

Heavy Element Production and Neutrino-Driven Jets in Neutron-Star Merger

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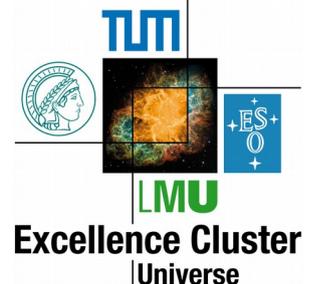
**RIKEN-RESCEU Joint Seminar
Tokyo, Japan, July, 27th, 2016**

*With: H.-Th. Janka, M. Obergaulinger, S. Goriely,
A. Bauswein, R. Ardevol, and others*

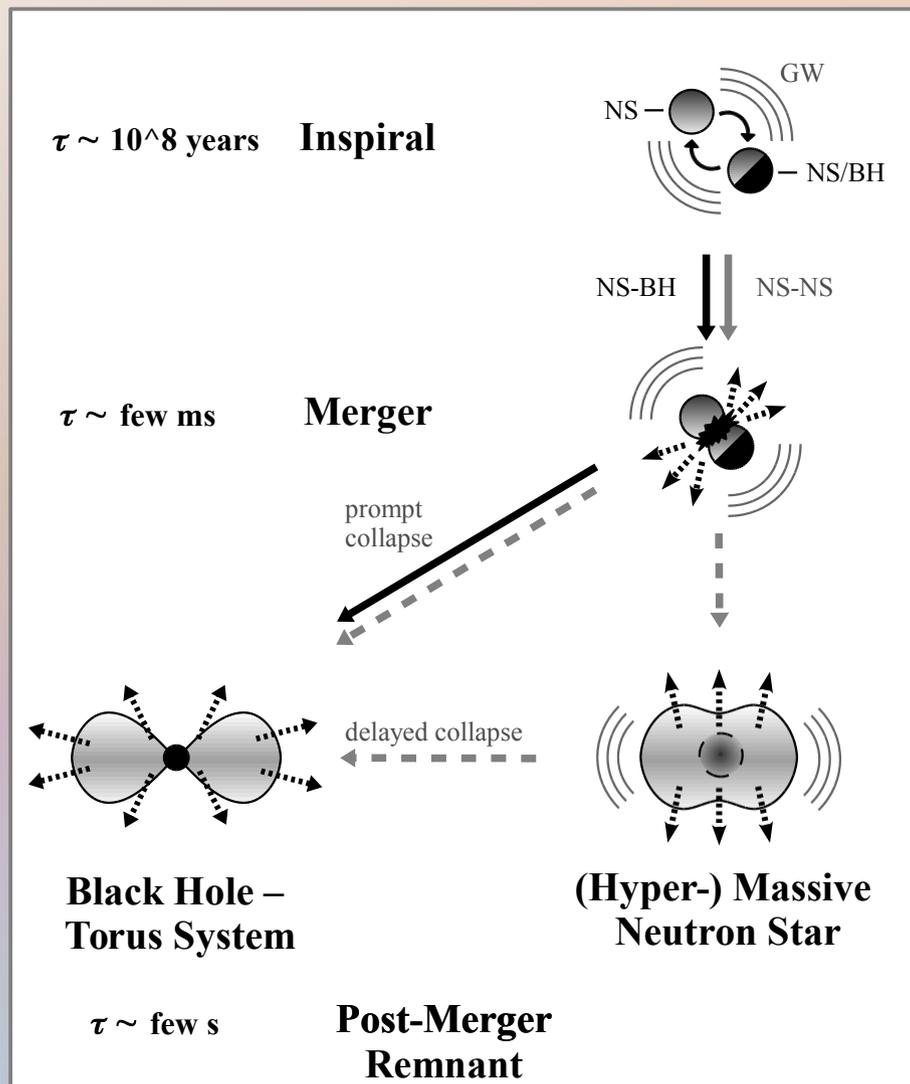


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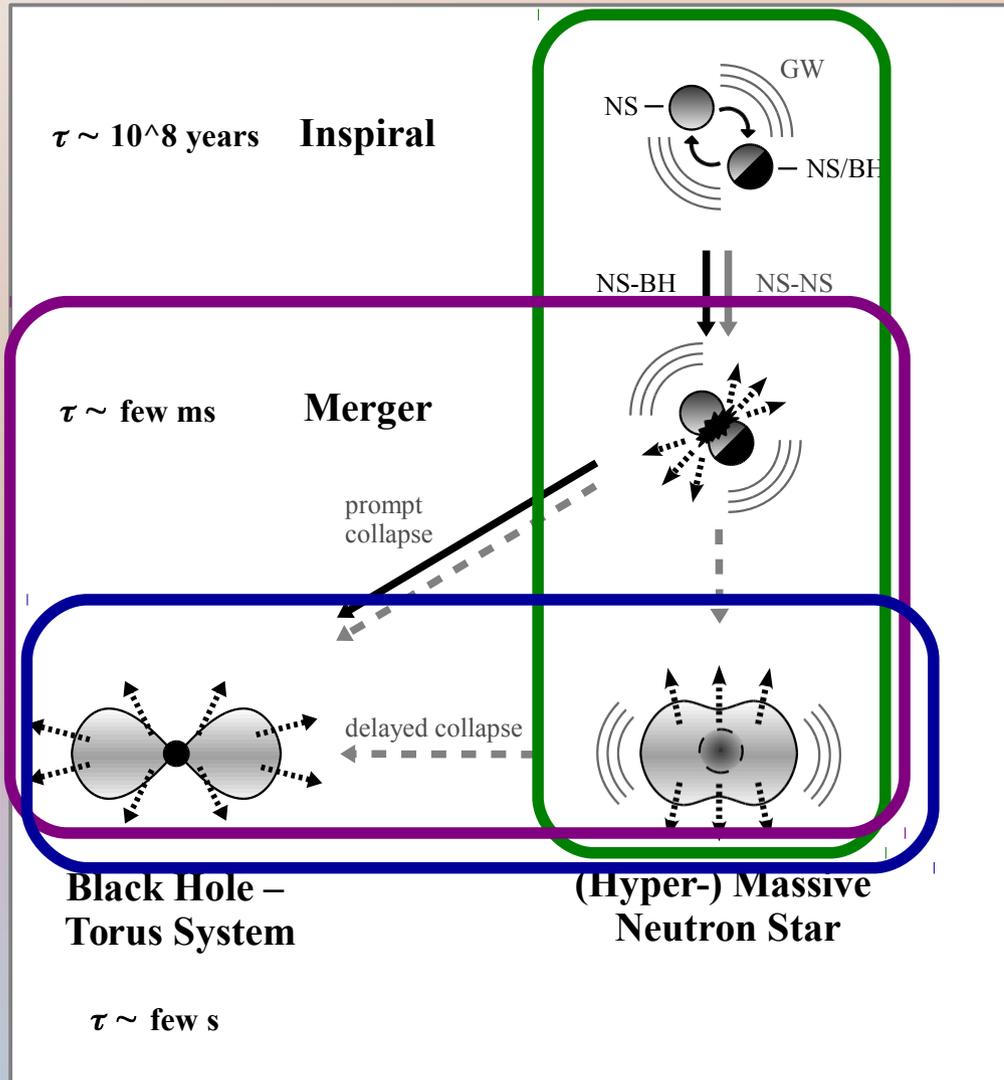
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Evolution Channels, Multi-Messenger Astronomy



