

July 28th 2011

Resummed Propagators

. . . with the eikonal approximation

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Introduction

Curvature perturbations

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Density/velocity perturbations

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Propagator: evolution of perturbations

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Propagator: evolution of perturbations

Fluid description \leadsto Vlasov equation

Fluid equations

{ continuity equation

$$\frac{1}{H} \dot{\delta} + \theta = 0$$

Fluid equations

$$\left\{ \begin{array}{l} \text{continuity equation} \\ \frac{1}{H} \dot{\delta} + \theta = 0 \\ \frac{1}{H} \dot{\theta} + \left(2 + \frac{\dot{H}}{H^2} \right) \theta + \frac{3}{2} \Omega_m \delta = 0 \\ \text{Euler equation} \end{array} \right.$$

Fluid equations

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The continuity equation is labeled with a curved arrow pointing to the first equation. The Euler equation is labeled with a curved arrow pointing to the second equation.

growing mode: $\delta(\mathbf{x}, t) = D_+(t)\delta_+(\mathbf{x}) \quad f_+ = d \ln D_+ / d \ln a$

Fluid equations

$$\Psi_a \equiv \begin{pmatrix} \delta \\ \Theta \end{pmatrix} \xleftarrow{-\frac{\theta}{f_+}}$$

$$\frac{\partial}{\partial \eta} \Psi_a(\mathbf{k}) + \Omega_{ab} \Psi_b(\mathbf{k}) = \gamma_{abc}(\mathbf{k}, \mathbf{k}_1, \mathbf{k}_2) \Psi_b(\mathbf{k}_1) \Psi_c(\mathbf{k}_2)$$

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$$\begin{pmatrix} 0 & -1 \\ -\frac{3\Omega_m}{2f_+^2} & \frac{3\Omega_m}{2f_+^2} - 1 \end{pmatrix}$$

$$\begin{aligned} & \frac{(\mathbf{k}_1 + \mathbf{k}_2) \cdot \mathbf{k}_1}{k_1^2} \\ & \frac{(\mathbf{k}_1 + \mathbf{k}_2)^2 \mathbf{k}_1 \cdot \mathbf{k}_2}{2k_1^2 k_2^2} \end{aligned}$$

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Linear propagator $g_{ab}(\eta, \eta_0)$

$$\frac{\partial}{\partial \eta} g_{ab}(\eta, \eta') + \Omega_{ac}(\eta) g_{cb}(\eta, \eta') = 0$$

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Linear propagator $g_{ab}(\eta, \eta_0)$

$$\boxed{\Psi_a(\mathbf{k}, \eta) = g_{ab}(\eta, \eta_0) \Psi_b(\mathbf{k}, \eta_0) \quad \text{Scoccimarro (2001)} \\ + \int_{\eta_0}^{\eta} d\tilde{\eta} \ g_{ab}(\eta, \tilde{\eta}) \gamma_{bde}(\mathbf{k}, \mathbf{k}_1, \mathbf{k}_2) \Psi_d(\mathbf{k}_1, \tilde{\eta}) \Psi_e(\mathbf{k}_2, \tilde{\eta})}$$

Nonlinear propagator

$$G_{ab}(\eta, \eta_0) \equiv \left\langle \frac{\partial \Psi_a(\eta)}{\partial \Psi_b(\eta_0)} \right\rangle$$

