# Cosmic rays, photons and neutrinos at VHE: messengers of the extreme Universe

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# Outline



- 1. Overview: the CR-γ-ν connection and motivations
  - a. Introduction: detection and experimental facilities...
  - b. The CR-γ-ν connection
  - c. Motivations: extreme astrophysics, fundamental physics, dark matter physics etc.

## 2. Beyond PeV: the extreme energy frontier

- a. The ultra-high energy cosmic ray enigma
- b. A prototype source: gamma-ray burst
- c. Connection with neutrinos and photons

## 3. <u>Astroparticle physics in the GeV – PeV range</u>

- a. Supernovae as cosmic ray sources
- b. Connection to the origin of Ice Cube neutrinos
- c. Transport of cosmic rays and the anti-matter 'anomalies'

1.a Introduction: photon diffuse backgrounds





# **1.a Introduction: air showers**



# 1.a Introduction: cosmic ray spectrum





a window on the non-thermal Universe

# 1.a Introduction: GeV – PeV cosmic rays



# 1.a Introduction: >EeV cosmic rays



1.a Introduction: Akeno Giant Air Shower Array (AGASA)





# **1.a Introduction: fluorescence observations**

High Resolution Fly's Eye now followed by the Telescope Array (2007-)

# 1.a Introduction: HESS II Cerenkov gamma-ray telescope





# 1.a Introduction: gamma-ray observations of supernovae



10<sup>15</sup>

E [eV]



# 1.a Introduction: neutrino diffuse backgrounds





# 1.a Introduction: the Ice Cube neutrino detector





## The Ice Cube experiment:

- located at the South Pole
- effective volume 1km<sup>3</sup>
- about 5000 photomultipliers between 1500m and 2500m deep



# 1.a Introduction: the Ice Cube neutrinos



# Ice Cube has reported in 2013 the first detection of « astrophysical » neutrinos...

... a few dozens of events at  ${\sim} \text{PeV}$  energies

... consistent with isotropic arrival directions...

and with flavor ratio 1:1:1 expected from oscillations...



# 1.b CR-ν-γ connection: energy loss processes

<u>High energy electrons (and e+)</u>:  $\gamma_e \equiv E_e/(m_e c^2)$ 

synchrotron radiation, with typical frequency:  $u_{\rm syn} \sim 70 \,{\rm Mhz} \, E_{\rm GeV}^2 B_{\mu \rm G}$ 

inverse Compton on seed photon:  $e + \gamma \rightarrow e + \gamma$ 

$$E'_{\gamma} \simeq 2\gamma_e^2 E_{\gamma} \sim 10 \,\mathrm{keV} \, E^2_{\mathrm{GeV}} E_{\gamma,\mathrm{CMB}}$$

#### High energy protons:

p-p interaction:  $p + p \rightarrow p + (p \text{ or } n) + \pi + ...$ p- $\gamma$  interaction:  $p + \gamma \rightarrow (p \text{ or } n) + \pi + ...$ with threshold:  $E_p E_{\gamma} \gtrsim m_{\pi}^2 c^4$ 

hence neutrino production:  $\pi^+$ 

$$\begin{array}{rccc} \pi^+ & \to & \mu^+ + \nu_\mu \\ \mu^+ & \to & e^+ + \nu_e + \overline{\nu}_\mu \end{array}$$

large loss rates imply that most radiation is of leptonic origin (in general!)

astrophysical v: unambiguous signature of hadron acceleration

and gamma production:  $\pi^0 \rightarrow \gamma + \gamma$ 

(+pair production  $p+\gamma \rightarrow p+e+e$  with small inelasticity, possibly synchrotron radiation in strong magnetic fields, etc.)

## 1.b CR- $\nu$ - $\gamma$ connection: energy output



Concomittant production of neutrinos and photons in hadronic interactions imply a relationship between cosmic ray, neutrinos and photons spectra:

- ightarrow assume a source of cosmic ray (protons) with spectrum  $\,\mathrm{d}N_p\,\propto\,E_p^{-s_p}\,\mathrm{d}E_p$
- $\rightarrow$  write  $f_{\pi}$  the fraction of proton energy lost to pion production (either in source or during transport)

pp or py produce neutral and charged pions with approximate ratio 1:1

 $\rightarrow$  energy carried by individual photons/neutrinos:  $\langle E_{\gamma} \rangle \sim 0.5 E_{\pi} \sim 0.1 E_p$  $\langle E_{\nu} \rangle \sim 0.25 E_{\pi} \sim 0.05 E_p$ 

 $\rightarrow$  corresponding energy spectra:

$$\begin{split} E_{\gamma}^{2} \frac{\mathrm{d}N_{\gamma}}{\mathrm{d}E_{\gamma}} &\simeq f_{\pi \to \gamma} E_{p}^{2} \frac{\mathrm{d}N_{p}}{\mathrm{d}E_{p}} \\ E_{\nu}^{2} \frac{\mathrm{d}N_{\nu}}{\mathrm{d}E_{\nu}} &\simeq f_{\pi \to \nu} E_{p}^{2} \frac{\mathrm{d}N_{p}}{\mathrm{d}E_{p}} \end{split} \quad \text{with} \quad f_{\pi \to \gamma} \sim 0.5 f_{\pi} \sim f_{\pi \to \nu} \end{split}$$

⇒ gamma-rays and neutrinos offer complementary probes of the inner engines of cosmic rays ...