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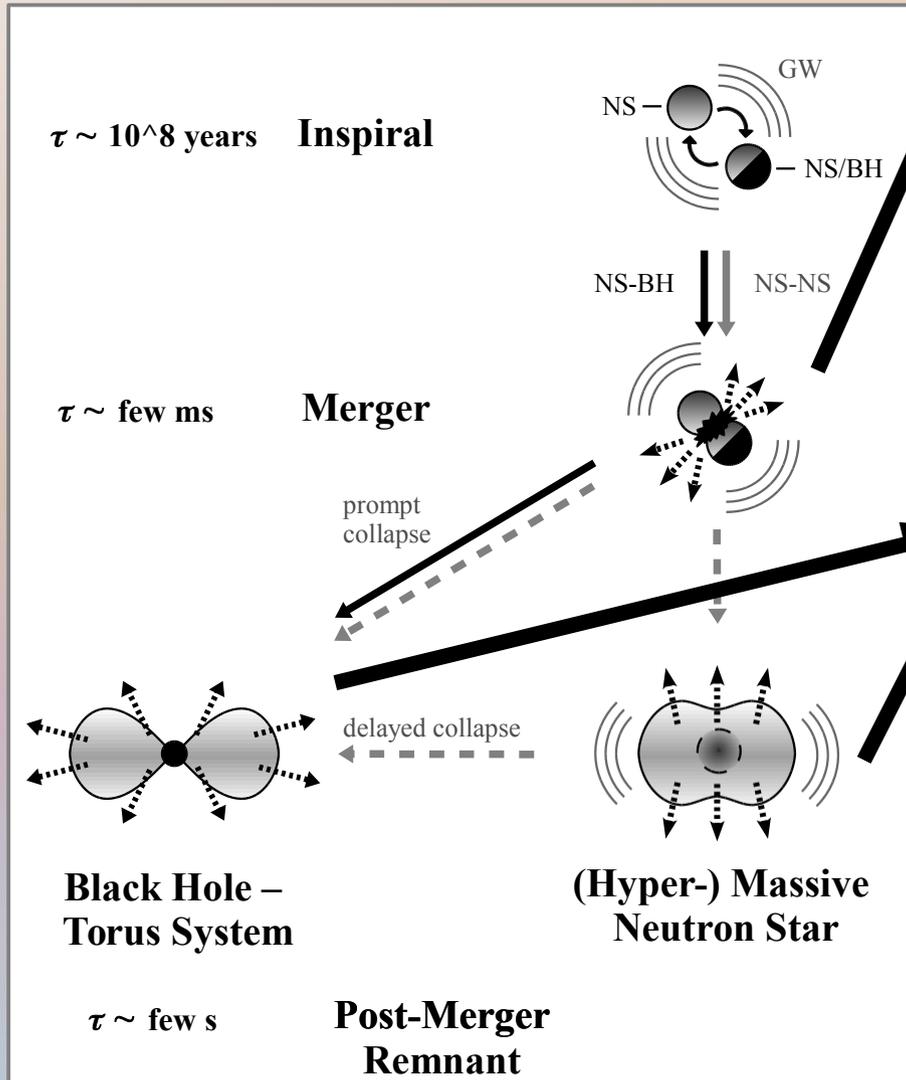


→ GW signal
 ... binary parameters?
 ... nuclear EOS?

→ short GRB
 ... from BH-torus system?
 ... from (H)MNS?

→ Massive Ejecta
 ... sources for r-process elements?
 ... observable electromagnetic signals (Macro-/Kilonova)?

Ejecta Components, Modeling Ingredients



dynamical/prompt ejecta

→ tidal tails

→ shock-heated

3D, GR, ν -transport, MHD

post-merger ejecta

→ neutrino-driven

→ viscous expansion

→ MHD turbulence

ν -tran, MHD/Vis, 3D, GR

Ingredients For Modeling NSMs (and CCSNe)

- ✓ **6D neutrino transport**
 - ✓ **nuclear EOS**
 - ✓ **detailed weak interaction cross sections**
 - ✓ **general relativistic effects**
 - ✓ **magnetic fields**
 - ✓ **sufficient (?) resolution**
 - ✓ **...**
- **all combined only feasible with exa-scale computers (or even much bigger)**
- **algorithms with reasonable compromise between accuracy and efficiency highly desired**

“ALCAR” Neutrino Transport Module (OJ, Obergaulinger, Janka '15, MNRAS, 453, 3386)

Radiation-hydro with Boltzmann solver too expensive!

Our approach:

→ Spectral Two-moment scheme with algebraic Eddington factor (**AEF or M1 scheme**)

$$E = \int d\Omega \mathcal{I}(\mathbf{x}, \mathbf{n}, \epsilon, t) \quad \leftarrow \text{energy density}$$

$$F^i = \int d\Omega \mathcal{I}(\mathbf{x}, \mathbf{n}, \epsilon, t) n^i \quad \leftarrow \text{momentum density}$$

$$P^{ij} = \int d\Omega \mathcal{I}(\mathbf{x}, \mathbf{n}, \epsilon, t) n^i n^j \quad \leftarrow \text{pressure}$$

$$Q^{ijk} = \int d\Omega \mathcal{I}(\mathbf{x}, \mathbf{n}, \epsilon, t) n^i n^j n^k$$

$$\left. \begin{aligned} \partial_t E + \nabla_j F^j + \nabla_j (v^j E) + (\nabla_j v_k) P^{jk} - (\nabla_j v_k) \partial_\epsilon (\epsilon P^{jk}) &= C^{(0)} \\ \partial_t F^i + c^2 \nabla_j P^{ij} + \nabla_j (v^j F^i) + F^j \nabla_j v^i - (\nabla_j v_k) \partial_\epsilon (\epsilon Q^{ijk}) &= C^{(1),i} \end{aligned} \right\} \text{evolution equations}$$

$$\left. \begin{aligned} P^{ij} &= P^{ij}(E, F^i) \\ Q^{ijk} &= Q^{ijk}(E, F^i) \end{aligned} \right\} \text{approximate algebraic closure relations (e.g. "M1 closure")}$$

