

Three body decay of Gravitino and the indirect detection

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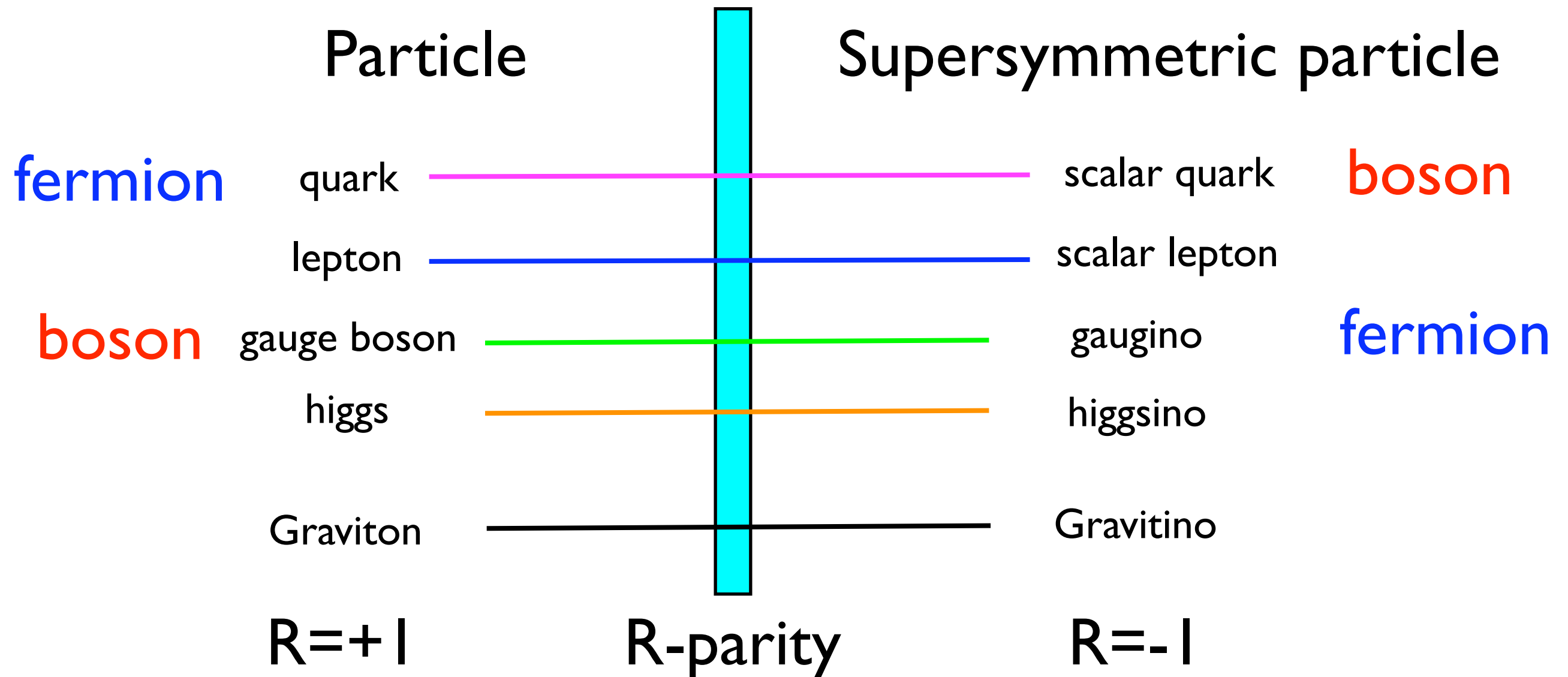
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1. Gravitino LSP and Dark Matter
2. R-parity violation and Gravitino decay
3. Gravitino three-body decay and indirect detection
4. Discussion

- Supersymmetry and Gravitino

Supersymmetry introduces the symmetry between bosons and fermions



► Lightest Spersymmetric Particle (LSP) is stable with R-parity

a good candidate for Dark Matter : Neutralino, sneutrino, Gravitino, Axino, etc

[Talk by L.Roszkowski and L. Covi]

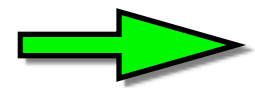
- Gravitino is massless in the supersymmetric limit

 become massive with SUSY breaking

$$m_{\tilde{G}} = \langle W e^{K/2} \rangle = \frac{F_X}{M_P}$$

from eV to TeV depending on the SUSY breaking mechanism
eg, GMSB, Gravity-MSB, gaugino-MSB, Anomaly-MSB,

