Large Scale Structures in Kinetic Gravity Braiding Model and Observational Implication

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Based on R. Kimura and K. Yamamoto JCAP 04 (2011) 025 R. Kimura, T.Kobayashi and K. Yamamoto, in preparation





Introduction

- Kinetic Gravity Braiding Model
- Large Scale Structures
- Observational Constraints
- Summary

Modified Gravity



What causes an accelerated expansion of the universe ??

- Cosmological constant ??
- Dark energy ??
- Modification of gravity ??

Modified gravity theory must satisfy ...

- \checkmark Modification at large distance
- \checkmark Recovery of GR at small scale
- \checkmark Accelerated expansion of the universe at present
- \checkmark Consistency with cosmological observations

Kinetic Gravity Braiding Model

KGB model (Deffayet et al. 2010, Kobayashi et al. 2010)

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\rm Pl}^2}{2} R + K(X) - G(X) \Box \phi + \mathcal{L}_{\rm m} \right]$$

- $\begin{aligned} X &= -g^{\mu\nu} \nabla_{\mu} \phi \nabla_{\nu} \phi / 2 \\ \Box \phi &= g^{\mu\nu} \nabla_{\mu} \nabla_{\nu} \phi \\ K(X), \ G(X) : \text{ arbitrary functions of the kinetic term} \end{aligned}$
- \checkmark Second order differential equation
- \checkmark No ghost instability
- \checkmark Self-accelerating solution
- ✓ GR at small scale (Vainshtein effect)

Kinetic Gravity Braiding Model

Example (Kimura and Yamamoto. 2011)

$$K(X) = -X$$
$$G(X) = M_{\rm Pl} \left(\frac{r_c^2}{M_{\rm Pl}^2} X\right)^n$$

 r_c : crossover scale ($\sim H_0^{-1}$) n: model parameter (n > 1/2)

Choosing the attractor solution, $\dot{\phi} = K_X/3G_XH$

 $\left(\frac{H}{H_0}\right)^2 = (1 - \Omega_0) \left(\frac{H}{H_0}\right)^{-\frac{2}{2n-1}} + \Omega_0 a^{-3}$ behave like dark energy Dvali-Turner's Model (2003)



Observational Constraints







How about large scale structures ?

Kinetic Gravity Braiding Model

Newtonian gauge

$$ds^{2} = -(1+2\Psi)dt^{2} + a^{2}(t)(1+2\Phi)\delta_{ij}dx^{i}dx^{j}$$

Perturbed Einstein equations

$$2M_{\rm Pl}^2 \left[-3H(\dot{\Phi} - H\Psi) + \frac{1}{a^2} \nabla^2 \Phi \right] = -\delta\rho - \delta\rho_\phi$$
$$2M_{\rm Pl}^2 \left(\dot{\Phi} - H\Psi \right) = \delta q + \delta q_\phi$$
$$2M_{\rm Pl}^2 \left[(3H^2 + 2\dot{H})\Psi + H\dot{\Psi} - \ddot{\Phi} - 3H\dot{\Phi} \right] = \delta p_\phi$$
$$\Psi + \Phi = 0$$

where

$$\begin{split} \delta\rho_{\phi} &= K_X \delta X + G_X \left(3\dot{\phi}^3 \dot{\Phi} - 12H\dot{\phi}^3 \Psi + 9H\dot{\phi}^2 \dot{\delta\phi} - \frac{\phi^2}{a^2} \nabla^2 \delta\phi \right) + 3G_{XX} H\dot{\phi}^3 \delta X \\ \delta q_{\phi} &= -K_X \dot{\phi} \delta\phi - G_X \dot{\phi}^2 \left(\dot{\phi} \Psi - \dot{\delta\phi} + 3H\delta\phi \right) \\ \delta p_{\phi} &= K_X \delta X + G_X \left(\dot{\phi}^3 \dot{\Psi} - \dot{\phi}^2 \ddot{\delta\phi} + 4\dot{\phi}^2 \ddot{\phi} \Psi - 2\dot{\phi} \ddot{\phi} \dot{\delta\phi} \right) - G_{XX} \dot{\phi}^2 \ddot{\phi} \delta X \end{split}$$

