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BAOs in 2D

**Modeling redshift-space power spectrum
from perturbation theory**

Atsushi Taruya
(RESCEU, Univ.Tokyo)

In collaboration with

Takahiro Nishimichi (IPMU), Shun Saito (Berkeley)

Contents

- Introduction: why 2D ?
- Modeling redshift-space distortion
- Impact on parameter estimations
- Summary

Introduction

Precise measurement of **B**aryon **A**coustic **O**scillations
(**BAOs**)

{ A major science goal of galaxy redshift surveys
{ Characteristic scale of BAO as a standard ruler
to trace cosmic expansion history

➔ nature of dark energy

Reducing systematics is a big issue :

Aim of this
talk

- non-linear evolution
- redshift distortion
- galaxy bias

Methodology

There are several techniques to reduce systematics
(still on-going subjects)

Fitting

Sophisticated parametric formula and/or hybrid fitting

Seo et al. ('08, '09); Padmanabhan & White ('09)

Reconstructing

Degradation of acoustic features by Zel'dovich approx.

Eisenstein et al. ('07); Huff et al. ('07); Padmanabhan et al. ('09)

Forward modeling

Perturbation theory (PT) based modeling of BAOs

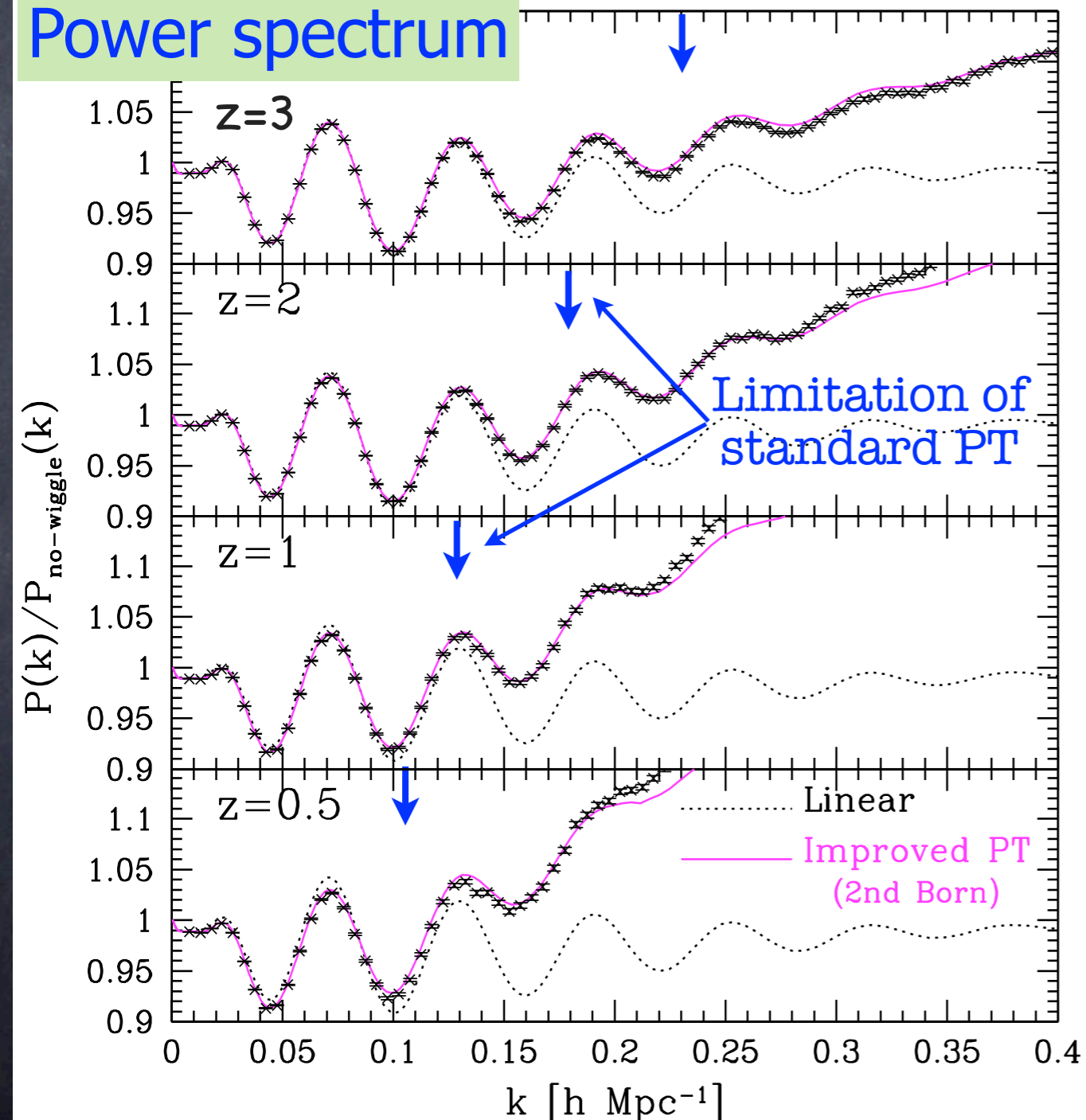
Crocce & Scoccimarro ('08); Jeong & Komatsu ('06,'09);
Matsubara ('08a,b); AT & Hiramatsu ('08); AT et al. ('09); etc. ...

Development of new analytic method

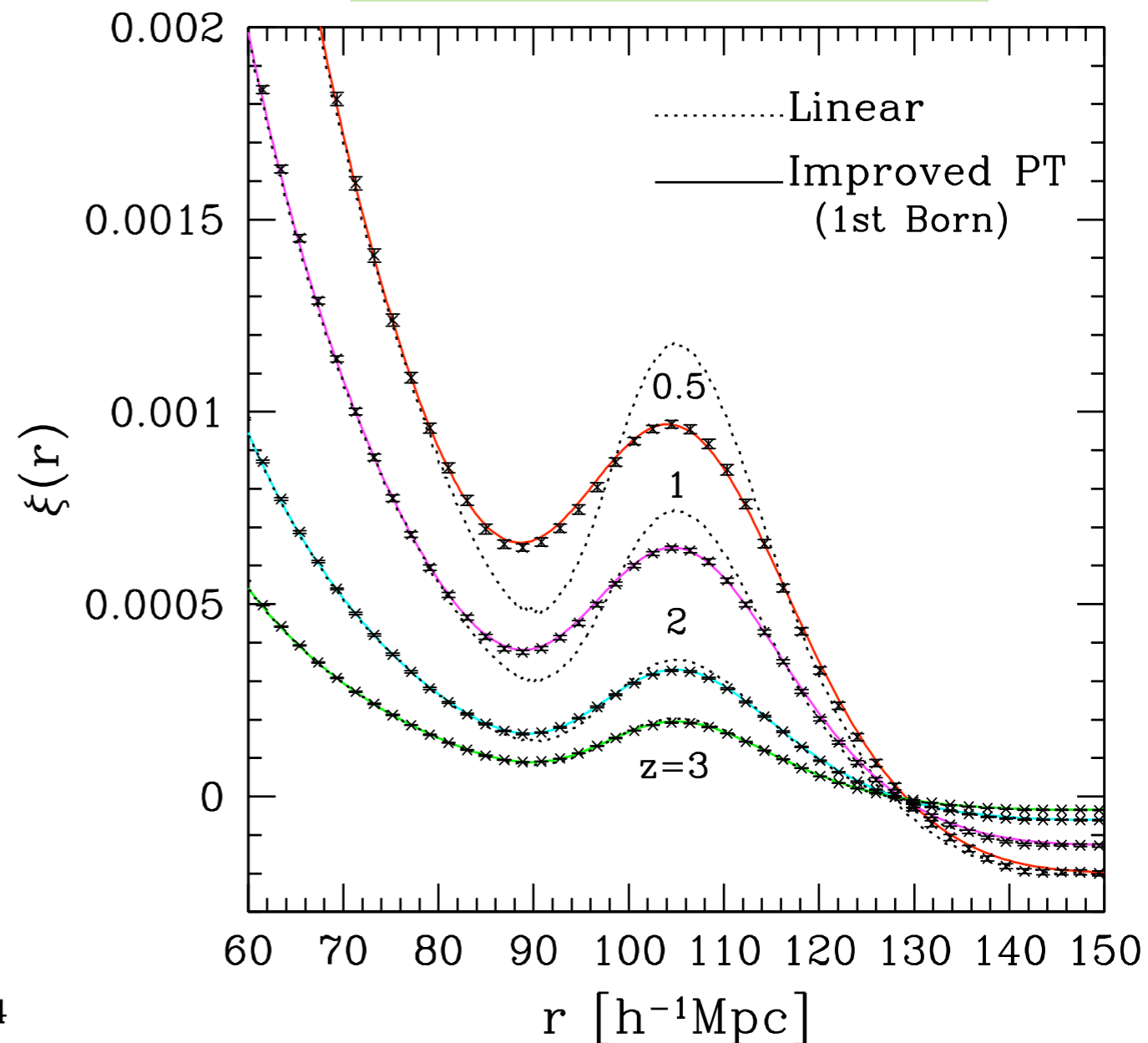
An improved treatment of perturbation theory (PT) to deal with non-linear gravitational evolution

