

# Magneto-genesis from Higgs Inflation

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# Table of contents

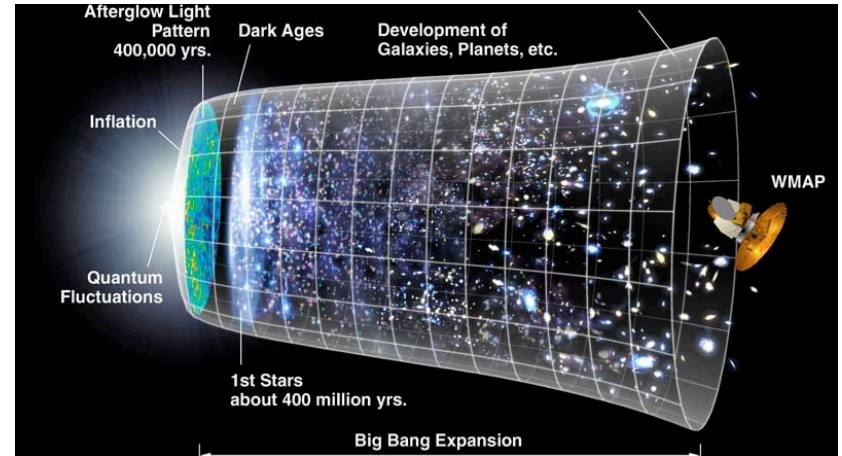
1. Introduction
2. Cosmology in Einstein-Electroweak theory
3. Dynamics of gauge fields
4. Generation of primordial magnetic fields
5. Conclusion

# Table of contents

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2. Cosmology in Einstein-Electroweak theory
3. Dynamics of gauge fields
4. Generation of primordial magnetic fields
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# The inflationary paradigm

The best idea explaining the early universe!



- There have been various models induced by high energy theories.

SUSY, super string, extra dimension...

However...

- There exists no signal of such theories in experiments.
  - some physicists want to explain the inflation without using them.
  - Higgs Inflation might meet their expectation!

# What is Higgs Inflation?

- Action is as follows [1] :

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} R + \xi \Phi^\dagger \Phi R - (\partial_\mu \Phi)^\dagger (\partial^\mu \Phi) - \lambda \left( \Phi^\dagger \Phi - \frac{v_0^2}{2} \right)^2 \right]$$

appears in the curved space-time

$\Phi$  : Higgs field (inflaton)

- **It explains the inflation in the framework of Standard Model!**
- for successful inflation, required conditions are:

$$n_s \simeq 0.97, \quad r \simeq 0.0033, \quad \xi \simeq 2 \times 10^4$$

—————> Planck's data strongly support! ~~————— against the BICEP2 data~~

# If there exist gauge fields in early universe...

- It is worth investigating a possible role of gauge fields in inflationary universe.



## Several models with gauge fields

- ❖ Anisotropic Inflation
- ❖ Gauge-flation
- ❖ Chromo-natural Inflation etc.

- **What happens if electroweak gauge fields are included in Higgs Inflation model?**

As a result... ↓

- Higgs Inflation can make sizable primordial magnetic fields.

# Motivation of generating primordial magnetic fields

- Magnetic fields in clusters of galaxies have been observed.



Very little astrophysical substances

They might be generated from inflation!

- M. S. Turner and L. M. Widrow, Phys. Rev. D **37**, 2743 (1988).
  - B. Ratra, Astrophys. J. **391**, L1 (1992).
  - K. Bamba and J. Yokoyama, Phys. Rev. D **69**, 043507 (2004) [astro-ph/0310824].
- Etc.

Observational constraint [2]:  $10^{-15}[\text{G}] \lesssim B_{obs} \lesssim 10^{-9}[\text{G}]$

Higgs Inflation might satisfy this constraint!

[2]A. Neronov and I. Vovk, Science 328, 73 (2010) [arXiv:1006.3504 [astro-ph.HE]].

D. G. Yamazaki, T. Kajino, G. J. Mathew and K. Ichiki, Phys. Rept. **517**, 141 (2012) [arXiv:1204.3669 [astro-ph.CO]].

Etc.