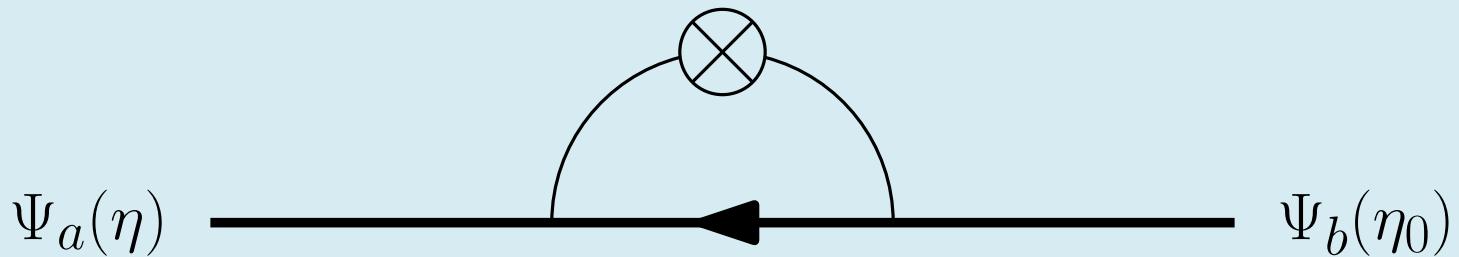
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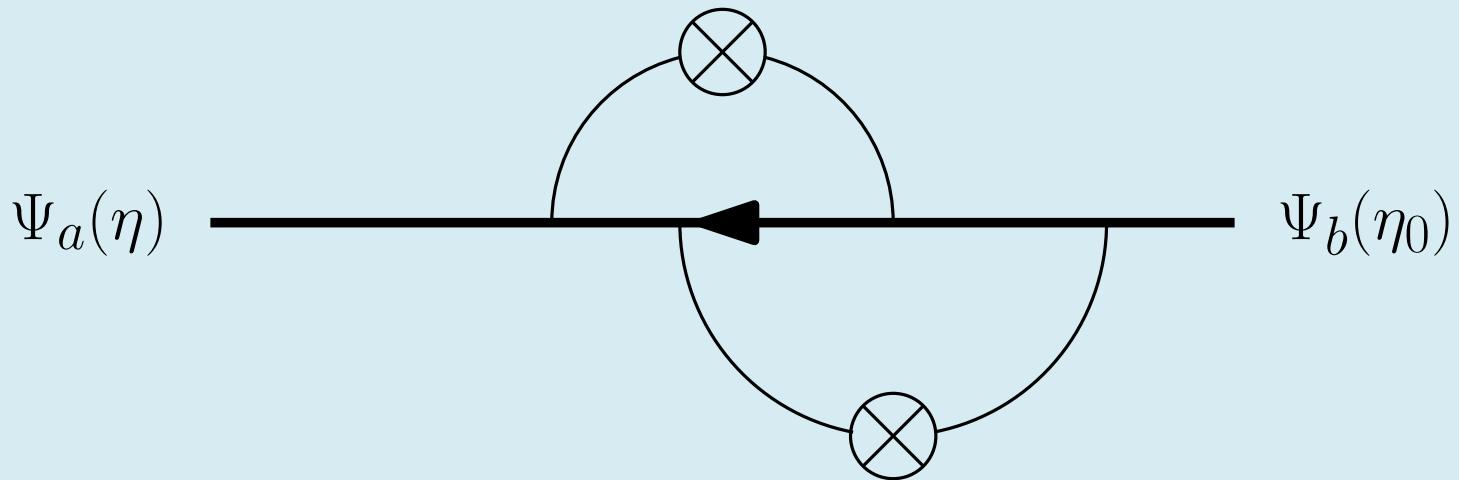
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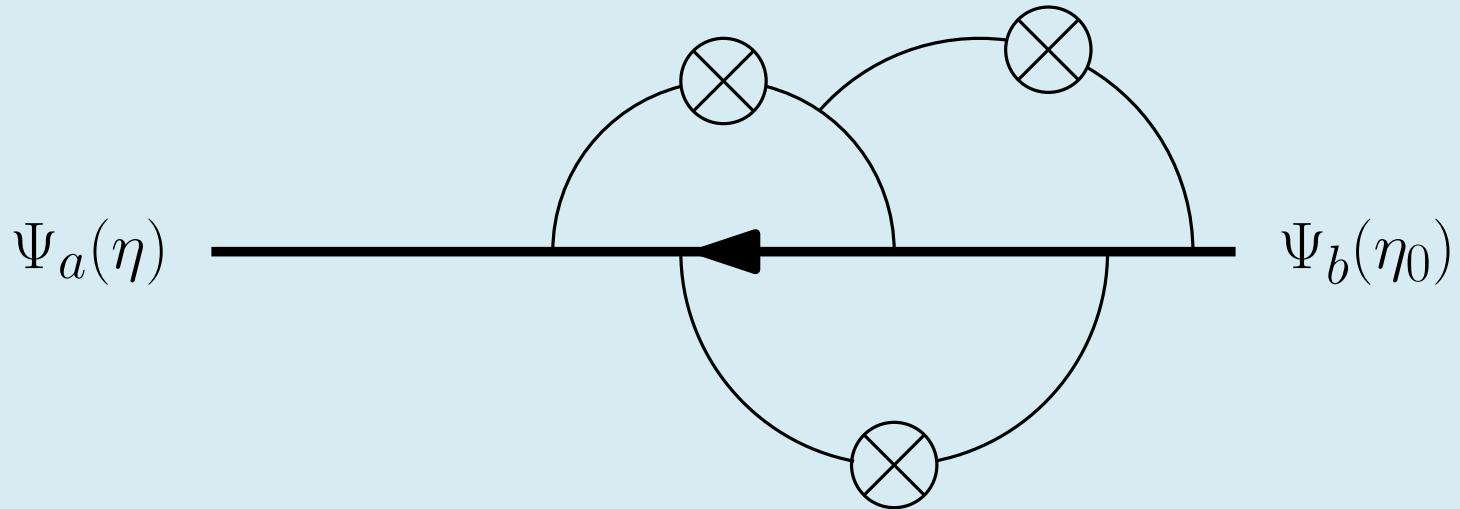
