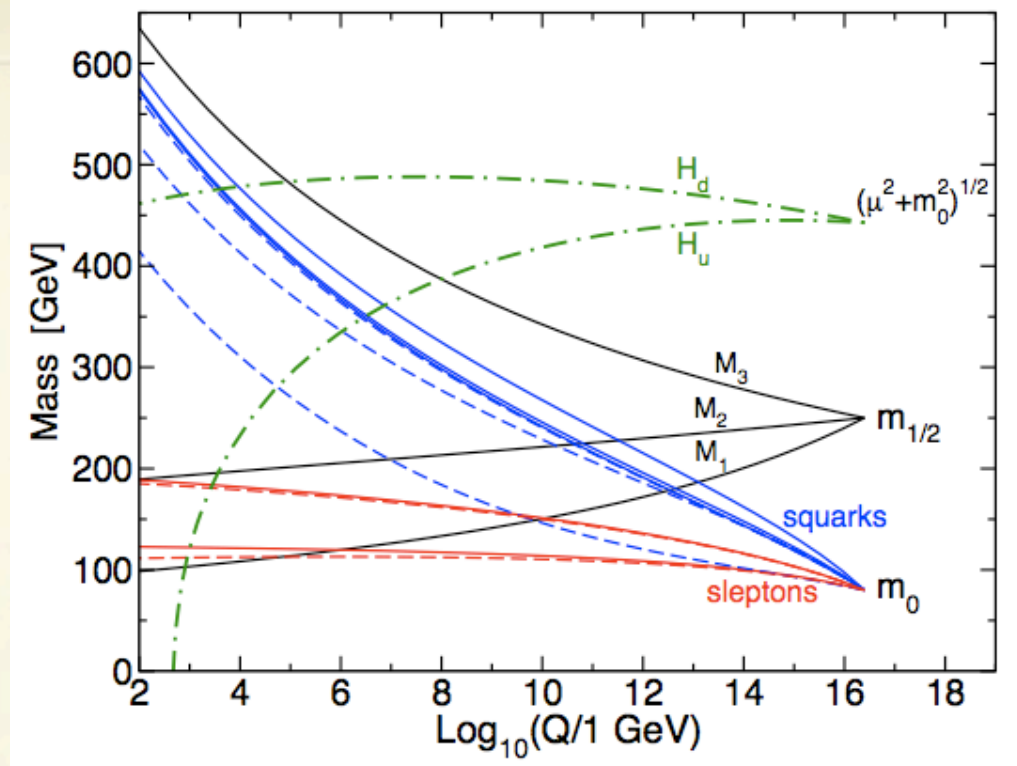


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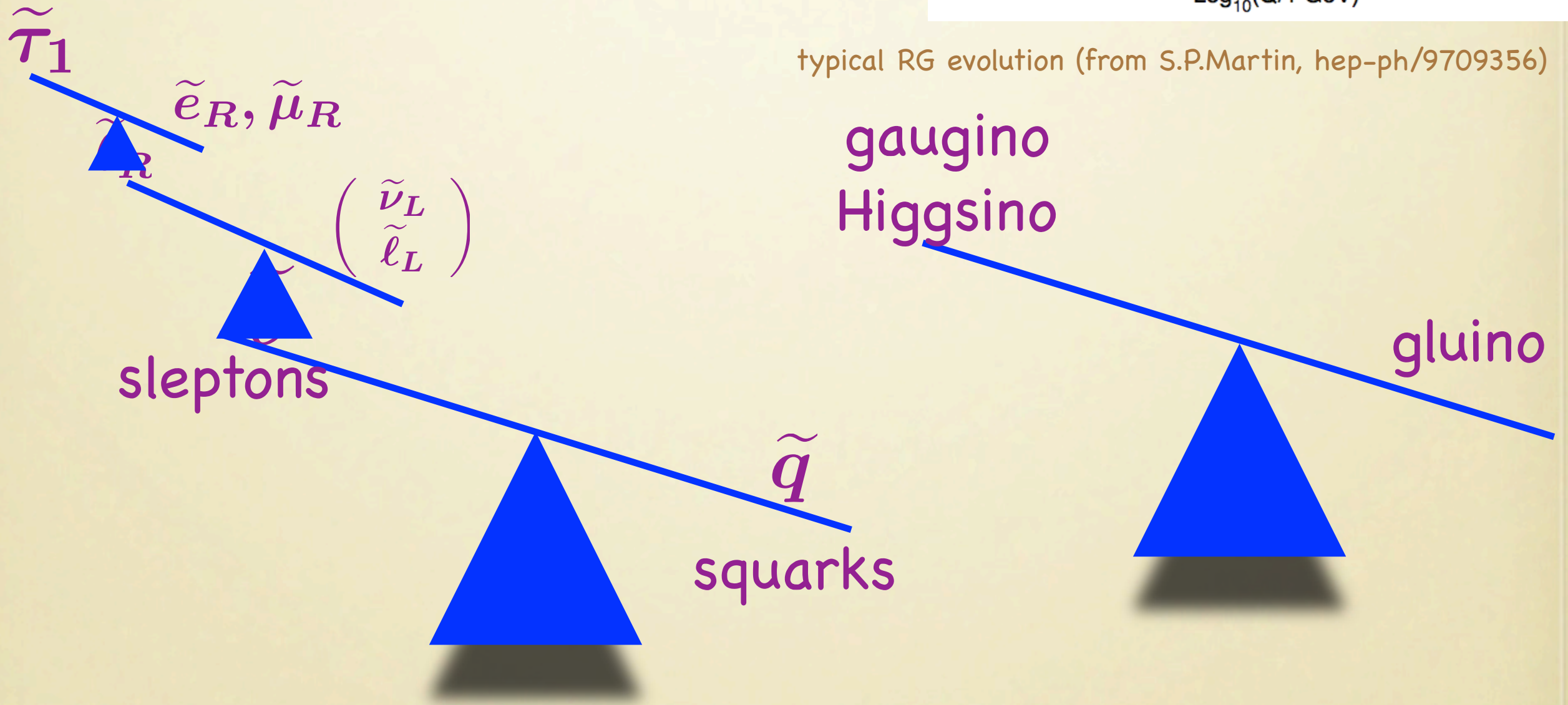


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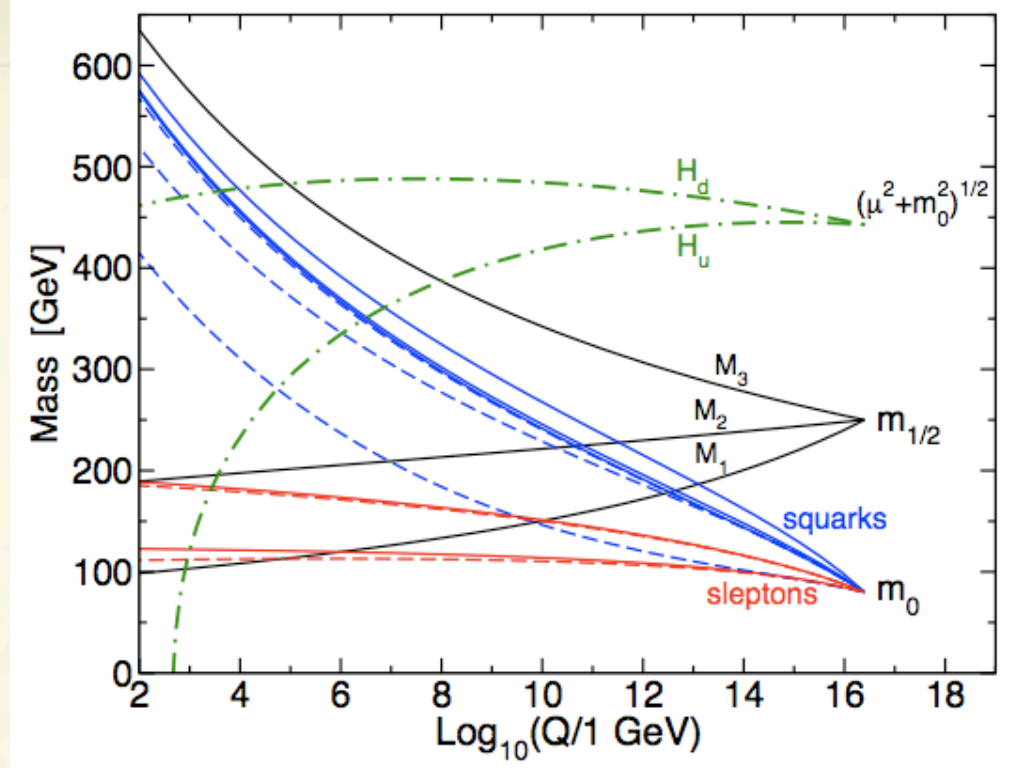


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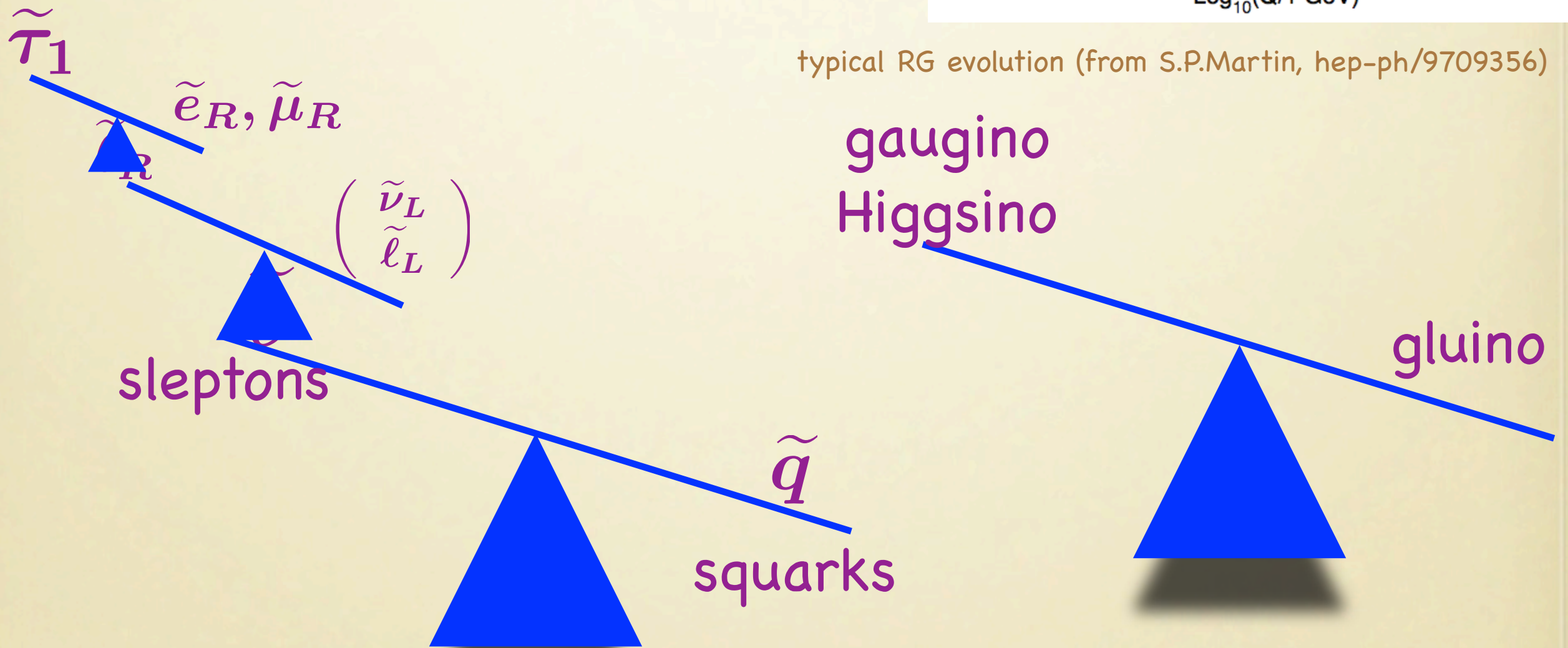


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- In most cases, either **Stau** or **Neutralino** is the NLSP

scenario: SUSY + gravitino DM + stau NLSP

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→ **stau = long-lived charged particle.**

important for cosmology and collider.

2

T_R determination at the LHC with long-lived staus.

.... in gravitino DM scenario with stau NLSP.

M.Endo, KH, K.Nakaji, in progress
+ S.Asai, KH, S.Shirai,'09

See also related works:
Choi, Roszkowski, Ruiz De Austri,'07
Steffen,'08

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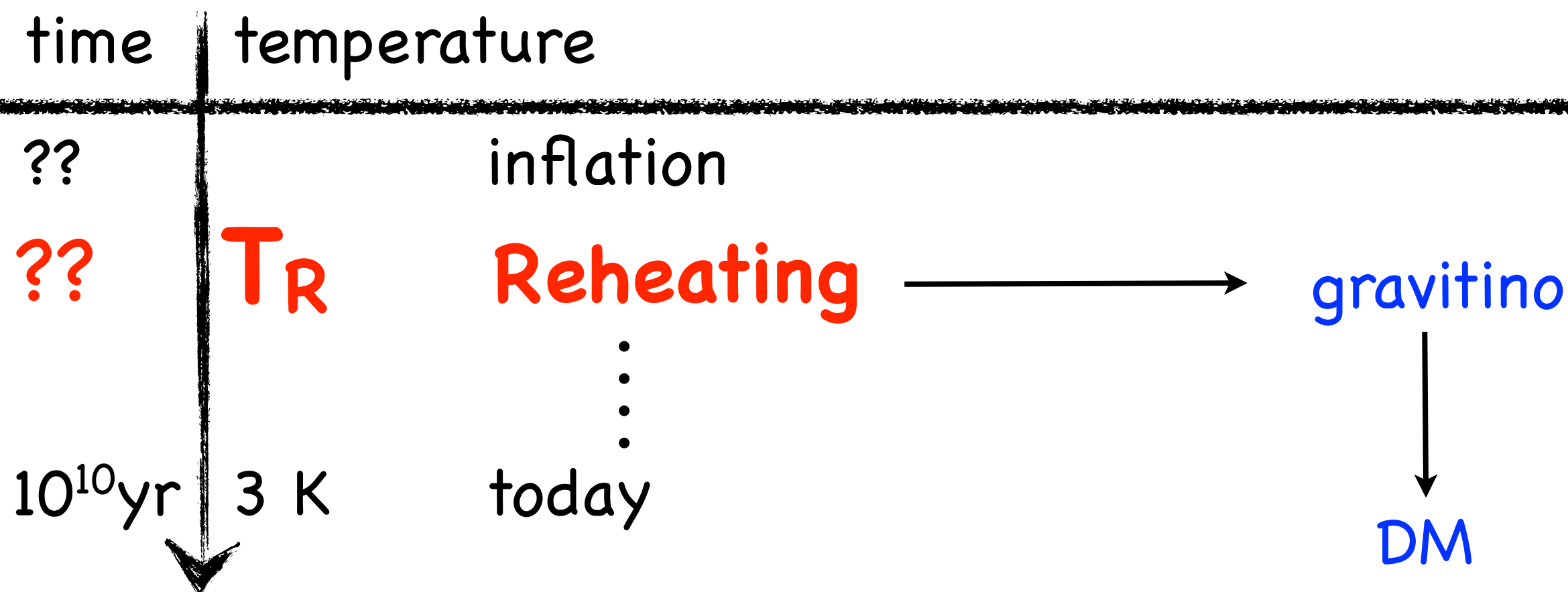
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POINT: gravitino abundance is determined by T_R



