



A New Method To Determine Large Scale Structure From The Luminosity Distance

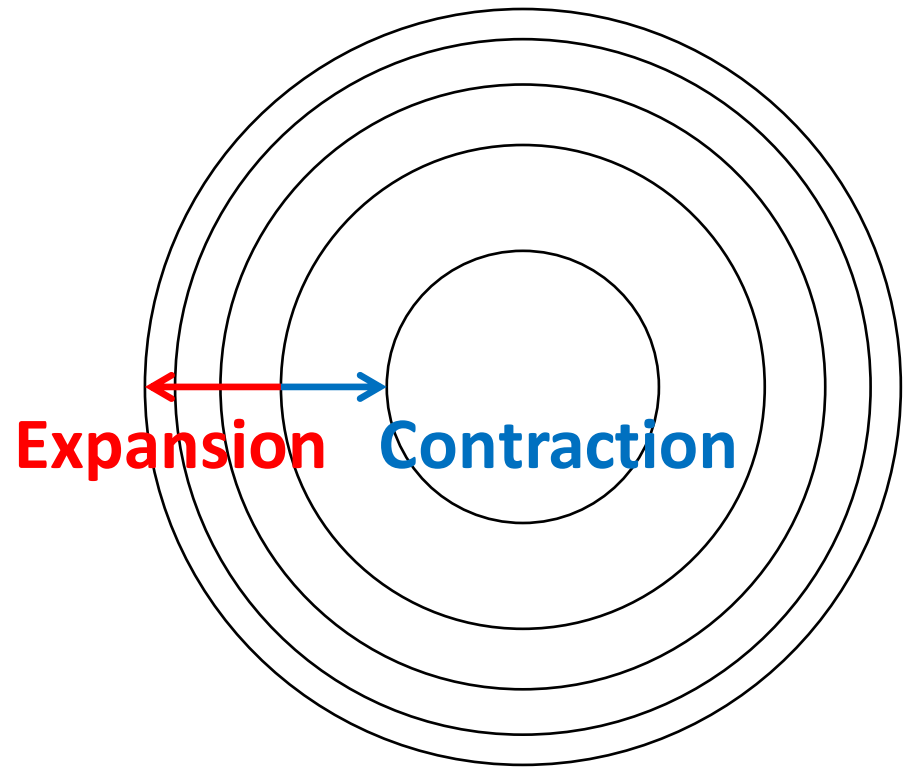
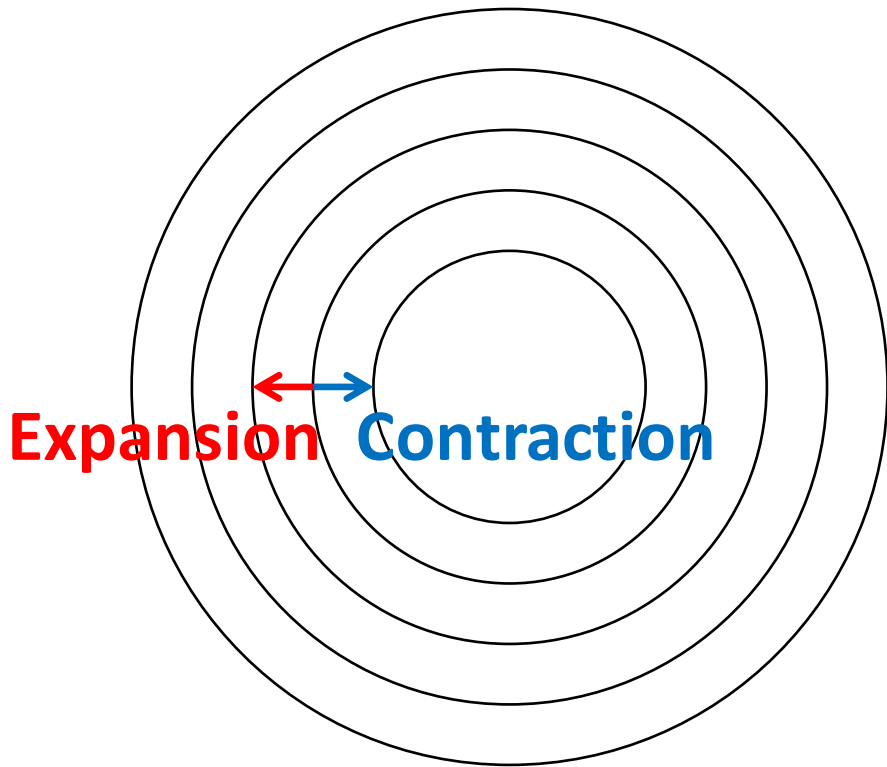
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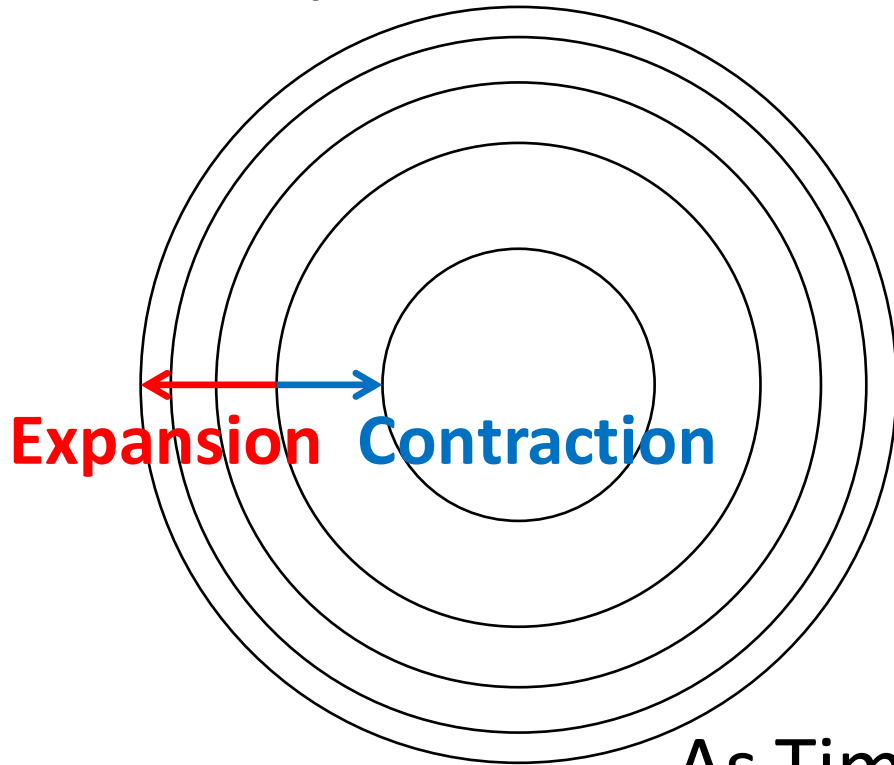
Accelerating Expansion from Large Scale Inhomogeneity