(also several other realizations by Kuroda '15, O'Connor '15, Skinner '15, Foucart '15)

Algorithm Details

→ Advantages:

- effective save up of the two angular degrees of freedom
- truly multidimensional instead of Ray-by-Ray(-plus)
- computational costs: \sim hydro times # energy bins (more expensive with energy-bin coupling interactions)
- since equations are hyperbolic they can be solved just like hydro-eq.
- describes shadows better than FLD
- better to parallelize than FLD, because advection is explicit in time

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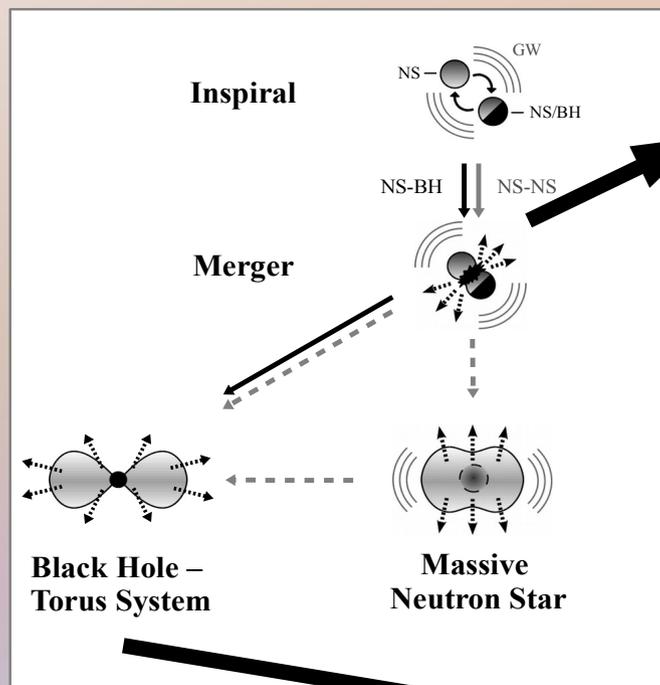
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→ Disadvantages:

- less accurate for more complex geometries of radiation sources
- AEF closure introduces non-linearity \rightarrow light fronts interact with each other, unphysical shocks can develop (“two-beam instability”)
- accuracy for CCSNe and NSMs only partially explored so far

What and how many heavy elements are ejected in which phase of a NS-NS/BH merger?

(OJ, Bauswein, Ardevol, Goriely, Janka '15, MNRAS 448, 541)



NS-NS and NS-BH merger phase modeled with relativistic 3D code (without neutrinos)

Nucleosynthesis analysis of prompt ejecta

Hydrodynamic simulations

Post-processing of ejecta

BH-torus phase modeled in 2D axisymmetry with spectral neutrino transport

Nucleosynthesis analysis of disk ejecta

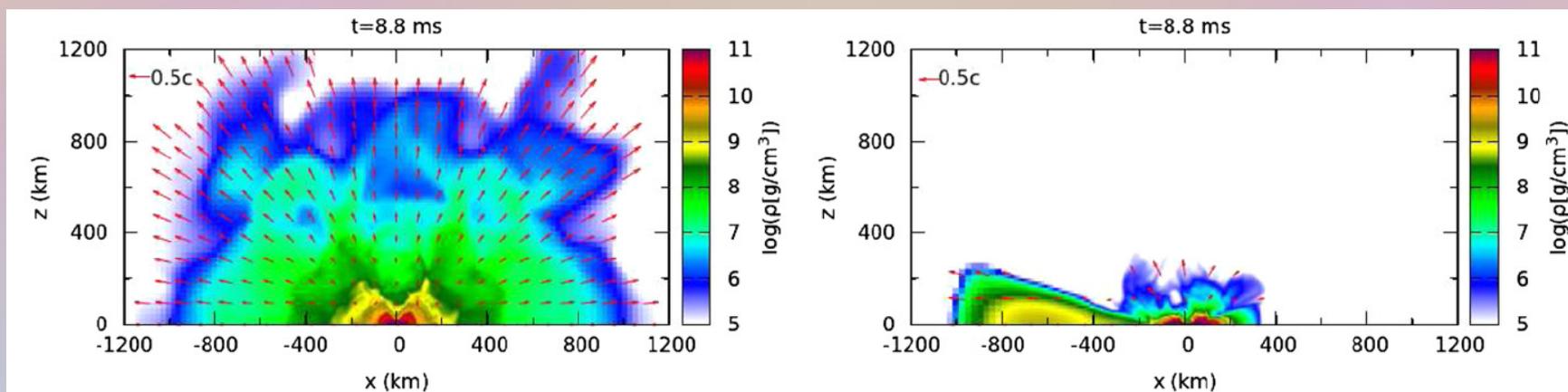
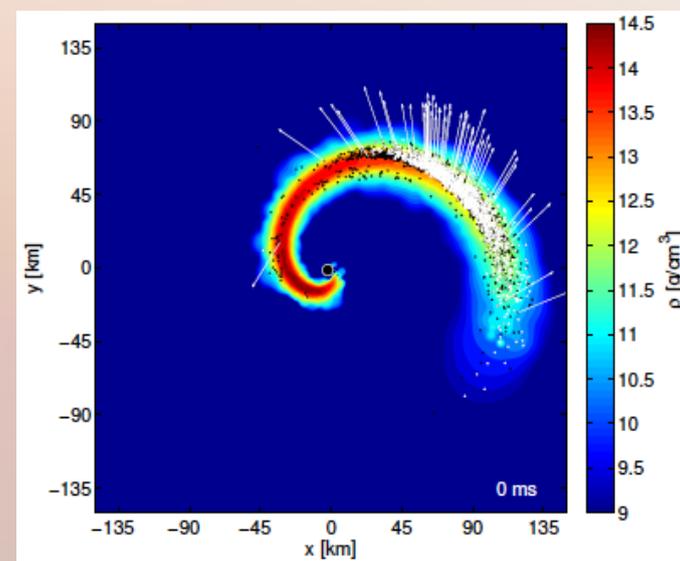
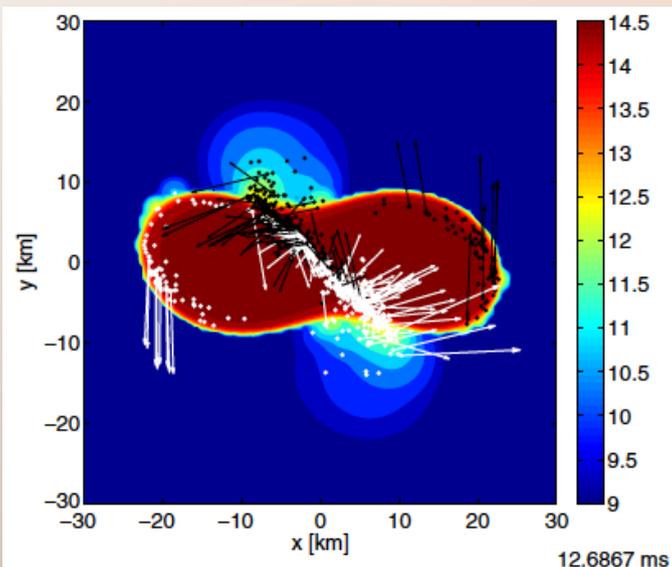
Dynamical Ejecta