AT & Hiramatsu (2008)
AT, Nishimichi, Saito & Hiramatsu (2009)

Power spectrum



Correlation function



Development of new analytic method

An improved treatment of perturbation theory (PT) to deal with non-linear gravitational evolution

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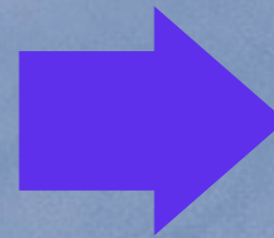
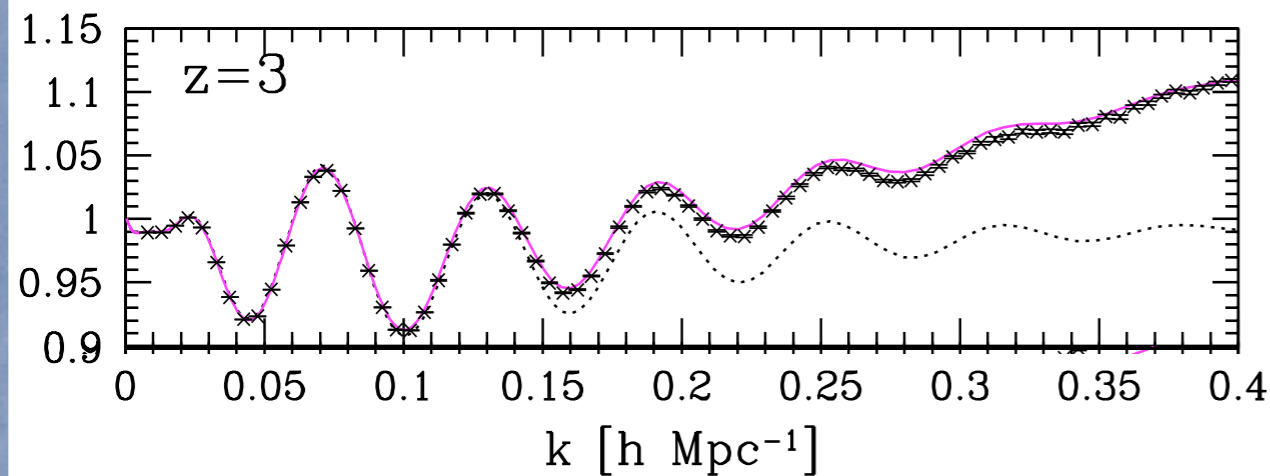
This talk

Based on improved PT of non-linear structure growth, further investigation on modeling BAOs