- A homogeneous universe (FRW model)
- An inhomogeneous universe (Void model)

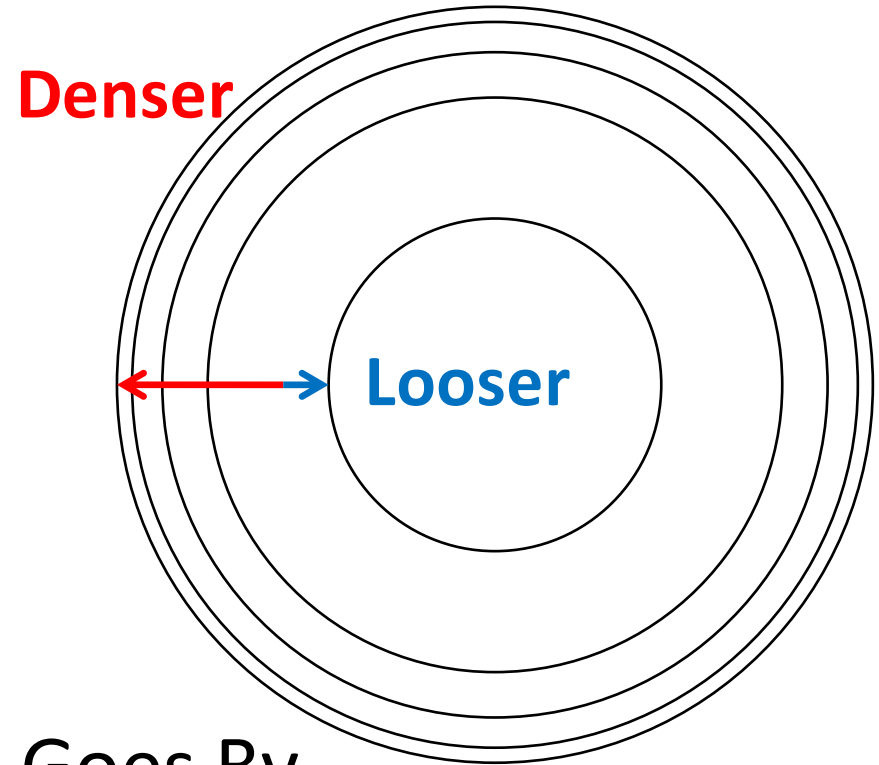


Accelerating Expansion from Large Scale Inhomogeneity

- Early time:



- Now:



As Time Goes By

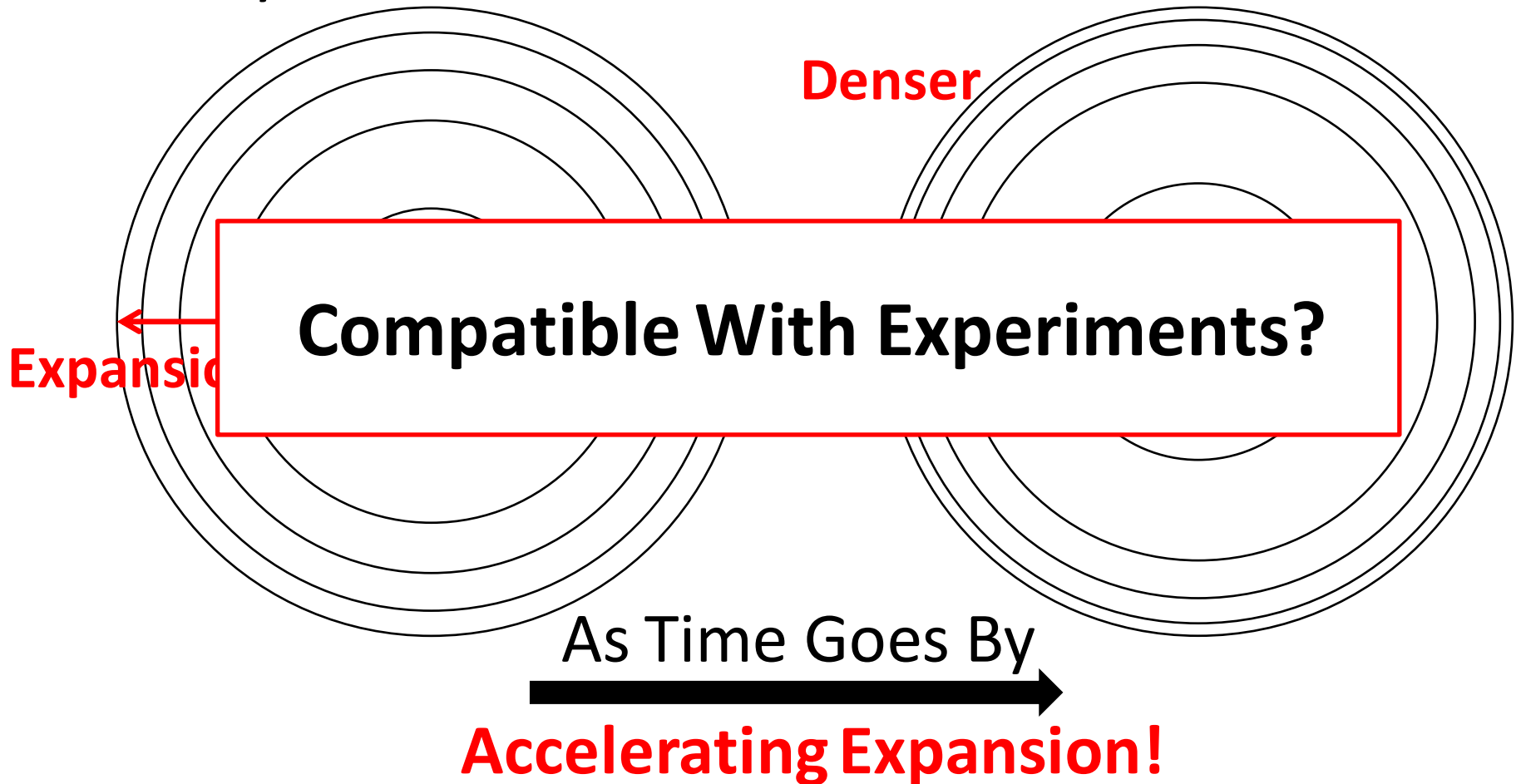


Accelerating Expansion!