NS-NS

Typical properties:

NS-BH

- outflow masses:
 $\sim 0.001 - 0.1 M_{\text{sun}}$
- electron fraction:
 $Y_e < 0.1$ (*)
- entropy per baryon:
 $s \sim 1 - 30 \text{ kB}$
- velocity:
 $v \sim 0.1 - 0.4 c$

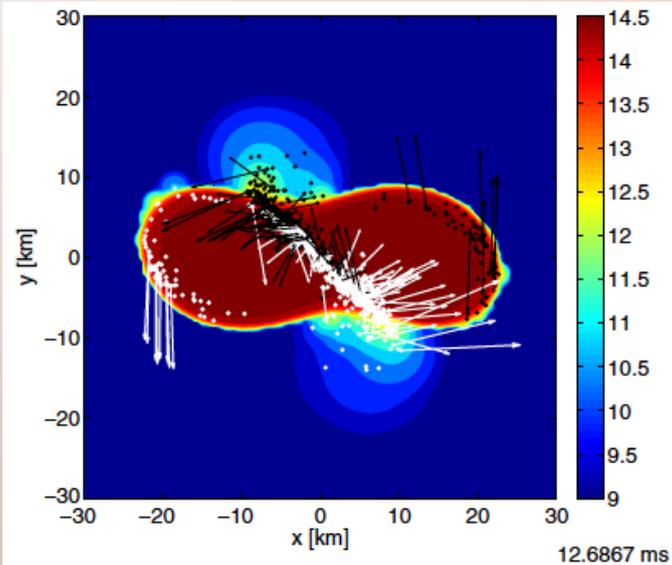


(Hotokezaka et. al. '13)

(* : Depends on neutrino treatment for NS-NS mergers)

Dynamical Ejecta

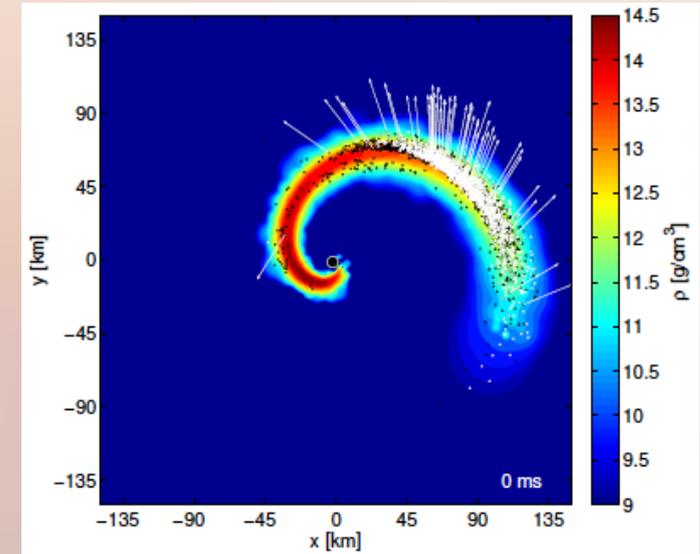
NS-NS



Typical properties:

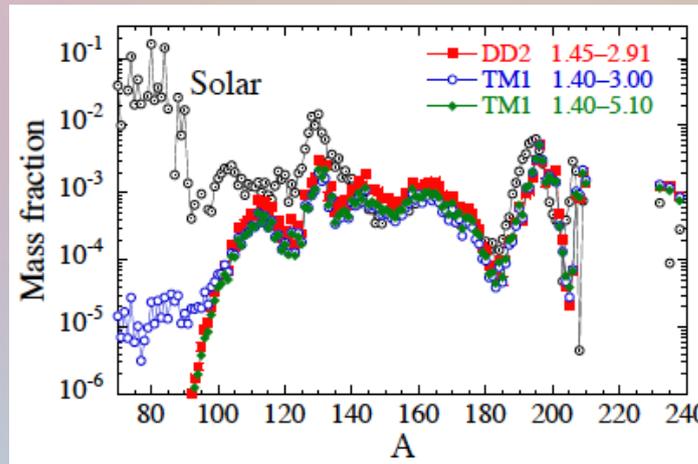
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NS-BH



Typical nucleosynthesis pattern:

→ sub-solar for $A < 140$ (*)