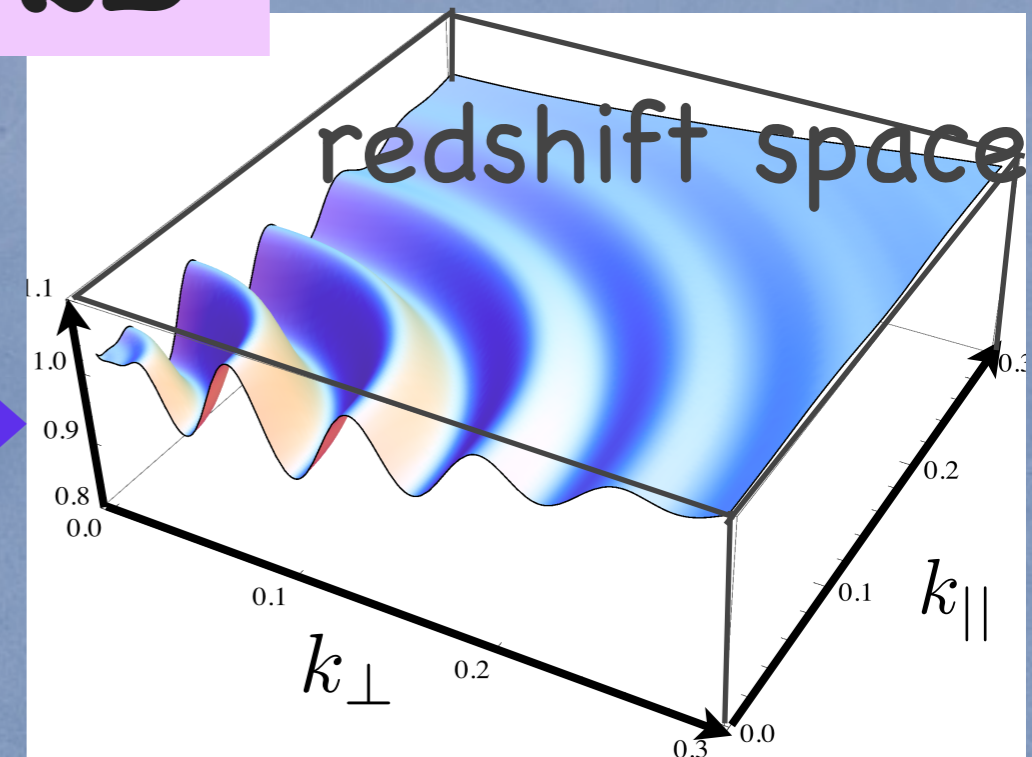
Key word

From 1D to 2D

real space

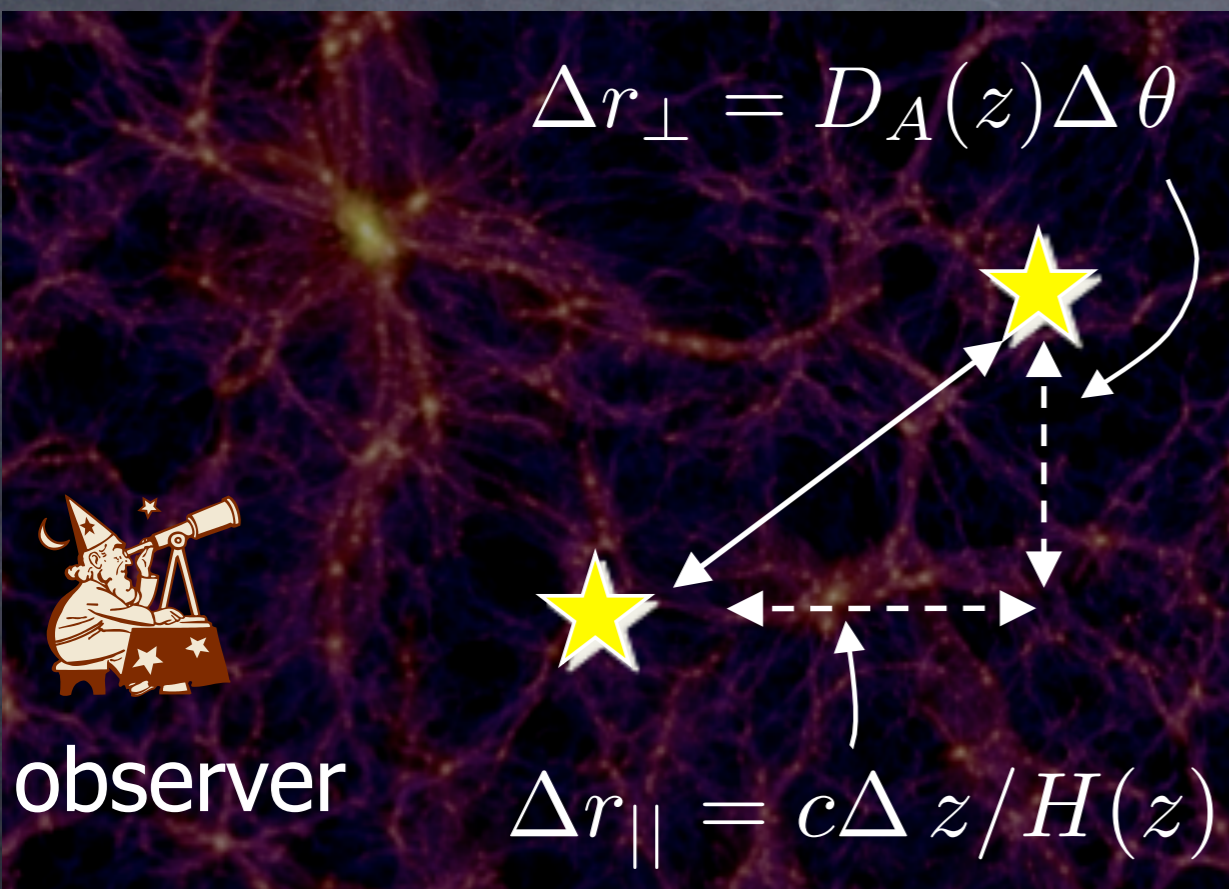


redshift space



Alcock-Paczynski (AP) effect

Different distance measurements between **parallel** & **transverse** directions cause apparent anisotropies in $P(k)$



Using BAO as standard ruler,

$H(z)$ & **$D_A(z)$**

can be determined simultaneously



strong constraints on
dark energy E.O.S

[e.g., Seo & Eisenstein ('03); Hu & Haiman ('03);
Blake & Glazebrook ('03); Shoji et al.('09)]

Note majority of recent works focuses on angle-averaged 1D BAOs

→ only constrain the combination, $D_A(z)^2/H(z)$

Gravity test

In modeling 2D, redshift distortion effect is inevitable, **but** it provides a way to test gravity on cosmological scales

On linear regime,

Kaiser ('87)

strength of redshift distortion

\propto growth-rate parameter :

$$f(z) \equiv \frac{d \ln D_+}{d \ln a}$$

$[D_+(z) : \text{linear growth rate}]$

In general relativity,

$$f(z) \simeq \{\Omega_m(z)\}^\gamma \quad ; \quad \gamma = 0.55$$

e.g., Linder ('05)

Deviation of γ from 0.55 implies a breakdown of GR on cosmological scales

Redshift distortion

Definition

redshift space real space

$$\vec{s} = \vec{r} + \frac{(\vec{v} \cdot \hat{z})}{a H(z)} \hat{z}; \quad \begin{cases} \vec{v} & : \text{peculiar velocity} \\ \hat{z} & : \text{observer's} \\ & \text{line-of-sight direction} \end{cases}$$

Observed clustering pattern is apparently distorted.

- Anisotropy (2D power spectrum)

$$P(k) \longrightarrow P^{(S)}(k, \mu); \quad \mu \equiv (\vec{k} \cdot \hat{z}) / |\vec{k}|$$

- Power spectrum amplitude

Enhancement

Kaiser effect (small-k)

Suppression

Finger-of-God effect (large-k)

Power spectrum in redshift space

Exact expression

$$\mathbf{x} = \mathbf{r} - \mathbf{r}'$$

$$P^{(S)}(\mathbf{k}) = \int d^3\mathbf{x} e^{i\mathbf{k}\cdot\mathbf{x}} \left\langle e^{-ik\mu \Delta u_z} \left\{ \delta(\mathbf{r}) - \nabla_z u_z(\mathbf{r}) \right\} \left\{ \delta(\mathbf{r}') - \nabla_z u_z(\mathbf{r}') \right\} \right\rangle$$



$$u_z = (\vec{v} \cdot \hat{\mathbf{z}}) / (aH)$$

$$\Delta u_z = u_z(\mathbf{r}) - u_z(\mathbf{r}')$$

(Popular) analytic model

1D velocity dispersion

e.g., Scoccimarro (2004)

$$P^{(S)}(k, \mu) = e^{-(k\mu \sigma_v)^2} \left[P_{\delta\delta}(k) - 2\mu^2 P_{\delta\theta}(k) + \mu^4 P_{\theta\theta}(k) \right]$$

Finger of God

(non-linear) Kaiser

fitting parameter

... physical, but still empirical formula

Missing terms, found

From low-k expansion of the exact formula for $P^{(S)}(k, \mu)$,

$$P^{(S)}(k, \mu) = e^{-(k\mu f\sigma_v)^2} \left[P_{\delta\delta}(k) - 2f\mu^2 P_{\delta\theta}(k) + f^2\mu^4 P_{\theta\theta}(k) + A(k, \mu) + B(k, \mu) \right]$$

Leading-order corrections to the mode-coupling btw velocity & density

Non-Gaussian correction

$$A(k, \mu) = -2k\mu \int \frac{d^3\mathbf{p}}{(2\pi)^3} \frac{p_z}{p^2} B_\sigma(\mathbf{p}, \mathbf{k} - \mathbf{p}, -\mathbf{k})$$

antiphase oscillation

$$\langle \theta(\mathbf{k}_1) \{ \delta(\mathbf{k}_2) - \mu_2^2 \theta(\mathbf{k}_2) \} \{ \delta(\mathbf{k}_3) - \mu_3^2 \theta(\mathbf{k}_3) \} \rangle = (2\pi)^3 \delta_D(\mathbf{k}_{123}) B_\sigma(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3)$$

Gaussian correction

$$B(k, \mu) = (k\mu)^2 \int \frac{d^3\mathbf{p}}{(2\pi)^3} F(\mathbf{p}) F(\mathbf{k} - \mathbf{p})$$