# 1.b CR- $\nu$ - $\gamma$ connection: the Waxman-Bahcall limit



The correspondence between neutrino and cosmic ray spectra can be used to set a limit on the diffuse neutrino flux from the observed CR spectra at UHE:

 $\rightarrow$  assuming a power-law with index s<sub>p</sub> = 2.0 in the range 10<sup>19</sup>-10<sup>20</sup> eV, one infers from the observed cosmic-ray flux dI<sub>p</sub>/dE<sub>p</sub> an emissivity

$$\frac{\mathrm{d}I_p}{\mathrm{d}E_p} \sim \frac{c}{4\pi} \tau_p \frac{\mathrm{d}\dot{n}_p}{\mathrm{d}E_p} \Rightarrow E_p^2 \frac{\mathrm{d}\dot{n}_{\mathrm{CR}}}{\mathrm{d}E_p} \sim 10^{44} \,\mathrm{erg/Mpc^3/yr}$$
  
$$\tau_p \,\mathrm{lifetime \, of \, cosmic-rays} \,(\rightarrow \text{lecture 2})$$

 $\rightarrow$  extrapolating this flux to lower energies, in the PeV range, one infers a neutrino limit flux of

$$E_{\nu}^2 \frac{\mathrm{d}I_{\nu}}{\mathrm{d}E_{\nu}} \simeq \frac{c}{4\pi} f_z \frac{f_\pi}{4} t_\mathrm{H} E_p^2 \frac{\mathrm{d}\dot{n}_p}{\mathrm{d}E_p} \sim 2 \times 10^{-8} \,\mathrm{GeV/cm^2/s/sr}$$

 $t_{H}$ : Hubble time, accounts for the fact that neutrinos up to  $z \sim 1$  are observed,  $f_{z} \sim 1$ : accounts for cosmological redshift effects

# 1.b CR- $\nu$ - $\gamma$ connection: the WB limit vs Ice Cube



The Waxman-Bahcall bound sets the standard for the detection of high energy neutrinos, predicting O(few) events per yr in a km<sup>3</sup> cube neutrino detector such as Ice Cube at PeV energies...



... the observed flux at Ice Cube matches nicely the WB bound: a connection to high energy cosmic ray physics?

# 1.c Motivations: gamma-ray bursts





Gamma-ray bursts:

- a burst of <1 to >100sec at MeV energies (and higher), non-repeating, with erratic behavior
- cosmological origin: likely associated with rare (~1/Gpc<sup>3</sup>/yr) events at redshift z ~ 1 and above, originating from the collapse of massive stars of coalescence of compact objects

# 1.c Motivations: Lorentz invariance violation probes

The large distance of gamma-ray bursts and the small time dispersion of photons arrival times are used to place constraints on modified photon dispersion relations at high energy:

e.g. 
$$E^2 = p^2 c^2 \left[ 1 + \sum_{n=1}^{\infty} s_{n,\pm} \left( \frac{E}{E_{\text{QG}}} \right)^{2n} \right]$$



# 1.c Motivations: extreme astrophysics, pulsars



Some supernovae remnants harbor a pulsar, surrounded by a pulsar wind nebula...



... origin of the nebular radiation: electrons with energies up to a PeV! acceleration at a relativistic collisionless shock wave moving with Lorentz factor 10<sup>6</sup>?

... origin of the radio to GeV emission of the pulsar, including pulses and flares?

# 1.c Motivations: the unexpected flaring of the Crab



AGILE and Fermi have detected flares at GeV energies from the Crab nebula....



## **Blazar:**

radiation emitted by ultra-relativistic particles, accelerated in a jet launched from a super-massive black hole...

high energy radiation beamed towards the (on-axis) observer



## Main questions:

- How are jets launched and accelerated?
- How are particles accelerated in the jets, where, and what are the main radiation processes? Do blazar contribute to ultra-high energy cosmic rays?
- What causes the extreme time variability observed?



# 1.c Motivations: the extreme variability of blazars



Some blazars, e.g. PKS2155-304, are observed at very high energies with variability down to a few minutes!



and relativistic bulk motion (which enhances the variability by the Lorentz factor)

# 1.c Motivations: the origin of magnetic fields



Recent observations of blazars have placed interesting lower limits on the strength of an all-pervading magnetic field...

... an interesting connection to the long-standing problem of the origin of cosmic magnetic (Hoyle 58)

... some scenarios of magnetogenesis:



# 1.c Motivations: electromagnetic cascades



The limits on an all-pervading magnetic field rests on the development of electromagnetic cascades following the absorption of VHE photons by diffuse backgrounds (infra-red for TeV photons)

 $\rightarrow$  Pair production:  $\gamma + \gamma_b \rightarrow e^{\scriptscriptstyle +} + e^{\scriptscriptstyle -}$ 

$$\lambda_{\gamma} \simeq 80 \,\mathrm{Mpc} \left( E_{\gamma,0} / 10 \,\mathrm{TeV} \right)^{-1}$$

(interaction/cooling length)

electron energy:  $E_e \simeq E_{\gamma} / 2$ 

 $\rightarrow$  Inverse Compton: e +  $\gamma_b \rightarrow$  e +  $\gamma$ 

e cooling length through inverse Compton on the cosmic microwave background:

 $\lambda_e \simeq 0.07 \,\mathrm{Mpc} \left(E_e/5 \,\mathrm{TeV}\right)^{-1}$ energy of upscattered photons:

$$E_{\gamma} \simeq \frac{4}{3} k_{\rm B} T_{\rm cmb} \left( E_e / m_e c^2 \right)^2$$
  
$$\simeq 80 \, {\rm GeV} \left( E_{\gamma,0} / 10 \, {\rm TeV} \right)^2$$



# ⇒ overall: all >TeV energy of far away sources is redeposited in the 1-100 GeV range through an electromagnetic cascade...

(Mpc)

log<sub>10</sub> Distance

deflection of e<sup>+</sup>e<sup>-</sup> in intervening magnetic fields also leads to angular deflection (halos) and time delay (echos) of impulsive point-like energy injection (Aharonian et al. 94)

# 1.c Motivations: magnetic field signatures in blazars



intrinsic

The energy of photons emitted at >TeV energies should be seen in the multi-GeV range... ... upper limits on multi-GeV flux allows to put a lower limit on intergalactic B

 $\rightarrow$  reprocessed component is deflected by  $\delta\theta$  , spread over halo of extension  $\Theta$  :

if  $\delta\theta$  exceeds the beaming angle of the source  $\theta_c$ , or if  $\Theta$  exceeds the PSF extension, reprocessed flux is depleted...