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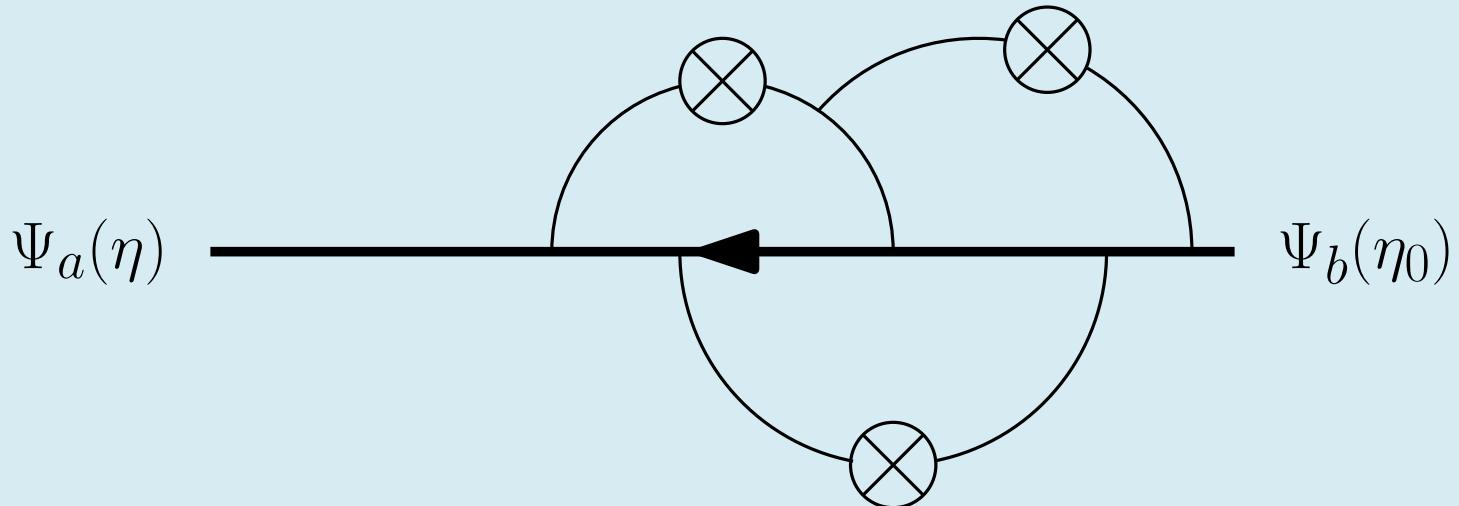
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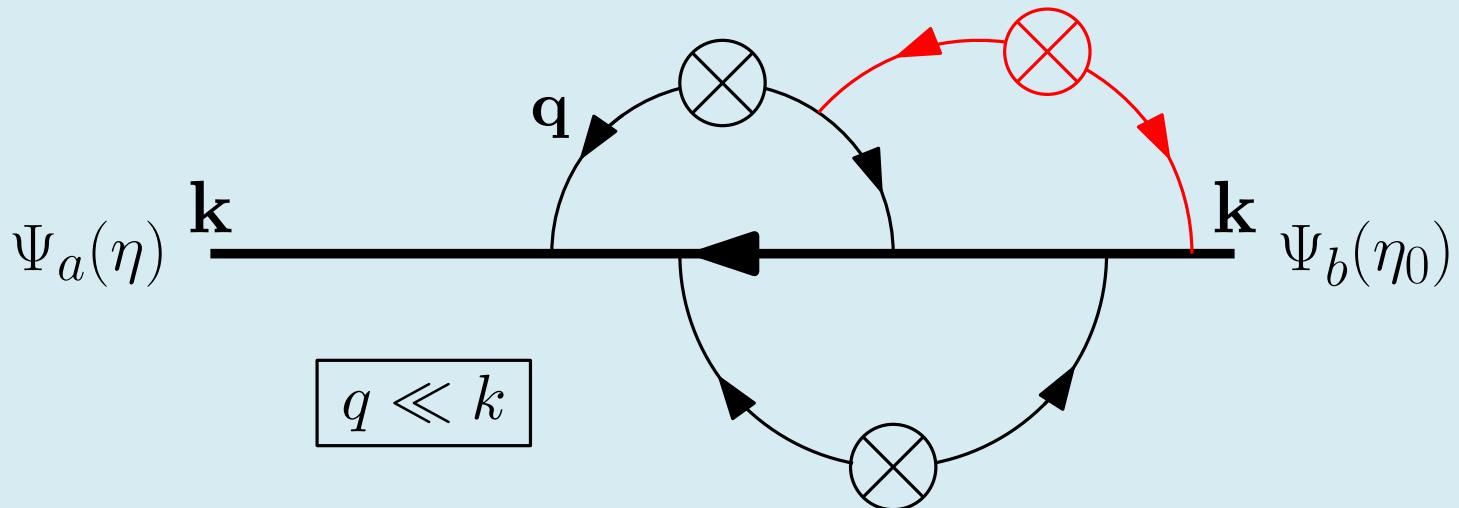