Perturbed scalar field equations

$$\begin{split} &-K_{X}\left[3\dot{\phi}\dot{\Phi}-\dot{\phi}\dot{\Psi}-2(\ddot{\phi}+3H\dot{\phi})\Psi+\ddot{\delta}\phi+3H\dot{\delta}\phi-\frac{1}{a^{2}}\nabla^{2}\delta\phi\right]\\ &-3G_{XXX}H\dot{\phi}^{3}\ddot{\phi}\delta X-G_{X}\left[3\dot{\phi}^{2}\ddot{\Phi}+6(\ddot{\phi}+3H\dot{\phi})\dot{\phi}\dot{\Phi}-9H\dot{\phi}^{2}\dot{\Psi}\right.\\ &-12\left\{(\dot{H}+3H^{2})\dot{\phi}^{2}+2H\dot{\phi}\ddot{\phi}\right\}\Psi-\frac{\dot{\phi}^{2}}{a^{2}}\nabla^{2}\Psi+6H\dot{\phi}\ddot{\delta}\phi\\ &+6\left\{H\ddot{\phi}+(\dot{H}+3H^{2})\dot{\phi}\right\}\dot{\delta}\phi-\frac{2}{a^{2}}(\ddot{\phi}+2H\dot{\phi})\nabla^{2}\delta\phi\right]\\ &-G_{XX}\left[3\dot{\phi}^{3}\ddot{\phi}\dot{\Phi}-3H\dot{\phi}^{4}\dot{\Psi}-3\left\{8H\dot{\phi}^{3}\ddot{\phi}+(\dot{H}+3H^{2})\dot{\phi}^{4}\right\}\Psi\\ &+3H\dot{\phi}^{3}\ddot{\delta\phi}+3\left\{5H\dot{\phi}^{2}\ddot{\phi}+(\dot{H}+3H^{2})\dot{\phi}^{3}\right\}\dot{\delta\phi}-\frac{\dot{\phi}^{2}\ddot{\phi}}{a^{2}}\nabla^{2}\delta\phi\right]=0\end{split}$$

Large Scale Structures

Large n solution

$$\dot{\delta X} + 3H\delta X = 0$$

$$\delta X = \dot{\phi}\dot{\delta\phi} - \dot{\phi}^2\Psi$$

$$\delta X = \operatorname{Const}/a^3 \to 0$$
$$c_s^2 \simeq 0$$

NO effects caused by the scalar field at linear order (= Λ CDM model)

• <u>Small n solution (+sub-horizon)</u>

$$\mathcal{O}(k^2 c_s^2/a^2) \gg \mathcal{O}(H^2)$$

 $\ddot{\delta} + 2H\dot{\delta} \simeq 4\pi G_{\text{eff}}\rho\delta$ $G_{\text{eff}} = G\left[1 + 4\pi G\frac{G_X^2\dot{\phi}^4}{\beta(a)}\right]$

The growth of density perturbations should be different from the ACDM model !!









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Constraints from the matter power spectrum





What could be a powerful tool to constrain on the KGB model ?

Galaxy-ISW Cross-correlation

ISW term in CMB anisotropy

$$\left(\frac{\Delta T(\vec{\gamma})}{T}\right)_{\rm ISW} = \int_{\eta_d}^{\eta_0} d\eta [\Psi'(\eta, \mathbf{x}) - \Phi'(\eta, \mathbf{x})]$$

Modified Poisson equation

$$\frac{\nabla^2}{a^2}\Psi = 4\pi G_{\rm eff}\rho\delta$$

Galaxy-ISW Cross-correlation



determine the sign of CCF

 $\Lambda CDM model$

Kinetic gravity braiding (small n)

Galaxy-LSS Cross-correlation



Summary



✓KGB model has a self-accelerating solution and passes solar system constraints

✓ The background evolution can mimic the ACDM model. However, the growth of LSS in KGB model has different signatures from the ACDM model

 \checkmark ISW-LSS cross-correlation is a powerful tool to constrain on modified gravity

✓ Small n value in the KGB model (correspond to the galileon model) is disfavored by Galaxy-ISW cross-correlation



Thank you !!