- Interaction is suppressed by Planck scale

 suffers from the Gravitino problem

$$\tau \sim 10^2 - 10^{12} \text{ sec}$$

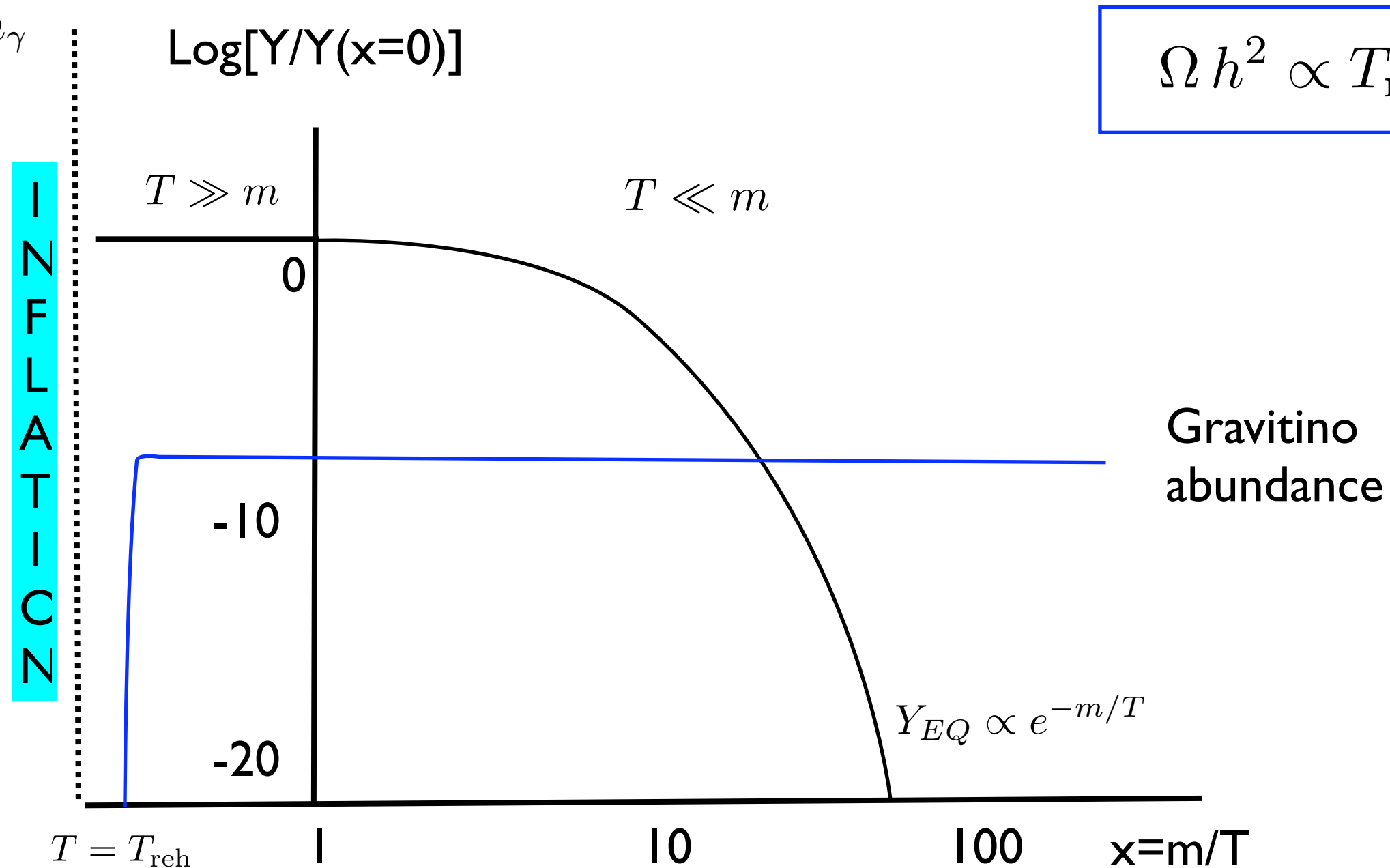
: disturb the standard Big Bang Nucleosynthesis

- Production of gravitinos

► **Thermal production:** from the scatterings of the particles in the thermal equilibrium

[Ellis, Kim, Nanopoulos, 1984]

$$Y \equiv \frac{n}{n_\gamma}$$



$$\Omega h^2 \propto T_{\text{reh}}$$

► Non-thermal production

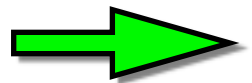
: from the decay of next LSP such as neutralino, stau, etc

and/or from the decay of inflaton, during preheating

or from the decay of moduli field

► Total relic density of gravitino

$$\Omega_{\tilde{G}} h^2 = \Omega_{\tilde{G}}^{TP} h^2 + \Omega_{\tilde{G}}^{NTP} h^2$$



Gravitinos can be produced to give correct amount
required for Dark Matter

- Supersymmetry with R-parity violation and Gravitino LSP

No symmetry to protect the decay of LSP

Gravitino LSP can decay

: However Gravitino lifetime can be much longer than the age of Universe with small R-parity breaking and Planck suppressed interaction

and still can be a good candidate for dark matter

- Bilinear R-parity violation and Gravitino DM [Takayama, Yamaguchi, 2000]