$$G_{ab}(\eta, \eta_0) = g_{ab}(\eta, \eta_0) \exp \left[-k^2 \sigma_d^2 (e^\eta - e^{\eta_0})^2 / 2 \right]$$

Crocce, Scoccimarro (2006)

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Eikonal approximation

$$\frac{\partial}{\partial \eta} \Psi_a(\mathbf{k}) + \Omega_{ab} \Psi_b(\mathbf{k}) = \gamma_{abc}(\mathbf{k}, \mathbf{k}_1, \mathbf{k}_2) \Psi_b(\mathbf{k}_1) \Psi_c(\mathbf{k}_2)$$

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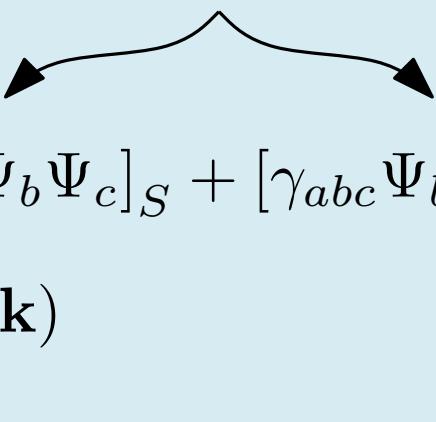
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$$\begin{aligned} & [\gamma_{abc} \Psi_b \Psi_c]_S + [\gamma_{abc} \Psi_b \Psi_c]_H \\ & \Xi_{ab}(\mathbf{k}, \cancel{\mathbf{k}_1}, \cancel{\mathbf{k}_2}) \Psi_b(\mathbf{k}) \\ & \Xi(\mathbf{k}) \Psi_a \end{aligned}$$

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$$\left(\frac{\partial}{\partial \eta} - \Xi(\mathbf{k}) \right) \Psi_a(\mathbf{k}) + \Omega_{ab} \Psi_b(\mathbf{k}) = [\gamma_{abc} \Psi_b \Psi_c]_H$$

Resummed propagator

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$$G_{ab}(\mathbf{k}, \eta) = \langle \xi_{ab}(\mathbf{k}, \eta) \rangle_{\Xi}$$

Multi-component fluid

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decaying modes

→ isodensity modes $\left(\delta = \sum \frac{\Omega_\alpha}{\Omega_m} \delta_\alpha = 0 \right)$

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$$\Xi_{ab} = \Xi_{ab}^{\text{ad}} + \Xi_{ab}^{\text{iso}}$$

Conclusion

- We describe a new method to derive the resummed propagators
- Applicable to non-standard cosmologies and initial conditions
- Applicable to multi-component fluids