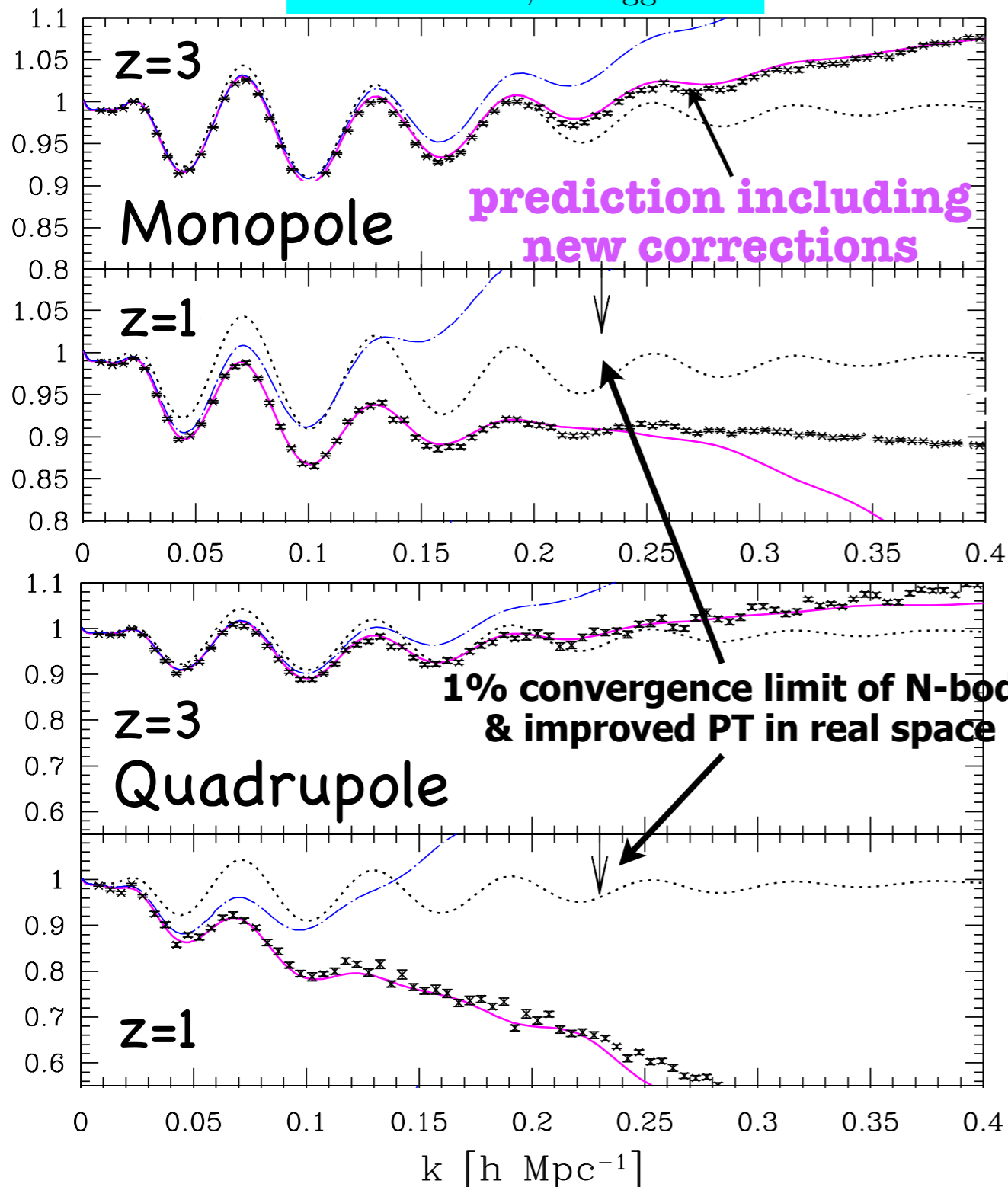
small in amplitude (<1-2%)

$$F(\mathbf{p}) \equiv \frac{p_z}{p^2} \left\{ P_{\delta\theta}(p) - \frac{p_z^2}{p^2} P_{\theta\theta}(p) \right\}$$

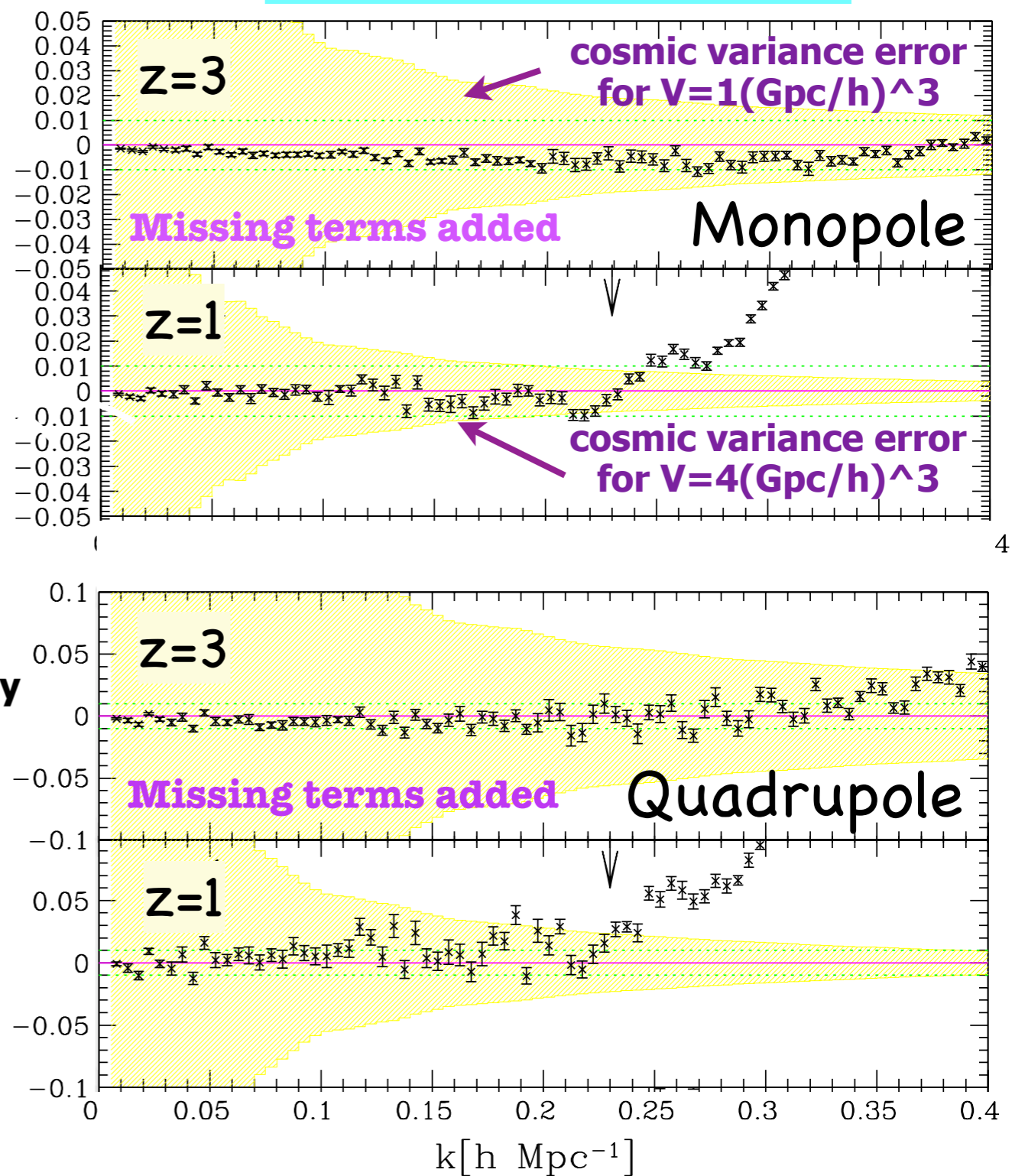
These also depend on 'f'