$$W = \mu H_1 H_2 + \sum_{i=1}^3 \mu_i L_i H_2$$

: No BBN problem due to early decay of NLSP

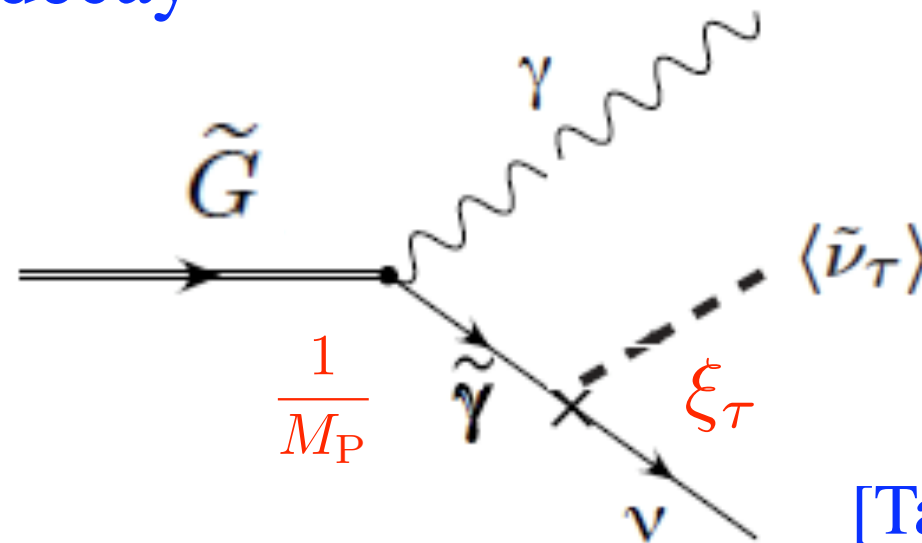
$$\mathcal{L}_{soft} = B_i \tilde{L}_i H_u + m_{L_i H_2}^2 \tilde{L}_i H_1^* + h.c.$$

: gives sneutrino VEV and mixing between neutrino and neutralino

$$\langle \tilde{\nu}_i \rangle = \frac{B_i \sin \beta + m_{L_i H_d}^2 \cos \beta}{m_{\tilde{\nu}_i}^2} v$$

- Gravitino two-body decay

$$\tilde{G} \rightarrow \gamma + \nu$$



[Takayama, Yamaguchi, 2000]

➡
$$\Gamma(\tilde{G} \rightarrow \gamma \nu) = (10^{23} \text{ sec})^{-1} \left(\frac{m_{\tilde{G}}}{100 \text{ GeV}} \right)^3 \left(\frac{\xi_\tau}{10^{-7}} \right)^2 |U_{\tilde{\gamma}\tilde{Z}}|^2$$

$$\xi_\tau \equiv \frac{\langle \tilde{\nu}_\tau \rangle}{v} \quad \langle \tilde{\nu}_i \rangle = \frac{B_i \sin \beta + m_{L_i H_d}^2 \cos \beta}{m_{\tilde{\nu}_i}^2} v$$

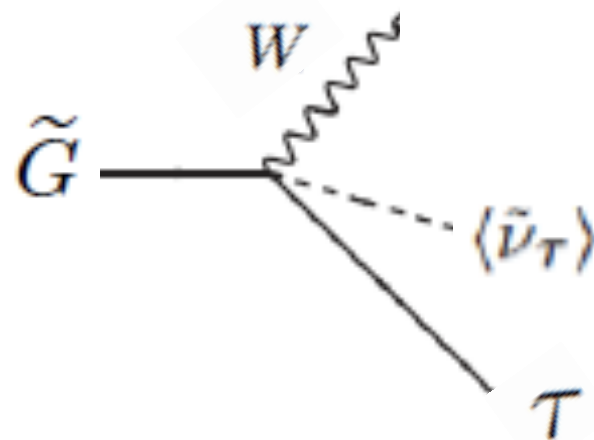
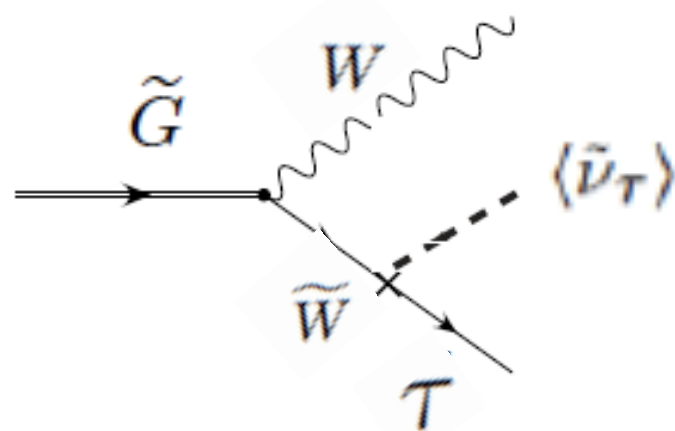
$$|U_{\tilde{\gamma}\tilde{Z}}| \simeq \frac{M_Z(M_2 - M_1)s_W c_W}{(M_1 c_W^2 + M_2 s_W^2)(M_1 s_W^2 + M_2 c_W^2)} \quad [\text{Ibarra, Tran, 2008}]$$

M_1, M_2 : gaugino mass s_W, c_W : weak mixing angle

- Gravitino two-body decay: If Gravitino is massive than W and Z bosons

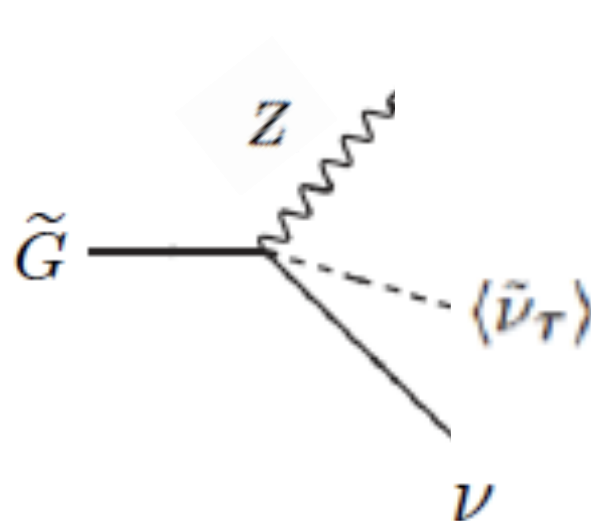
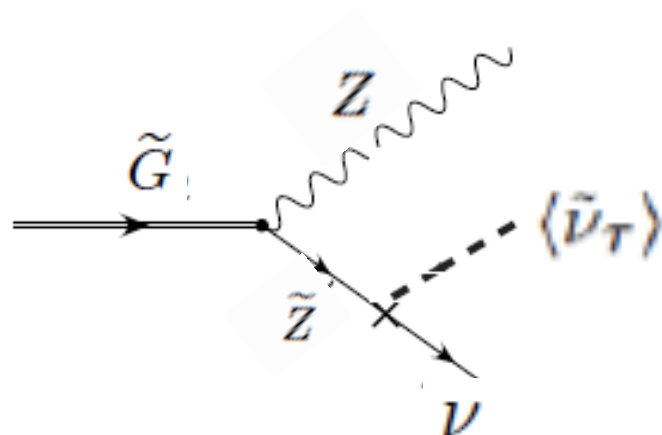
[Ibarra, Tran, 0709.4593]