→ solar-like for $A > 140$

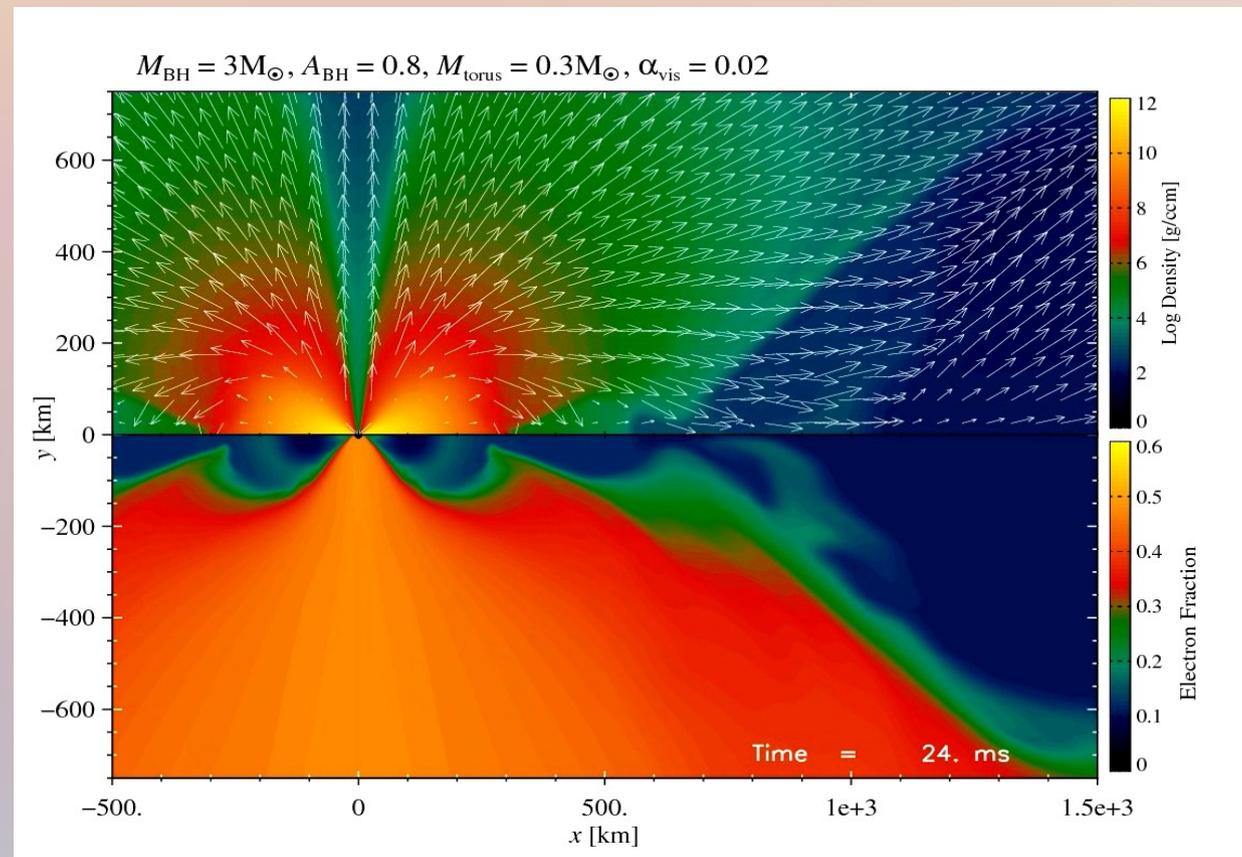
(* : Depends on neutrino treatment for NS-NS mergers)

Ejecta from BH-torus remnant

*comparable to mass
of dynamical ejecta!!!*

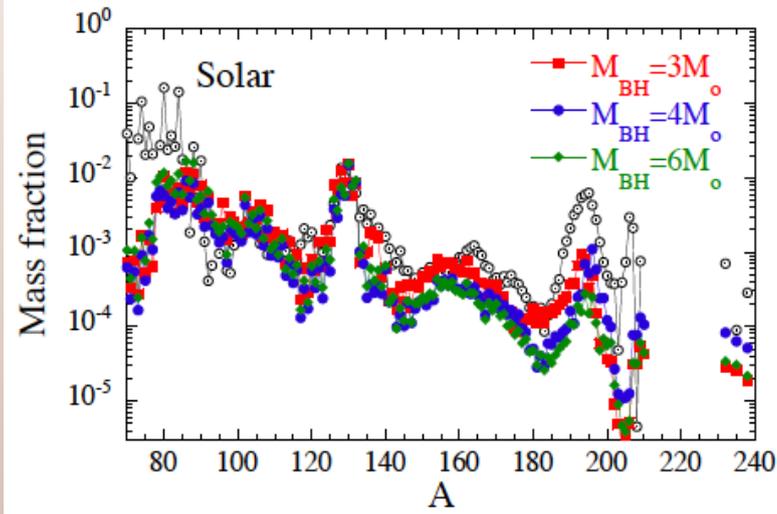
Typical properties:

- outflow masses:
~ 5-20% of initial
disk mass
- electron fraction:
 $Y_e \sim 0.1-0.3$
- entropy per baryon:
 $s \sim 10 - 30 \text{ kB}$
- velocity:
 $v \sim 0.05 - 0.1 c$
- **small** neutrino-driven
component
- **dominant** viscous
component

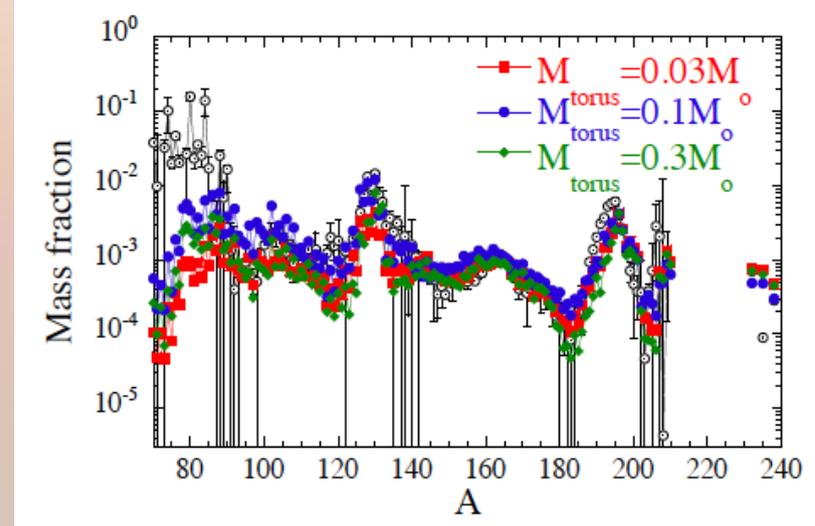


Combined nucleosynthesis yields

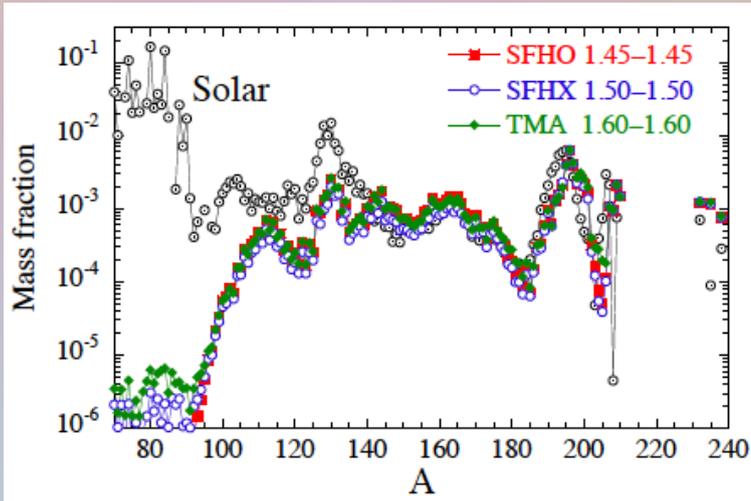
→ DISK ejecta (mainly $A \sim 90 - 140$)