Accelerating Expansion from Large Scale Inhomogeneity

- Early time:

- Now:



Outline

- Accelerating Expansion from Large Scale Inhomogeneity: LTB Model
- Mimicking Λ CDM Model: Central Spatial Curvature as “Free Parameter”
- Climbing over Apparent Horizon: An Unique Solution

(R = Areal Radius)

Modeling Large Scale Inhomogeneity

- Assuming spherical symmetry (LTB metric):

$$ds^2 = -dt^2 + a(t, r)^2 \left[\left(1 + \frac{r \partial_r a(t, r)}{a(t, r)} \right)^2 \frac{dr^2}{1 - k(r) r^2} + r^2 d\Omega^2 \right],$$

$$\rho = \frac{2 \partial_r M}{a^2 r^2 (r \partial_r a + a)}, \quad H^2 = \left(\frac{\partial_t a}{a} \right)^2 = \frac{-k(r)}{a^2} + \frac{2M(r)}{a^3 r^3}$$

- Fix gauge freedom of r by setting $M(r) = \frac{1}{6} \rho_0 r^3$

- Conformal time $\eta = \int_0^t \frac{d\tau}{a(\tau, r)} + t_b(r)$

$$a(\eta, r) = \frac{\rho_0}{6k(r)} \left[1 - \cos(\sqrt{k(r)} \eta) \right] \quad \text{From now on we set } t_b = 0$$

$$t(\eta, r) = \frac{\rho_0}{6k(r)} \left[\eta - \sqrt{k(r)}^{-1} \sin(\sqrt{k(r)} \eta) \right] + t_b(r)$$

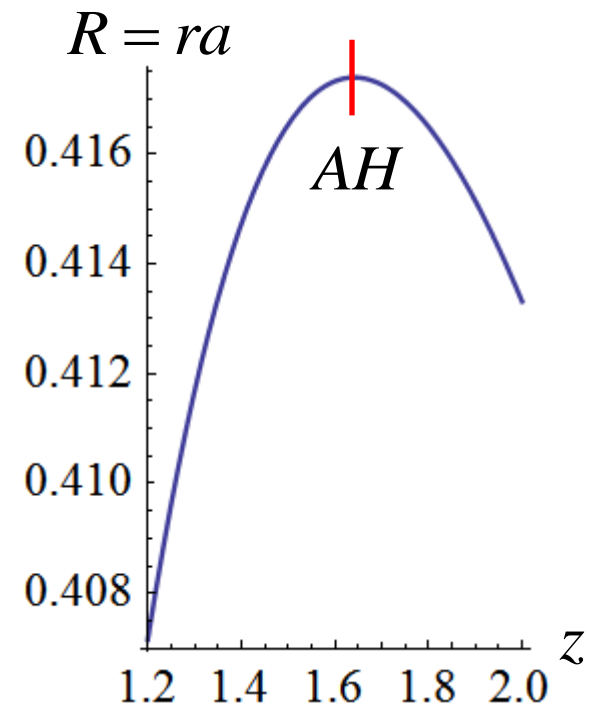
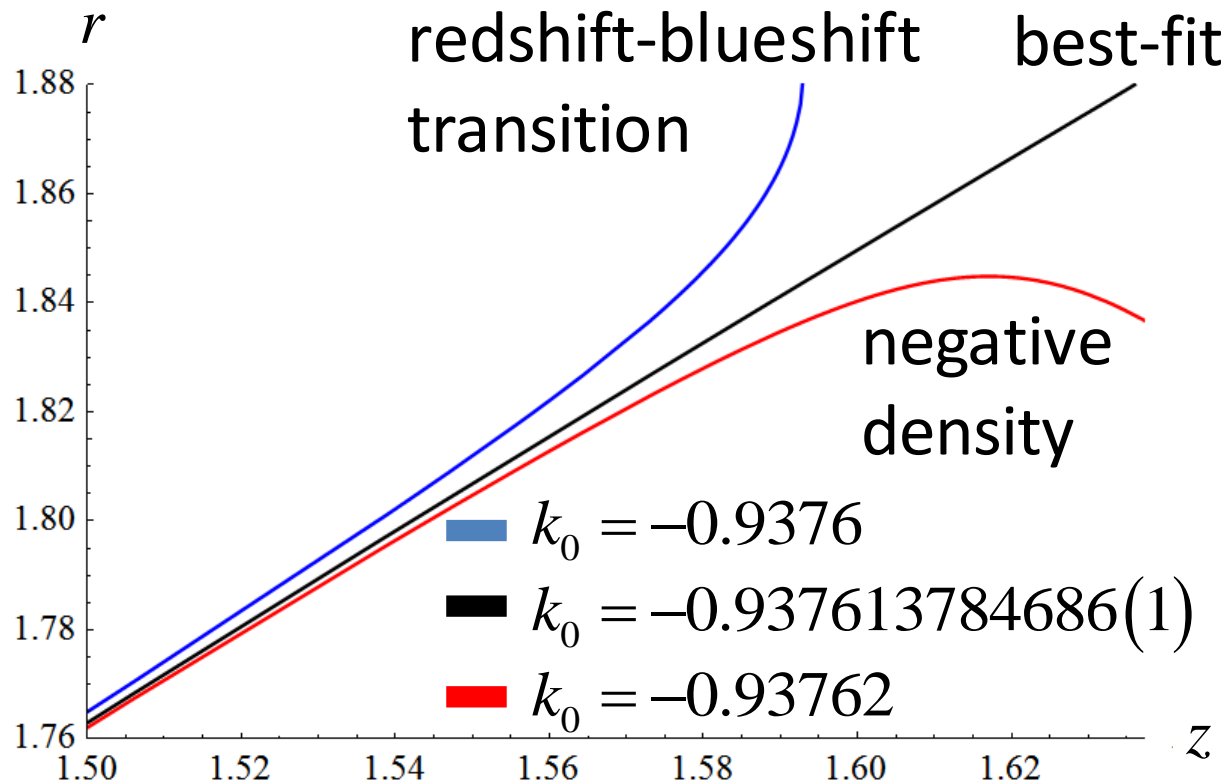
Fixing Initial Condition

