



Daisuke Yamauchi, JGRG 22(2012)111420

“Weak lensing generated by vector perturbations and
detectability of cosmic strings”

**RESCEU SYMPOSIUM ON
GENERAL RELATIVITY AND GRAVITATION**

JGRG 22

November 12-16 2012

Koshiba Hall, The University of Tokyo, Hongo, Tokyo, Japan



Weak lensing generated by vector perturbations and detectability of cosmic strings

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- JCAP 1201 (2012) 007, arXiv:1110.1718 [astro-ph.CO]
- JCAP 1210 (2012) 030, arXiv:1205.2139 [astro-ph.CO]
- arXiv:12???.????

Vector mode in cosmology

Scalar mode

- dominant component, gravitational clustering,...

Vector mode

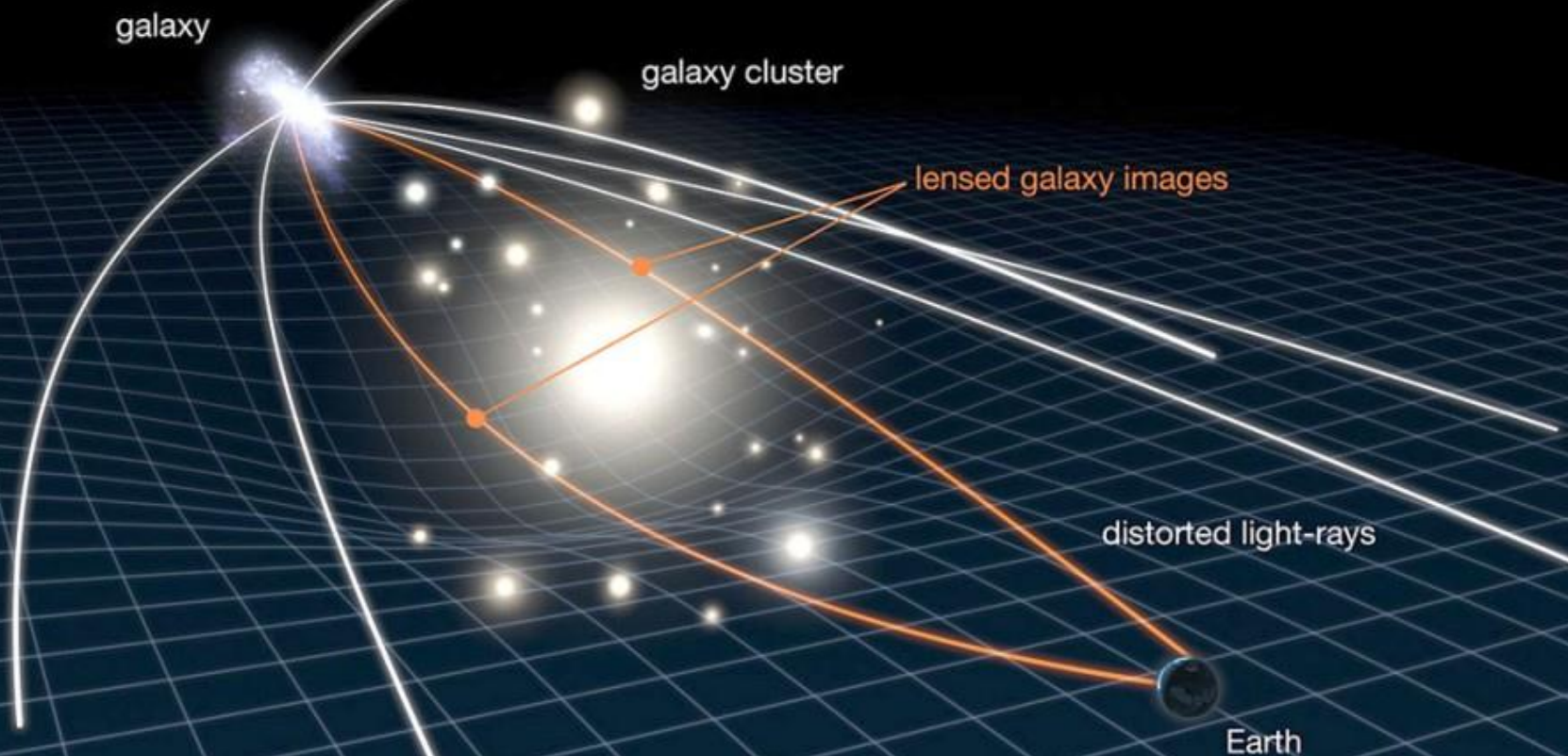
- very minor component and does not serve as a seed of structure formation
- In absence of sources, it decays away, and rapidly becomes negligible ...