Results including corrections

$$P_\ell^{(S)}(k)/P_{\ell,\text{no-wiggle}}^{(S)}(k)$$

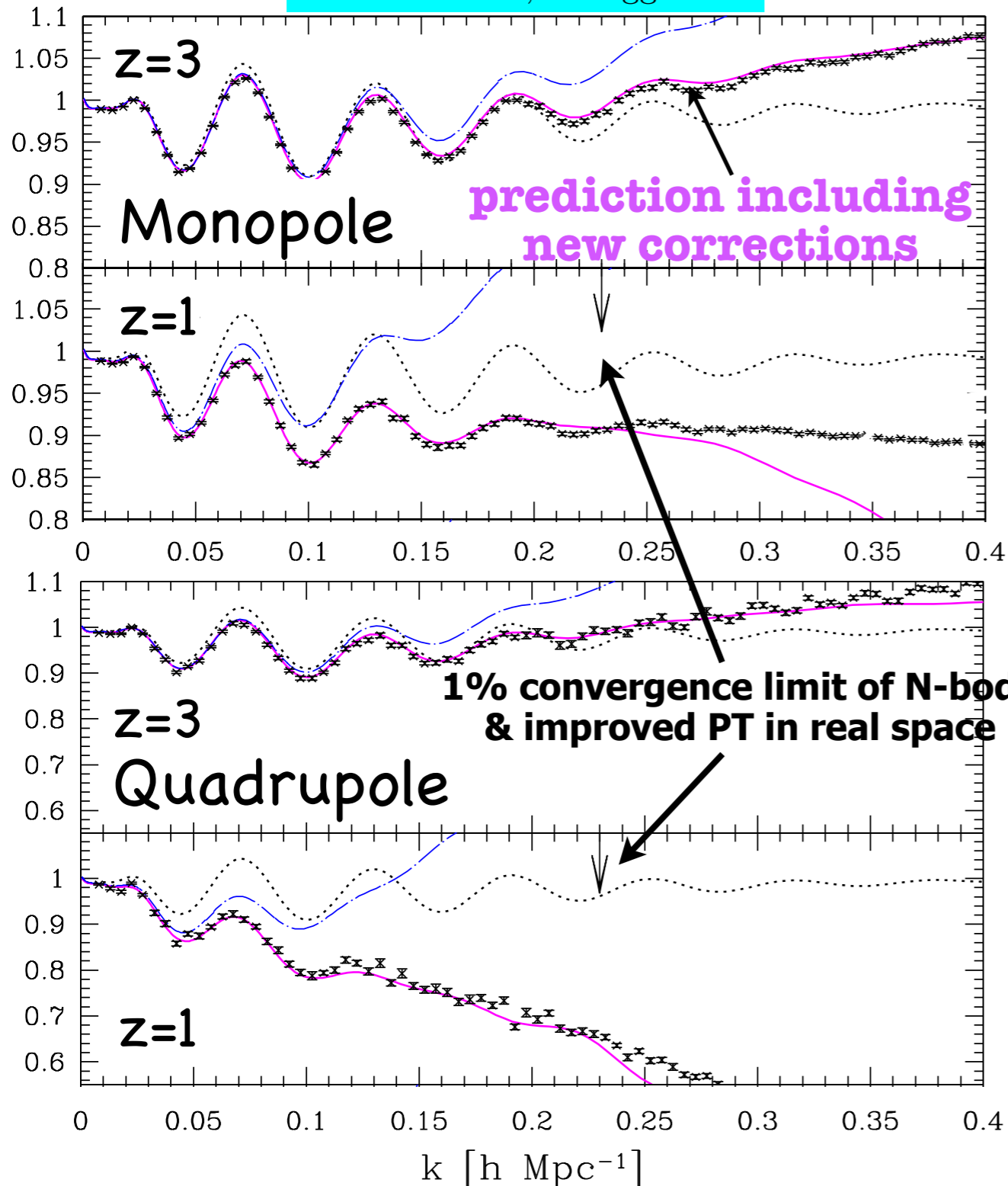


fractional residuals

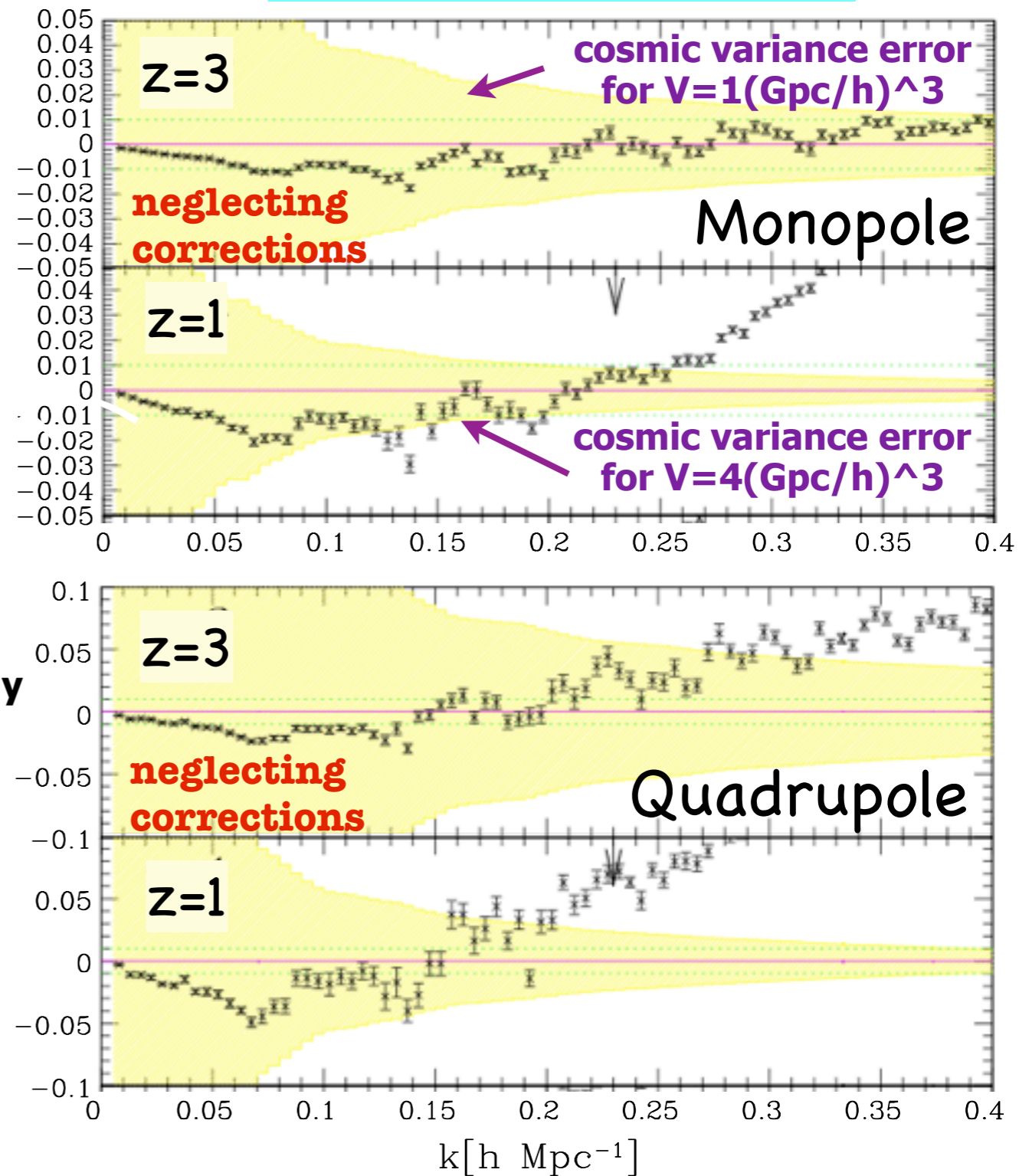


Results including corrections

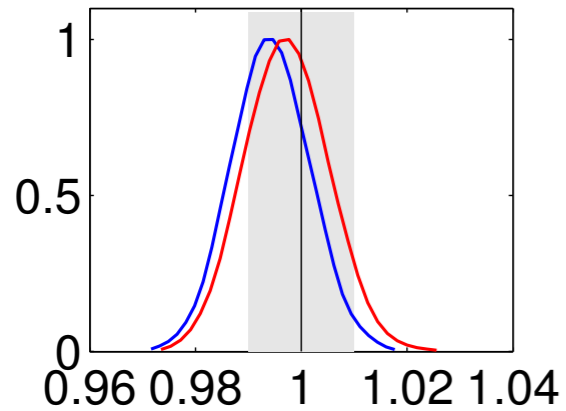
$$P_l^{(S)}(k)/P_{l,\text{no-wiggle}}^{(S)}(k)$$



fractional residuals

