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Results

Light Propagation through exact non-linear inhomogeneities in ACDM cosmology

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28.07.2011 / RESCEU/DENET Summer School Work with Marco Bruni: arXiv:1103.0501 and arXiv:1107.4433

Outline



- 2
 - Our model
 - The metric and density deviations, arXiv:1103.0501
 - Light tracing equations, arXiv:1107.4433

Results 3

- Compensated δ
- Non-compensated δ

Summary



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Summary

Current status of cosmology



$$\begin{split} &\mathsf{WMAP7+BAO+}H_0\\ &\Omega_\Lambda = 0.725 \pm 0.016\\ &\Omega_m h^2 = 0.1126 \pm 0.0036\\ &w = -1.10 \pm 0.14\\ &\mathsf{Komatsu\ et\ al.\ (2011)} \end{split}$$

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The distance measure

IN FLRW:

 $\frac{d^2(d_A)}{d\lambda^2} = -\frac{1}{2}\bar{E}^2\bar{\rho}d_A,$



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The distance measure

IN FLRW:

$$\frac{d^2(d_A)}{d\lambda^2} = -\frac{1}{2}\bar{E}^2\bar{\rho}d_A,$$

In general:

$$egin{array}{rcl} rac{d^2(d_{\mathcal{A}})}{d\lambda^2} &=& -\left[rac{1}{2}E^2ar
ho(1+\delta)+|\sigma|^2
ight]d_{\mathcal{A}}, \ &rac{d\sigma}{d\lambda} &+& 2rac{rac{d(d_{\mathcal{A}})}{d\lambda}}{d_{\mathcal{A}}}\sigma=\Psi_0, \end{array}$$

where d_A is the angular diameter distance *E* is photon energy σ is shear of light bundle

 Ψ_0 is Weyl scalar

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Our metric

Consider here the flat ACDM sub-case

$$ds^2 = -dt^2 + S(t)^2 \left[dx^2 + dy^2 + Z(\mathbf{x}, t)^2 dr^2
ight]$$

where

$$Z(\mathbf{x},t) = A(\mathbf{x}) + F(r,t)$$

and x stands for all three Cartesian spatial coordinates

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The metric deviation Z splits

We find

$$\boldsymbol{A} = \boldsymbol{1} + \boldsymbol{B}\beta_{+}(\boldsymbol{r})\left\{ [\boldsymbol{x} + \gamma(\boldsymbol{r})]^{2} + [\boldsymbol{y} + \omega(\boldsymbol{r})]^{2} \right\}$$

and

$$F(r,t) = \beta_+(r)f_+(t) + \beta_-(r)f_-(t)$$

Results

The metric deviation Z splits

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$$F(r,t) = \beta_+(r)f_+(t) + \beta_-(r)f_-(t)$$



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Density deviation example



Here $\gamma(r) = \omega(r) = \beta_{-}(r) = 0$ and $\beta_{+} \propto sin(kr)$ for $k = 2\pi/8 \text{Mpc}^{-1}$.

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The null geodesic equations

The very simplified case for light rays along *r*-axis

$$-\frac{E'}{E}=\frac{S'}{S}+\frac{F'}{1+F},$$

and

$$r'=rac{2}{3}rac{1}{H_0\sqrt{\Omega_\Lambda}SZ},$$

Motivation

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The Sachs optical equations

$$\begin{aligned} d''_{A} + d'_{A} \frac{E'}{E} &= \left(-\frac{2}{9} \frac{\rho}{H_{0}^{2} \Omega_{\Lambda}} - \frac{4}{3} \frac{|\tilde{\sigma}|^{2}}{E^{2}} \right) d_{A}, \\ \tilde{\sigma}' &+ 2 \frac{d'_{A}}{d_{A}} \tilde{\sigma} = 0 \end{aligned}$$

Here, the shear can be set to zero.

Results

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Inhomogeneities periodic on 8 Mpc scale



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Inhomogeneities periodic on 100 Mpc scale



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red: over-densities, blue: under-densities, black: FLRW background

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- we have developed an exact inhomogeneous ACDM model
- can model non-linear structure growth and mode coupling

- compensated density has negligible effect on d_A
- non-compensated profiles have large effect
- identifying correct background is crucial

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Results

Thank you

Motivation	Our model	Results 00000	Summary
In our model The shocker slide			

$$-\frac{E_{\tau}}{E} = S^2 \left(\frac{3\Lambda}{4}\right) \left\{ r_{\tau}^2 (F+A) \left[(F+A) \frac{S_{\tau}}{S} + (f_+)_{\tau} \beta_+ \right] + \frac{S_{\tau}}{S} (x_{\tau}^2 + y_{\tau}^2) \right\},\tag{1}$$

$$\mathbf{x}_{\tau\tau} + \left(2\frac{\mathbf{S}_{\tau}}{\mathbf{S}} + \frac{\mathbf{E}_{\tau}}{\mathbf{E}}\right)\mathbf{x}_{\tau} - (\mathbf{F} + \mathbf{A})\beta_{+}(\mathbf{x} + \gamma)\mathbf{r}_{\tau}^{2} = \mathbf{0},$$
(2)

$$y_{\tau\tau} + \left(2\frac{S_{\tau}}{S} + \frac{E_{\tau}}{E}\right)y_{\tau} - (F + A)\beta_{+}(y + \omega)r_{\tau}^{2} = 0, \qquad (3)$$

$$r_{\tau\tau} + r_{\tau}^{2} \frac{(F+A)_{r}}{F+A} + 2r_{\tau} \left[\frac{S_{\tau}}{S} + \frac{1}{2} \frac{E_{\tau}}{E} + \beta_{+} \frac{(f_{+})_{\tau} + (x+\gamma)x_{\tau} + (y+\omega)y_{\tau}}{F+A} \right] = 0.$$
(4)

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Motivation	Our model	Results	Summary
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Non-FLRW contributions in the Sachs equations



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angular diameter distance at high z



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