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$$\Omega_{\tilde{G}} h^2 \simeq 0.1 \left(\frac{3 \text{ GeV}}{m_{\tilde{G}}} \right) \left(\frac{m_{\text{gluino}}}{1 \text{ TeV}} \right)^2 \left(\frac{T_R}{10^8 \text{ GeV}} \right)$$

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assumption

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step 2

measure stau mass


step 3

measure stau lifetime

step 4

measure gluino mass

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T_R determination at the LHC with long-lived staus.

$$\Omega_{\tilde{G}} h^2$$

heavy charged particle
(like muon)
low velocity (\neq muon)

cf. Talk by S.Asai

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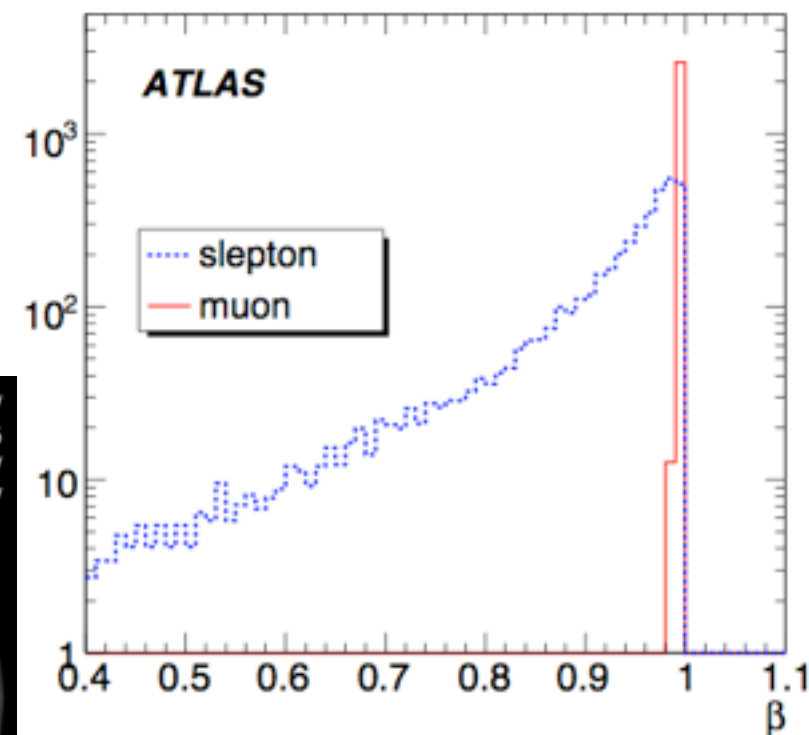
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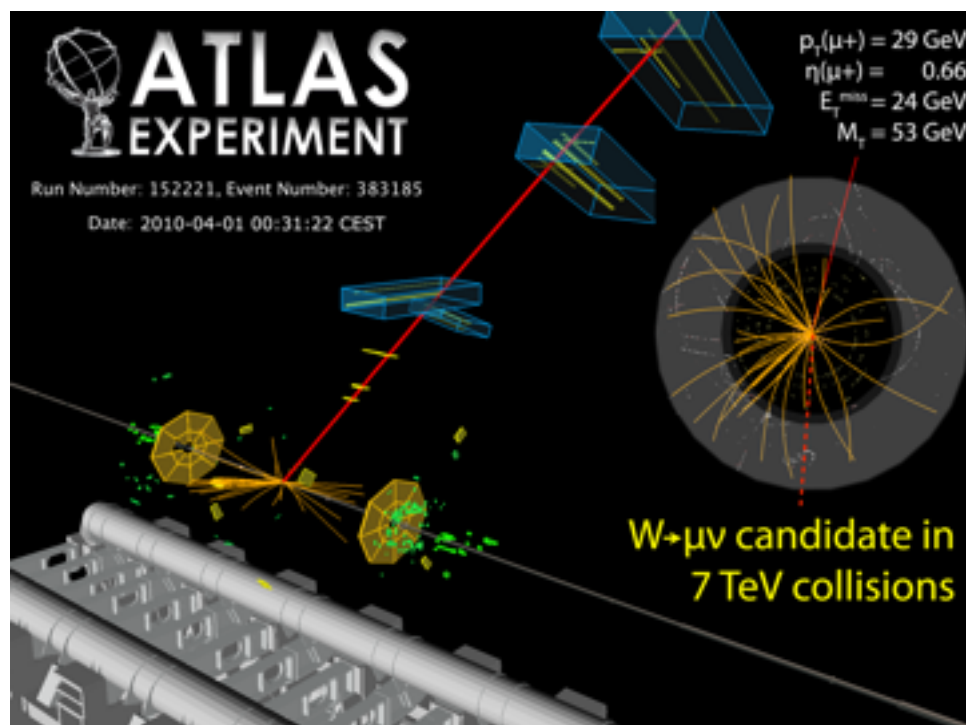
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ATLAS, 0901.0512



<http://atlas.web.cern.ch/Atlas/public/EVTDISPLAY/events.html>

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T_R determ

by velocity measurement
(+ momentum measurement)

$$\text{mass} = p/(\beta\gamma)$$

5.

Ω

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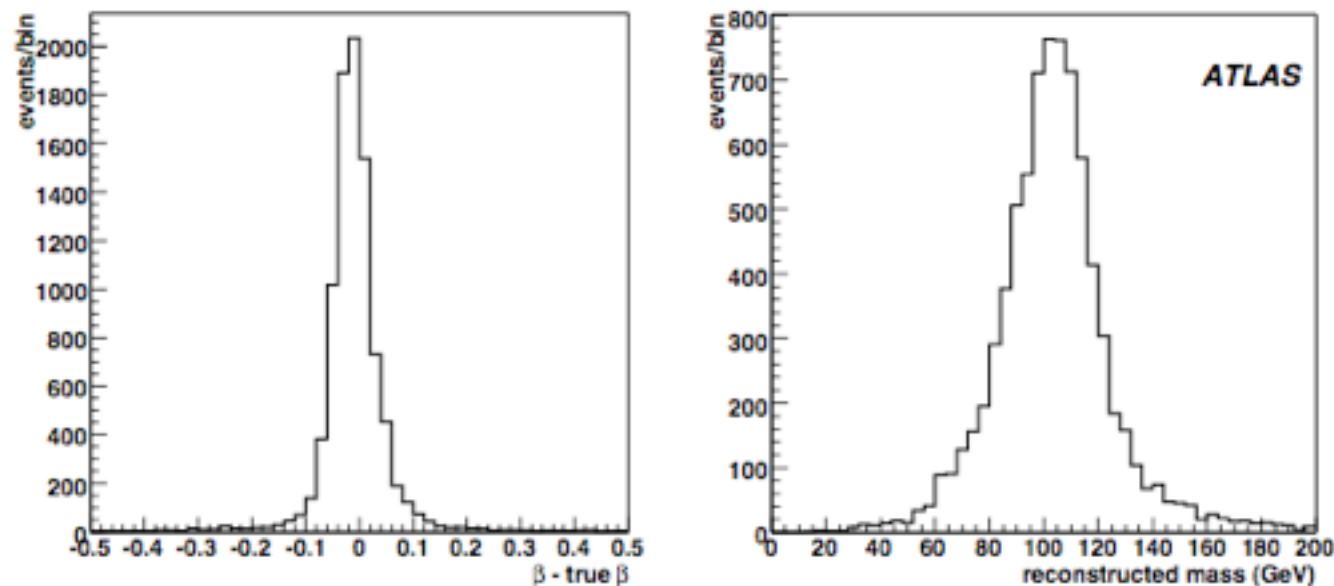


Figure 18: β resolution and reconstructed mass for sleptons from the GMSB5 sample.

ATLAS, 0901.0512

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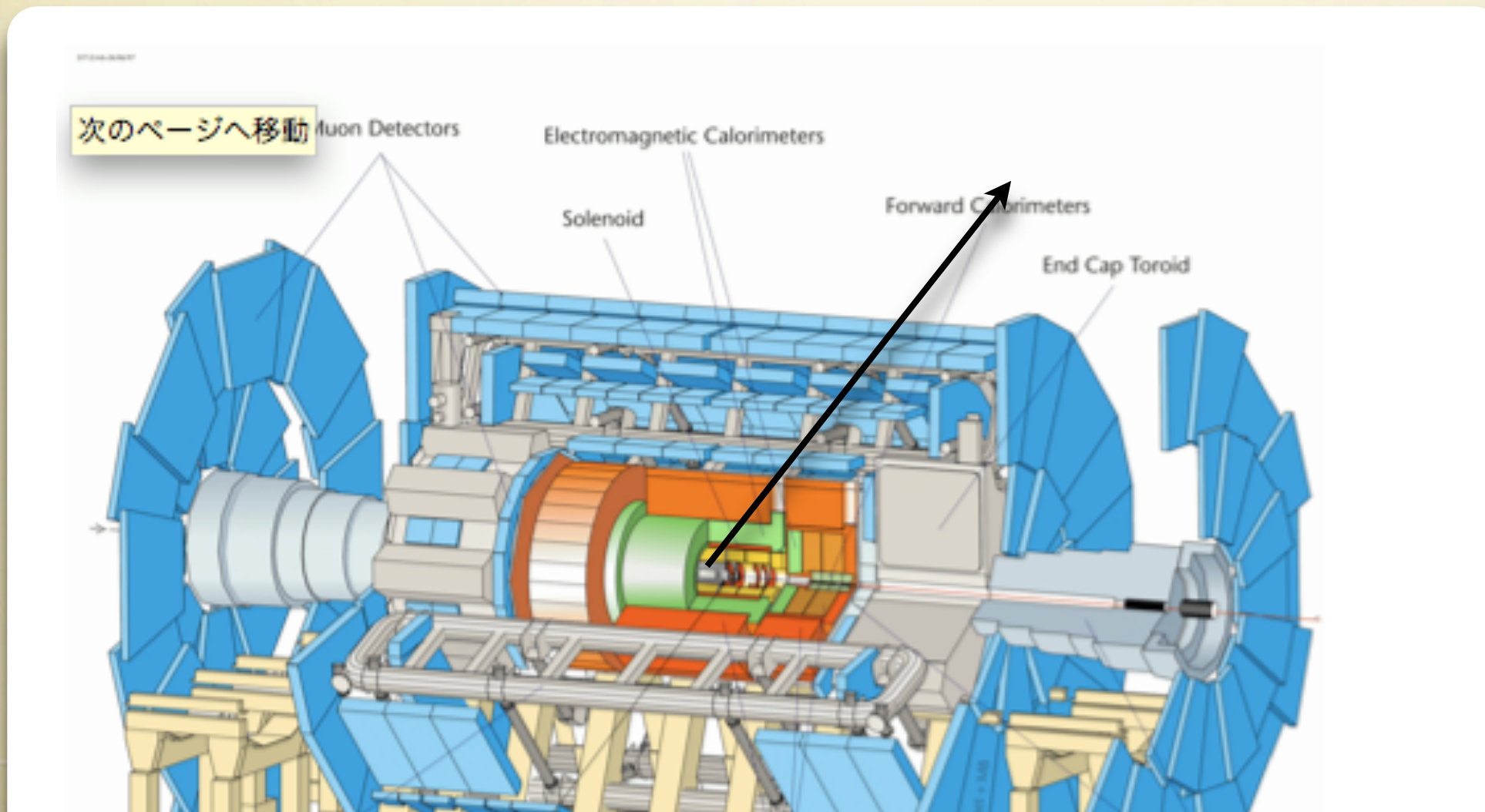
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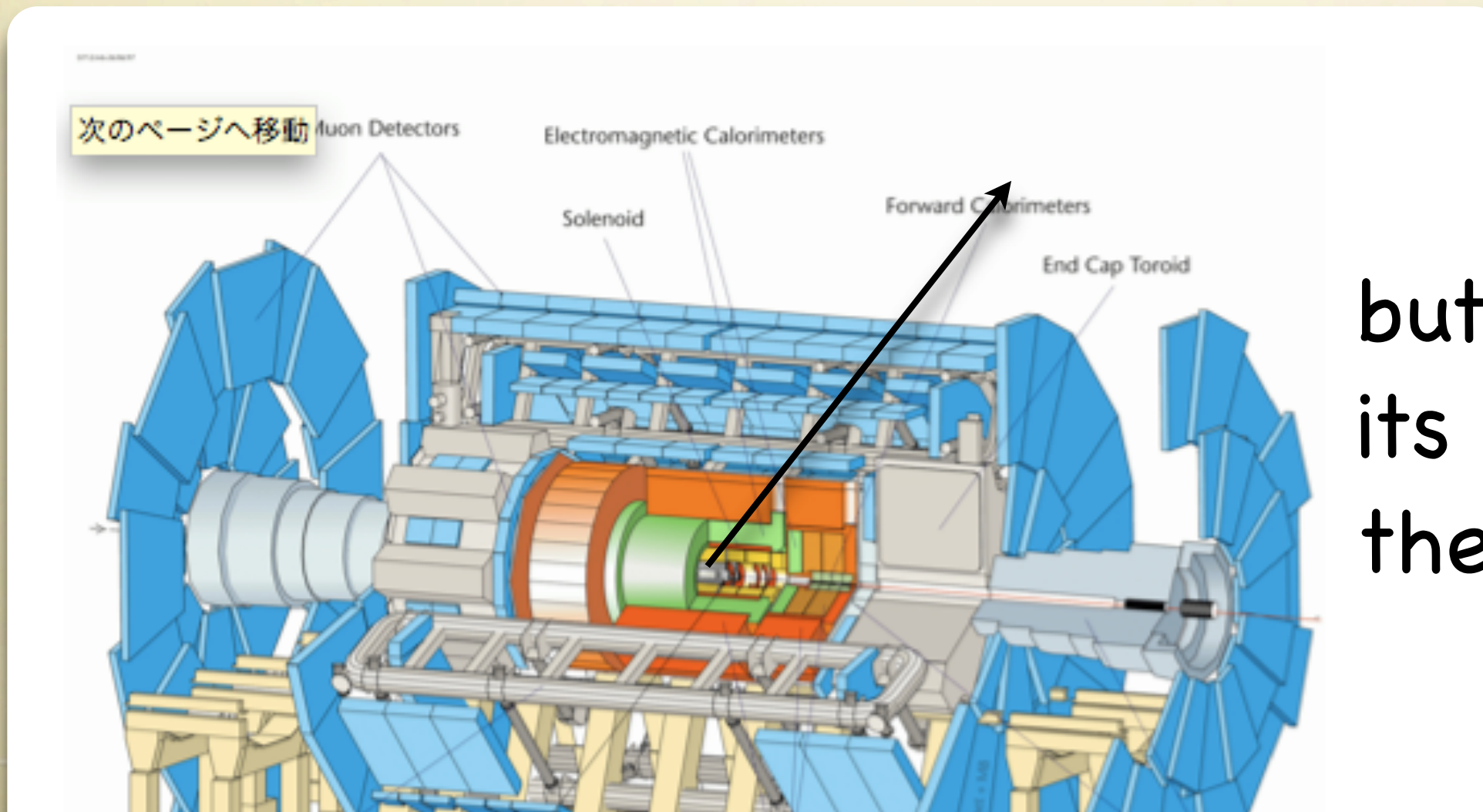
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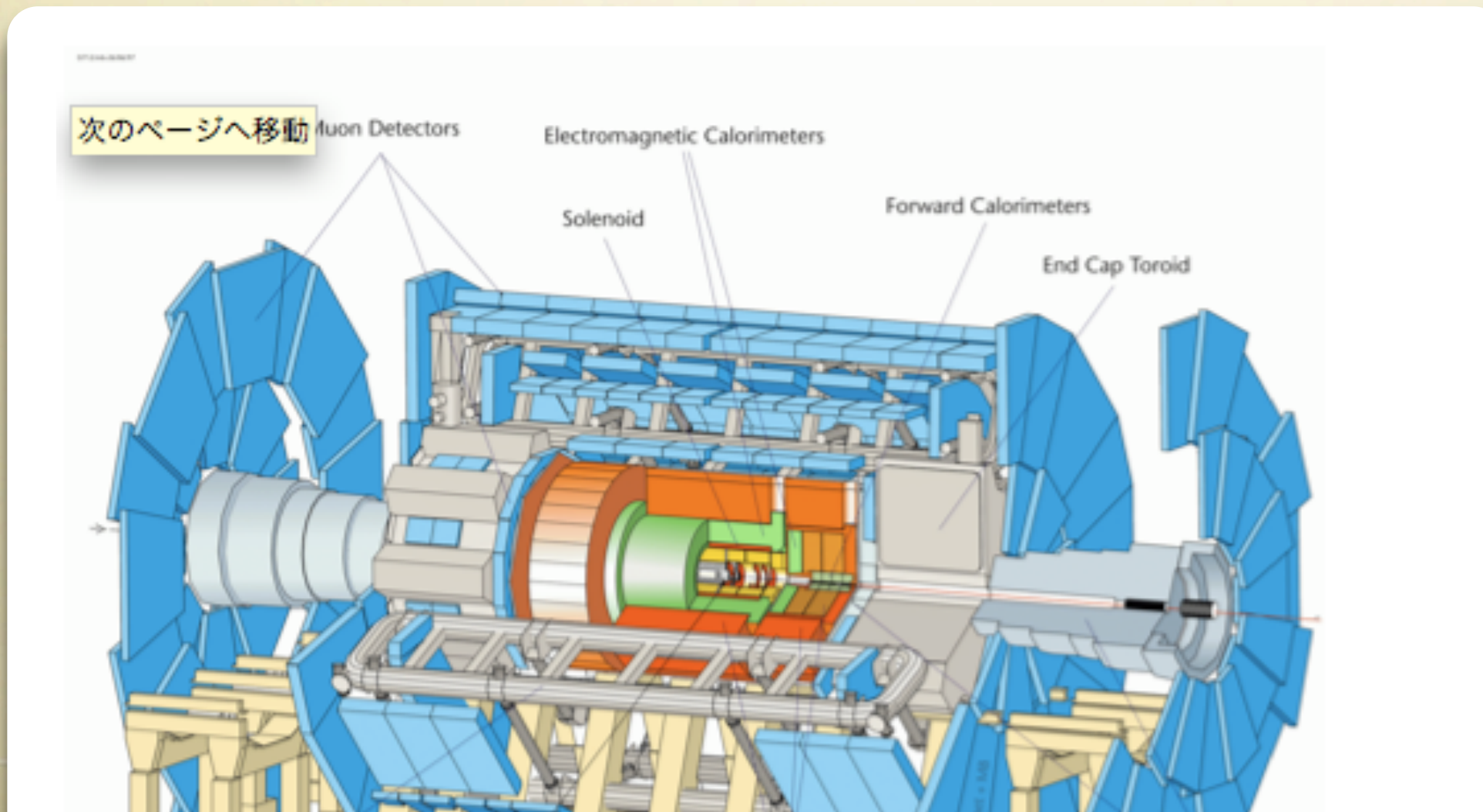
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but we can't see
its decay in
these events....