→ DISK + PROMPT ejecta



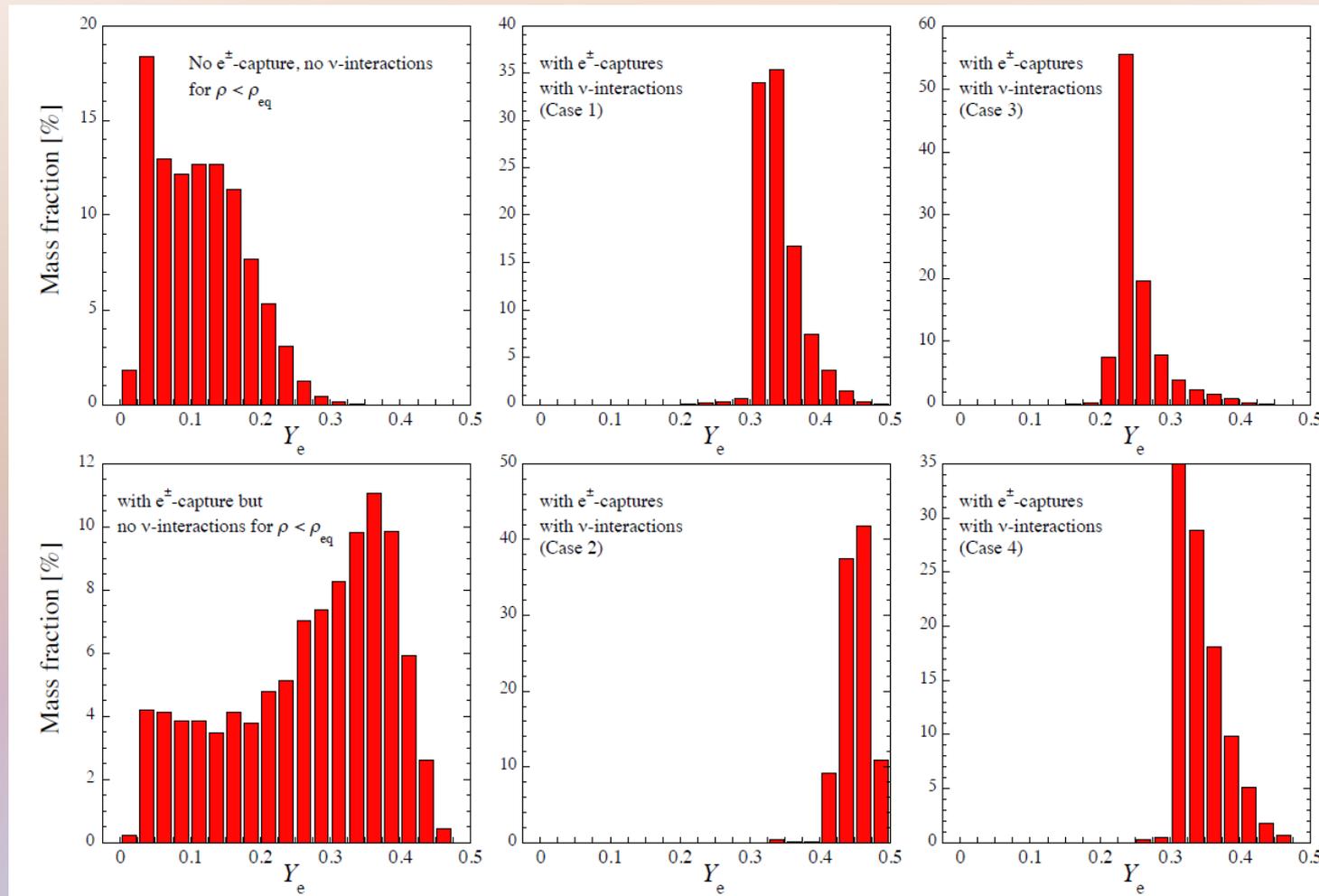
→ PROMPT ejecta (mainly $A \sim 140 - 210$)



- nicely recovers the full mass range $A > 90$
- disk ejecta could be significant sources of intermediate mass elements with $90 < A < 140$
- NS-remnant ejecta give similar pattern (Metzger+14, Perego+'14, Martin+'15)

Impact of Weak Interaction on Dynamical Ejecta in NS-NS Mergers?

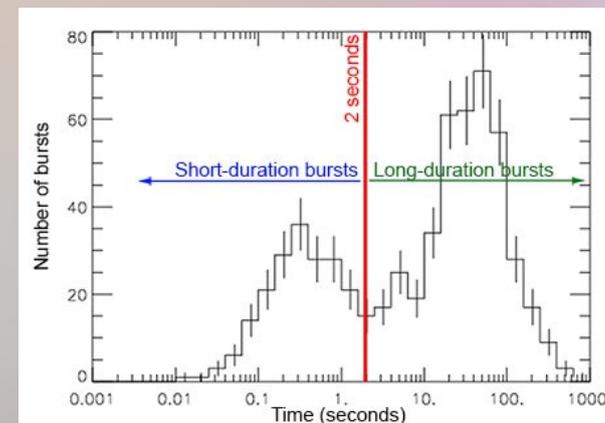
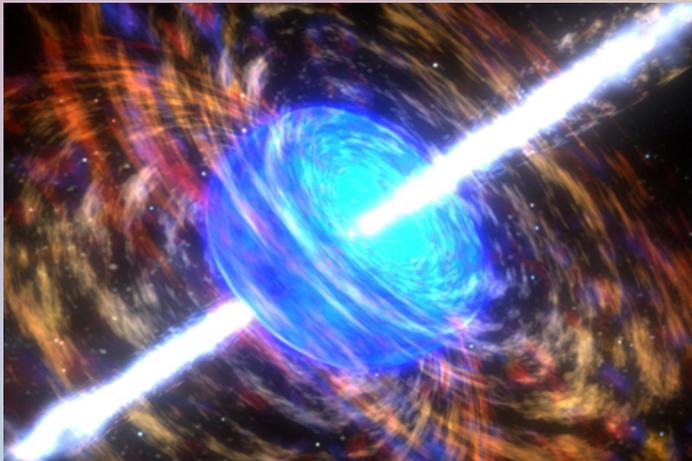
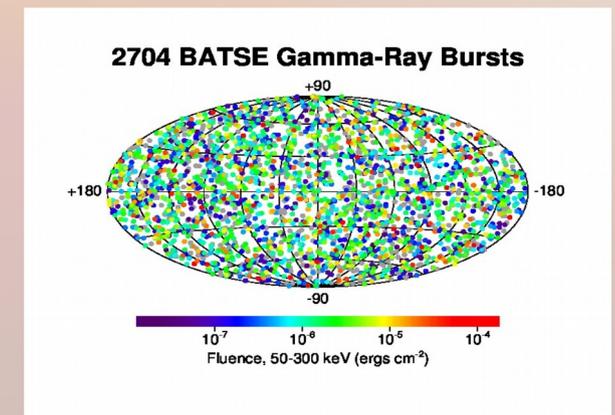
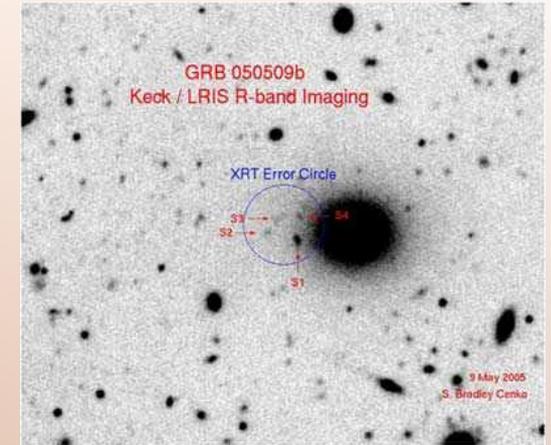
(Goriely, Bauswein, OJ, Pllumbi, Janka '15, MNRAS 452, 389)