- Assuming a central observer, we have
 $r(z=0)=0$, $a(z=0)=a_0$, $\eta(z=0)=\eta_0$, $k(z=0)=k_0$
and $H^{LTB}(z=0)=H_0$, where $H^{LTB} = \frac{\partial \ln a(t, r)}{\partial t}$.
- Since $a(\eta=\eta_0, r=0)=a_0$ and $H^{LTB}(\eta=\eta_0, r=0)=H_0$
are given, η_0 , k_0 and ρ_0 can be determined up
to k_0 .
- Central spatial curvature as “free parameter”
- Luminosity distance on past null geodesic
 $D_L^{obs}(z) = (1+z)^2 a(t(z), r(z)) r = D_L^{FRW}(z)$ is the input.

(R = Areal Radius)

Numerical Results

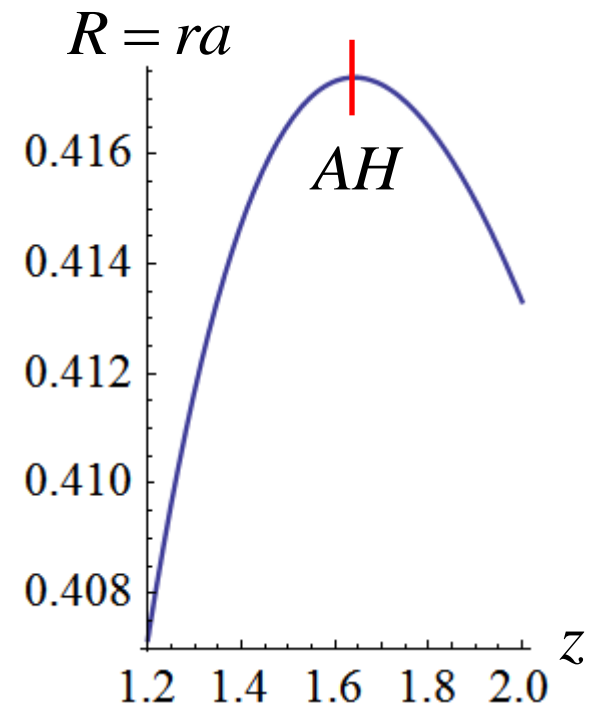
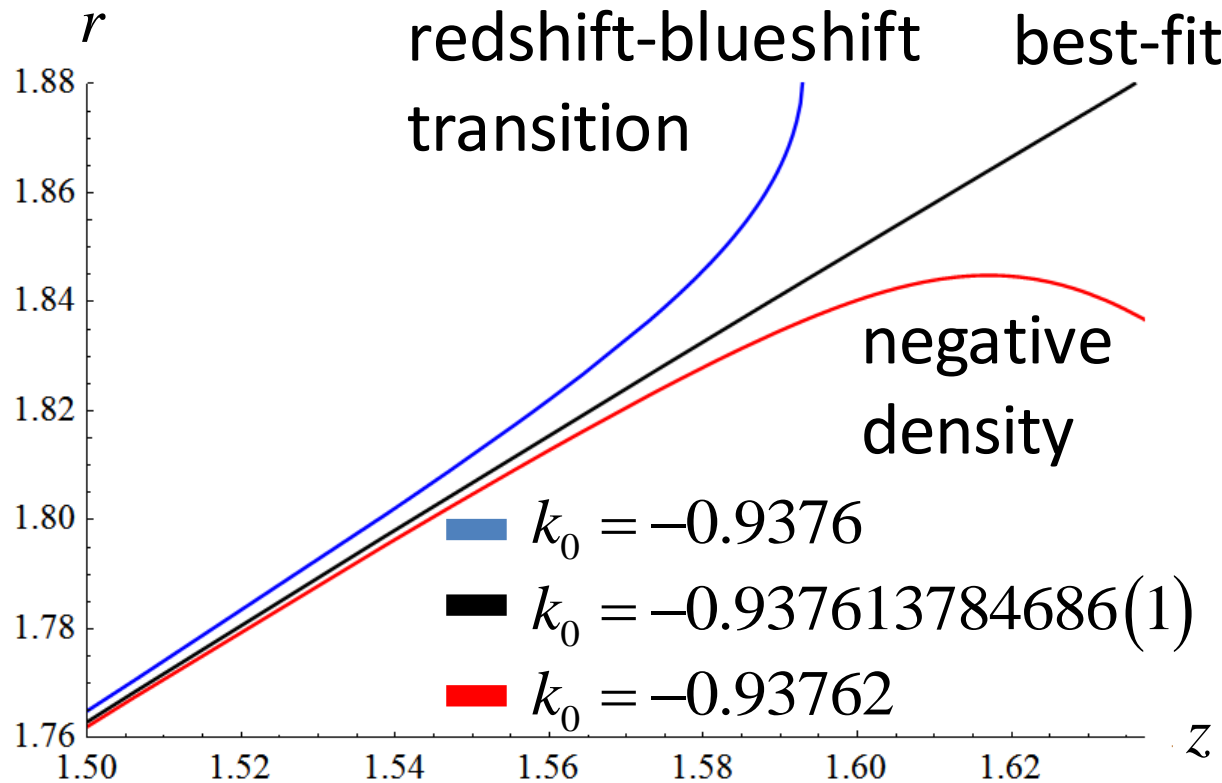
- Not so stable around apparent horizon at $z \sim 1.6$
- Overcome AH through extrapolation?!