 $<sup>\</sup>Rightarrow$  assuming homogeneous B, B > 10<sup>-15</sup>G

Note: constraint assumes steady source, if limited activity timescale  $T_s$ ,  $\delta t < T_s$  implies a weaker limit B > 10<sup>-17</sup>G... (Dermer et al. 11, Taylor et al. 11)

Neronov & Vovk (10), Tavecchio et al. (10)

# 1.c Motivations: dark matter in the Galactic center?



Sky map at very high energies with annihilating dark matter (w/o astrophysical foregrounds)



High energy photons, neutrinos and cosmic rays can be used to search for indirect signatures of dark matter annihilation, e.g. for photons:

$$\frac{\mathrm{d}I_{\gamma}}{\mathrm{d}E_{\gamma}} = \frac{\langle \sigma_{\chi\chi}v\rangle}{8\pi m_{\chi}^2} \frac{\mathrm{d}N_{\gamma}}{\mathrm{d}E_{\gamma}} \int_{\mathrm{l.o.s}} \mathrm{d}s \,\rho_{\chi}^2$$

# 1.c Motivations: dark matter in the Galactic center?





likely origin:

- Planck: synchrotron emission of accelerated electrons
- Fermi: inverse Compton on CMB of electrons...

# 1.c Motivations: dark matter in the Galactic center?





recently, a GeV excess has been detected in the Galactic center, with an (interesting) spherical morphology extending over 10deg... Hooper+ 08

... interestingly, the level of emission is compatible with standard halo models of dark matter, and with the canonical annihilation rate  $3x10^{-26}$  cm<sup>3</sup>/s...

 $\Omega_{\chi} \sim \frac{10^{-26} \,\mathrm{cm}^3/\mathrm{s}}{\langle \sigma_{\chi\chi} v \rangle}$ 

... however, extraction of astrophysical foregrounds in the Galactic center is complex!

	DB00-58	X-ray Thread	
— \$NR 0.9+0.1	Arches Cluster		
Sagittarius B2 1E 1743.	.1-2843 Sagittarius A	and the second	
a film and the second			1E 1740.7-2942
Quir	ntuplet	Sagittarius C	
Sagittarius B1 CI	üster		
DB00-6			
	Cold Gas Cloud		
	& Radio Arc	DB01-4	2
Chandra			





 $10^{22}$ 

ν [Hz]

# 1.c Motivations: dark matter signatures in cosmic rays?



Recent data of cosmic-ray anti-matter has been claimed to present an excess, which might constitute an indirect signature of dark matter annihilation...



# 1.c Motivations: dark matter signatures in cosmic rays?



In the anti-proton sector of cosmic rays, the most recent AMS data has also been claimed to be in excess of standard expectations: needs for an exotic source?



## Summary



- 1. <u>Overview: the CR-γ-ν connection and motivations</u>
  - a. Introduction: detection and experimental facilities...

... Cosmic rays are seen from GeV to  $10^{20}$  eV, photons up to 100TeV and now neutrinos with PeV energies: a new window on the extreme Universe

#### b. The CR-γ-ν connection

... Charged leptons generically produce high energy photons through inverse Compton (or synchrotron etc.) interactions ... Charged hadrons produce photons and neutrinos in similar proportions through

hadronic pp or pgamma interactions: opens a connection between photons, neutrinos and cosmic rays which can be used to probe the inner engines of cosmic ray sources

... application: the Waxman-Bahcall bound connects PeV neutrinos to UHECRs

c. Motivations: extreme astrophysics, fundamental physics, dark matter physics etc.

... VHE photons, neutrinos and cosmic rays are useful probes of extreme astrophysics (highly relativistic outflows, compact objects, powerful sources etc.), of fundamental physics (e.g. Lorentz invariance violation), and of dark matter annihilation

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# 2.a UHECR: the enigma of extreme energies



## <u>World record (1993, Fly's Eye)</u>: $E \sim 3 \pm 0.6 \ 10^{20} \text{ eV}$

 $\rightarrow$  center of mass interaction energy (with atmosphere):

$$\sqrt{s} = \sqrt{2m_pc^2(E+m_pc^2)} \simeq 800 \,\mathrm{TeV}$$

... requires extrapolation of known hadronic physics!

→ energy: 
$$E \simeq 48 \,\text{J}$$
 ... macroscopic!  
→ velocity:  $\gamma = \frac{E}{m_p c^2} \simeq 3 \times 10^{11} \Rightarrow v = c (1 - \epsilon)$ ,  $\epsilon \simeq 5 \, 10^{-24}$ 

Two main questions:

 $\rightarrow$  how (+ where) to accelerate particles to such extreme energies?

→ why do not we see (could we?) counterparts to arrival directions of the highest energy events?

10<sup>20</sup>eV giant air shower: a few km<sup>2</sup> on the ground!
# 2.a UHECR: the energy spectrum, with a cut-off





<u>Greizen-Zatsepin-Kuzmin cut-off</u>: the Universe becomes opaque to cosmic rays with  $E > 6 \ 10^{19} eV$  because of pion production on the cosmic microwave background:  $p + \gamma_{cmb} \rightarrow \pi + p/n$ 

⇒  $10^{20}$  eV protons (or Fe nuclei) come from within ~ 100 Mpc... ...  $10^{19}$  eV protons (Fe) come from within ~ 1000 Mpc!

note: the observed cut-off may correspond to the maximal energy at the source...

# 2.a UHECR: cosmological spectrum



For a continuous distribution of sources, the spectrum can be calculated as follows:



with:  $E_g = E_g(t; E, t_0)$  injection energy at t such that  $E_g=E$  at  $t=t_0$ : accounts for energy losses

$\frac{\mathrm{d}\dot{N}}{\mathrm{d}E_g} \sim \frac{L_{\mathrm{cr}}}{E_{\mathrm{min}}} \left(\frac{E}{E_{\mathrm{min}}}\right)$	$ \frac{1}{2} \int_{-\gamma}^{-\gamma} E_{min}: minimum E at source  \gamma (\geq 2): injection index  L_{cr}: CR luminosity (ergs/s) $
$\frac{\mathrm{d}I_{\mathrm{cr}}}{\mathrm{d}E} = \frac{c}{4\pi} \frac{\mathrm{d}n_{\mathrm{cr}}}{\mathrm{d}E}$	observed differential flux

Note: - L<sub>cr</sub> may depend on t

- n<sub>s</sub> generally depends on t, e.g. / SFR (star formation rate)
 ... higher by a few at z > 3-4...