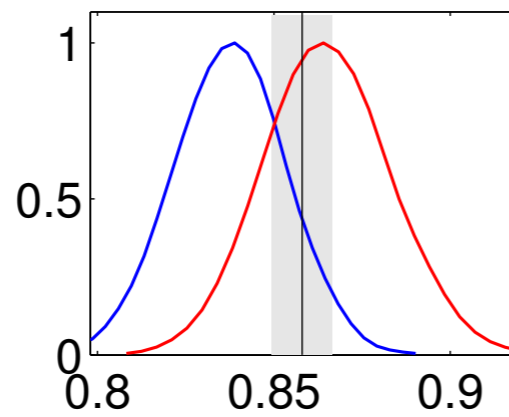
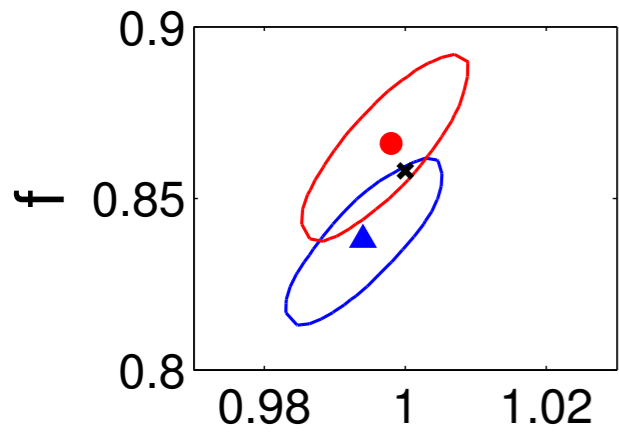


Check: recovery of D_A , H & f



* Fiducial

- — New model of redshift distortion
- ▲ — Phenomenological model (w/o corrections)

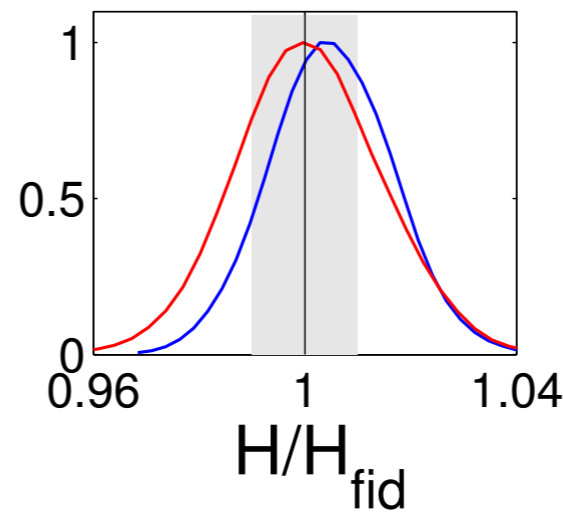
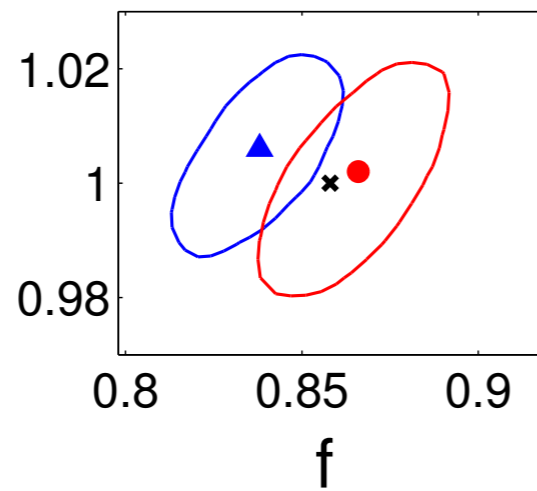
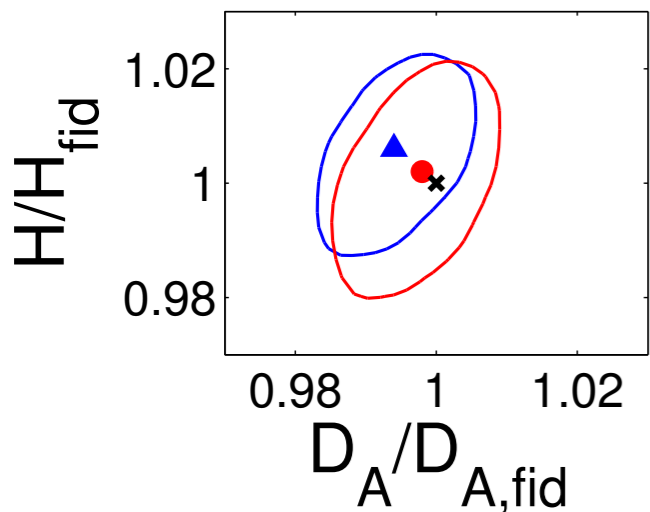