(R = Areal Radius)

Numerical Results

Only One Valid k_0 , Why?



Expansion Around Apparent Horizon

- Staring at the geodesic equations, we found a common denominator $r\sqrt{k(1-s^2)} - s\sqrt{1-kr^2}$, where $s = \sqrt{H_0ka(1+k_0)^{-1}}$.

Unstable when $R \approx (1+k_0)r^3/H_0$ or $s \approx \sqrt{kr^2}$

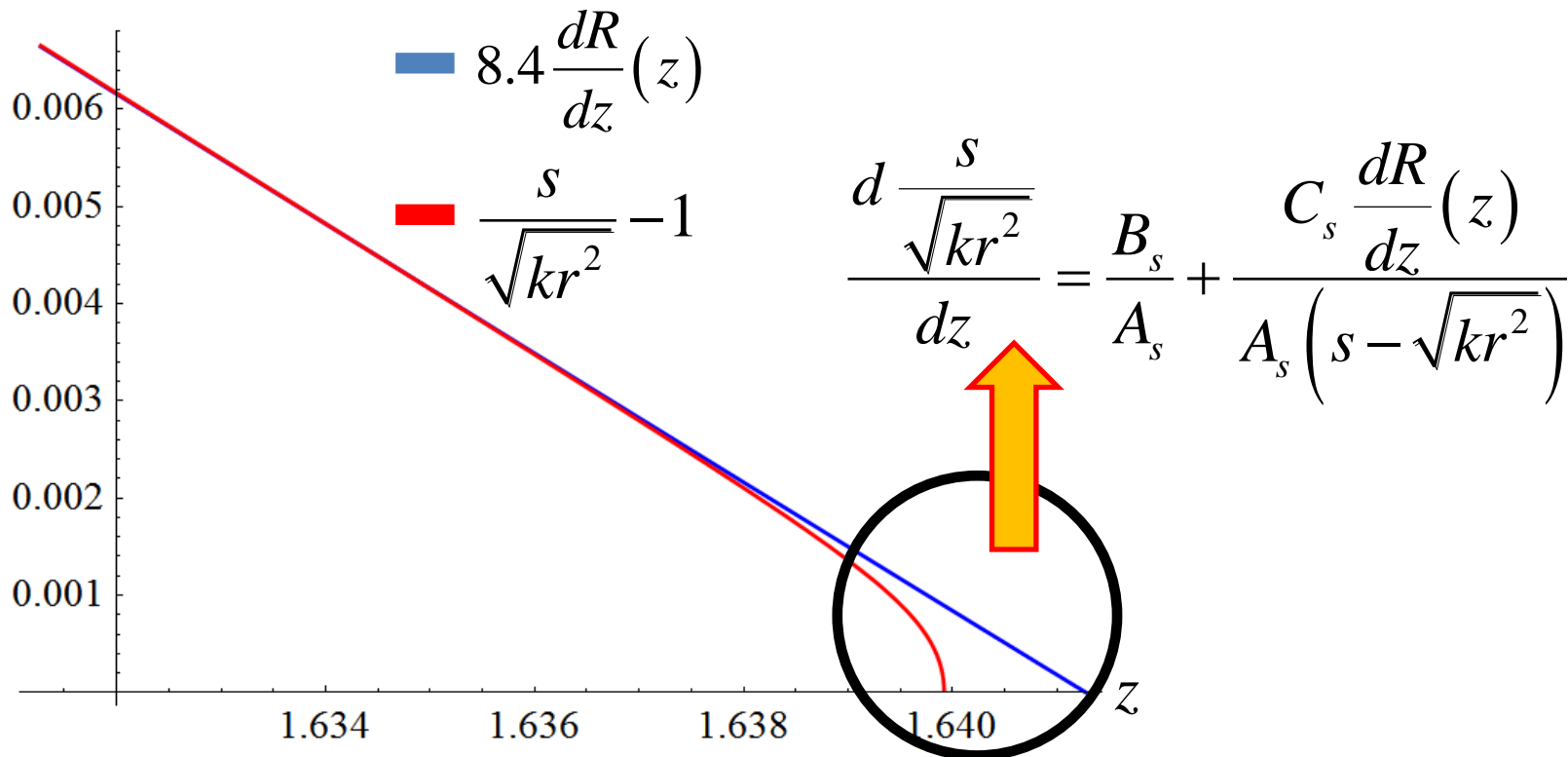
- Expand the numerator around $R = (1+k_0)r^3/H_0$