### 2.a UHECR: Greisen – Zatsepin – Kuz'min cut-off



Greisen 66, Zatsepin & Kuzmin 66

 $\rightarrow$  threshold energy: p( $\gamma,\pi$ )N reaction is permitted when

$$E'_{\gamma} \geq m_{\pi}c^2\left(1+\frac{m_{\pi}}{2m_p}\right)$$
 in nucleus rest frame

or, in the cosmic rest frame (frame where CMB is isotropic):



# 2.a UHECR: Maximum distance scale



Consequence 1:> 10<sup>20</sup> eV particles must lie within 100 Mpc> 4 10<sup>19</sup> eV particles must lie within 1000 Mpc

<u>Consequence 2:</u> above 10<sup>19</sup> eV, the cosmic ray lifetime decreases rapidly with energy... recall:  $\frac{\mathrm{d}I_p}{\mathrm{d}E_p} \sim \frac{c}{4\pi} \tau_p \frac{\mathrm{d}\dot{n}_p}{\mathrm{d}E_p} \Rightarrow E_p^2 \frac{\mathrm{d}\dot{n}_{\mathrm{CR}}}{\mathrm{d}E_p} \sim 10^{44} \,\mathrm{erg/Mpc}^3/\mathrm{yr}$ 

# 2.a UHECR: main principles of acceleration



#### Kinematics:

 $\rightarrow$  Lorentz force:

$$rac{\mathrm{d} oldsymbol{p}}{\mathrm{d} t} \,=\, q \left(oldsymbol{E} + oldsymbol{v} imes oldsymbol{B}
ight)$$

 $\rightarrow$  in astrophysical plasmas: **E** = 0 (generic) in the rest frame of the plasma... (consequence of very large conductivity of plasmas)

 $\rightarrow$  in the lab reference frame, in which the plasmas moves with velocity  $v_p$ :

 $E_{\perp}\,=\,-v_{
m p} imes B$ 

Consequence 1: acceleration generically takes place in motional electric fields

Consequence 2: transport across B is necessary, which requires an external force

B uniform  $\Leftrightarrow$  particle confined in a cylinder of radius  $r_L$ 

Consequence 3: typical acceleration timescale 
$$t_{\rm acc} \simeq \left| \frac{1}{p} \frac{\mathrm{d} \boldsymbol{p}}{\mathrm{d} t} \right|^{-1} \simeq \frac{c}{v_{\rm p}} t_{\rm L}$$

... in astrophysical sources, acceleration competes against losses and escape: reaching VHE requires a fast moving plasma!

# 2.a UHECR: main principles of acceleration

Kinematics:

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$$E_{\perp}\,=\,-v_{\mathbf{p}} imes B$$

Examples: - turbulent Fermi acceleration

- Fermi acceleration at shock waves

- acceleration in sheared velocity fields

- magnetized rotators

- reconnection





### 2.a UHECR: acceleration to UHE



The typical acceleration timescale can be used to constrain the source of ultra-high energy cosmic rays, e.g. consider acceleration in an outflow moving towards the observer, with

• acceleration timescale (comoving frame):  $t_{
m acc} = \mathcal{A} t_{
m L}$ 

A >> 1, A ~ 1 at most: wind - for non-relativistic Fermi I,  $A \sim g/\beta_{sh}^2$  with g > 1 • time available for acceleration (comoving frame):  $t_{\rm dyn} \approx \frac{R}{\beta \Gamma c}$ • maximal energy:  $t_{\rm acc} \leq t_{\rm dyn} \Rightarrow E_{\rm obs} \leq \mathcal{A}^{-1} Z e B R / \beta$ • 'magnetic luminosity' of the source:  $L_B = 2\pi R^2 \Theta^2 \frac{B^2}{2} \Gamma^2 \beta c$ • lower bound on total luminosity:  $L_{\rm tot} \geq 0.65 \times 10^{45} \,\Theta^2 \Gamma^2 \mathcal{A}^2 \beta^3 Z^{-2} E_{20}^2 \,{\rm erg/s}$ 10<sup>45</sup> ergs/s is robust: for  $\beta \rightarrow 0$ ,  $\mathcal{A}^2 \beta^3 \geq 1/\beta \geq 1$ for  $\Theta\Gamma \rightarrow 0$ ,  $L_{\rm tot} \geq 1.2 \times 10^{45} \mathcal{A}\beta \frac{\kappa}{r_{\rm L}c} Z^{-2} E_{20}^2 \, {\rm erg/s}$  $L_{\rm tot} > 10^{45} Z^{-2} \, {\rm erg/s}$ Lower limit on luminosity of the source:

low luminosity AGN:  $L_{bol} < 10^{45}$  ergs/s high luminosity AGN:  $L_{bol} \sim 10^{46}$ - $10^{48}$  ergs/s gamma-ray bursts:  $L_{bol} \sim 10^{52}$  ergs/s

⇒ only most powerful AGN jets, gamma-ray bursts or very energetic

pulsars for UHE protons...

# 2.a UHECR: acceleration in giant radio-galaxies



#### Faranoff-Riley II radio-galaxy Cygnus A



# 2.a UHECR: arrival directions of Auger events





... arrival directions do not point to most powerful giant-radio galaxies...

... interpretation: UHECRs are accelerated in sources which camouflage in galaxies (e.g. gamma-ray bursts), or angular deflection from the source is significant, because UHECRs are heavy nuclei?

# 2.b A prototype model: gamma-ray bursts



#### Gamma-ray bursts: (in short)

- impulsive event, from <1sec to >1000sec
- compact (<10<sup>17</sup> cm), ultra-relativistic ( $\Gamma$  > 100)
- tremendous energy release: about 0.01M<sub>sun</sub> c<sup>2</sup>
- associated with the collapse of giant stars or with the coalescence of neutron stars
- typical distance: radius of observable Universe



# 2.b A prototype model: gamma-ray bursts





Gamma-ray burst:

output energy:  $E \sim 10^{46} \,\mathrm{J} \frac{\Delta \Omega}{4\pi} \sim 0.1 M_{\odot} c^2 \frac{\Delta \Omega}{4\pi}$ 

• rate:  $\dot{n} \sim 10^{-9} \, {
m Mpc}^{-3} \, {
m yr}^{-1}$  in principle enough to match the flux of UHECRs!