Fitting to P_0 & P_2 of N-body data,
 (D_A, H, f) are estimated

using MCMC

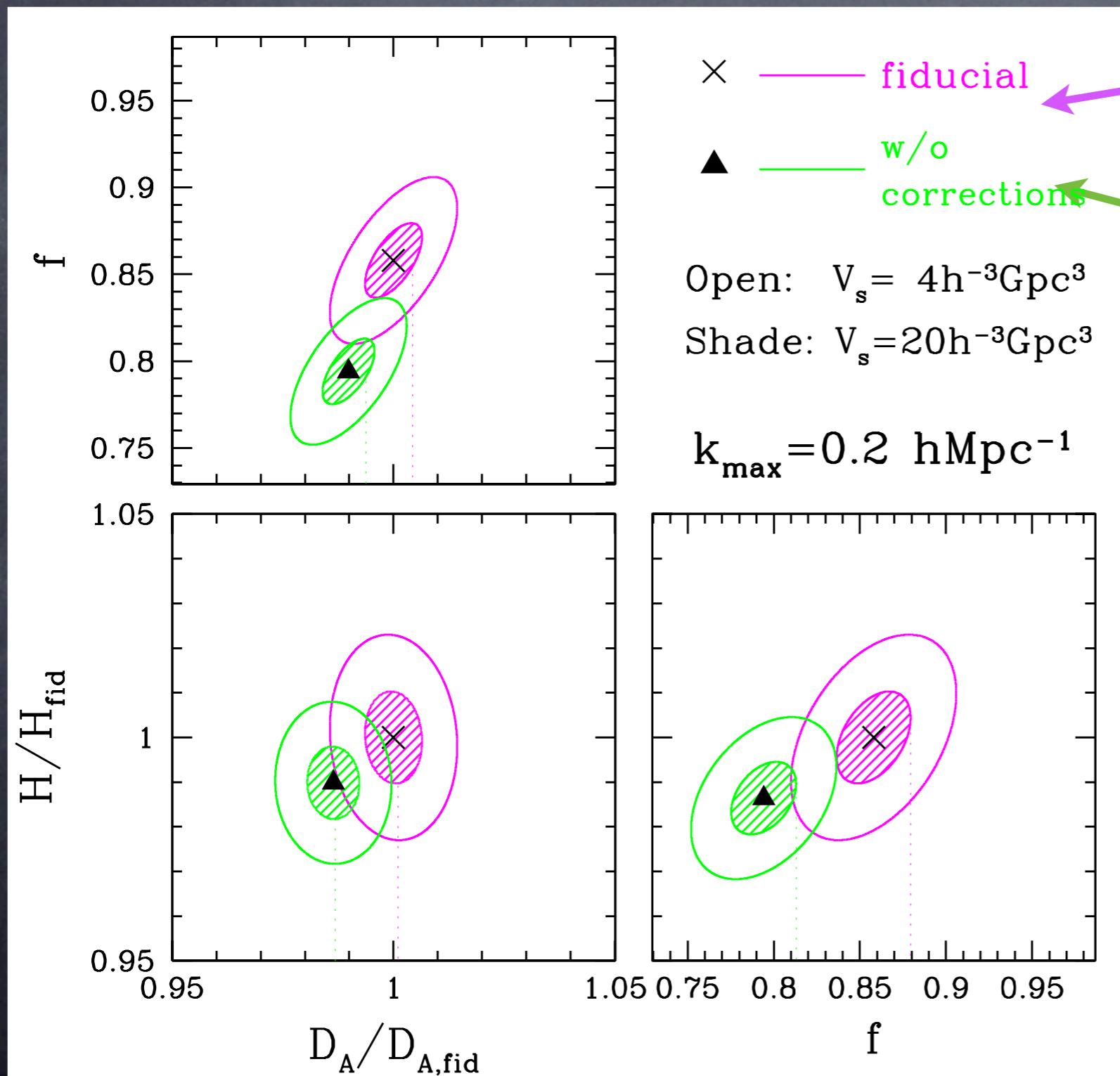


New model correctly recovers the input values



Impacts on future observations

Fisher matrix analysis using full 2D information



New model of redshift distortion

Phenomenological model (w/o corrections)

Assumptions

$z = 1$, linear bias ($b = 2$),
 $n_g = 5 \times 10^{-4} h^3 \text{Mpc}^{-3}$

Systematic biases:

- $D_A, H : 1 \sim 2 \%$
- $f : \sim 5 \%$

cannot be negligible even for stage III-class surveys

Summary

Modeling BAOs in 2D taking account of the effects of both non-linear clustering & redshift distortion



A new model of redshift distortion

Fisher-matrix analysis on systematic bias

- With improved PT, a (sub-)percent precision is achieved for predictions, and the model can correctly recover D_A , H & f
- Systematic bias caused by incorrect model assumption of redshift distortion would produce

$$\begin{cases} D_A(z) \ \& \ H(z) & : \text{small, but non-negligible (1}\sim\text{2\%)} \\ f(z) & : \text{significant } (\sim\text{5\%)} \end{cases}$$

crucial for stage III surveys

Appendix

Modeling non-linear P(k)

Analytic calculation based on perturbation theory (PT)

would be very useful complementary to N-body simulations

Perturbative expansion of (CDM+baryon) system described as pressureless fluid:

Density fluctuation

$$\delta = \delta^{(1)} + \delta^{(2)} + \delta^{(3)} + \dots$$



Power spectrum

$$\langle \delta(\vec{k}) \delta(\vec{k}') \rangle = (2\pi)^3 \delta_D(\vec{k} + \vec{k}') P(k)$$

◆ Standard PT

Jeong & Komatsu (2006)

◆ Improved perturbation theory

{ Renormalized Perturbation Theory (**RPT**)

Crocce &
Scoccimarro (2008)

{ **Closure Approximation (CLA)**

AT & Hiramatsu (2008)

Standard PT vs. Improved PT

$$P^{(mn)}(k) \sim \langle \delta^{(m)} \delta^{(n)} \rangle$$

Standard PT

$$P(k) = \underbrace{P^{(11)}(k)}_{\text{Linear (tree)}} + \underbrace{[P^{(13)}(k) + P^{(22)}(k)]}_{\text{1-loop}} + \underbrace{[P^{(33)}(k) + P^{(24)}(k) + P^{(15)}(k)]}_{\text{2-loop}} + \dots$$

$\propto D^2(t)$ $\propto D^4(t)$ $\propto D^6(t)$

Straightforward calculation based on naïve expansion

Improved PT (CLA)

$$P(k;t) = \underbrace{G^2(k | t, t_0)}_{\text{propagator}} P(k; t_0) + \underbrace{\int \int ds dt G(k | s, t_0) G(k | t, t_0) \Phi_{1\text{-loop}} [P(k); s, t]}_{\text{Mode-coupling term}}$$