# Table of contents

1. Introduction
2. Cosmology in Einstein-Electroweak theory
3. Dynamics of gauge fields
4. Generation of primordial magnetic fields
5. Conclusion



# Higgs Inflation with gauge fields

- Action is as follows:

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} R + \xi \Phi^\dagger \Phi R - \frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} - \frac{1}{4} G_{\mu\nu}^a G^{a\mu\nu} - (D_\mu \Phi)^\dagger (D^\mu \Phi) - \lambda \left( \Phi^\dagger \Phi - \frac{v_0^2}{2} \right)^2 \right]$$

$SU(2)_L \times U(1)_Y$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c, \quad D_\mu \Phi = \left( \partial_\mu - i \frac{g}{2} \tau^a A_\mu^a - i \frac{g'}{2} B_\mu \right) \Phi$$

$$G_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu .$$

- We assume a homogeneous and closed universe:

$$ds^2 = -dt^2 + a(t)^2 d\Omega_3^2$$

$$= -dt^2 + a(t)^2 ((\sigma^1)^2 + (\sigma^2)^2 + (\sigma^3)^2) .$$

# Configurations of gauge fields

- If only  $SU(2)_L$  gauge field exists...

$$\left( \begin{array}{l} ds^2 = -dt^2 + a(t)^2 d\Omega_3^2 \\ \quad = -dt^2 + a(t)^2 ((\sigma^1)^2 + (\sigma^2)^2 + (\sigma^3)^2) . \\ \mathbf{A} = \frac{1}{2g} f(t) (\sigma^1 \tau^1 + \sigma^2 \tau^2 + \sigma^3 \tau^3) , \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v(t) \end{pmatrix} . \end{array} \right.$$

$S^3$



isotropic

- Considering  $U(1)_Y$  gauge field...

$$\mathbf{B} = h(t) \sigma^3$$

$$\left( \begin{array}{l} ds^2 = -dt^2 + a_1(t)^2 ((\sigma^1)^2 + (\sigma^2)^2) + a_3(t)^2 (\sigma^3)^2, \\ \quad \text{where } a_1(t) \equiv e^{\alpha(t)+\beta(t)} , a_3(t) \equiv e^{\alpha(t)-2\beta(t)} . \\ \mathbf{A} = \frac{1}{2g} [f_1(t) (\sigma^1 \tau^1 + \sigma^2 \tau^2) + f_3(t) \sigma^3 \tau^3] , \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v(t) \end{pmatrix} . \end{array} \right.$$



anisotropic!