- ✓ BUT! It can be generated via a variety of mechanisms :
cosmic strings, magnetic fields, massive neutrinos,
2nd-order density perturbations, modification of gravity,...

...The detection/measurement of them offers an important clue to probe the unknown physics and history of the very early universe!

Gravitational Lensing

= method to “see” invisibles



WEAK LENSING observations can provide a direct evidence for the intervening “**VECTOR MODES**” along a line of sight by measuring the spatial patterns of the deformation of the photon path.

Main Goal

1. We derive the full-sky formulas for

- vector-induced cosmic shear for galaxy survey
- vector-induced deflection angle for CMB-lensing

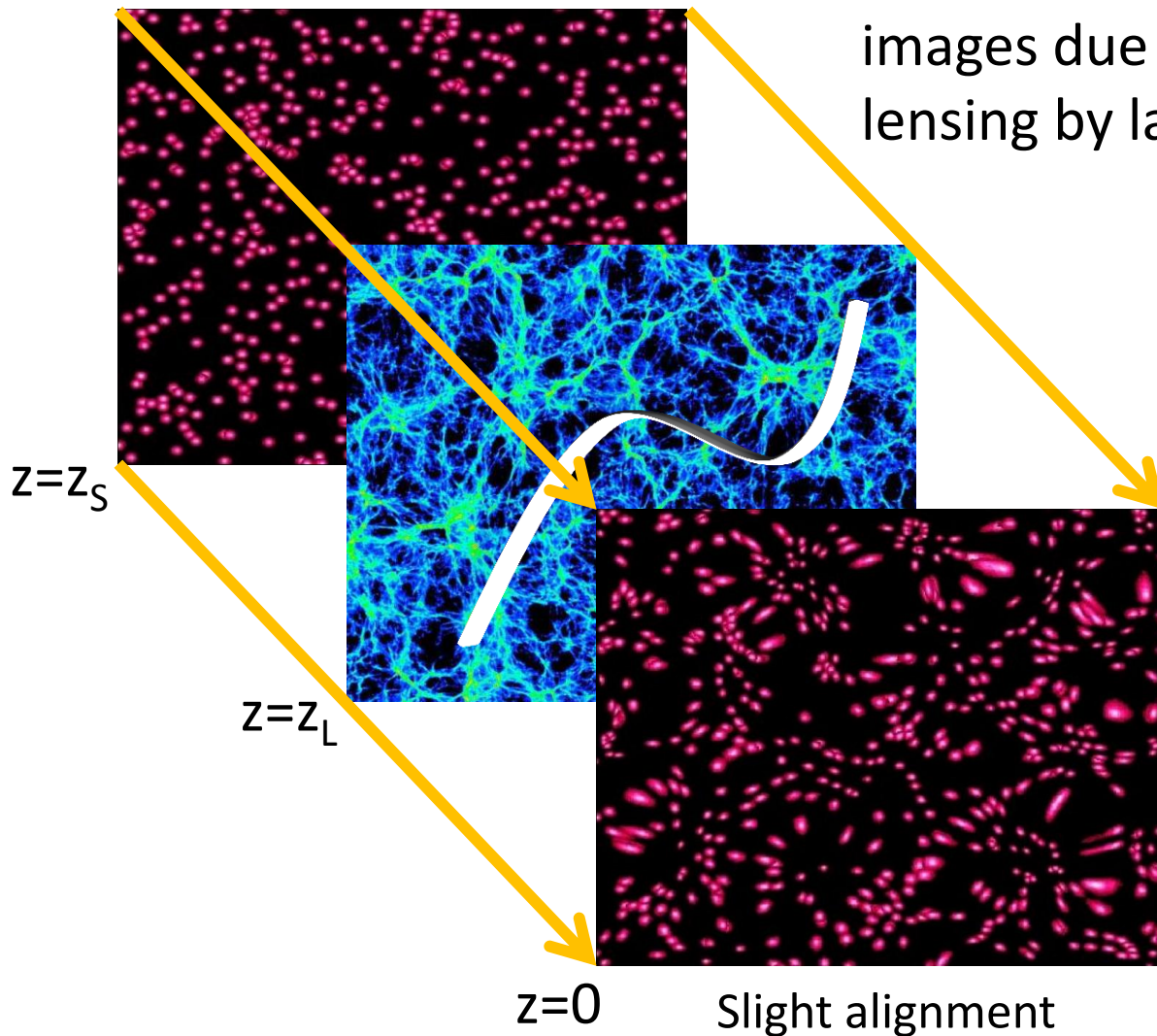
2. As a prospect for detecting shear/deflection, we consider a cosmic string network as a possible source for seeding vector perturbations.