$$\tilde{G} \rightarrow W + \tau$$



$$\Gamma \propto \frac{\xi_\tau^2}{M_{\text{P}}^2} (A|U_{\tilde{W}\tilde{W}}|^2 + BU_{\tilde{W}\tilde{W}} + C)$$

$$\tilde{G} \rightarrow Z + \nu \quad : \text{similar diagrams to W-boson production}$$



- Gravitino two-body decay and the indirect detection

Leptogenesis and gamma ray signature [[Buchmuller, Covi, Hamaguchi, Ibarra, Yanagida, 2007](#)]

Gamma rays from decaying dark matter [[Bertone, Buchmuller, Covi, Ibarra, 2007](#)]
[[Ibarra, Tran, 0709.4593](#)]

Antimatter signatures of Gravitino dark matter decay [[Ibarra, Tran, 0804.4596](#)]

High energy cosmic rays from the decay of Gravitino DM decay
[[Ishiwata, Matsumoto, Moroi, 0805.1133](#)]

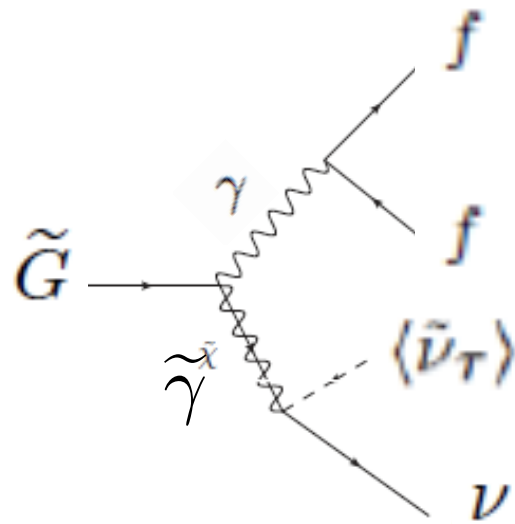
Unstable Gravitino dark matter and neutrino flux [[Covi, Grefe, Ibarra, Tran, 0809.5030](#)]

Probing Gravitino dark matter with PAMELA and Fermi
[[Buchmuller, Ibarra, Shindou, Takayama, Tran, 2009](#)]

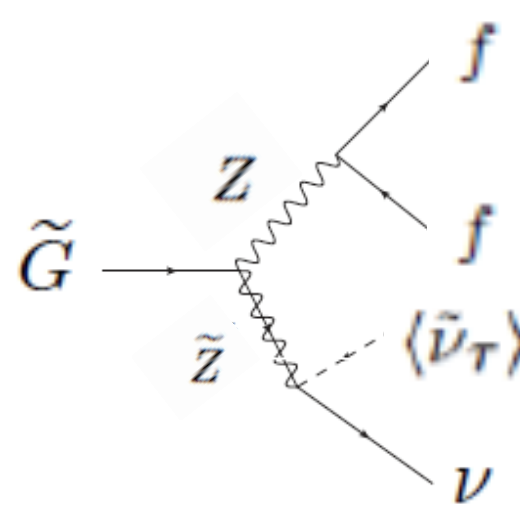
Gamma-ray detection from Gravitino DM decay in the MuNuSSM
[[K.Y.Choi, Lopez-Fogliani, Munz, Ruiz de Austri, 2010](#)]

• Three-body decays of Gravitino

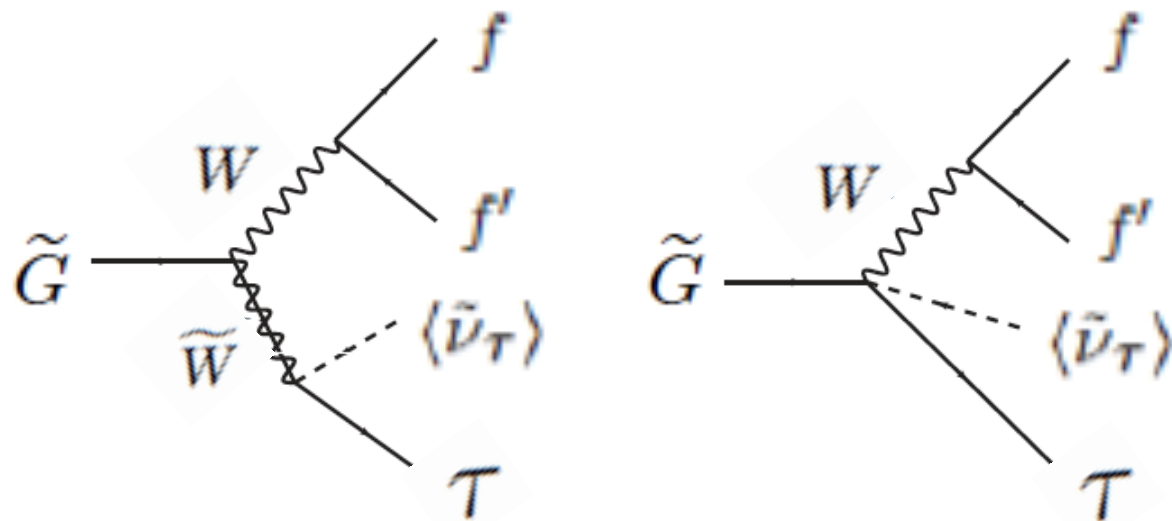
[KYChoi, Yaguna, 2010]