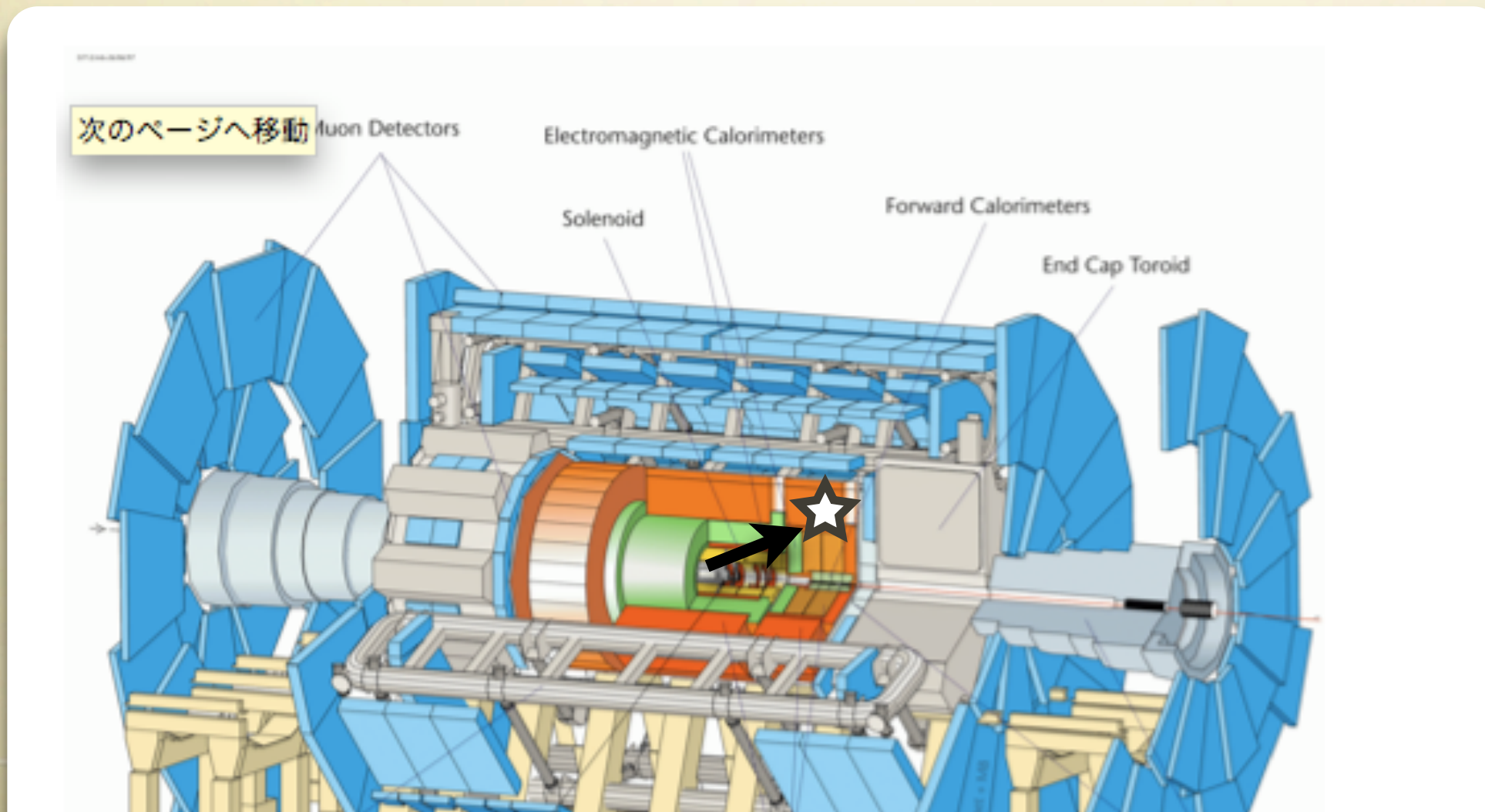
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- but some of them have sufficiently small velocity and stop at calorimeters.



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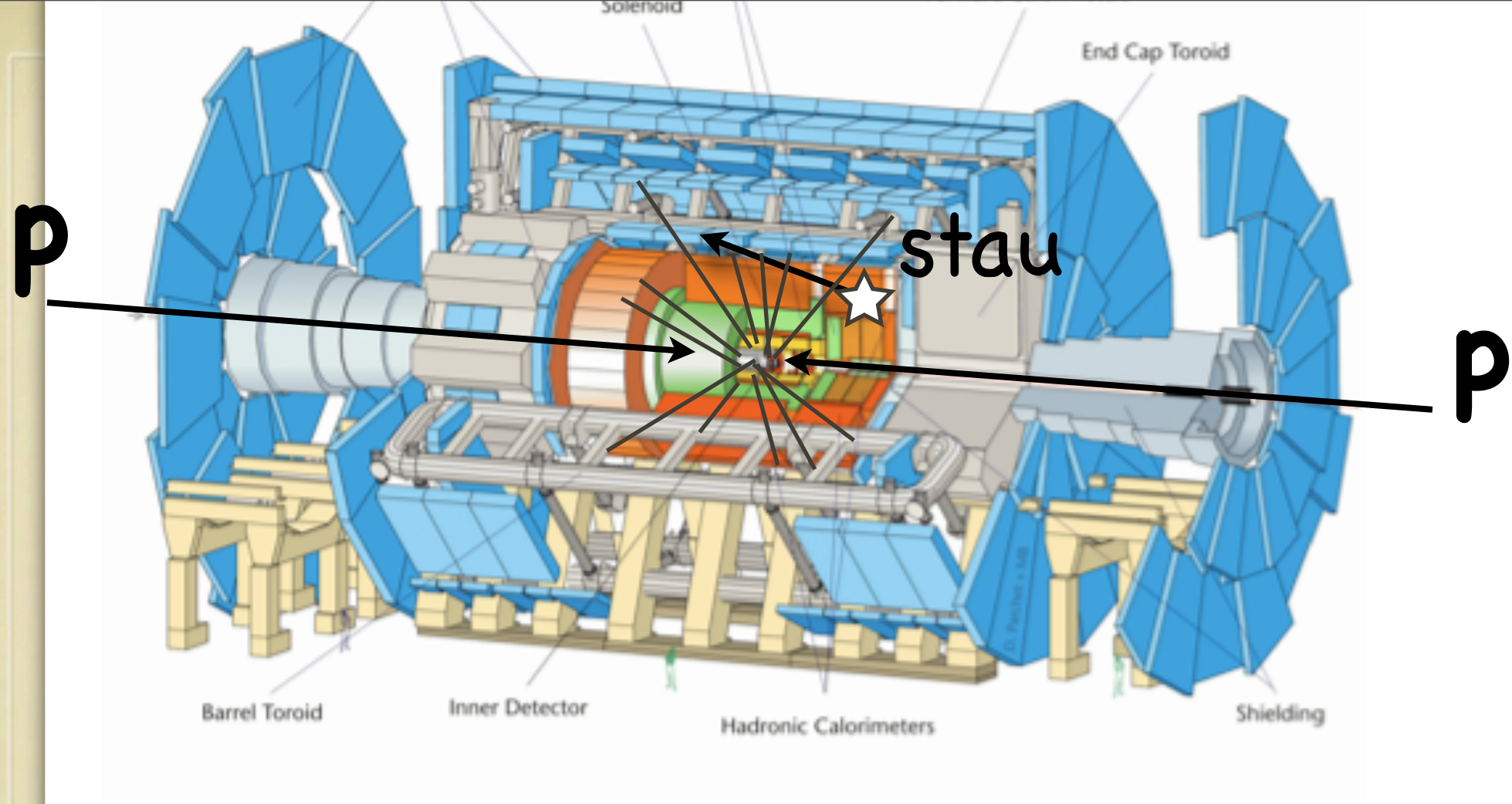


Figure 1-1 Overall layout of the ATLAS detector.

- but their late-time decay has **wrong timing** and **wrong direction**;
- **difficult** to reject **backgrounds**
- **difficult** to **trigger**.

..... during pp collision.