- ...can be quite significant!!! Dynamical ejecta may also produce lighter elements (consistent with Wanajo, Sekiguchi+'14,'15)
- more neutrino-transport simulations needed!

Gamma-Ray Bursts

- source is moving highly relativistically
- natural suggestion: jet from rotating compact object
- long bursts ($T > 2\text{s}$): connection to death of massive stars
- short bursts ($T < 2\text{s}$) still mysterious, most likely from NS mergers

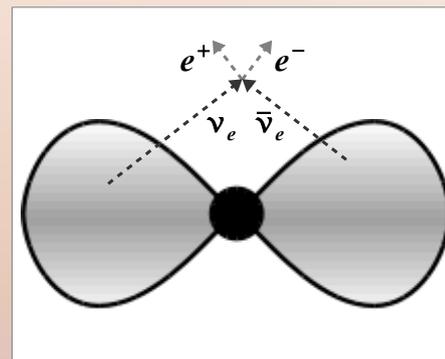


(NASA)

Popular central engine scenarios

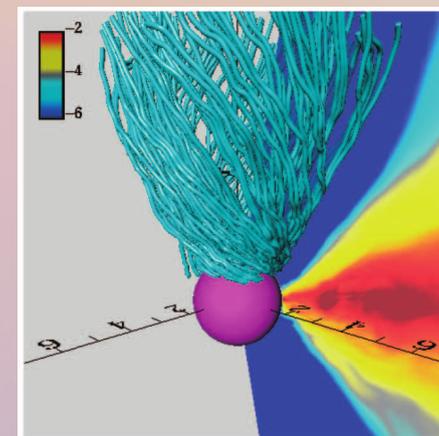
→ neutrino-pair annihilation

- neutrinos tap **gravitational energy of disk**
- e^+e^- pairs thermalize → thermal fireball
- **efficiency** of converting gravitational energy into jet energy?
- **baryon loading** in the funnel?



→ Blandford-Znajek process

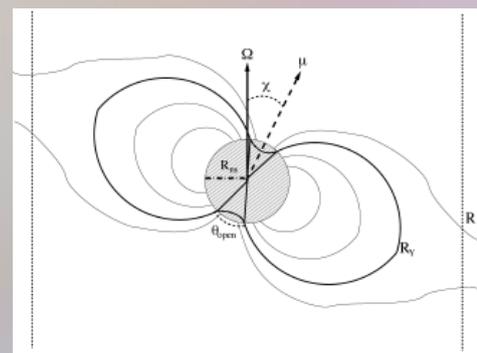
- B-field taps **rotation energy of central BH**
- Poynting-dominated jet
- efficient only for large-scale poloidal B-fields
- can **large-scale fields** be produced and sustained? MRI? Dynamo?



(Hirose+ '04)

→ magnetar spin-down emission

- B-field taps **rotation energy of central NS**
- Poynting dominated jet
- is **dipole model** appropriate?
- consistent with short burst timescale?



(Metzger+ '11)

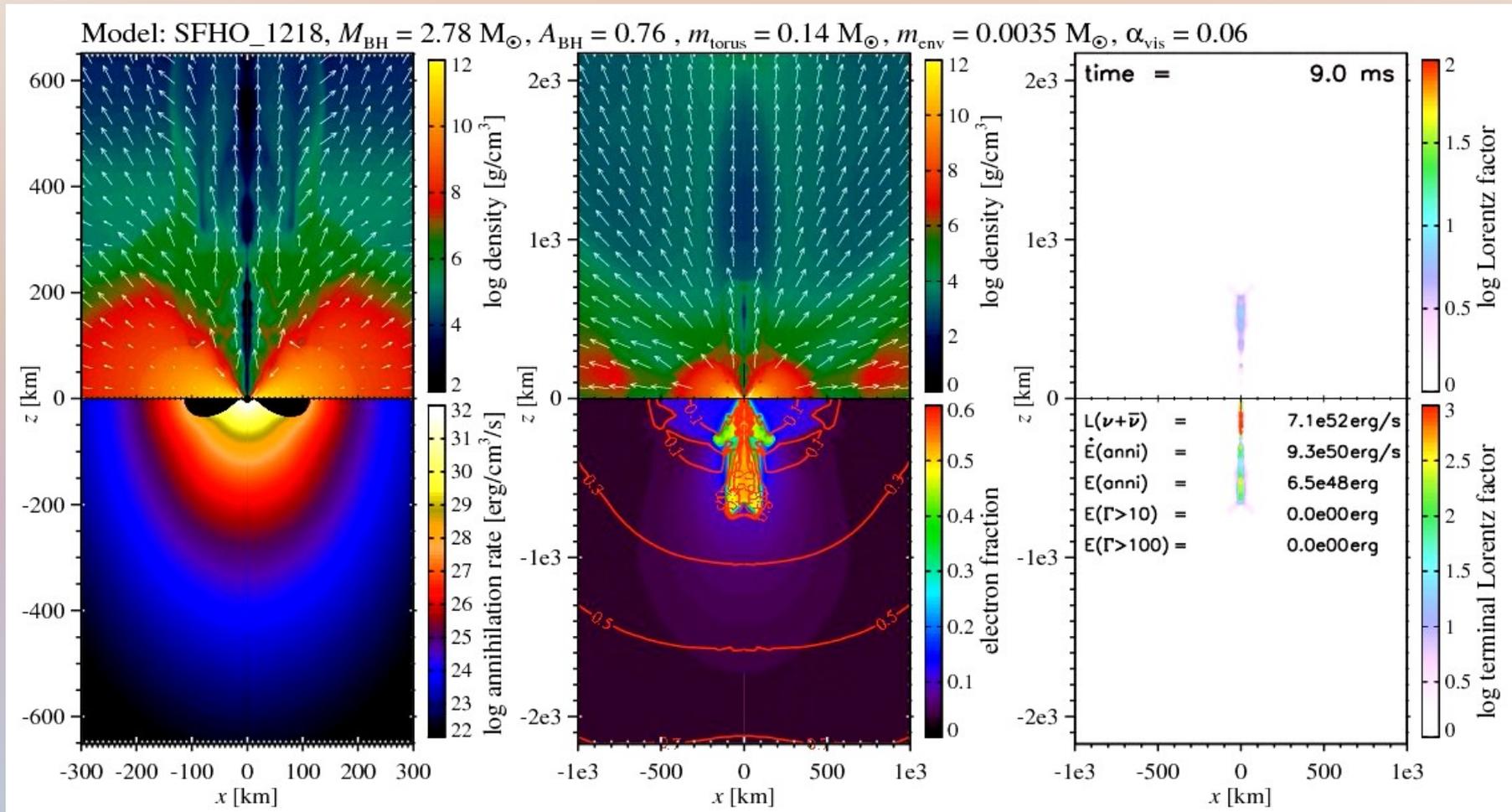
Is neutrino annihilation alone powerful enough to explain all sGRBs?