$$\tilde{G} \rightarrow \gamma^* + \nu \rightarrow f + f + \nu$$



$$\tilde{G} \rightarrow Z^* + \nu \rightarrow f + \bar{f} + \nu$$



$$\tilde{G} \rightarrow W^* + \tau \rightarrow f + \bar{f}' + \tau$$

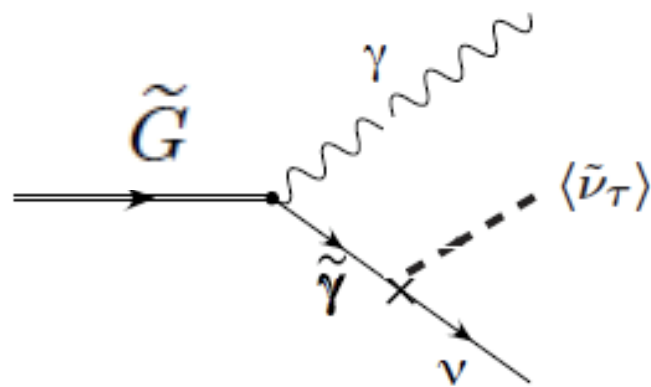
► Three-body decays:

usually suppressed compared to two-body decays

► However unusual situation happens

: just below the threshold of new decay modes $m_{\tilde{G}} < m_W$

► Two-body decay: Gravitino to photon and neutrino

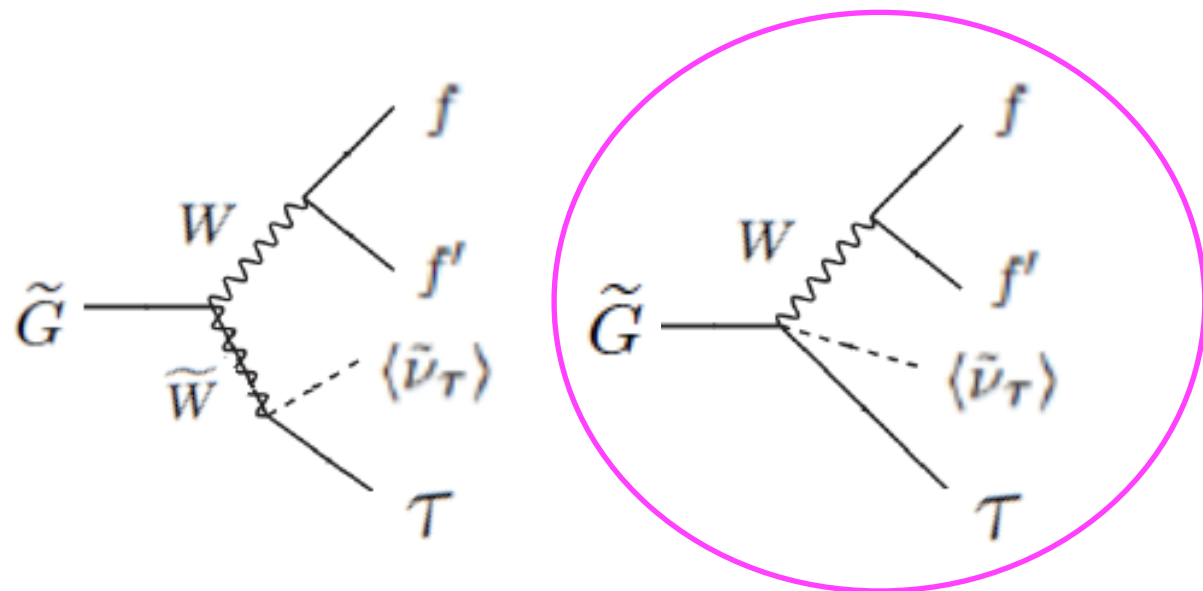


$$\Gamma \propto \frac{\xi_\tau^2}{M_P^2} |U_{\tilde{\gamma}\tilde{Z}}|^2 \propto \left(\frac{M_Z}{M_1}\right)^2$$

: suppressed by the mass of gaugino

* Gravitino two-body decay to W or Z boson is kinematically disallowed

► Three-body decays mediated by W- and Z-boson



$$\Gamma \propto \frac{\xi_\tau^2}{M_P^2} \left(A |U_{\tilde{W}\tilde{W}}|^2 + B U_{\tilde{W}\tilde{W}} + C \right) \propto \left(\frac{M_Z}{M_1}\right)^2 \propto \left(\frac{M_Z}{M_1}\right)$$