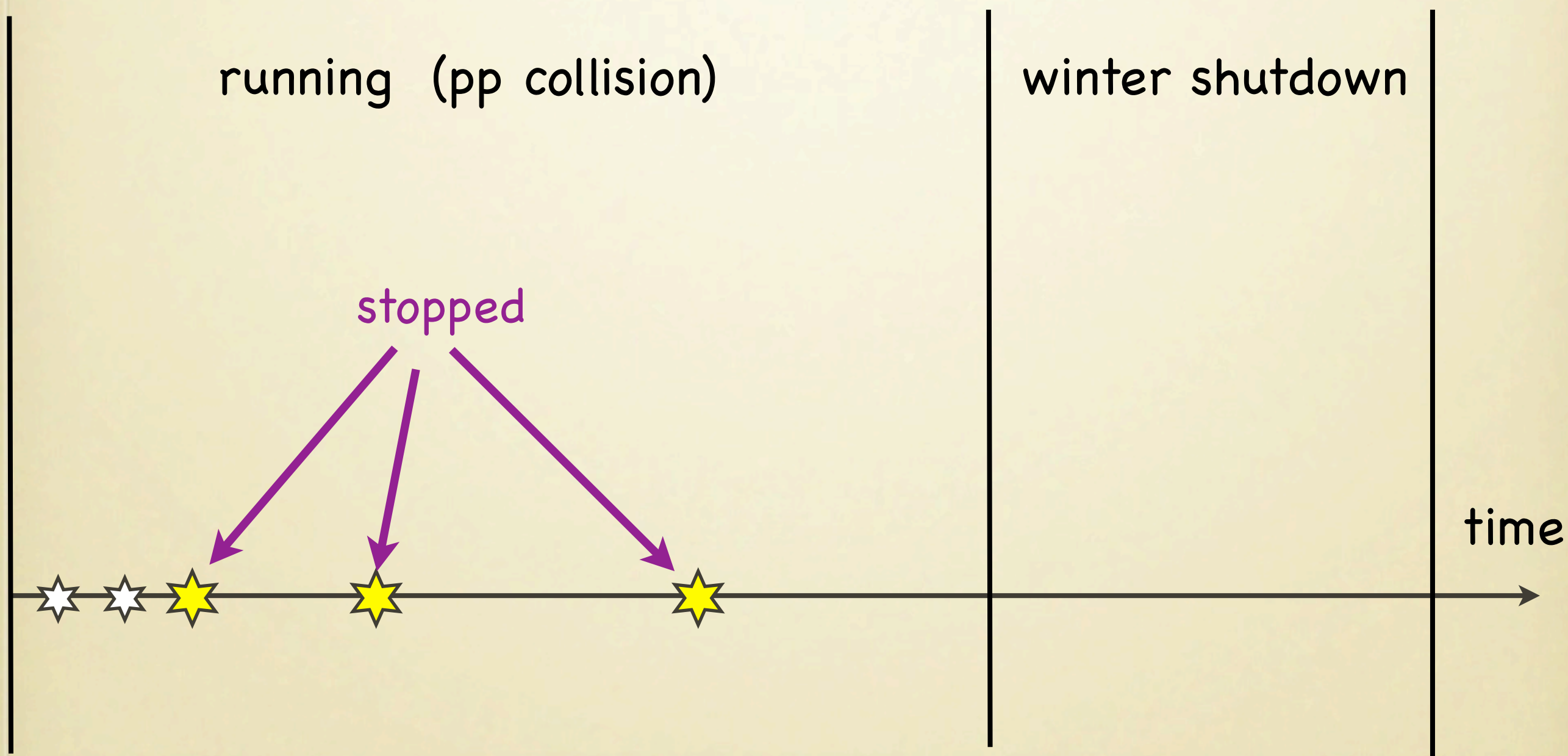
Idea:

use periods of no pp collision !!

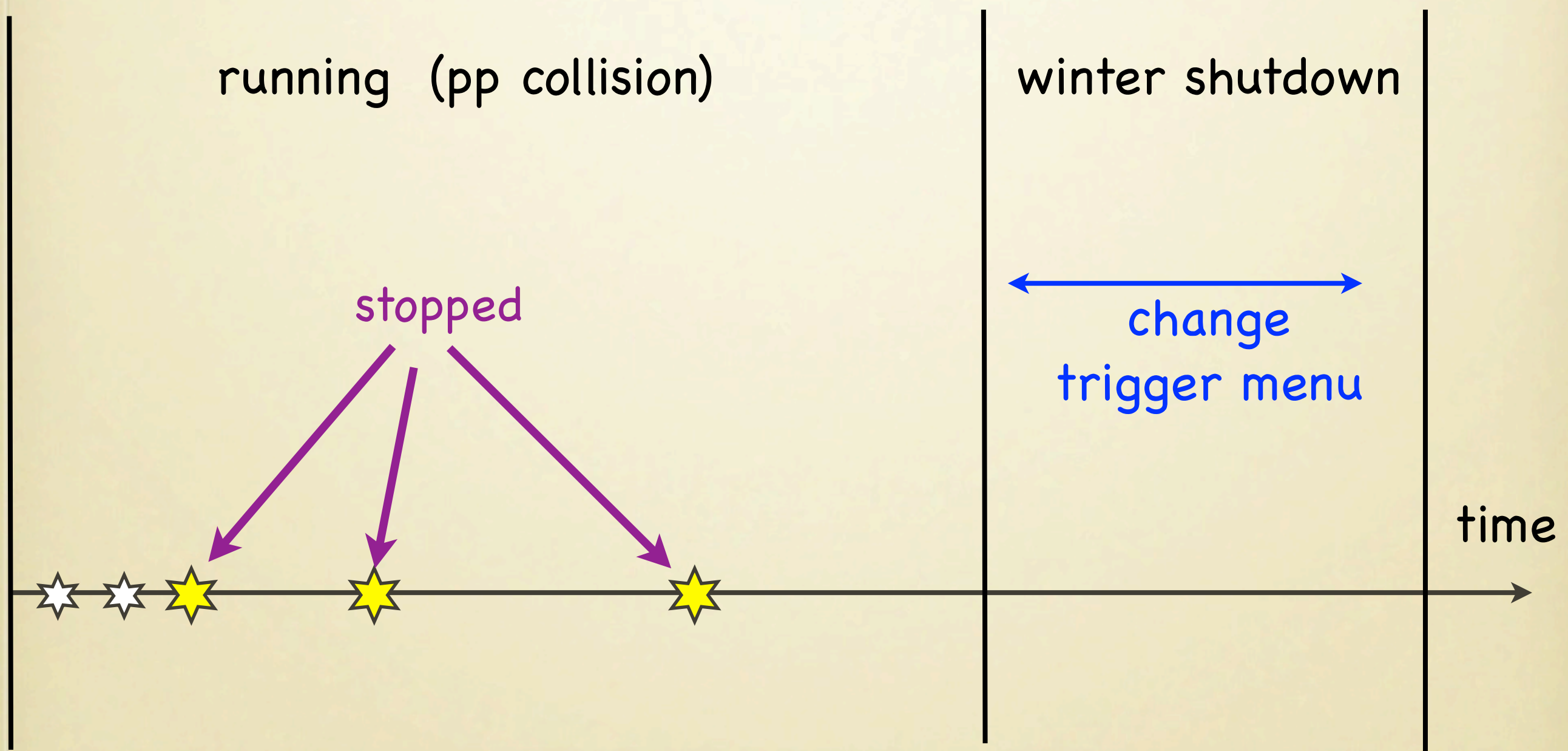
two possible strategies:

- for long lifetime: use shutdown time.
- for short lifetime: use beam-dump signal.
(or use empty bunch.)

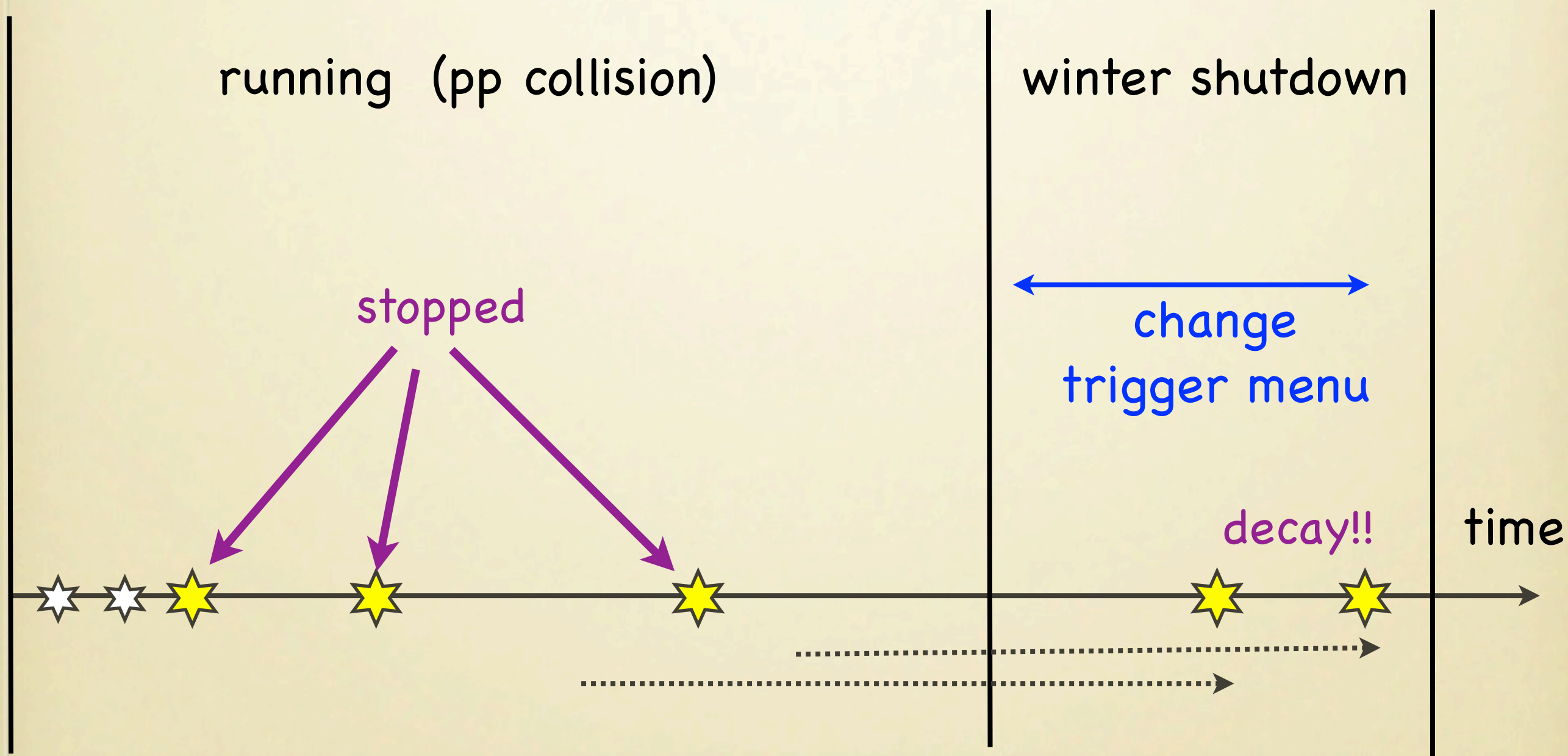
- for long lifetime: use shutdown time



- for long lifetime: use shutdown time

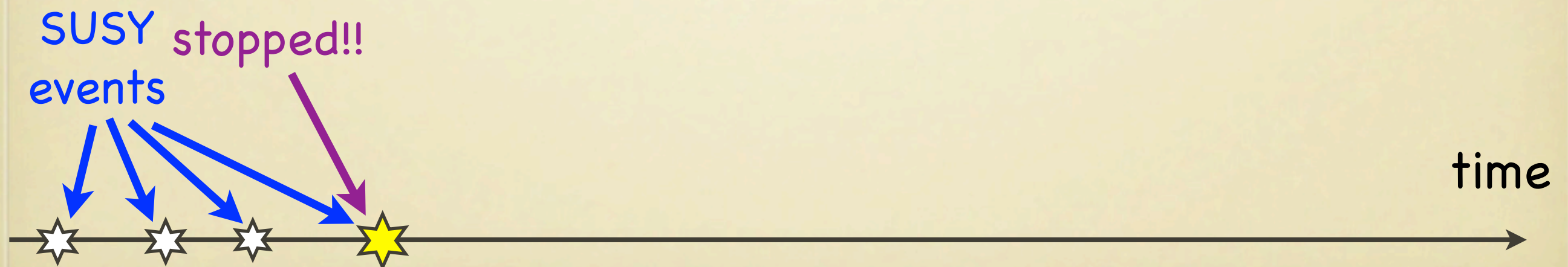


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- for short lifetime: use **beam-dump signal**.

(I) select the stopping event by **online Event Filter**.



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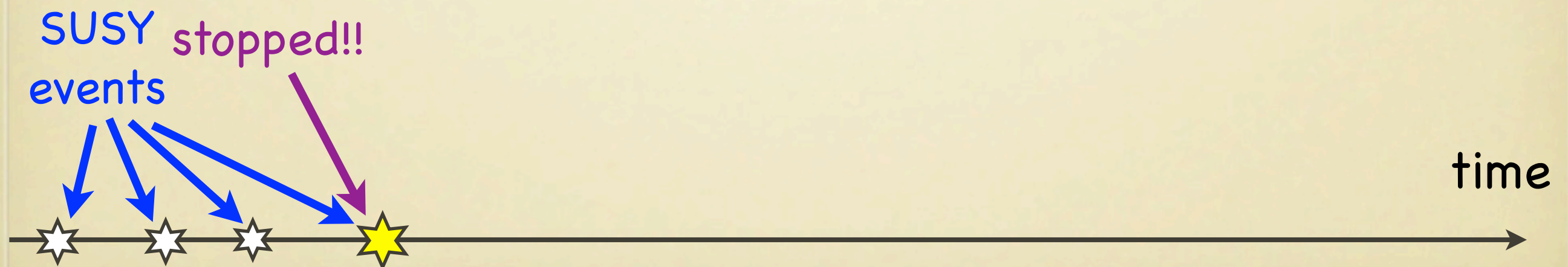
(1) missing $ET > 100 \text{ GeV}$

(2) 1 jet $PT > 100 \text{ GeV}$ + 2 jets $PT > 50 \text{ GeV}$

(3) isolated track with $PT > 0.1 \text{ m(stau)}$.

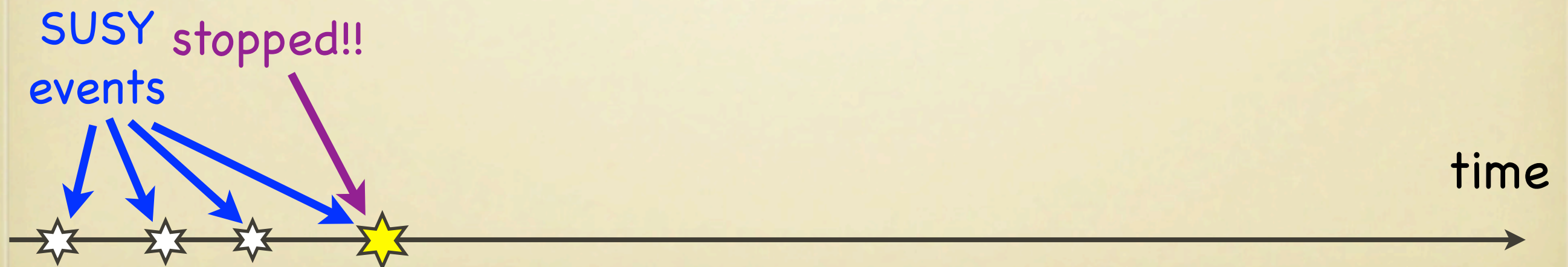
(4) extrapolate the track to calorimeter and energy deposit $< 0.2 \text{ p(stau)}$.