(Just+ '16, ApJL 816, 30)

- properties of **observed sGRB jets**:
energies $\sim 10^{48} - 10^{50}$ erg and Lorentz factors $\Gamma > 10 - 100$
- Are **necessary condition(s)** fulfilled?:
 - sufficient energy provided by nu-annihilation *and / or*
 - sufficiently small energy loss during expansion
- What is the impact of the **dynamical ejecta** on the jet?

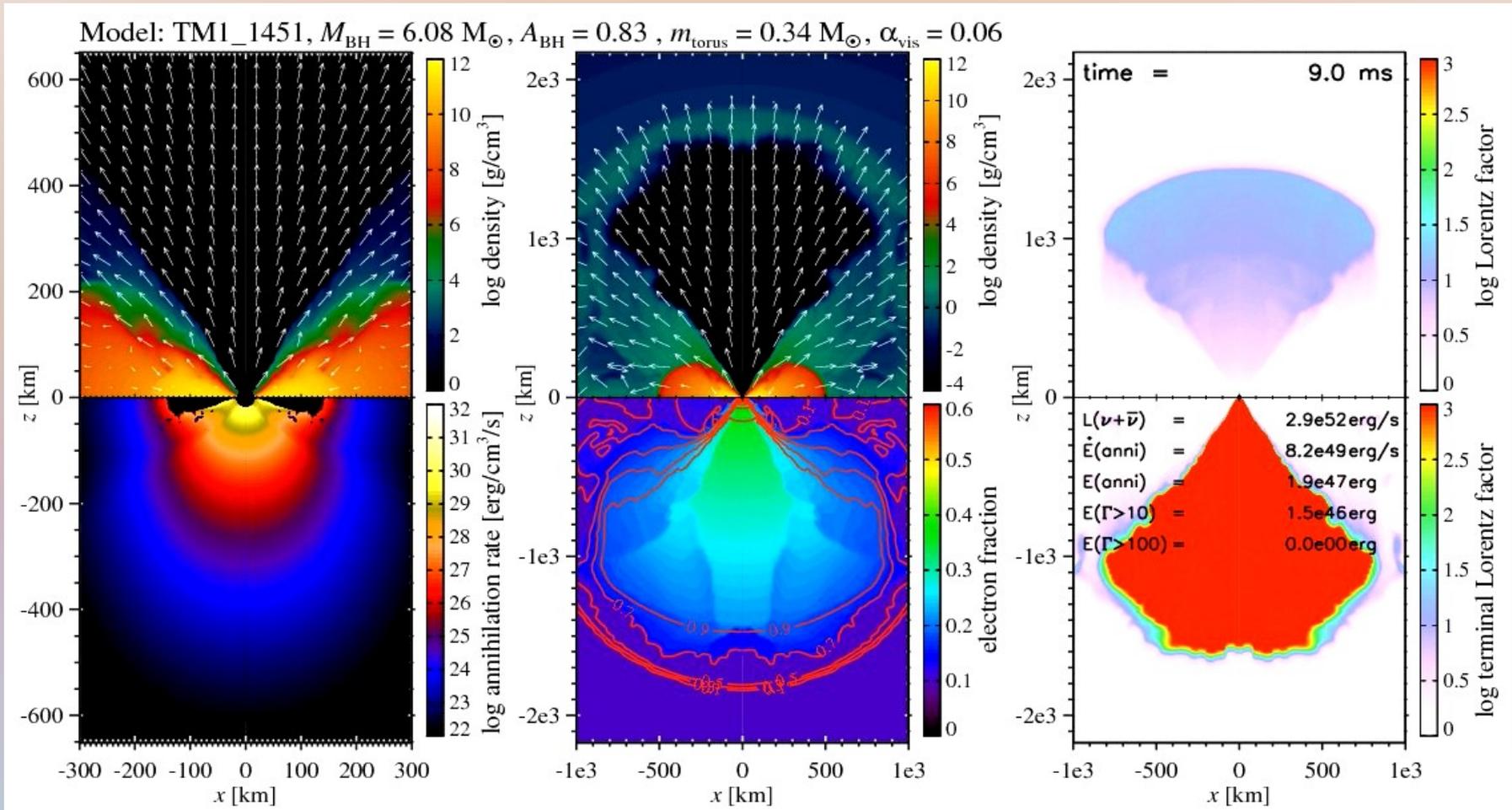
NS-NS Merger

- jet is **successfully launched**, but then **dissipates most of its kinetic energy** into cloud of dynamical ejecta



NS-BH Merger

- no dynamical ejecta in polar regions → jet can expand freely
- however, energy too low to explain majority of sGRBs



Summary

- several types of ejecta with **large spread in Y_e , s , v**
- each component with **characteristic abundance pattern** and Kilonova signal
- each component potentially relevant for **chemogalactic evolution**
- so far **only tidal ejecta** have robustly predictable nucleosynthesis
- **large uncertainties** for other components
- post-merger ejecta probably dominantly produce **intermediate mass elements**
- sGRB jet **probably not** launched by neutrino-annihilation alone
- *better modeling of neutrino transport combined with MHD, GR needed!*

Thank you for your attention!