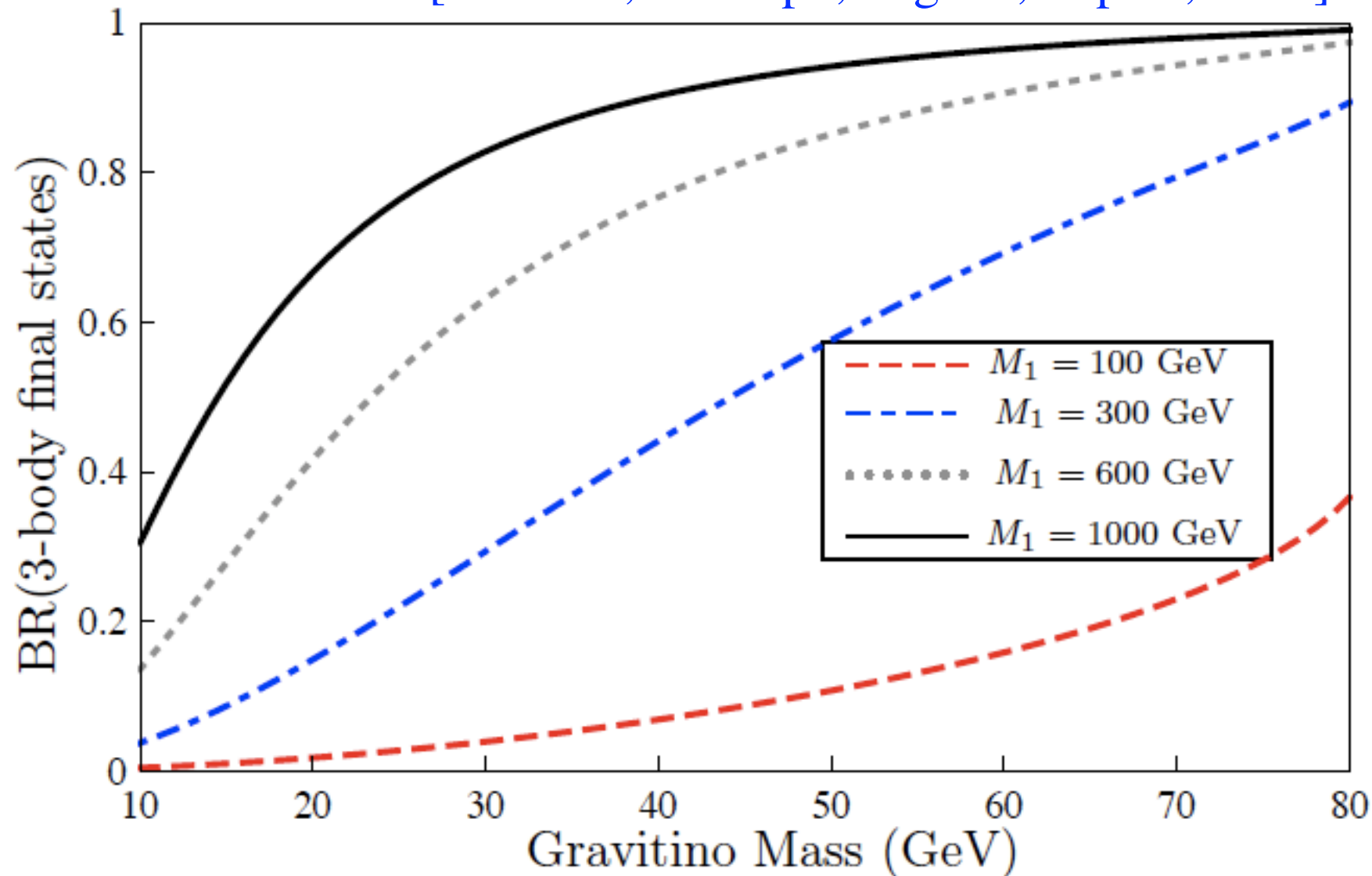
* Similar to Z-boson mediated ones

► 4-vertex interaction

: No suppression due to the gaugino mass

- Branching ratio of three-body decays of Gravitino
 - ▶ With large gaugino mass, 3-body decay dominates 2-body decay

[KYChoi, Restrepo, Yaguna, Zapata, 2010]