(5) extrapolate the track to muon system and no muon track.



- for short lifetime: use **beam-dump signal**.

(I) select the stopping event by **online Event Filter**.



- for short lifetime: use **beam-dump signal**.

(I) select the stopping event by **online Event Filter**.

(II) send a **beam-dump signal**, which immediately stops the pp collision.

trigger



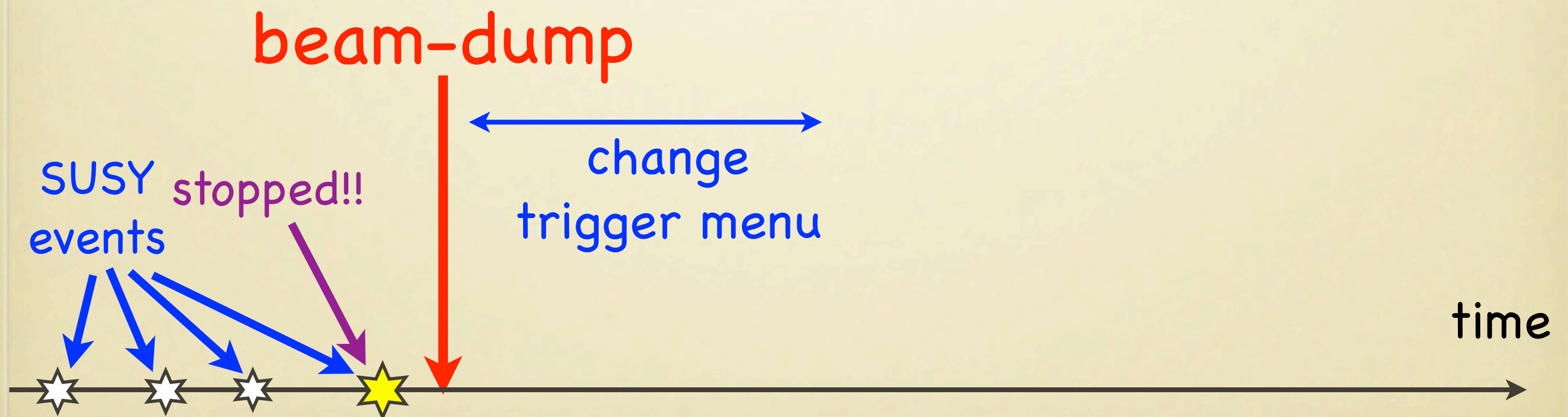
- for short lifetime: use **beam-dump signal**.

(I) select the stopping event by **online Event Filter**.

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(III) **change the trigger menu** to the one optimized for stau decay.

trigger



- for short lifetime: use **beam-dump signal**.

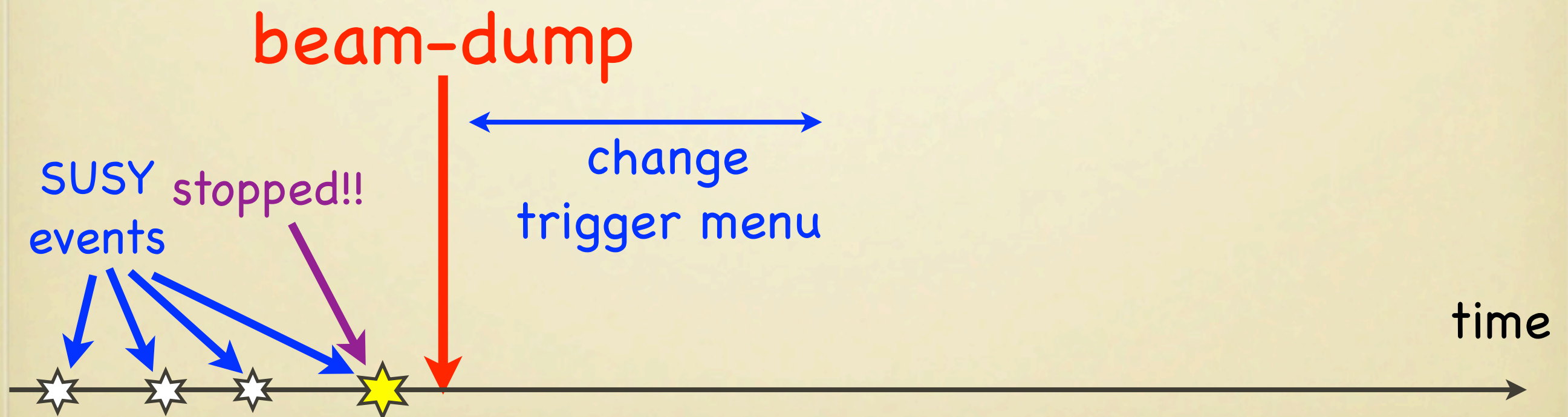
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(IV) **wait** for stau decay.

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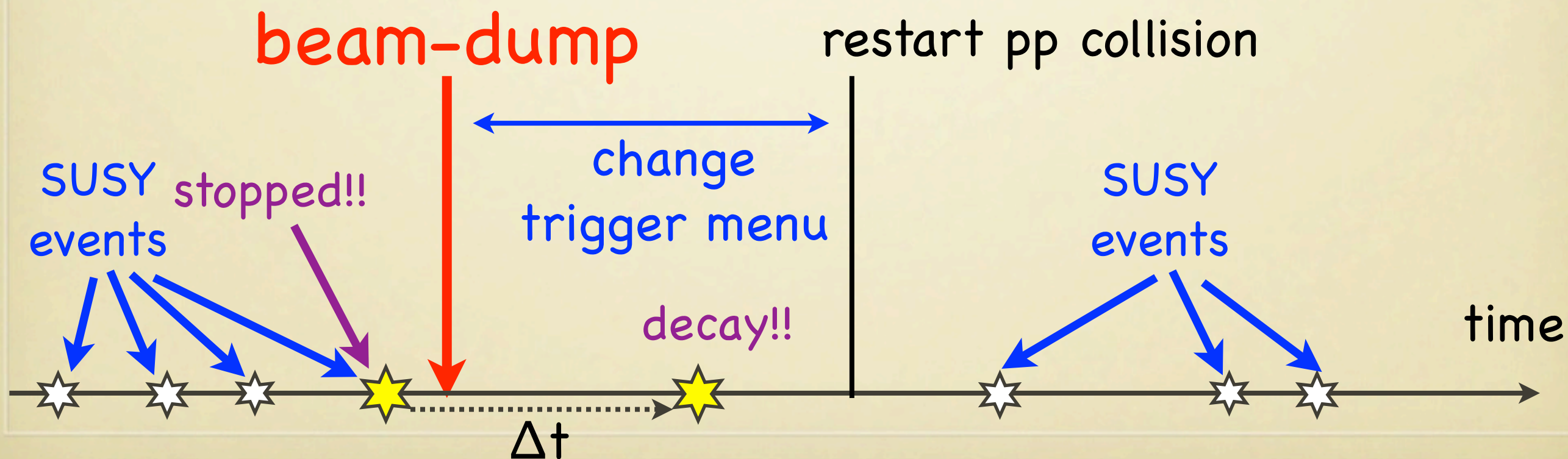
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trigger



lifetime measurement: Result

TABLE III: Expected statistical errors for each lifetime. $\langle N_D \rangle$ is the expected number of staus' decays in the corresponding period. (SPS7 point)

lifetime	10 fb^{-1}		100 fb^{-1}	
	$\langle N_D \rangle$	σ	$\langle N_D \rangle$	σ
0.1 sec	0.01	$\pm 0.1 \text{ sec}$	0.1	$\pm 0.1 \text{ sec}$
0.2 sec	1.8	$\pm 0.15 \text{ sec}$	18	$\pm 0.05 \text{ sec}$
0.5 sec	35	$\pm 0.1 \text{ sec}$	352	$\pm 0.03 \text{ sec}$
1 sec	96	$\pm 0.1 \text{ sec}$	956	$\pm 0.04 \text{ sec}$
10 sec	235	$\pm 0.7 \text{ sec}$	2353	$\pm 0.2 \text{ sec}$
100 sec	257	$\pm 7 \text{ sec}$	2574	$\pm 2.0 \text{ sec}$
1000 sec	217	$^{+180}_{-140} \text{ sec}$	2168	$\pm 51 \text{ sec}$
10 day	26	$\pm 2.2 \text{ day}$	262	$\pm 0.7 \text{ day}$
100 day	143	$^{+49}_{-25} \text{ day}$	1430	$^{+20}_{-13} \text{ day}$
10 year	14	$^{+7}_{-3} \text{ year}$	138	$^{+1.6}_{-1.2} \text{ year}$
50 year	2.8	$^{+110}_{-21} \text{ year}$	28	$^{+21}_{-12} \text{ year}$
300 year	0.5	—	5	$^{+224}_{-88} \text{ year}$

short

assumption

dead time: 1 sec
waiting time: 30 min.

running: 200 days
shutdown: 100 days

long

$O(0.1 \text{ sec} \dots 100 \text{ years})$ can be probed!!

T_R determination at the LHC with long-lived staus.

$$\Omega_{\tilde{G}} h^2 \simeq 0.1 \left(\frac{3 \text{ GeV}}{m_{\tilde{G}}} \right) \left(\frac{m_{\text{gluino}}}{1 \text{ TeV}} \right)^2 \left(\frac{T_R}{10^8 \text{ GeV}} \right)$$

step 1

see staus at the LHC

step 2

measure stau mass

step 3

measure stau lifetime

step 4

measure gluino mass

T_R determination at the LHC with long-lived staus.

$$\Omega_{\tilde{G}} h^2 \simeq 0.1 \left(\frac{3 \text{ GeV}}{m_{\tilde{G}}} \right)$$

by invariant mass method
[cf. Ito, Kitano, Moroi, '09]

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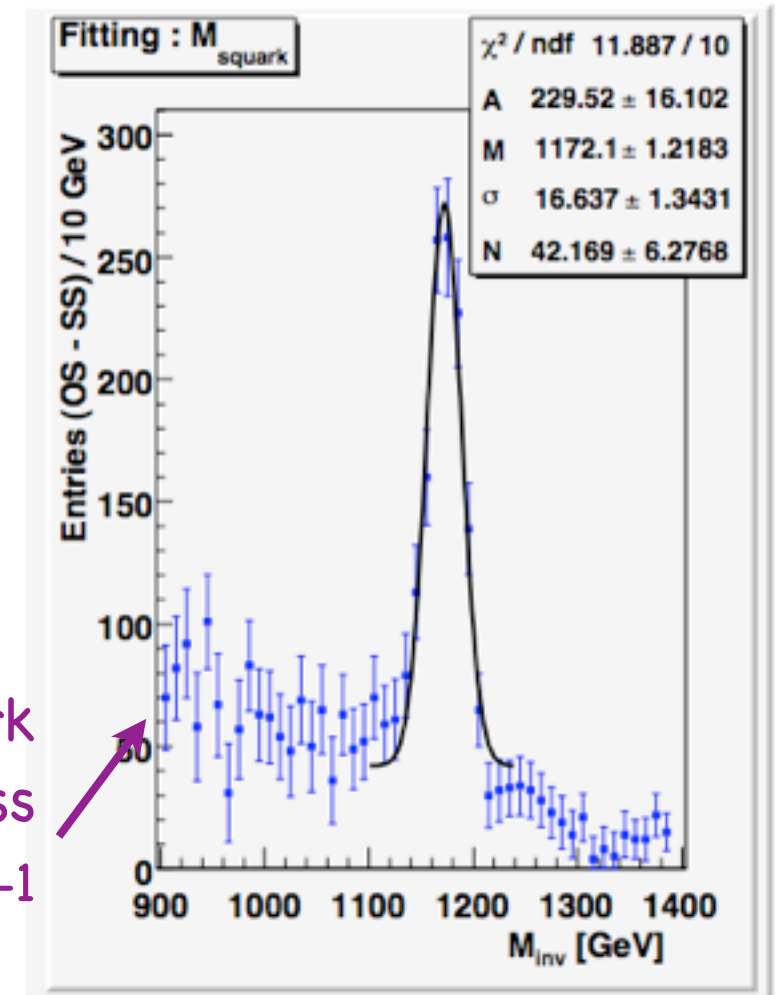
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Gluino mass is more difficult but should be possible at high luminosity

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$$\tau_{\tilde{\tau}} = \frac{48\pi M_{\text{pl}}^2 m_{\tilde{G}}^2}{m_{\tilde{\tau}}^5}$$

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step 3

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- Which range of parameters is accessible?
- Which range of T_R can be tested?
- Is it possible at the early stage (7TeV) ?

3

Probing **high T_R scenario**
at the LHC with long lived stau.

M.Endo, KH, K.Nakaji, arXiv:1008.2307

See also earlier works:
Choi, Roszkowski, Ruiz De Austri,'07
Steffen,'08

Probing **high T_R scenario** at the LHC with long lived stau.