# 2.b A prototype model: gamma-ray bursts



#### Fermi acceleration in mildly relativistic internal shocks:

Waxman 95,01; Rachen & Meszaros 96

$$\rightarrow$$
 internal energy density:  $u' = \epsilon_e^{-1} \frac{L_{\gamma}}{4\pi r^2 \Gamma^2 c}$ 

$$\rightarrow$$
 magnetic field:  $B' = \sqrt{8\pi\epsilon_B u'}$ 

$$\rightarrow$$
 acceleration timescale:  $t_{\rm acc} \approx \mathcal{A} t_{\rm L}$ 

(assumes fast conversion of fraction  $\varepsilon_{e} \sim 0.1$  of u into gamma-rays)

( $\epsilon_{\rm B} \sim 0.1$  assumes build-up of B through instabilities... to match obs. flux)

(assumes ~ Bohm scaling in mildly relativistic shocks)

$$\rightarrow$$
 age constraint:  $t_{\rm acc} < \frac{r}{\Gamma c} \Rightarrow E_{\rm obs,20} \lesssim 7 \times \epsilon_{B,-1}^{1/2} \epsilon_{e,-1}^{-1/2} \Gamma_{2.5} L_{\gamma,52}^{1/2} \mathcal{A}^{-1}$ 

 $\rightarrow$  synchrotron losses:  $t_{\rm acc} < t_{\rm syn}(E) \Rightarrow E_{\rm obs,20} \lesssim r_{12}^{1/3} \Gamma_{2.5}^{2/3} \mathcal{A}^{-2/3}$ 

(when combined with former bound)

 $\Rightarrow$  acceleration to ~ 10<sup>20</sup> eV if dissipation takes place at radii > 10<sup>12</sup>cm... as expected from observations

photo-pion production: 
$$t_{\gamma\pi}^{-1} = \frac{c}{2\gamma_p^2} \int_{\epsilon_{\rm th}}^{+\infty} \mathrm{d}\epsilon \,\sigma_{\gamma\pi}\epsilon\xi_{\gamma\pi} \int_{\epsilon/2\gamma_p}^{+\infty} \mathrm{d}\epsilon_{\gamma} \,\epsilon_{\gamma}^{-2} \frac{\mathrm{d}n_{\gamma}}{\mathrm{d}\epsilon_{\gamma}}$$

$$\Rightarrow \frac{r/(\Gamma c)}{t_{\rm pi}} \sim \frac{L_{\gamma,52}}{\Gamma_{2.5}^4 \Delta t_{-2}}$$

 $\rightarrow$ 

suggests that protons can escape the flow, avoid adiabatic losses, through conversion to n in  $\gamma$ - $\pi$  reactions, leading to a neutrino signal with  $E_{\nu} \sim E_{UHECR}$ 

# 2.b A prototype model: connection to neutrinos!

#### Neutrino production:

Ice Cube

 $\rightarrow$  the energy lost in pion production through photo-hadronic interactions is close to unity for the Waxman 95 model

→ predicts a neutrino flux at energies >PeV, close to the Waxman-Bahcall limit:



E, [GeV] Ice Cube 11, He et al. 12

# 2.b A prototype model: acceleration scenarios





#### Summary:

d b

G

d

 $\rightarrow$  acceleration in internal shocks may lead to a neutrino signal at the Waxman-Bahcall limit, now probed by Ice Cube... detection of PeV neutrinos would imply acceleration of p to >10<sup>17</sup> eV... absence of detection would not rule out acceleration to UHE... as other acceleration scenarios can be envisaged...

 $\rightarrow$  radiative signatures of proton acceleration to ultra-high energies? (Asano et al. 09, 10, Razzaque et al. 10)

### $\rightarrow$ a 'difficulty' for GRB model is production rate:

flux of UHECR above  $10^{19}$  eV requires an energy input rate: ~  $10^{44}$  erg/Mpc<sup>3</sup>/yr with a GRB rate  $\dot{n}_{\text{GRB}}$  this requires:  $E_{\text{UHECR/GRB}} \approx 10^{53} \text{ erg} \left(\frac{\dot{n}_{\text{GRB}}}{1 \text{ Gpc}^{-3} \text{ vr}^{-1}}\right)^{-1}$ i.e.,  $E_{UHECR/GRB} / E_{\gamma/GRB} \sim 10 - ...$ ? (Eichler & Pohl 11, Waxman 11)

 $\rightarrow$  do not expect association of UHECRs with observed GRBs! Time delay imparted by extra-galactic magnetic fields:  $\delta t \sim 10^4 - 10^5$  yr ... (Waxman & Miralda-Escude 96)

- p

# 2.c v and $\gamma$ from UHECRs: possible anisotropies



Pierre Auger Collaboration, Science 2007

... Anisotropies are seen, but at 95% confidence level only...

... interpretation: a mixture of light and heavy nuclei? intergalactic magnetic fields?

... no obvious counterpart in photons...

# 2.c v and $\gamma$ from UHECRs: expected angular deflection



The expected angular deflection increases rapidly with decreasing energy, because of decreasing magnetic rigidity and increasing source distance scale...





# 2.c v and $\gamma$ from UHECRs: High energy gamma-rays



In optimistic cases, counterparts in very high energy gamma-rays could be seen through the electromagnetic cascades seeded in the environment of the source by UHECRs, e.g.

$$p + \gamma_{\rm CMB} \rightarrow p + e^- + e^+$$
  
 $e + \gamma_{\rm CMB} \rightarrow \text{ e.m. cascade}$ 

 $\rightarrow$  detection with upcoming instruments requires a huge CR luminosity above 10<sup>19</sup>eV:  $L_{cr} \sim 10^{46}$  erg/s for a distance 1Gpc...

note: halo  $\Leftrightarrow$  smoking gun signature of p acceleration to >10<sup>19</sup>eV





# 2.c v and $\gamma$ from UHECRs: GZK neutrinos spectra



Secondary neutrinos are expected at UHE as a consequence of pion production on the CMB:



(nearly all GZK neutrinos are produced in immediate vicinity of source)



# 2.c v and $\gamma$ from UHECRs: GZK neutrino predictions



... The diffuse GZK neutrino (and photon) backgrounds are within reach of next generation experiments





. Lev

... electromagnetic cascades injected at ultra-high energies are reprocessed down to <TeV due to the opacity of the diffuse backgrounds to photon pair production



Berezinsky et al. 10

### Summary



#### 2. Beyond PeV: the extreme energy frontier

a. The ultra-high energy cosmic ray enigma

...Two essential questions: how to accelerate particles to extreme energies? Why are such powerful sources not seen in the arrival directions of the highest energy cosmic rays?