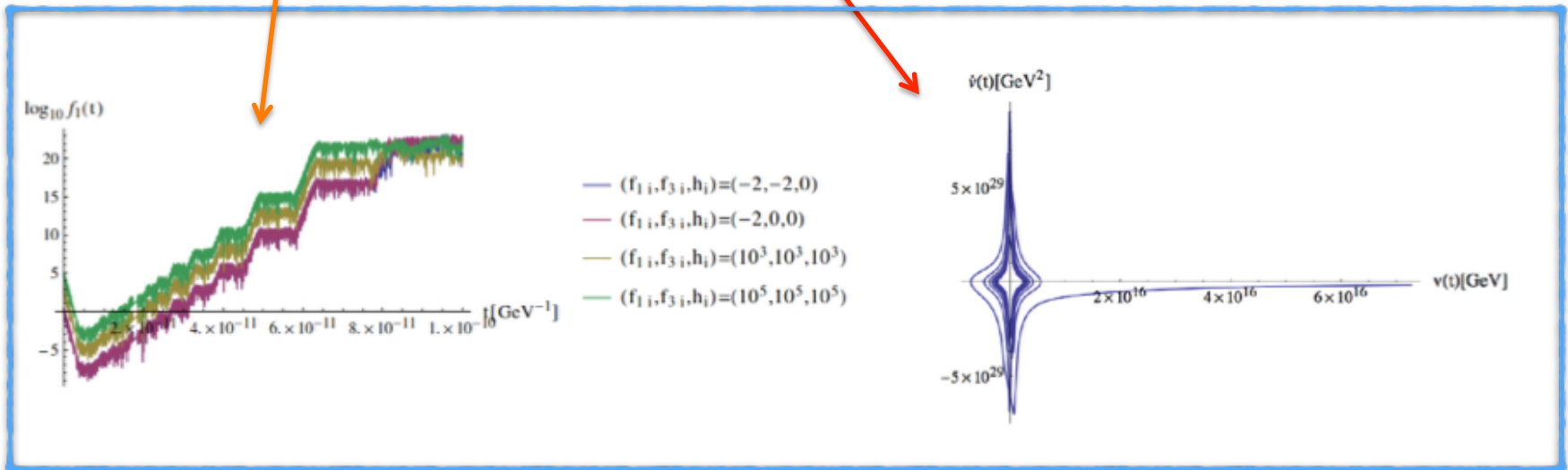
# Table of contents

1. Introduction
2. Cosmology in Einstein-Electroweak theory
3. Dynamics of gauge fields
4. Generation of primordial magnetic fields
5. Conclusion

# Solving calculation...

- Solving six parameters:

$$(f_1(t), f_3(t), h(t), v(t), \alpha(t), \beta(t))$$

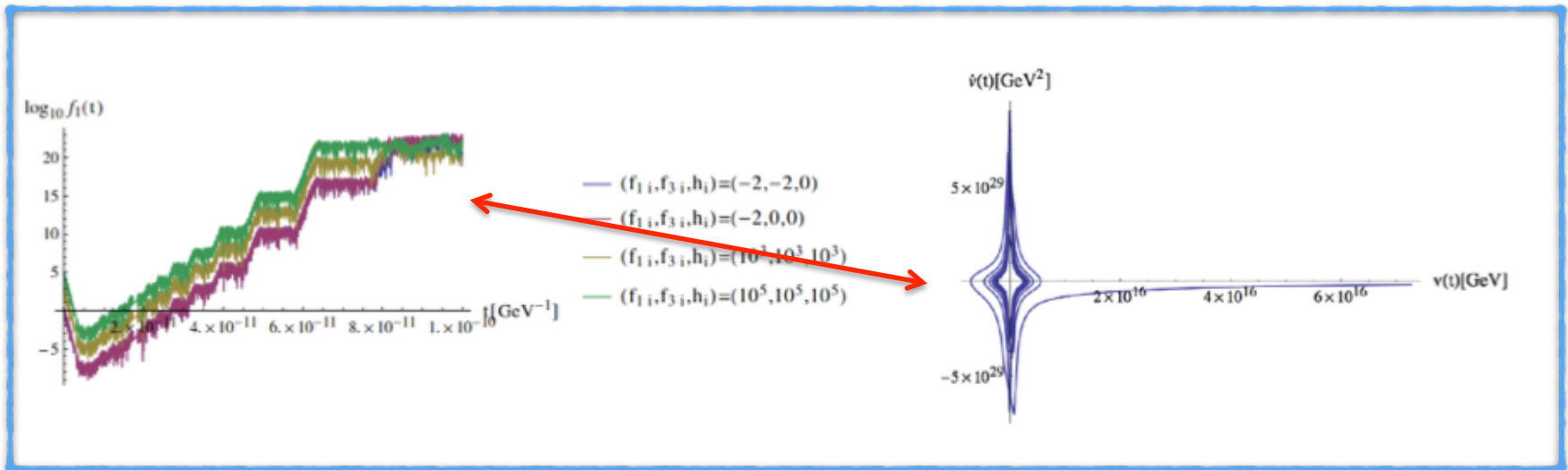


(Other two gauge components evolve like this one.)