M.Endo, KH, K.Nakaji,
arXiv:1008.2307

Logic

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(1) For a given **stau mass**

→ upper bound on **gravitino mass**

$$m_{\tilde{G}} \leq m_{\tilde{G}}^{\max}(m_{\tilde{\tau}})$$

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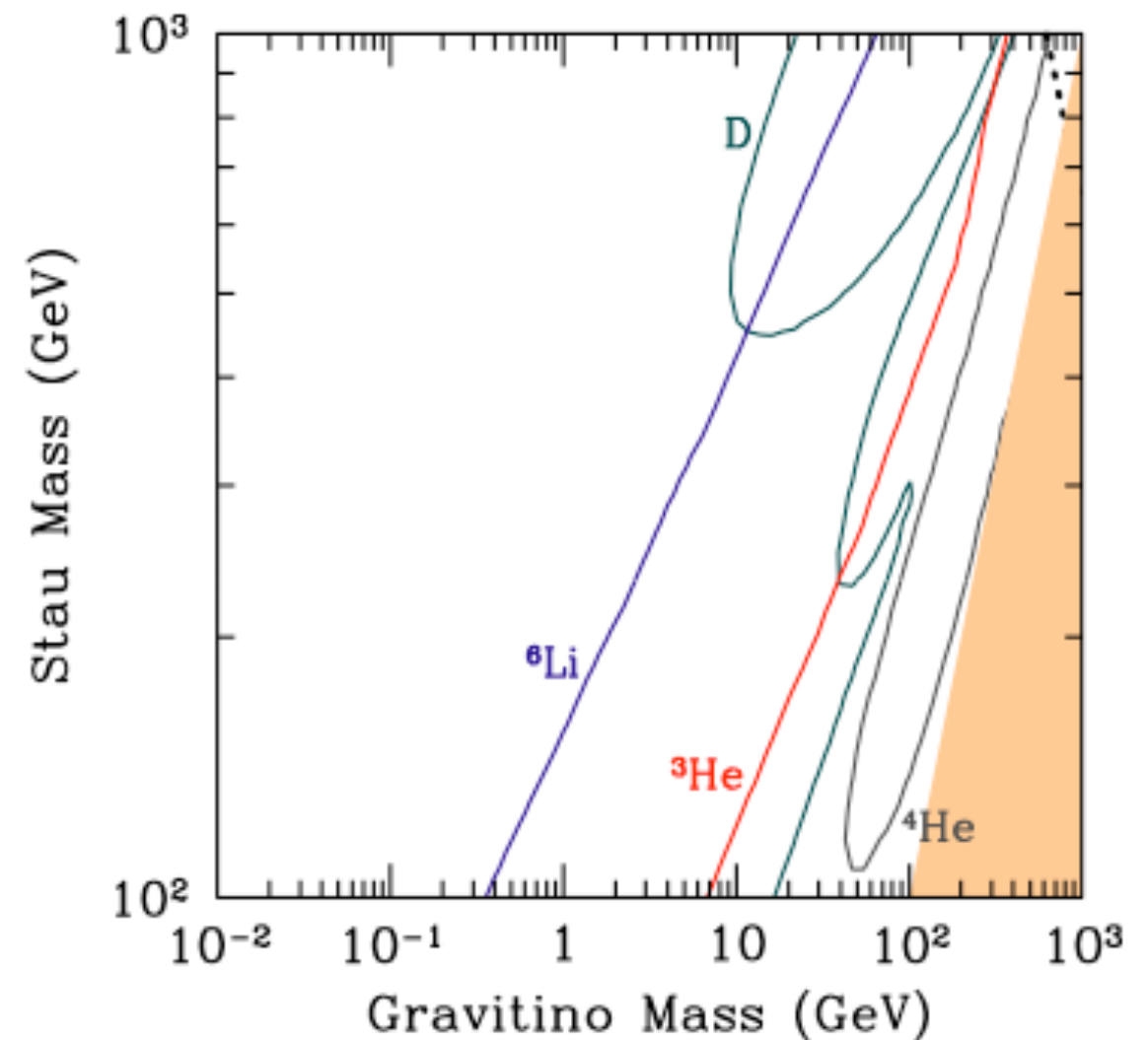
$$m_{\tilde{G}} \leq m_{\tilde{G}}^{\max}(m_{\tilde{\tau}})$$

BBN : constraint on $(Y_{\tilde{\tau}}, \tau_{\tilde{\tau}})$

$$Y_{\tilde{\tau}} = Y_{\tilde{\tau}}(m_{\tilde{\tau}})$$

$$\tau_{\tilde{\tau}} = \tau_{\tilde{\tau}}(m_{\tilde{\tau}}, m_{\tilde{G}})$$

⇒ constraint on $(m_{\tilde{\tau}}, m_{\tilde{G}})$



Kawasaki, Kohri, Moroi, Yotsuyanagi, '08

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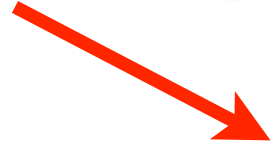
M.Endo, KH, K.Nakaji,
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(2) + for a given **T_R**

→ upper bound on **gluino mass**

Fujii, Ibe, Yanagida, '04

Probing **high T_R scenario** at the LHC with long lived stau.

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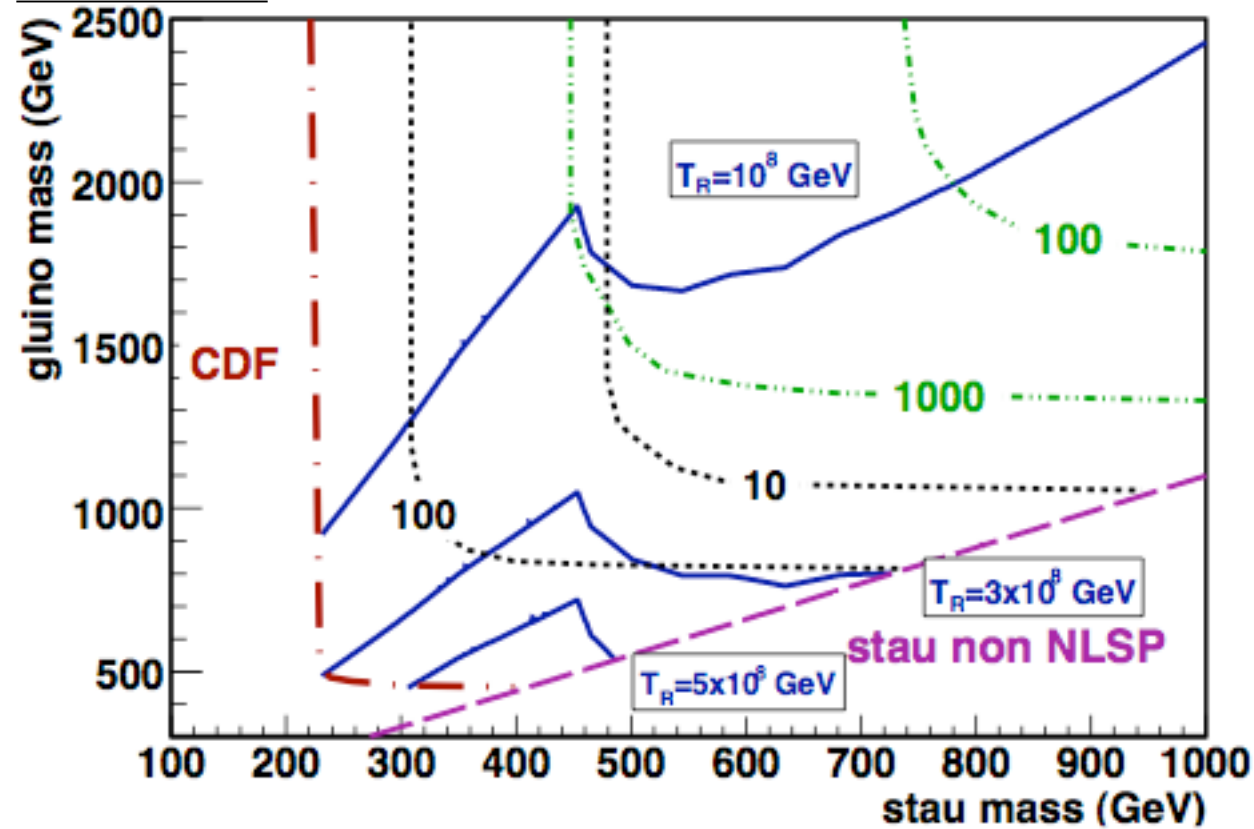
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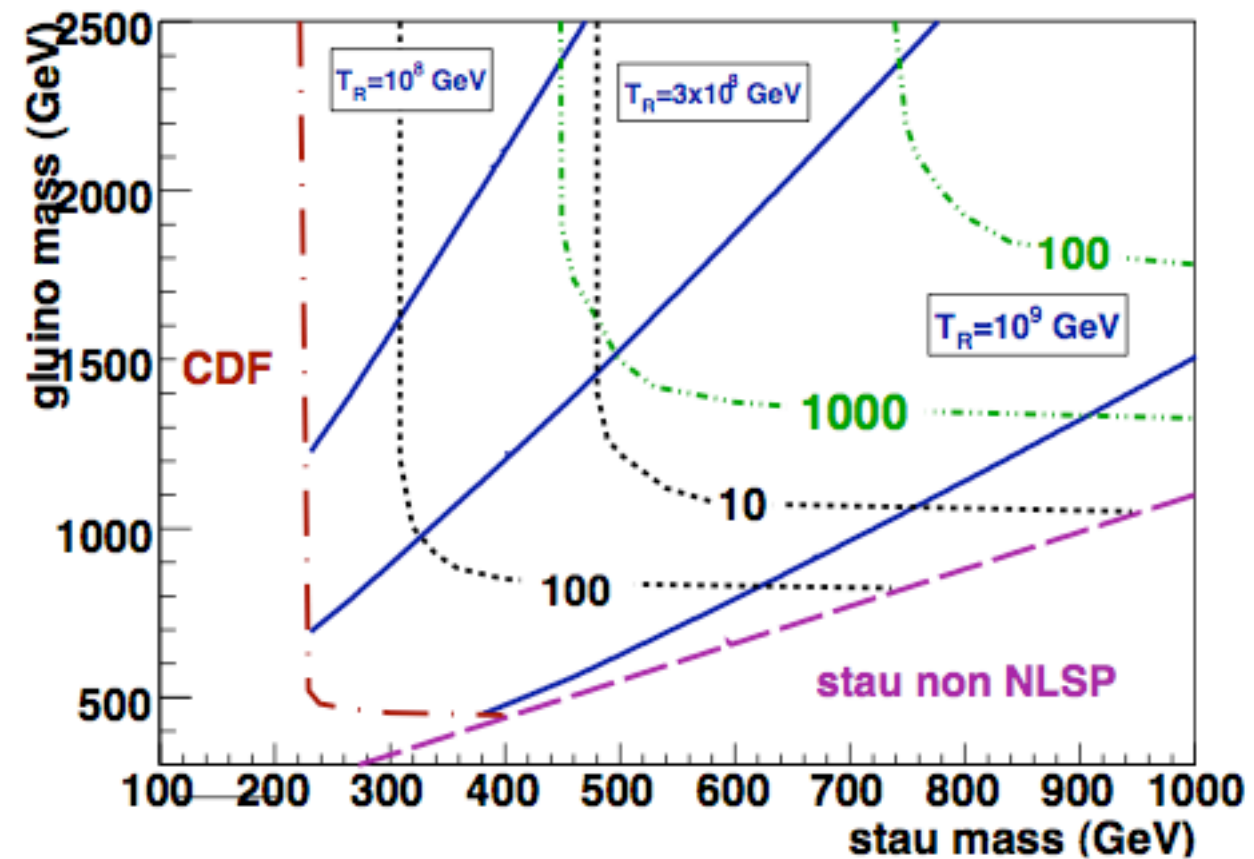
M.Endo, KH, K.Nakaji,
arXiv:1008.2307

Result



(a) = normal case

(in most cases) Stau annihilation is dominated by Electroweak process



(b) = reduced Y_{stau}

Stau annihilation is dominated by enhanced Higgs coupling

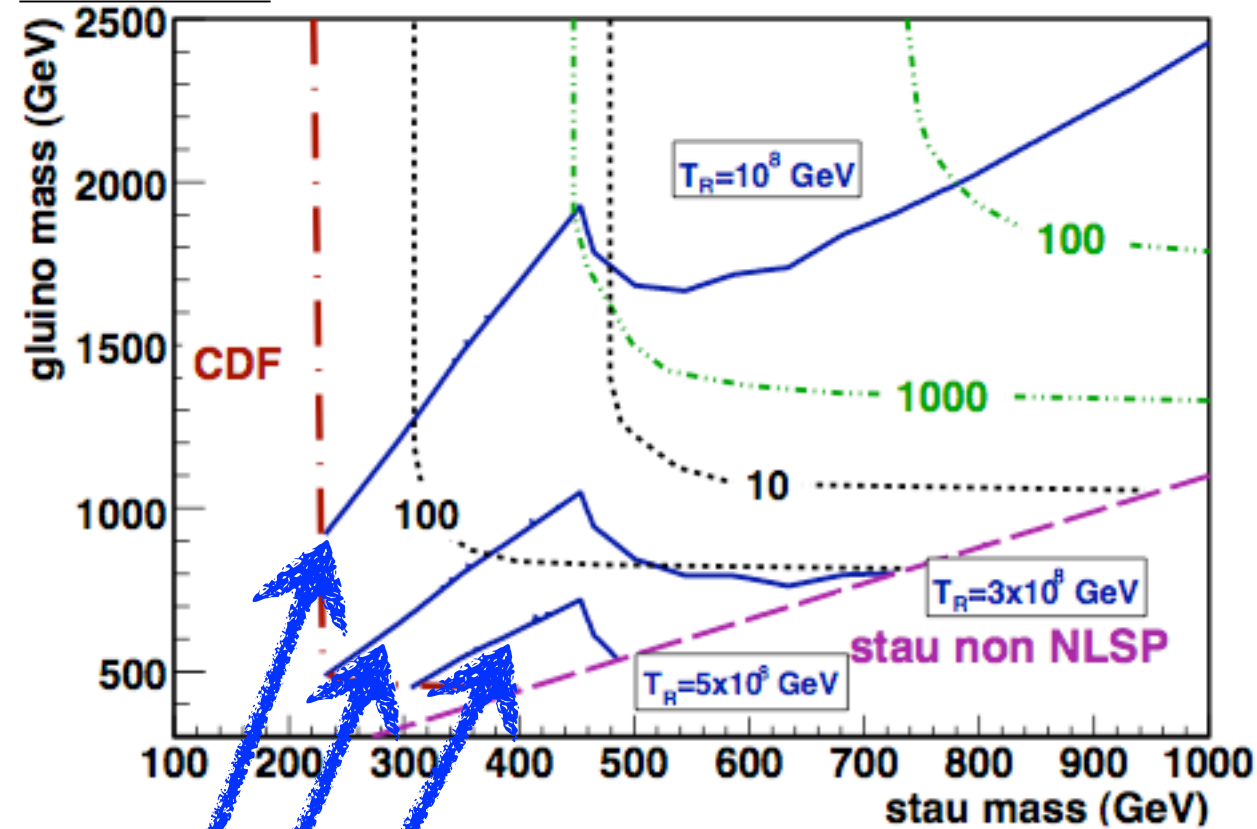
see, Ratz, Schmidt-Hoberg, Winkler,'08
Pradler, Steffen,'08

Note: take $m(\text{bino})=m(\text{wino})=1.1m(\text{stau})$
to have conservative bound on T_R .