M_1 : Bino mass

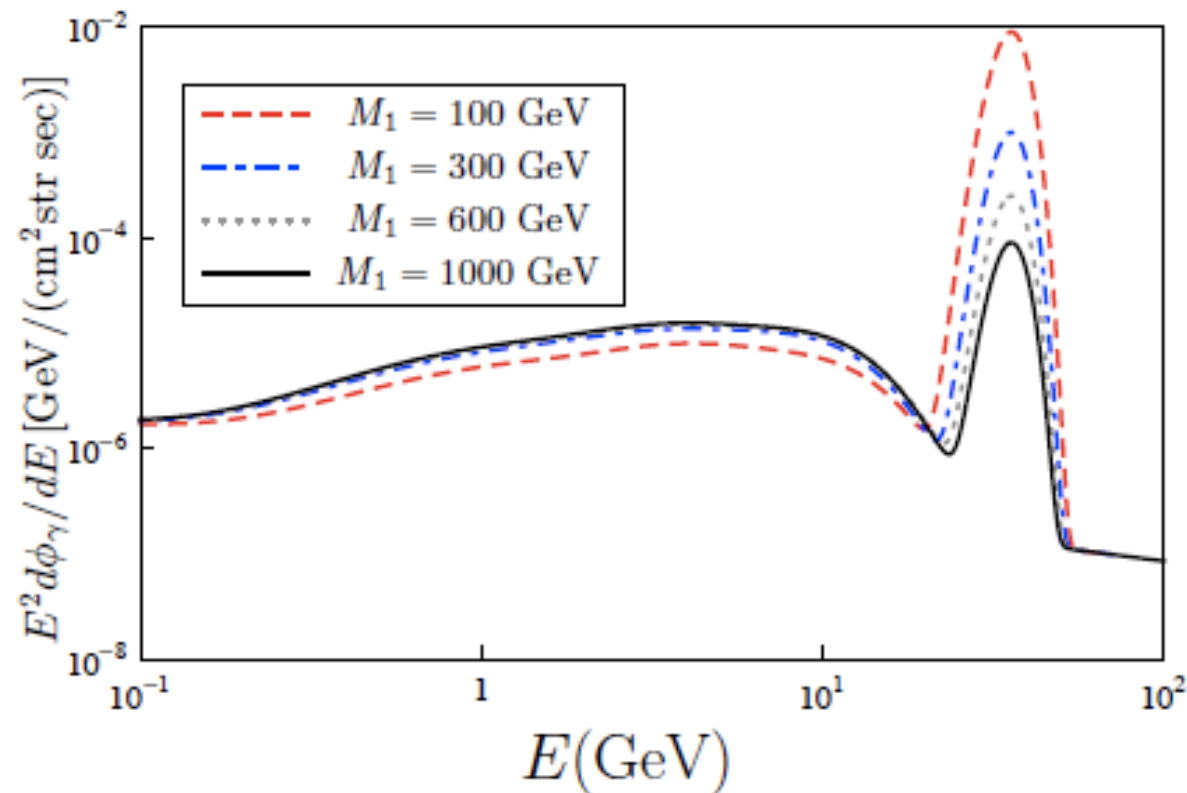
We assumed GUT unification of gaugino masses

- Gamma rays from Gravitino decay

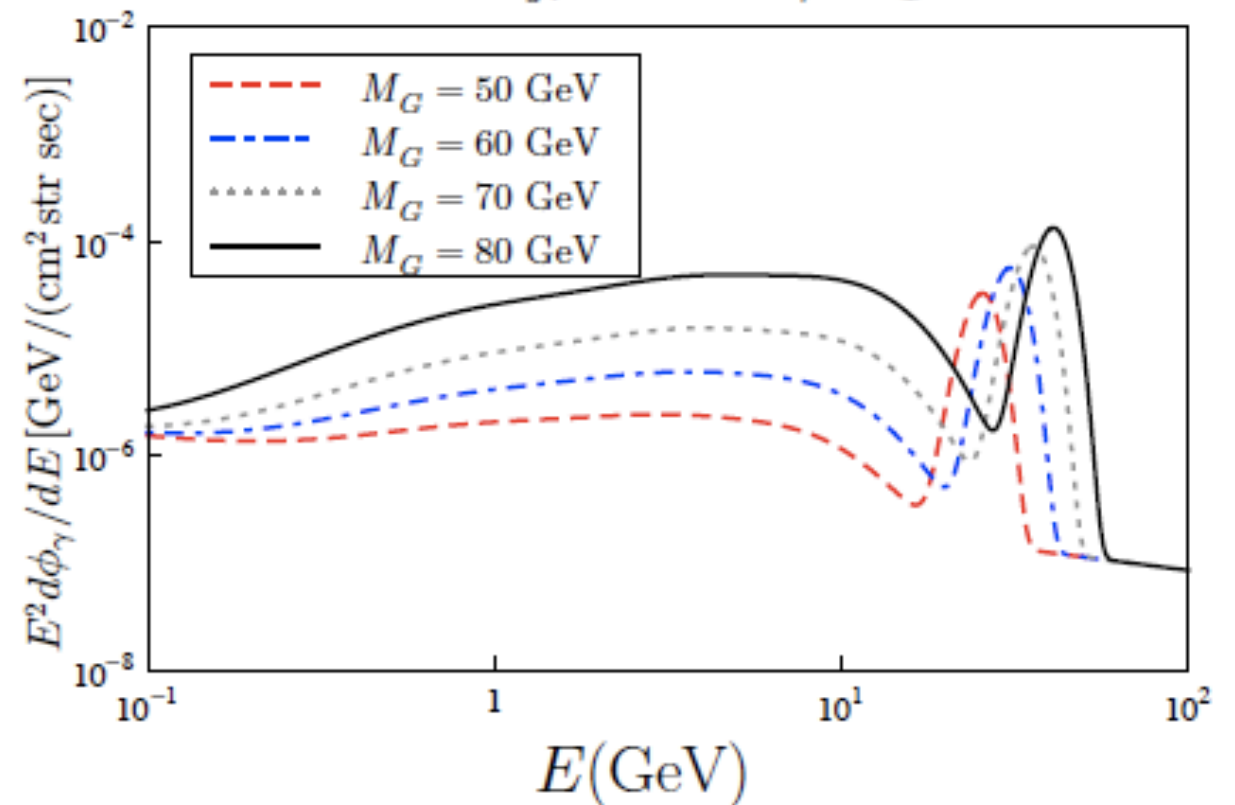
► **Monochromatic** from 2-body decay + **continuum** from 3-body decay