...Acceleration requires extreme conditions: relativistic outflows etc., corresponding to a source luminosity in excess of 10<sup>45</sup> erg/sec... (for protons, less for charged nuclei)

#### a. A prototype source: gamma-ray burst

...gamma-ray bursts seemingly present the requisite conditions for accelerating particles to the observed energies and possibly enough energy in CR for matching the flux ...the 'standard' model implies a neutrino luminosity at a PeV at the Waxman-Bahcall limit, which is not seen by Ice Cube... a more definite answer within a few years... ...non detection would require another acceleration scenario or another source...

#### b. Connection with neutrinos and photons

...direct counterparts are not generally expected because cosmic rays are slowed down and deflected by magnetic fields... another connection to the distribution of magnetic fields on large scales...

...in some optimistic cases, gamma-ray counterparts could be seen at <TeV energies ...diffuse photon and neutrino backgrounds are expected from GZK interactions, with possible detection in the next decade...

## Outline



- 1. Overview: the CR-γ-ν connection and motivations
  - a. Introduction: detection and experimental facilities...
  - b. The CR-γ-ν connection
  - c. Motivations: extreme astrophysics, fundamental physics, dark matter physics etc.

#### 2. Beyond PeV: the extreme energy frontier

- a. The ultra-high energy cosmic ray enigma
- b. A prototype source: gamma-ray burst
- c. Connection with neutrinos and photons

#### 3. <u>Astroparticle physics in the GeV – PeV range</u>

- a. Supernovae as cosmic ray sources
- b. Connection to the origin of Ice Cube neutrinos
- c. Transport of cosmic rays and the anti-matter 'anomalies'

# 3.a CR sources: diffuse spectrum at all energies



## 3.a CR sources: supernovae remnants

Supernovae remnants are efficient particle accelerators at young age, with R  $\sim$  1 pc,  $v_{sh} \sim 0.01$  - 0.1 c

<u>Energetics</u>: if each supernova (Galactic rate  $r_{SN} \sim 2 / 100yr$ ) injects 10% of its shock kinetic energy ( $E_{SN} \sim 10^{51}$  ergs) in cosmic rays, one recovers the observed low energy cosmic ray flux, accounting for the Galactic volume (V $\sim 10^{67}$  cm<sup>3</sup>) and escape timescale ( $\tau_{esc} \sim 10^7$  yrs):  $u_{cr} \sim \frac{r_{SN} \tau_{esc} 0.1 E_{SN}}{V} \sim 1 \, eV \, cm^{-3}$ 



#### Spectrum:

- in the test particle approximation, predicted spectral index s=2.0 (an oversimplification!)
 ... roughly consistent with observed spectral index 2.7 once transport effects are included...?

#### Maximal energy:

- comparing the age of the remnant and the acceleration timescale, one finds

 $E_{max} \sim 10 \text{ TeV Z} (B/3 \mu G)$  (assuming Bohm scaling  $t_{scatt} \sim t_L$ )

- understanding the cosmic ray spectrum requires to accelerate protons at least up to 2000 TeV !

- likely solution: magnetic field in the shock vicinity is amplified by the back-reaction of cosmic rays on the turbulence... a non-trivial problem in plasma physics!









## 3.a CR sources: test particle Fermi acceleration





 $\rightarrow$  generates a power-law with spectral index s: balance between energy gain and escape

 $dN/dE \propto E^{-s}$ , with  $s = 1 + \ln(1-P_{esc})/\ln(1+g) = 2.0$  in strong shocks

→ maximal energy: balance between acceleration timescale and loss/escape timescales

age limit:  $E_{max} \sim 10$  TeV Z  $B_{-6} \beta_{sh,-1.5} R_{3pc}$  (generally applies to ions)

synchrotron limit:  $E_{max} \sim 600 \text{ TeV } B_{-6} \, {}^{-1/2} \beta_{sh,-1.5}$  (for electrons)

e.g. Bell 78, Axford et al. 77, Krimsky 77, Blandford & Ostriker 78, Drury 83, Blandford & Eichler 87

# 3.a CR sources: test particle Fermi acceleration







#### Microscopic derivation of diffusive shock acceleration:

 $\rightarrow$  energy gain per Fermi cycle: dowstream to upstream then back to downstream

$$p_{\rm down}^{(1)} = \gamma_{\rm up/down}^2 \left(1 + \cos\theta_{\rm up}\beta_{\rm up/down}/\beta\right) \left(1 - \cos\theta_{\rm down}\beta_{\rm up/down}/\beta\right) p_{\rm down}^{(0)}$$

note: a double Lorentz transform from downstream rest frame to upstream rest frame, where random (elastic) pitch angle deflection takes place, then back to downstream... <u>this is a trick:</u> it avoids solving for the trajectory of the particles in the motional electric fields... of course, acceleration is provided by the motional electric fields...