# Solving calculation...

- After inflation, the equation for gauge fields approximate the following ones:

$$\ddot{f}_1 + \frac{g^2}{4} v(t)^2 f_1 \simeq 0, \quad (\ddot{f}_3 - g' \ddot{h}) + \frac{g^2 + g'^2}{4} v(t)^2 (f_3 - g' h) \simeq 0 .$$

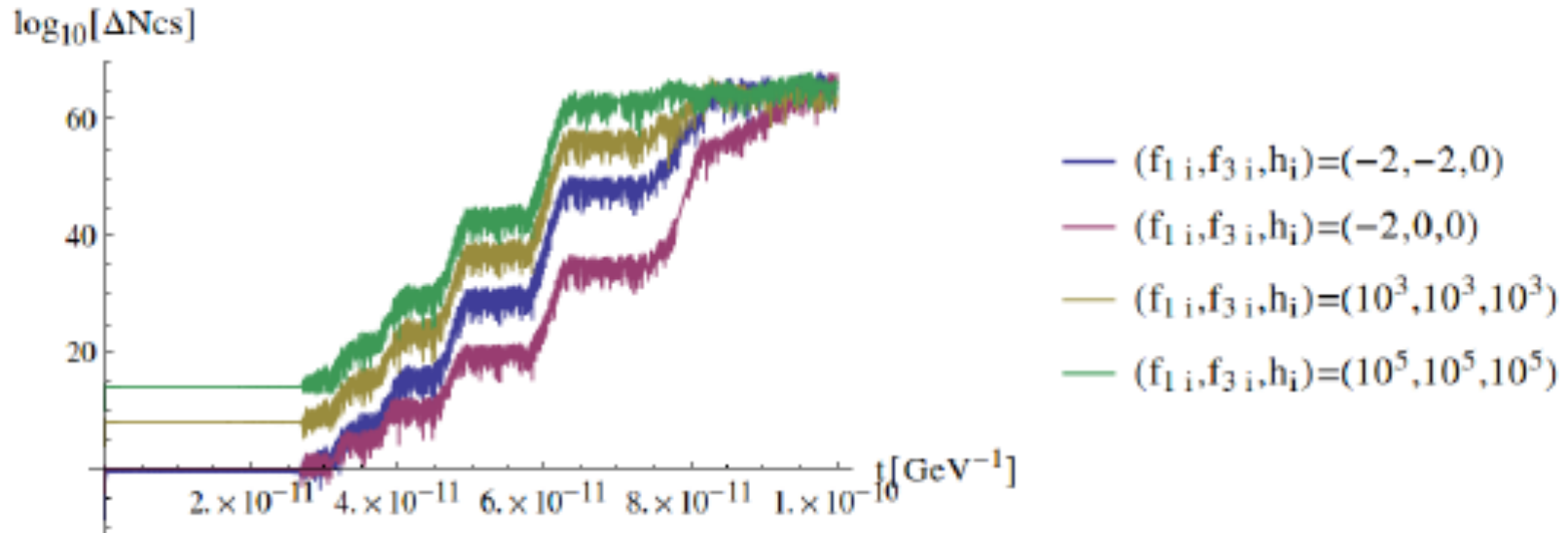


resonance happens!

# Topological Configuration

- The change of Chern-Simons number is

$$\begin{aligned}\Delta N_{CS} &\equiv \frac{g^2}{16\pi^2} \int_0^t dt \sqrt{-g} \text{Tr}(F_{\mu\nu} \tilde{F}^{\mu\nu}) \\ &= -\frac{1}{8} \int_0^t \frac{d}{dt} [f_1^2 f_3 + 2f_1^2 + f_3^2] dt.\end{aligned}$$