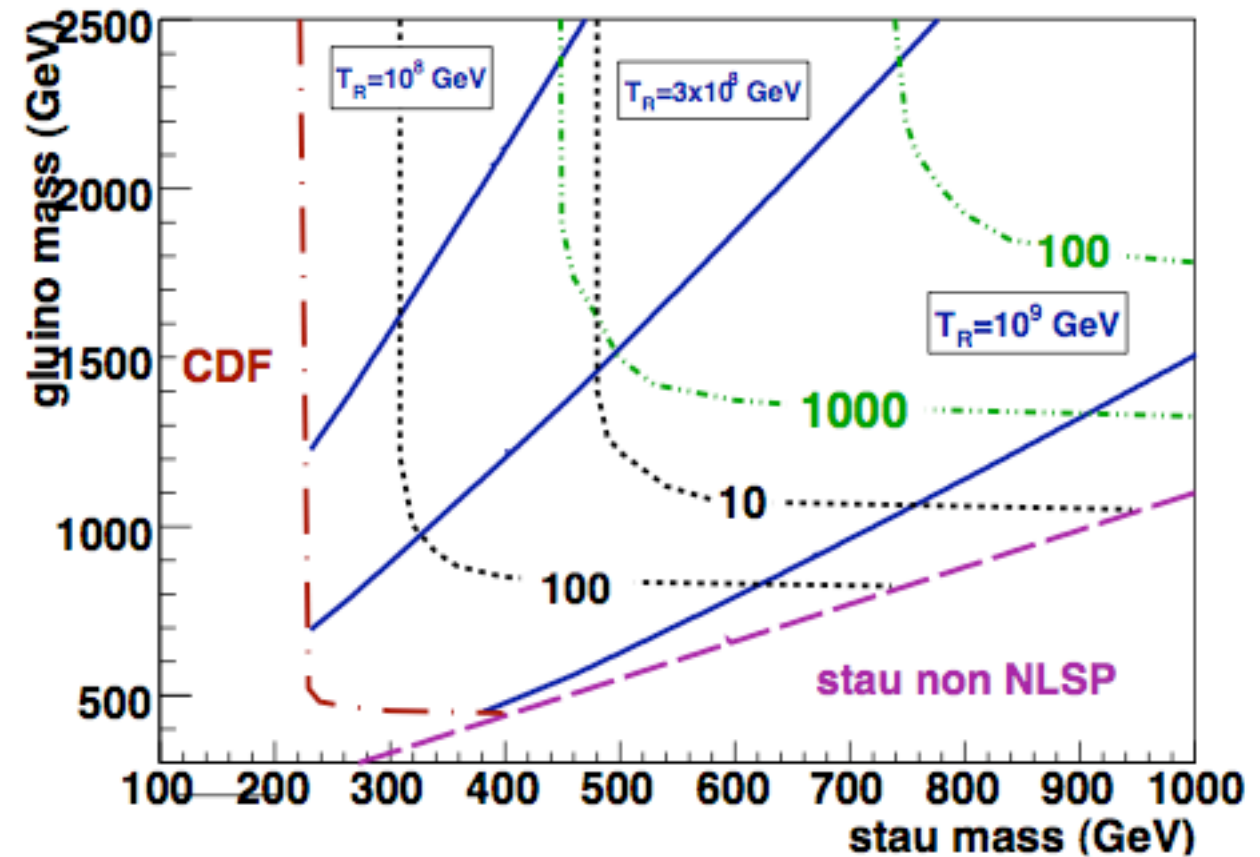
Probing **high T_R scenario** at the LHC with long lived stau.

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arXiv:1008.2307

Result



(a) = normal case



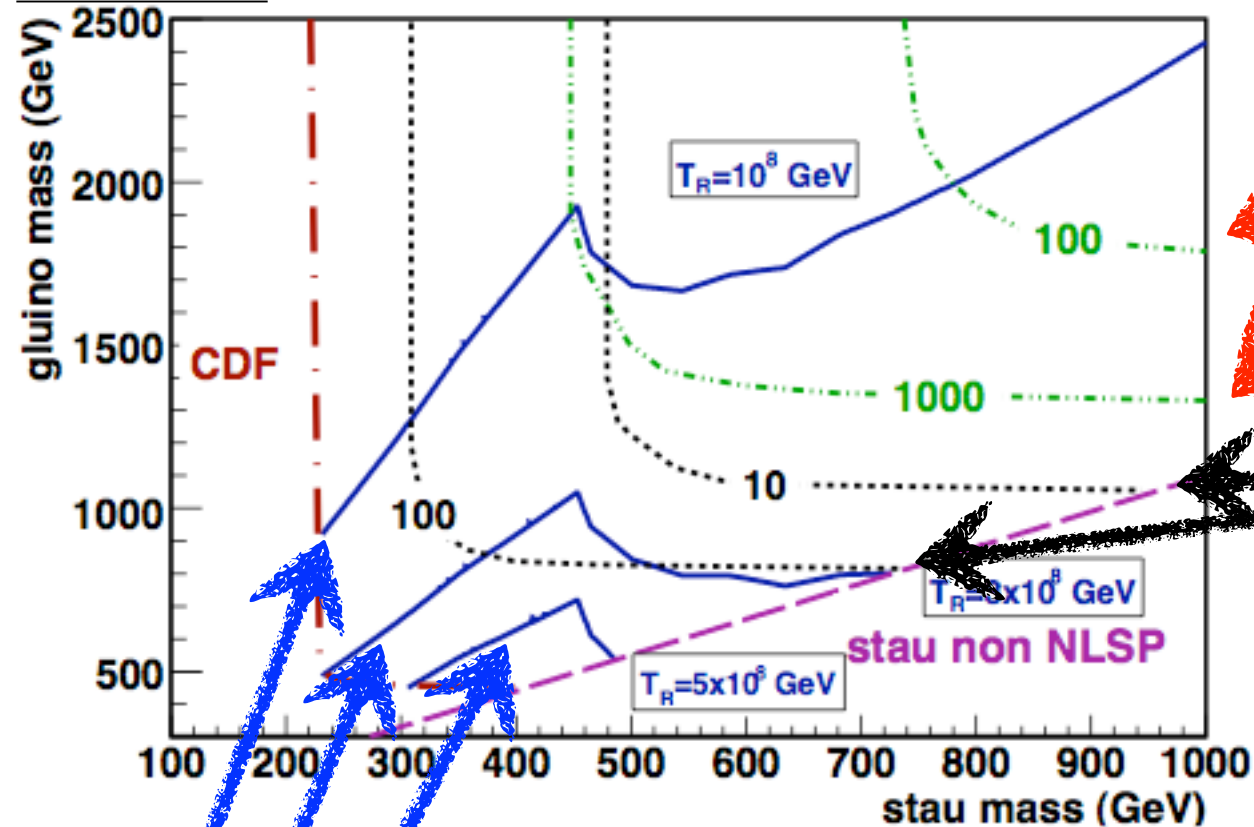
(b) = reduced Y_{stau}

upper bound on the gluino mass for given T_R

Probing high T_R scenario at the LHC with long lived stau.

M.Endo, KH, K.Nakaji,
arXiv:1008.2307

Result



of produced staus
at 14 TeV 10fb^{-1}

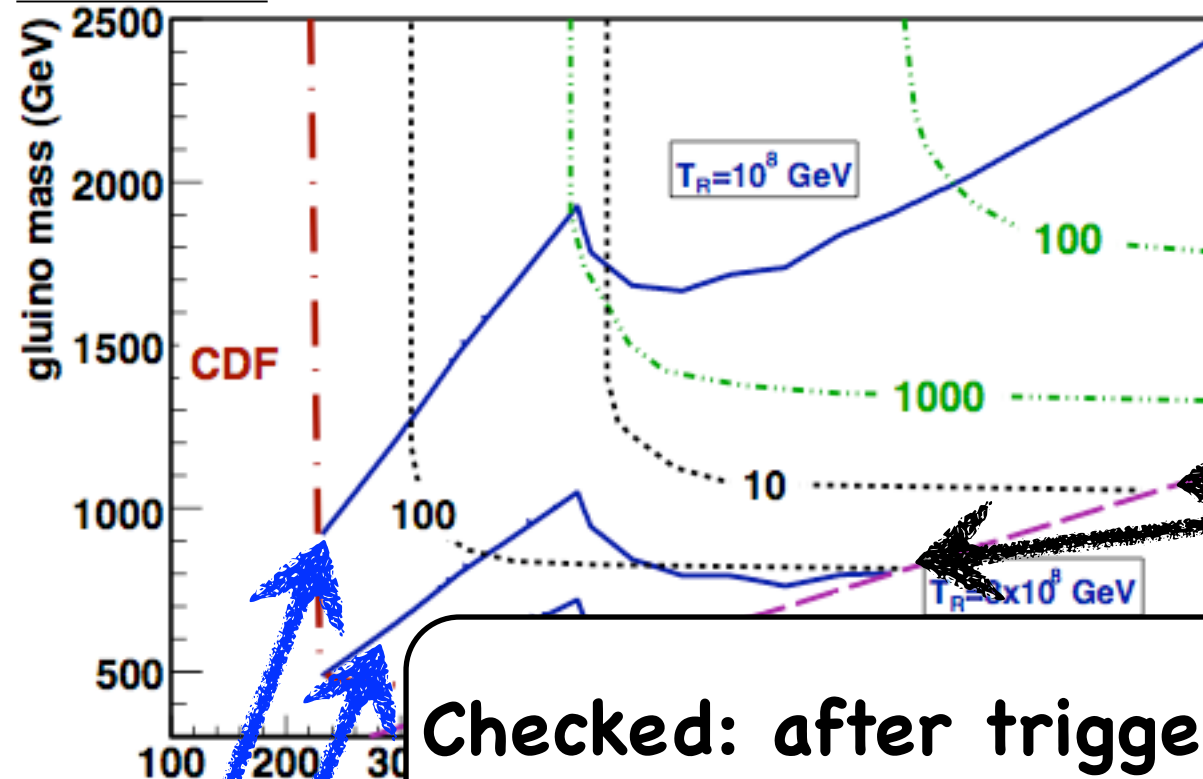
of produced staus
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upper bound on the gluino mass for given T_R

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Result



of produced staus at 14 TeV 10fb^{-1}

of produced staus at 7 TeV 1fb^{-1}

Checked: after triggers and cuts, 20-50% events remain.

trigger assumption:

- >=1 isolated e ($p_T > 20\text{GeV}$), or
- >=1 isolated mu ($p_T > 40\text{GeV}$), or
- >=1 isolated tau ($p_T > 100\text{GeV}$), or
- >=1 isolated stau ($p_T > 40\text{ GeV}$ and $\beta > 0.7$, $\text{eta} < 1.0$ or $\beta > 0.8$, $\text{eta} < 2.8$), or
- >=2 staus ($p_T > 40\text{ GeV}$ and $\beta > 0.7$, $\text{eta} < 1.0$ or $\beta > 0.8$, $\text{eta} < 2.8$)

stau cuts assumptions:

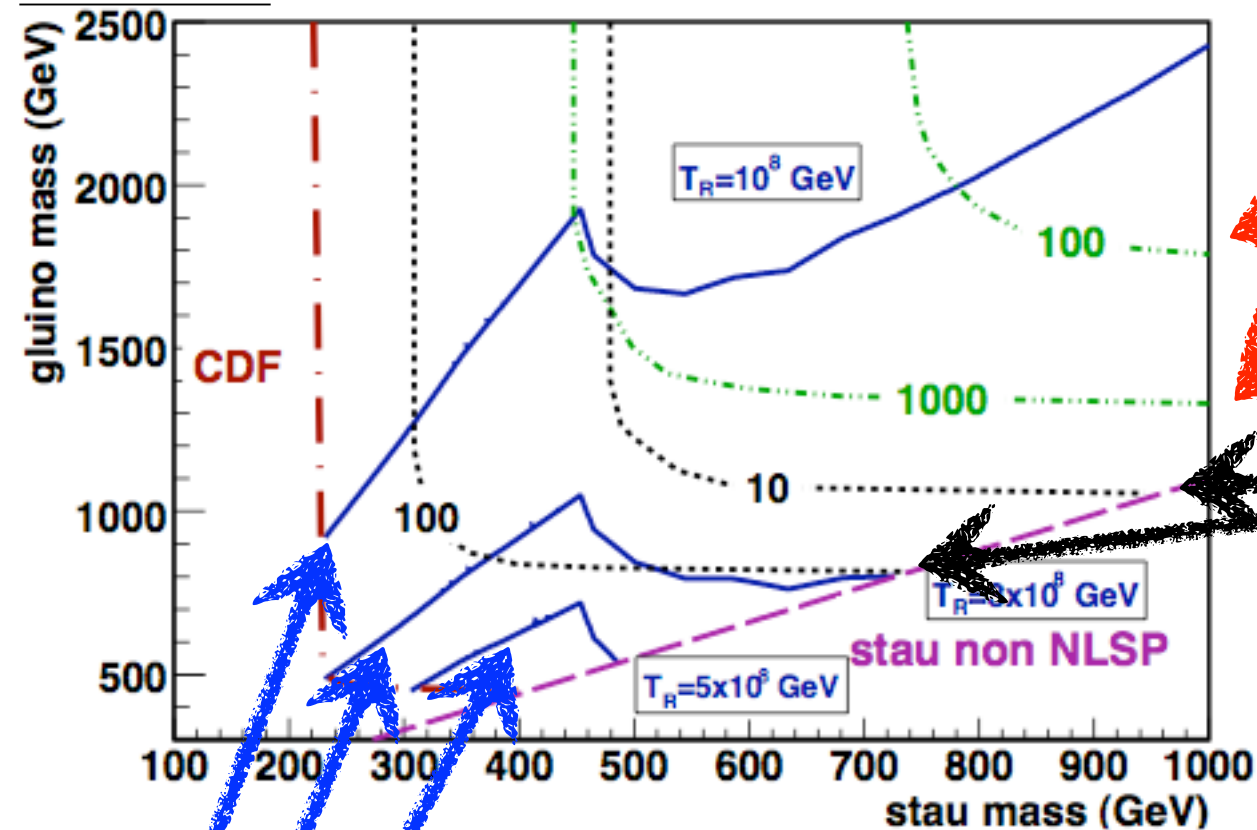
$p_T > 20\text{ GeV}$ & $\text{eta} < 2.5$ & $0.5 < \beta < 0.9$ -> almost background free!

upper bound

Probing high T_R scenario at the LHC with long lived stau.