 $\rightarrow$  flux average:

$$\left\langle p_{\rm up}^{(1)}/p_{\rm up}^{(0)} \right\rangle = 2 \int_0^1 \mathrm{d}\cos\theta_{\rm up} \,\cos\theta_{\rm up} \, 2 \int_{-1}^0 \mathrm{d}\cos\theta_{\rm down} \,\cos\theta_{\rm down} \, p_{\rm up}^{(1)}/p_{\rm up}^{(0)}$$
$$\simeq \frac{4}{3} \frac{\beta_{\rm up/down}}{\beta} \simeq \frac{4}{3} \frac{u_1 - u_2}{v} \quad (u_1 - u_2 \ll c)$$

#### **References:**

A. R. Bell, Month. Not. Roy. Astron. Soc., **182**, 147 (1978) A. R. Bell, Month. Not. Roy. Astron. Soc., **182**, 443 (1978)

# 3.a CR sources: diffusive shock acceleration

 $\rightarrow$  escape probability: (escape occurs downstream by advection)

$$P_{\rm esc} = \frac{\text{flux} \to +\infty}{\text{flux through shock}} = \frac{nu_2}{2\int_0^1 d\cos\theta_{\rm up}\cos\theta_{\rm up}nv} = \frac{4u_2}{v}$$

 $\rightarrow$  spectrum of accelerated particles:  $N(>p) = \int_{p}^{+\infty} dp' 4\pi p'^{2} f_{2}(p')$ 

 $N(>p_n) \propto {
m probability}$  of going through n Fermi cycles leading to momentum  $p_n$ 

$$\rightarrow \text{after n Fermi cycles,} \quad \langle p^{(n)} \rangle = p^{(0)} \prod_{i=1}^{n} \left( 1 + \frac{4}{3} \frac{u_2 - u_1}{v_i} \right)$$
  
i.e.  $\ln \left( \langle p^{(n)} \rangle / p^{(0)} \right) \simeq \sum_{\substack{i=1 \ n}}^{n} \frac{4}{3} \frac{u_1 - u_2}{v_i}$ 

→ probability of reaching n-cycles:  $P_n = \prod_{i=1}^{n} (1 - P_{\text{esc},i})$ i.e.  $\ln P_n \simeq \frac{-3u_2}{(u_2 - u_1)} \ln \left( \langle p^{(n)} \rangle / p^{(0)} \right)$ 

leads to: 
$$N(>p) \propto p^{-3/(r-1)} \Leftrightarrow \mathrm{d}N/\mathrm{d}p \propto p^{-(r+2)/(r-1)}$$

 $r \equiv \frac{u_2}{u_1}$ 

# 3.a CR sources: main questions beyond test particle



The above test particle approximation fails to explain the origin of cosmic rays with energy up to a PeV... Modern developments focus on the back-reaction of particles on the turbulence and how particles are accelerated in the self-amplified field...



# 3.a CR sources: connection to laboratory astrophysics





particle-in-cell simulation of a collisionless shock wave (Spitkovsky 05)

#### Collisionless shock wave:

- mean free path to Coulomb collisions >> thickness of the shock
- structure of the shock: a famous problem in plasma physics!
- effective coupling between particles: electromagnetic collective modes of the plasma...



## 3.a CR sources: searching for the source

-10

-20

RA (hours)

TeV image of SN1006 (HESS)

15h02m

(fed) Dec (deg) Dec 41.6

-41.8

-42

-42.2

-42.4

PSF

15h04m



energies in supernovae remnants...

however, the brightest TeV emitters so far rather indicate that the photons are emitted by accelerated electrons... where is the hadronic contribution?

### 3.a CR sources: searching for the source



#### hadronic signature of the emission:

• pion production cross-section ~ independent of energy above threshold (~m<sub>p</sub>), hence  $\frac{\mathrm{d}I_{\gamma}}{\mathrm{d}E_{\gamma}} \propto \frac{\mathrm{d}N_p}{\mathrm{d}E_p} \propto E_{\gamma}^{-s_\mathrm{p}} \quad (E_{\gamma} \gg 0.1\,\mathrm{GeV}) \qquad \text{with s}_{\mathrm{p}} \sim 2.0-2.4$ 

in apparent conflict with the data (Fermi data at GeV energies is essential here!)

### 3.a CR sources: searching for the source



So far, two convincing cases of hadronic signature in shell-type supernovae:

Cassiopae-A (SN1680):

emission compatible with ~1% of kinetic energy going into hadrons, for B ~ 120  $\mu$ G (amplification!)

Tycho (SN1572):

emission compatible with ~10% of kinetic energy going into hadrons, for B ~ 200  $\mu$ G (amplification!), so  $E_p$  possibly as large as 500TeV!



# 3.a CR sources: evidence for hadronic acceleration



Most robust evidence so far for hadronic production: pion production bump in W44 and IC443

$$p + p \rightarrow \pi^0 + \dots, \quad \pi^0 \rightarrow \gamma + \gamma$$



#### Finding a PeVatron:

... difficult, because galactic supernovae are rare (2-3 /century), invisible in optical because of absorption, and acceleration to PeV should take place during the first thousand years... a prime objective for future Cerenkov arrays!
### 3.b Ice Cube v: a connection to CR physics?



The Waxman-Bahcall bound sets the standard for the detection of high energy neutrinos, predicting O(10) events per yr in a km3 cube neutrino detector such as Ice Cube at PeV energies...



... the observed flux at Ice Cube matches nicely the WB bound: a connection to high energy cosmic ray physics?

# 3.b Ice Cube v: cosmic ray production in starbursts?



 $\rightarrow$  In starburst galaxies, the high density of the interstellar medium leads to a short pion interaction timescale:

$$t_{\pi} = \frac{1}{n_p \langle \sigma_{pp} c \rangle} \sim 2 \times 10^5 \,\mathrm{yr}$$

 $\rightarrow$  whereas the diffusive/convective escape timescale is t<sub>esc</sub>  $\sim 10^6$  yrs at 1 GeV, implying that most of CR energy is lost to pion production...

 $\rightarrow$  Fermi observations of starburst galaxies find  $L_{\gamma} \sim 10^{40}$  erg/s for an estimated  $L_{cr} \sim 10^{41}$  erg/s, i.e. a conversion factor of 10%

 $\rightarrow$  gamma spectrum must trace the injection spectrum of cosmic rays if t<sub> $\pi$ </sub> sets the life-time, and Fermi observations indicate dN/dE  $\propto$  E<sup>-2.2</sup>, OK with cosmic-ray phenomenology...

 $\rightarrow$  if gamma-rays are indeed mostly hadronic, then:

 $\rightarrow$  starburst galaxies provide a substantial contribution to the diffuse gamma-ray background

 $\rightarrow$  starburst galaxies output a neutrino flux

 $E_{\nu}\phi_{\nu} \sim 2 \times 10^{-8} \left(E_{\nu}/100 \,\text{TeV}\right)^{-2.2} \,\text{GeV}\,\text{cm}^{-2}\,\text{s}^{-1}\,\text{sr}^{-1}$  very close to the Ice Cube measurements...