$$\frac{dk}{dz} = \frac{B_k}{A_k} + \frac{C_k \frac{dR}{dz}(z)}{A_k (s - \sqrt{kr^2})}, \quad \frac{dr}{dz} = \frac{B_r}{A_r} + \frac{C_r \frac{dR}{dz}(z)}{A_r (s - \sqrt{kr^2})}, \quad \frac{d\eta}{dz} = \frac{B_\eta}{A_\eta} + \frac{C_\eta \frac{dR}{dz}(z)}{A_\eta (s - \sqrt{kr^2})}$$

- Needs $s = \sqrt{kr^2}$ and $\frac{dR}{dz}(z) = 0$ at same spot

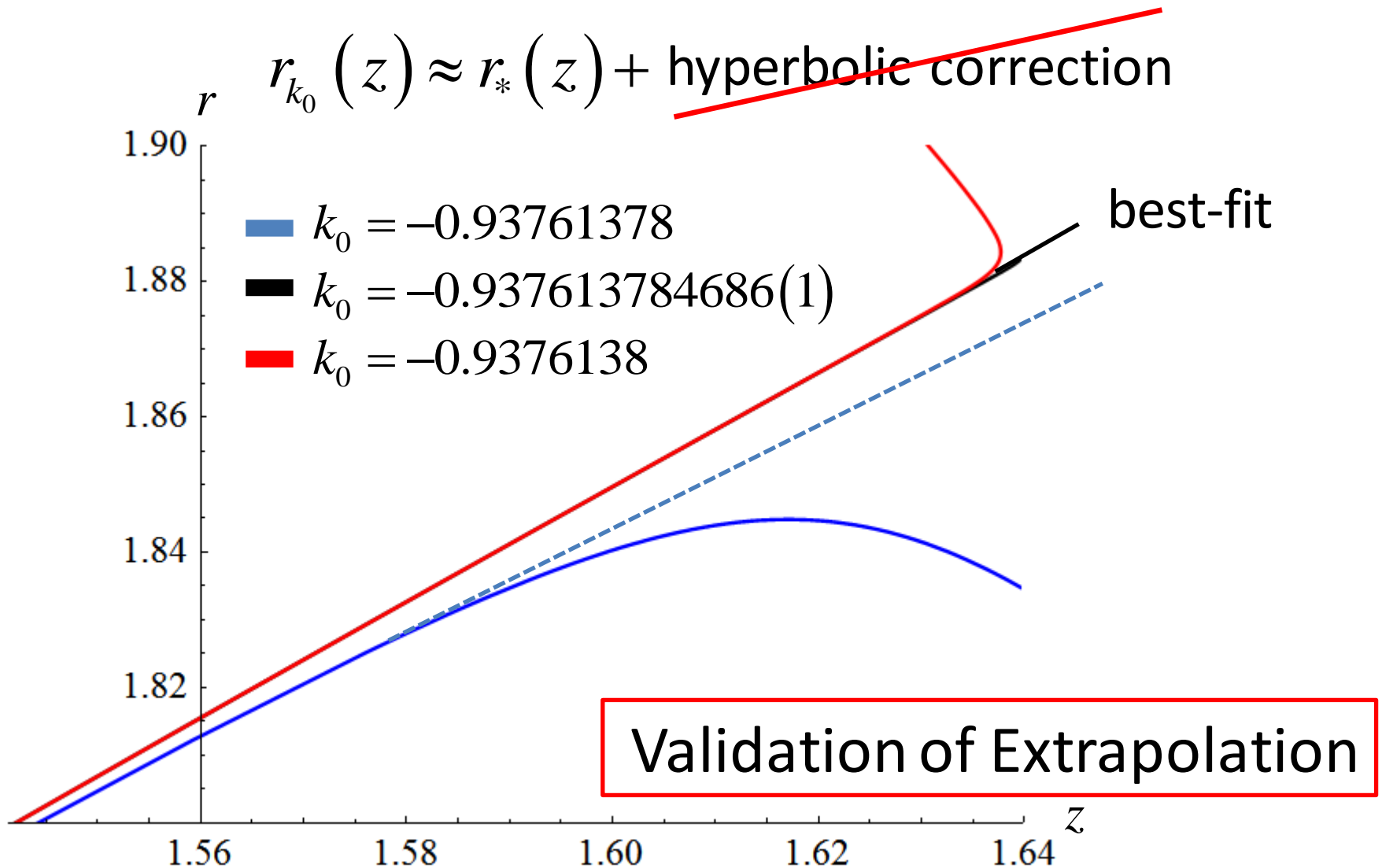
Expansion Around Apparent Horizon

- Needs $s = \sqrt{kr^2}$ and $\frac{dR}{dz}(z) = 0$ at same spot
- Also $\frac{s}{\sqrt{kr^2}} - 1 \propto (z - z_{AH})$



Expansion Around Apparent Horizon

$$r_{k_0}(z) \approx r_*(z) + \text{hyperbolic correction}$$

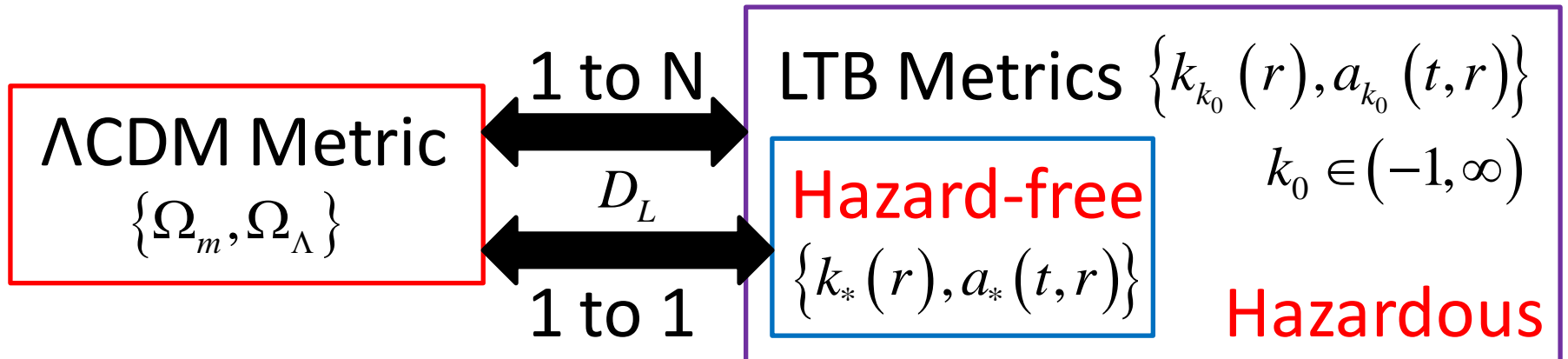


Uniqueness of k_0

- Sudden jump happens only at $R_{AH} \approx (1+k_0)r_{z_{AH}}^3/H_0$
- The existence of solution extended beyond AH is indicated by transit of cause of stop of integrator between $\frac{dr}{dz} = 0$ and $\frac{dz}{dr} = 0$.
- We scanned over parameter space $k_0 \in (-1, \infty)$ and found 1 solution.
- Not a rigorous proof yet

Conclusion

- There exists 1 to N correspondence between Λ CDM metric with certain parameter, and LTB metrics with specific setups that mimic the luminosity distance of that Λ CDM metric.
- But only 1 LTB metric can go beyond apparent horizon without hazards like negative density.



Conclusion

- The error of extrapolation used to overcome apparent horizon is marginal as long as k_0 used in simulation is close to the best fit k_0 value.
- Best extrapolation method is 1st order Taylor expansion.

The End