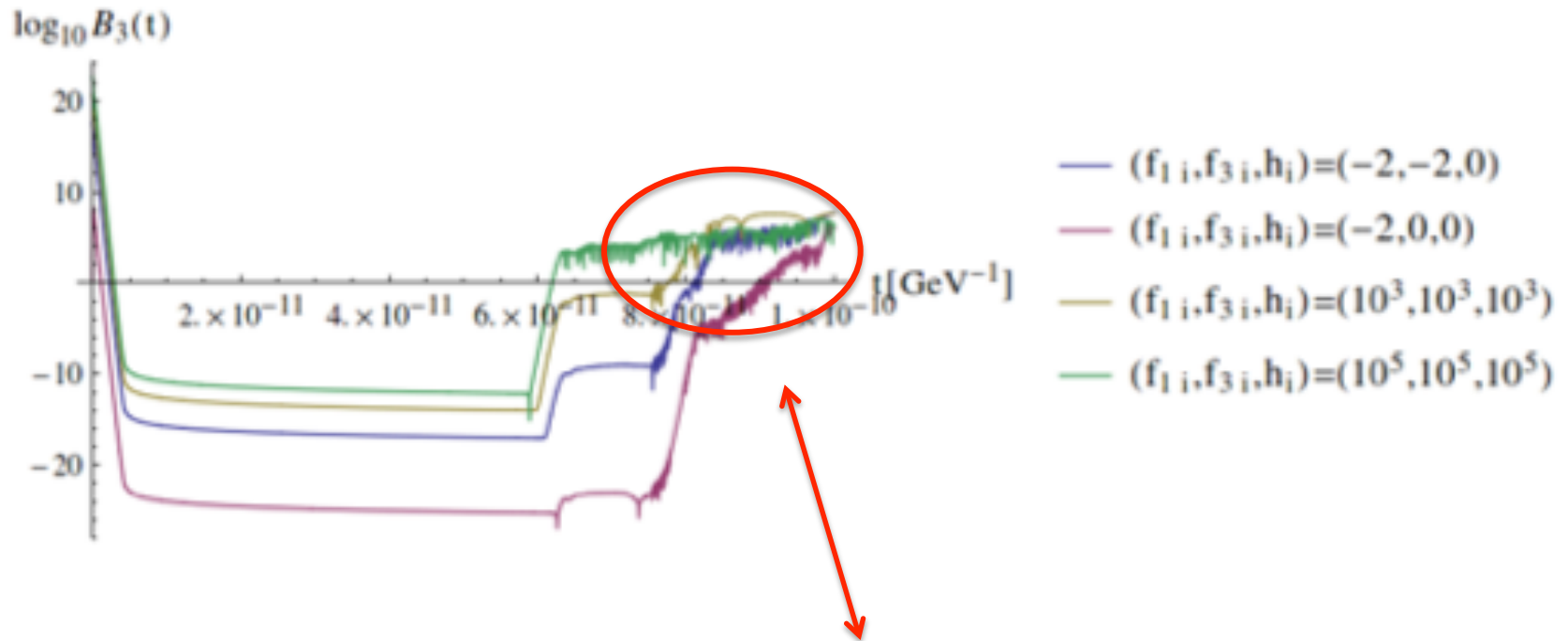


# Table of contents

1. Introduction
2. Cosmology in Einstein-Electroweak theory
3. Dynamics of gauge fields
4. Generation of primordial magnetic fields
5. Conclusion