[KYChoi, Restrepo, Yaguna, Zapata, 1007.1728]

$$\xi_\tau = 10^{-7}, m_{\tilde{G}} = 70 \text{ GeV}$$



$$\xi_\tau = 10^{-7}, M_1 = 1 \text{ TeV}$$

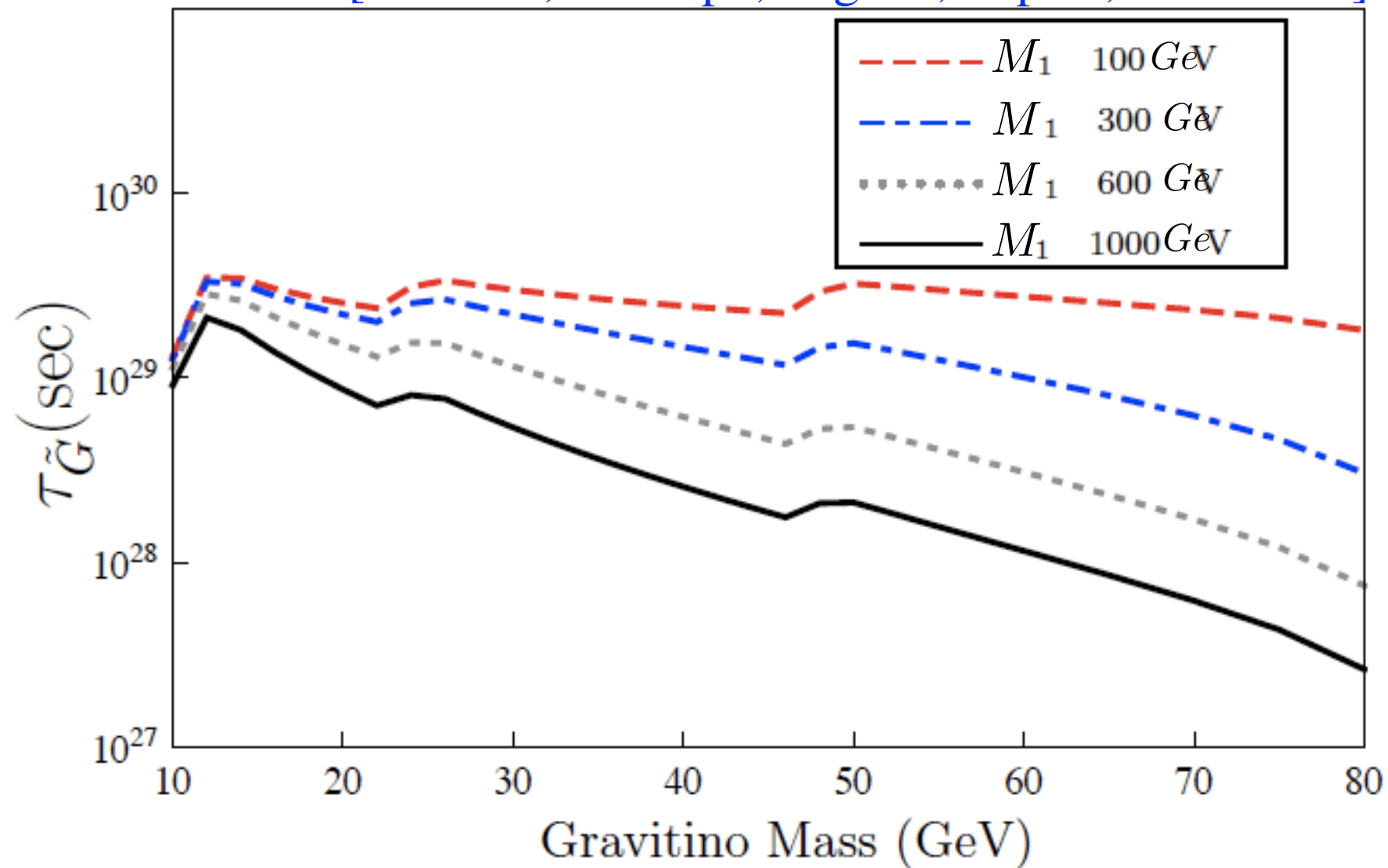


- Line signal is sensitive to the gaugino mass
- For larger gaugino mass, the line signal becomes smaller and continuum becomes larger

- Larger signal for larger gravitino mass
- The position of line signal depends on the mass of gravitino

- Constraint on the lifetime from diffuse gamma ray of Fermi

[KYChoi, Restrepo, Yaguna, Zapata, 1007.1728]



- The constraint on lifetime changes by one order of magnitude depending on the mass of gaugino

- Antimatter searches from Gravitino decay

- ▶ Three-body decay give rise to non-zero positron and antiproton flux

With compatible range of parameters from gamma ray observation it does not yield a visible signature for **positron flux** from gravitino decay

However for **antiproton flux**, it is marginally consistent with observation and may show some signature

[KYChoi, Restrepo, Yaguna, Zapata, 1007.1728]

Summary

- **Gravitino** can be the Lightest Supersymmetric Particle and become a good candidate for dark matter
- With slight **R-parity violation**, Gravitino can be Dark Matter and show Indirect Detection signatures
- **Three-body decay of Gravitino** can dominate the two-body decay and its lifetime and the signatures from the indirect detection are severely modified