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Result



of produced staus
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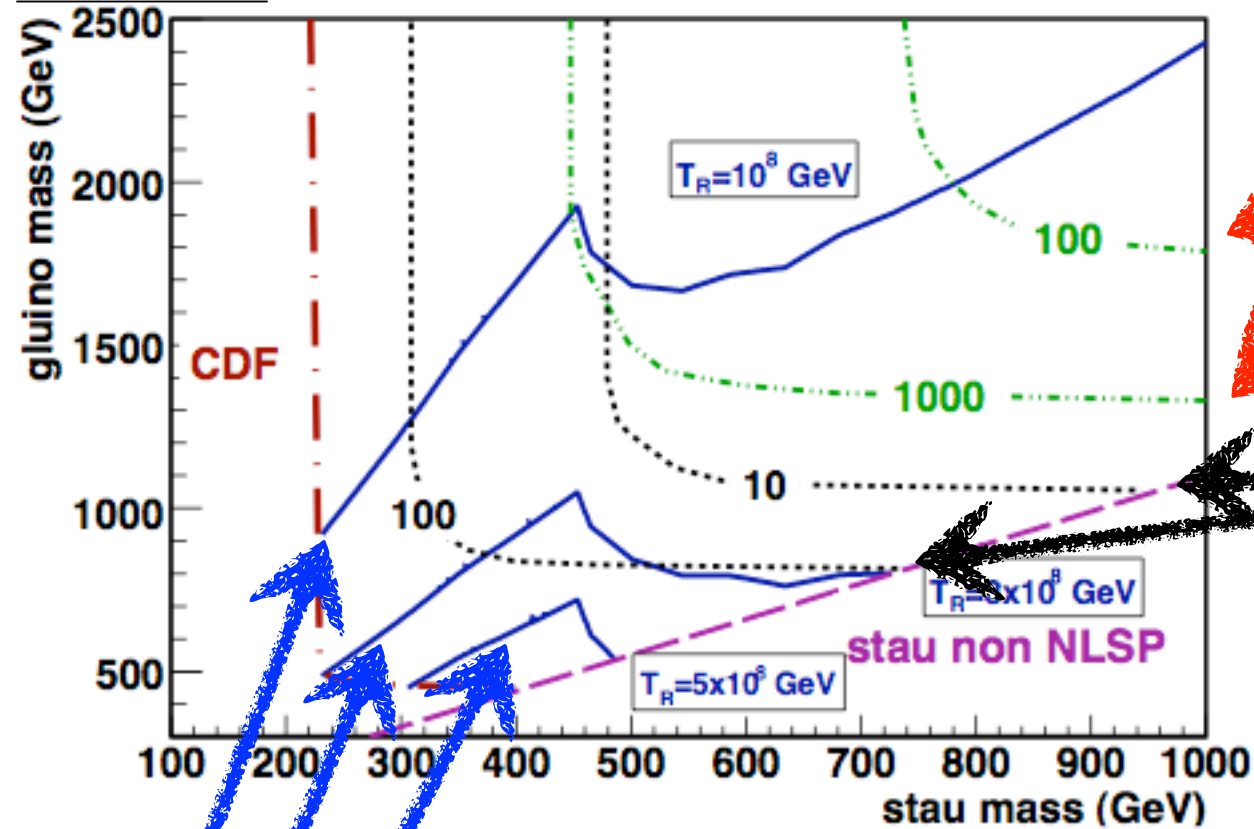
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Result



of produced staus
at 14 TeV 10fb^{-1}

of produced staus
at 7 TeV 1fb^{-1}

$T_R > \text{a few } 10^8 \text{ GeV}$
can be probed
at 7 TeV 1fb^{-1} !!!

upper bound on the gluino mass for given T_R

SUMMARY

Main message of this talk:

If **SUSY + gravitino DM + stau NLSP**
is realized in nature,

The **T_R** can be determined at the **LHC**.

At **7 TeV 1fb^{-1}** (\approx within 2 years),
 $T_R > \text{a few } 10^8 \text{ GeV}$ can be tested
in most of the parameter space!

L

H

C

!

Let's **H**ave a **C**offee **!**



additional slides

SUMMARY

If **metastable heavy charged particle** will be observed at the LHC,

→ **SUSY + gravitino DM + stau NLSP**

is the leading candidate, (which can be tested)

the **T_R** can be determined at the **LHC**,

(assuming no dilution, and $\Omega_{DM} = \Omega_G$)

at **7 TeV 1fb^{-1}** (\approx within 2 years),

$T_R > \text{a few } 10^8 \text{ GeV}$ can be tested
in most of the parameter space!

Reheating Temperature T_R

~ highest temperature in the rad. dom. universe.

= one of the most important parameters of cosmology.

- determined by **inflaton** decay rate:

$$\Gamma_\phi \sim m_\phi^3 / M_P^2$$

$$\longrightarrow T_R \sim \sqrt{\Gamma_\phi M_P} \sim 10^{10} \text{ GeV} (m_\phi / 10^{13} \text{ GeV})^{3/2}$$

- important for **baryogenesis**:

(e.g. thermal Leptogenesis $\longrightarrow T_R > \mathcal{O}(10^9 \text{ GeV})$)

- typically most of staus have large velocity and escape from detector.
- but some of them have sufficiently small velocity and stop at calorimeters.

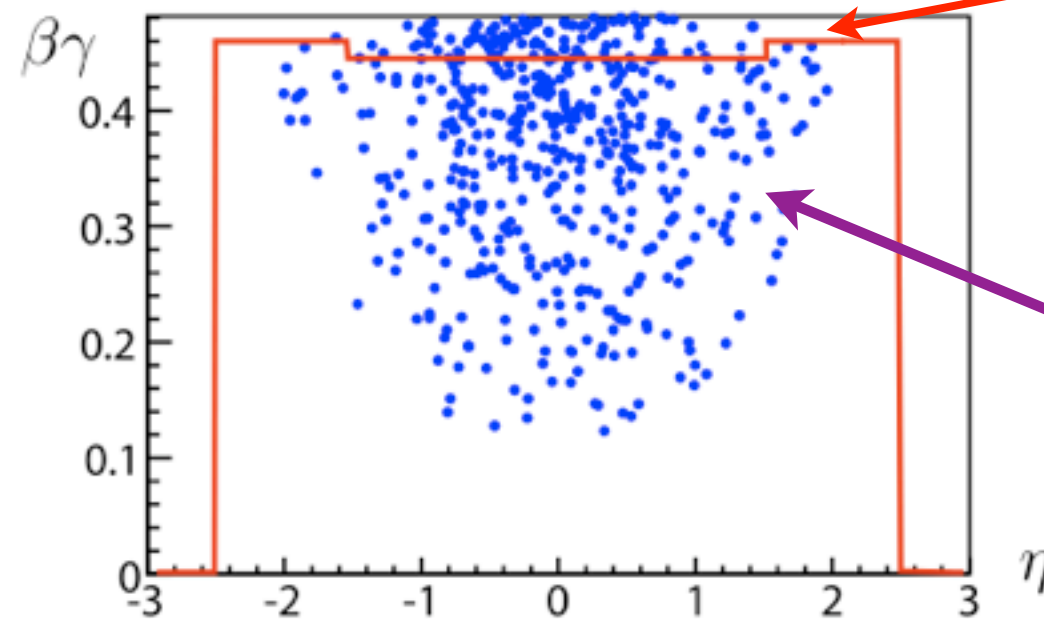
example of SUSY
model point SPS7
($\sigma_{\text{SUSY}} = 3.5 \text{ pb}$)

from Asai, KH,
Shirai '09

(See related work
"stopping gluino",
Arvanitaki et.al.)

TABLE II: The number of stopping staus for 10 fb^{-1} .

with cuts	without cuts
400	805



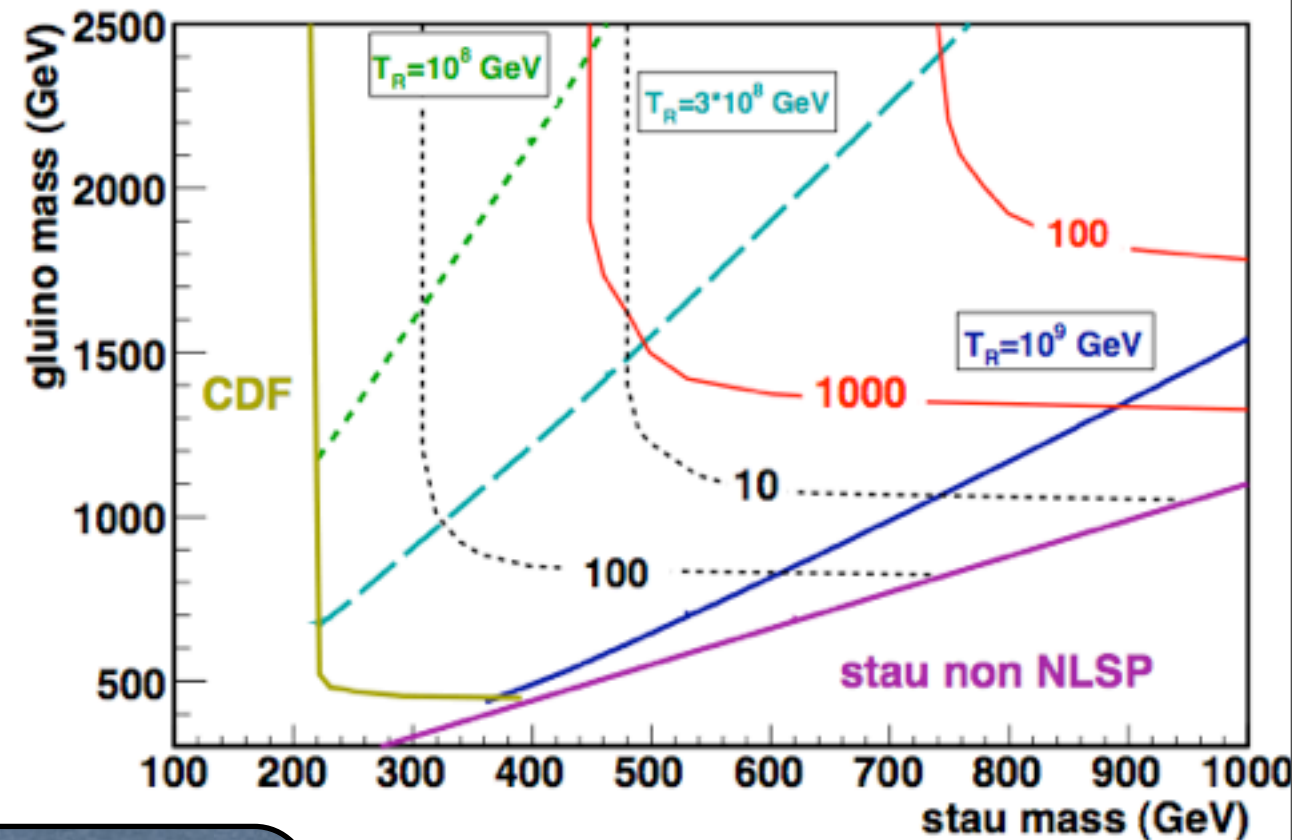
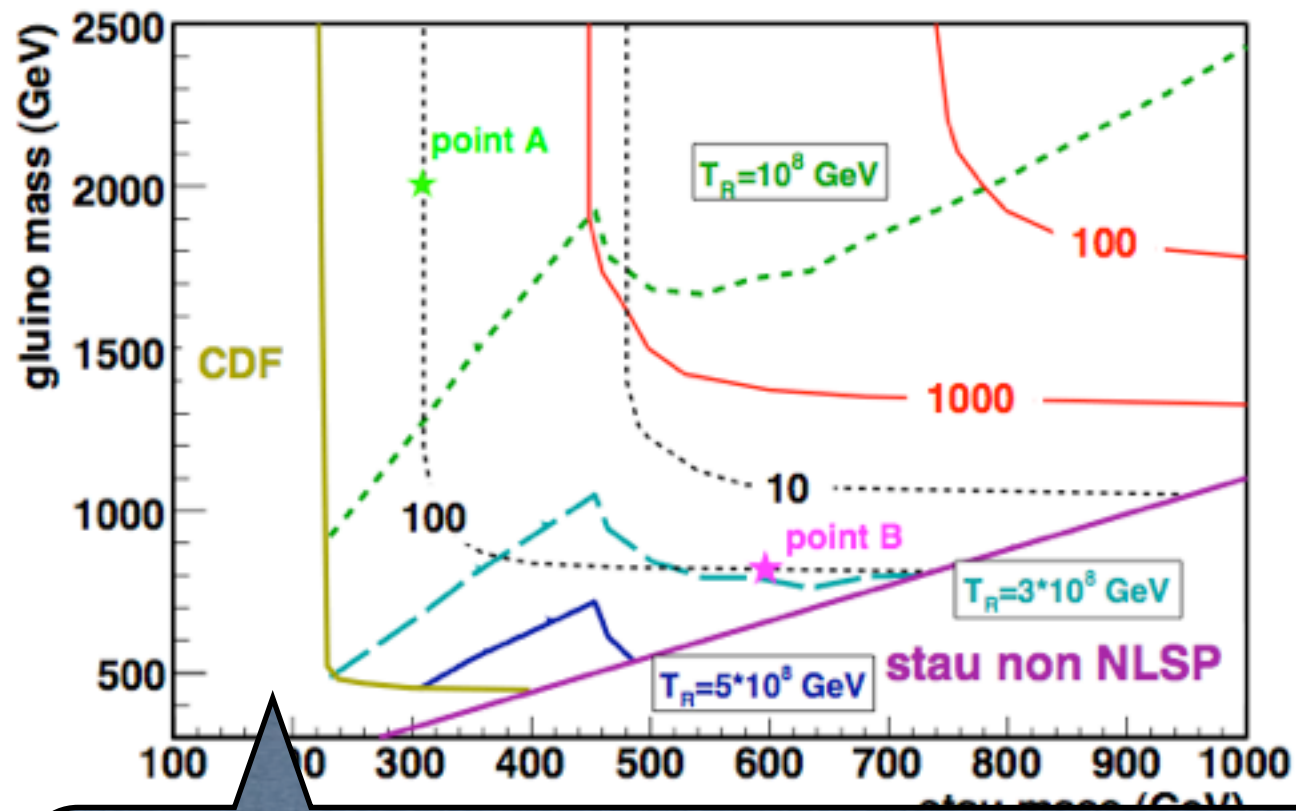
assume (ATLAS):
Fe 1440mm (barrel)
Cu 1400mm (end-cap)

- stopped events
- about 1% of total SUSY events
 - a few per day (for $10^{33} / \text{cm}^2 \text{ s}$)

FIG. 1: $\eta - \beta\gamma$ distribution of the staus. The red line shows the limit for the stau to stop in the detector.

Probing high T_R scenario at the LHC with long lived stau.

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(b) = reduced Y_{stau}

CDF boundusing [PRL 103, 021802 (2009)]
production cross section of CHAMPS with

- $\eta < 0.7$
- $p_T > 40$ GeV
- $0.4 < \beta < 0.9$
- $\text{sum}(E_T, R < 0.4) / p_T < 0.1$

should be less than 10 fb (2 sigma)