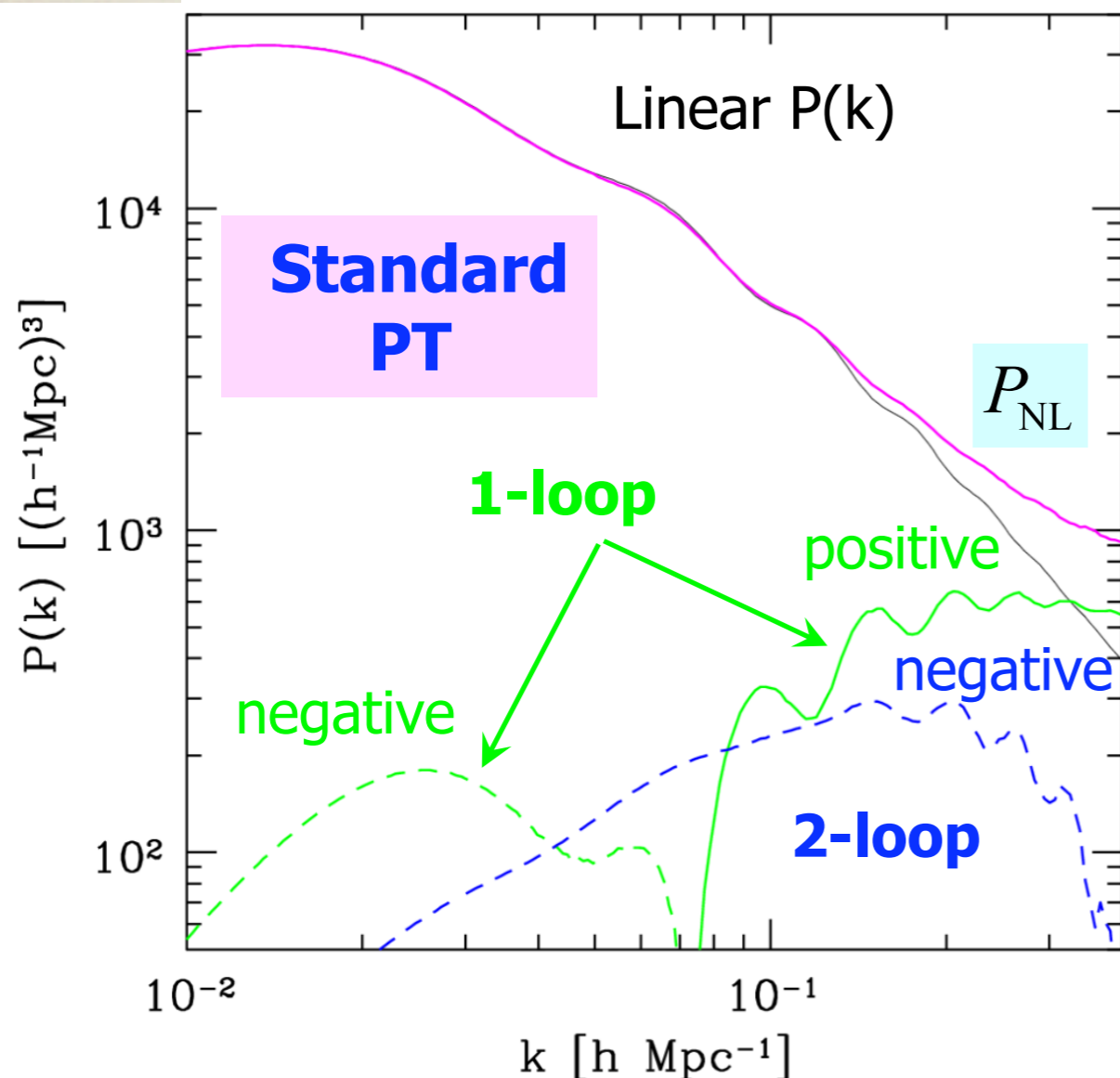
Initial P(k)

Non-perturbative effects is incorporated through **propagator**

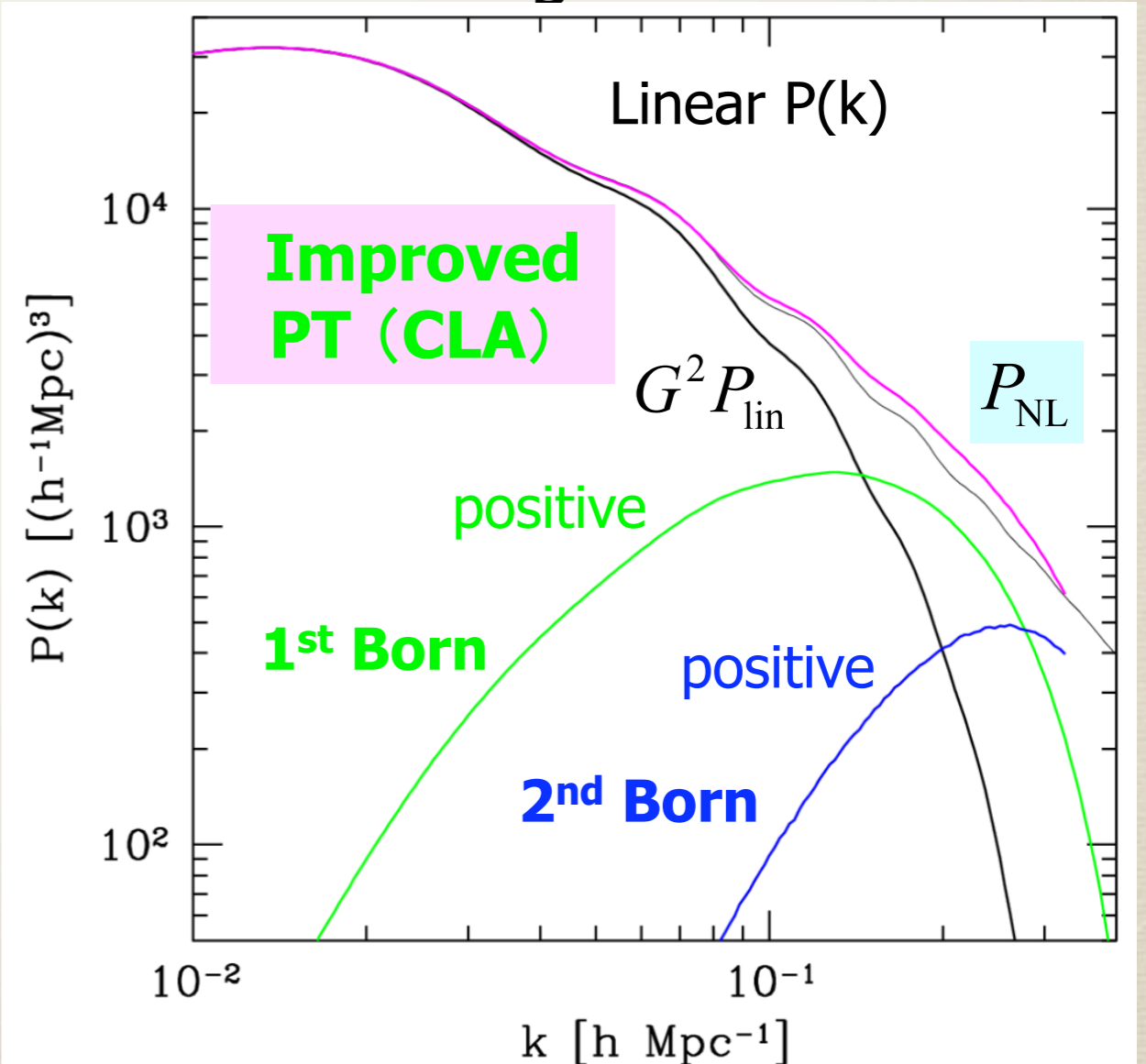
Iteratively evaluate **mode-coupling term** by Born approximation

Convergence of PT expansions

All contributions become comparable at low- z , positivity is not guaranteed



Contributions are all positive, shifted to higher k as increasing order of PT



Phenomenological models

$$P^{(S)}(k, \mu) = \underbrace{D[k\mu f \sigma_v]}_{\text{Damping func.}} \underbrace{(1 + f \mu^2)^2}_{\text{Linear Kaiser}} P_{\delta\delta}(k); \quad f \equiv d \ln D_+ / d \ln a$$

Damping func.

Linear Kaiser

fitting parameter
(1D velocity dispersion)

[Peacock & Dodds ('94); Cole et al.('95);
Ballinger et al. ('96); Magira et al. ('00)]

$$P^{(S)}(k, \mu) = \underbrace{D[k\mu f \sigma_v]}_{\text{Damping func.}} \underbrace{\{ P_{\delta\delta}(k) - 2f \mu^2 P_{\delta v}(k) + f^2 \mu^4 P_{vv}(k) \}}_{\text{Non-linear Kaiser}}$$

Damping func.

Non-linear Kaiser

[Scoccimarro ('04); Percival & White ('09);
Shoji et al. ('09)]

Damping func.

Finger-of-God
effect

$$D[x] = \begin{cases} 1/(1 + x^2) & \text{Lorentzian} \\ \exp\{-x^2\} & \text{Gaussian} \end{cases}$$

Comparison with N-body simulation

$$P^{(S)}(k, \mu) = \sum_{\ell=\text{even}} P_{\ell}^{(S)}(k) \mathcal{P}_{\ell}(\mu)$$