 $\rightarrow$  however, maximum energy of neutrinos imply acceleration and confinement up to 10^{17} eV?





... the observed flux at Ice Cube matches the prediction of the starburst model, but the maximal observed neutrino energy exceeds the prediction...?

3.b Ice Cube v: connection to the gamma-ray background





... a significant part of the diffuse gamma-ray background observed by Fermi must result from pion production in starburst galaxies, if these galaxies explain the signal observed by Ice cube... with already some tension...

### 3.c CR transport: general considerations



#### Transport of cosmic rays in galactic magnetic fields:



#### **References:**

V. S. Berezinsky, S. V. Bulanov, V. A. Dogiel, V. L Ginzburg, V. S. Ptsukin, « Astrophysics of cosmic rays » (1990, North-Holland)

### 3.c CR transport: a generic simplified model

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<u>A simple model</u>: particles are injected in sources, propagate in interstellar medium, then escape the galaxy after some (momentum dependent) time  $\tau_{\rm esc}$ 

ightarrow assume stationary + homogeneous source distribution:  $N_{SN} au_{
m esc} \gg 1$ 

 $\rightarrow$  consider stable species, neglect collisions and decays

$$\rightarrow$$
 average distribution:  $\frac{\mathrm{d}n}{\mathrm{d}p} \simeq \frac{\mathcal{N}}{p_0} \left(\frac{p}{p_0}\right)^{-s} \dot{n}_{\mathrm{s}} \tau_{\mathrm{esc}}$ 

 $\rightarrow$  connection to diffusion model: neglect all terms but diffusion and source, integrate over a finite volume... escaping flux balances injection,

$$au_{
m esc} \sim rac{H^2}{D}$$
 H scale height for escape D diffusion coefficient

ightarrow momentum dependence:  $\tau_{
m esc} \propto p^{-lpha} \Rightarrow {
m d}N/{
m d}p \propto p^{-s-lpha}$ 

theory - acceleration:  $s\simeq 2$  measurements:  $s+\alpha\simeq 2.7$ 

 $\Rightarrow$  suggests  $\alpha \simeq 0.7$ 

### 3.c CR transport: a generic simplified model



<u>A simple model</u>: particles are injected in sources, propagate in interstellar medium, then escape the galaxy after some (momentum dependent) time  $\tau_{\rm esc}$ 

 $\rightarrow$  momentum dependence:  $\tau_{\rm esc} \propto p^{-\alpha} \Rightarrow {\rm d}N/{\rm d}p \propto p^{-s-\alpha}$ 

theory - acceleration:  $s\simeq 2$  measurements:  $s+\alpha\simeq 2.7$ 

 $\Rightarrow$  suggests  $\alpha \simeq 0.7$ 

 $\rightarrow$  note: isotropic diffusion in Kolmogorov turbulence leads to  $D\propto p^{0.3}...$  which does not match the inferred scaling for  $\tau_{\rm esc}...$ 

... solutions: diffusion in interstellar space is ill-understood... or does acceleration lead to a slope  $s\sim 2.3-2.4$  ?

 $\rightarrow$  check of the model: secondary species are not injected in acceleration sites but produced through inelastic collisions of primary species, e.g. Li, Be, B from interactions of cosmic ray C, N, O nuclei with interstellar p and He

### 3.c CR transport: model validation

observed dependence:

 $\mathrm{d}N_{\mathrm{B}}/\mathrm{d}N_{\mathrm{C}} \propto p^{-0.6}$ 

indeed confirms simple model...



# 3.c CR transport: connection to anti-matter anomalies...



Simplified transport model neglects energy losses and decays/spallation of the secondary species... however it can be used to set an upper bound on the e+ secondary spectrum, the actual spectrum being suppressed by a factor  $f_{loss}$  characterizing energy losses (e.g. inverse Compton)



This analysis thus suggests that the excess exists with respect to transport models pre-AMS, not in itself, i.e. positrons could be pure secondary (no DM origin!) provided  $f_{loss}$  has the right dependence on energy

### 3.c CR transport: possible pulsar positron contribution



in the case of positrons, it had been predicted that nearby recent pulsars (e.g. Geminga) could possibly provide a primary source of positrons, leading an increasing positron fraction at high energies (pulsar wind:  $e^+ \sim e^-$ )... as observed!







current data cannot distinguish between a pulsar contribution, a contribution of leptophilic dark matter and a pure secondary origin....

### 3.c CR transport: anti-proton data



... the latest anti-proton data is also consistent with a pure secondary origin...

#### Summary



#### 3. <u>Astroparticle physics in the GeV – PeV range</u>

#### a. Supernovae as cosmic ray sources

... Supernovae remnants appear as the best candidate for the source of cosmic rays below 10<sup>17</sup>eV: energy output in principle sufficient...

... crucial questions: can supernovae accelerate particles to Z x a few PeV?

... a major question for theorists and a prime objective for future observatories: find a PeVatron

... current data indicate that the bulk of the VHE emission in some SNR is dominated by leptons, while in others by hadrons...

#### b. Connection to the origin of Ice Cube neutrinos

... starbursts galaxies behave as calorimeters for protons energies below a few PeV: all the energy going into hadrons is transferred to gamma-rays and neutrinos through pion production...

... a direct connection to the supernova origin of cosmic rays, which nicely explains the origin of Ice Cube neutrinos...

... or a subset thereof? in particular, where is the expected turn-over at <PeV energies?

#### c. Transport of cosmic rays and the anti-matter 'anomalies'

... transport of cosmic rays in the Galaxy involves many ill-known processes (diffusion, advection, re-acceleration etc.)...

... so far, positron and anti-proton spectra are consistent with a secondary origin, with possibly a nearby pulsar contribution to the positron spectra: no need for exotic dark matter!

### Some experimental perspectives



<u>Neutrinos</u>: a 10x Ice Cube is being projected, also KM3NET ( $\sim$ a few times Ice Cube) in the North

<u>UHECR:</u> situation not clear, developments projected at Auger, Telescope Array, possible deployment of the giant EUSO on the space station, with a clear need for better statistics and accurate chemical composition measurement...