# Generation of magnetic fields

- The magnetic field is given by:  $B \sim \frac{1}{a_1^2}(f_3 + h)$



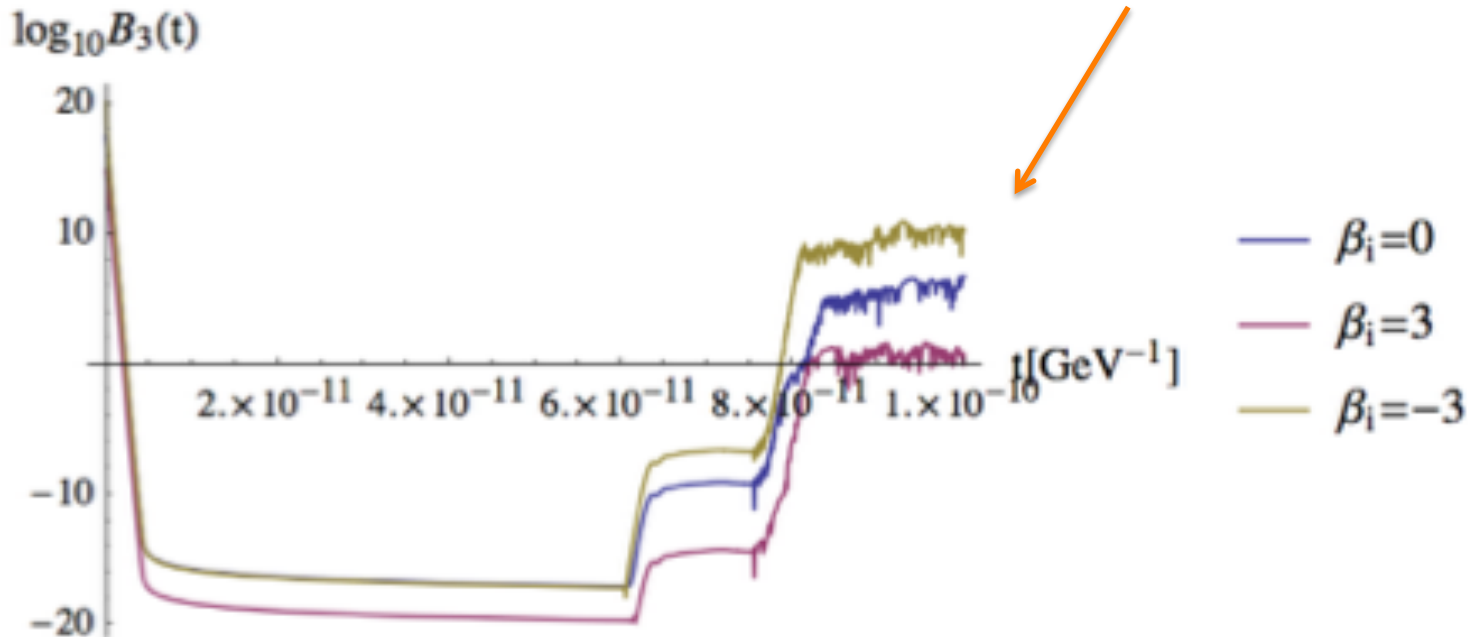
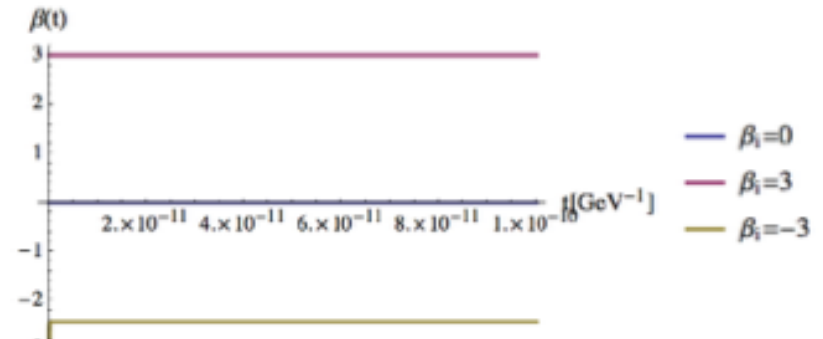
## Observational constraint

$$10^{-15}[\text{G}] \lesssim B_{obs} \lesssim 10^{-9}[\text{G}] \longrightarrow \underline{10^{11} \left( \frac{T_{reh}}{10^{10}[\text{GeV}]} \right)^2 [\text{GeV}^2] \lesssim B_{reh} \lesssim 10^{17} \left( \frac{T_{reh}}{10^{10}[\text{GeV}]} \right)^2 [\text{GeV}^2]}$$



# Anisotropy from Higgs Inflation

- The anisotropy survives during and after Higgs Inflation.



# Table of contents

1. Introduction
2. Cosmology in Einstein-Electroweak theory
3. Dynamics of gauge fields
4. Generation of primordial magnetic fields
5. Conclusion

# Conclusion and outlook

- We studied cosmological dynamics of electroweak gauge fields during and after Higgs Inflation.
- Gauge fields rapidly grow due to parametric resonances after inflation and then sizable primordial magnetic fields which might be compatible with observation.
- Remarkably, the resultant amplitude of magnetic fields depends on the anisotropy survived during inflation.
- To obtain precise results on magnetic fields, we need to introduce matter fields and investigate reheating processes. Moreover, we have to study evolution of inhomogeneity during and after Higgs Inflation.