Cosmic shear

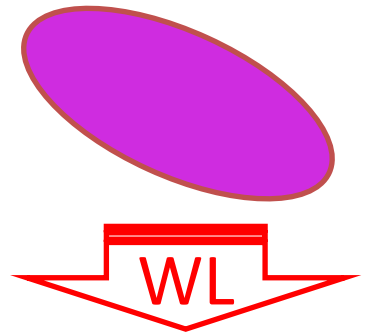
Galaxies randomly distributed

:the deformation of distant-galaxy images due to weak gravitational lensing by large-scale structure

...+ exotic objects



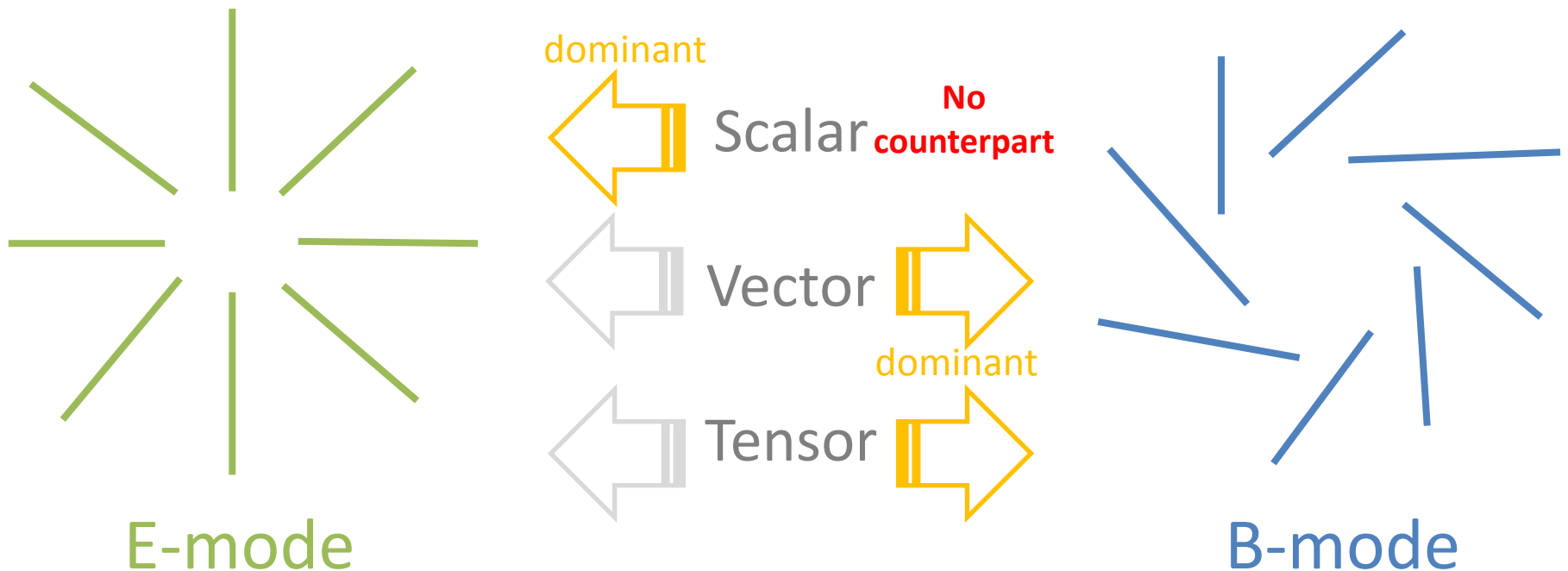
Intrinsic shape of a background galaxy



Galaxy shape actually seen after WL

E-/B-mode decomposition in WL

The spatial pattern of lensing shear fields can be decomposed into even-parity part (**E-mode**) and odd-parity part (**B-mode**).



The non-vanishing B-mode shear signal would be a direct evidence for non-scalar metric perturbations .

(See [Dodelson+(‘03), Schmidt+(‘12),...] for primordial GW)

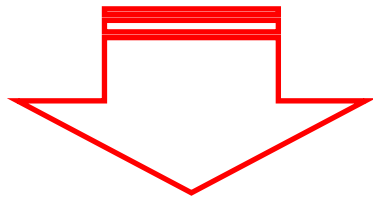
Formula for vector-induced B-mode shear

[DY+Namikawa+Taruya, 1205.2139]

➤ vector perturbations

$$ds^2 = a^2(\eta) \left[-d\eta^2 + 2B_i d\eta dx^i + (\delta_{ij} + 2H_{(i|j)}) dx^i dx^j \right]$$

$$\sigma_{g,i} \equiv \dot{H}_i - B_i : \text{gauge-invariant vector perturbations}$$



Solving the geodesic deviation equation,...

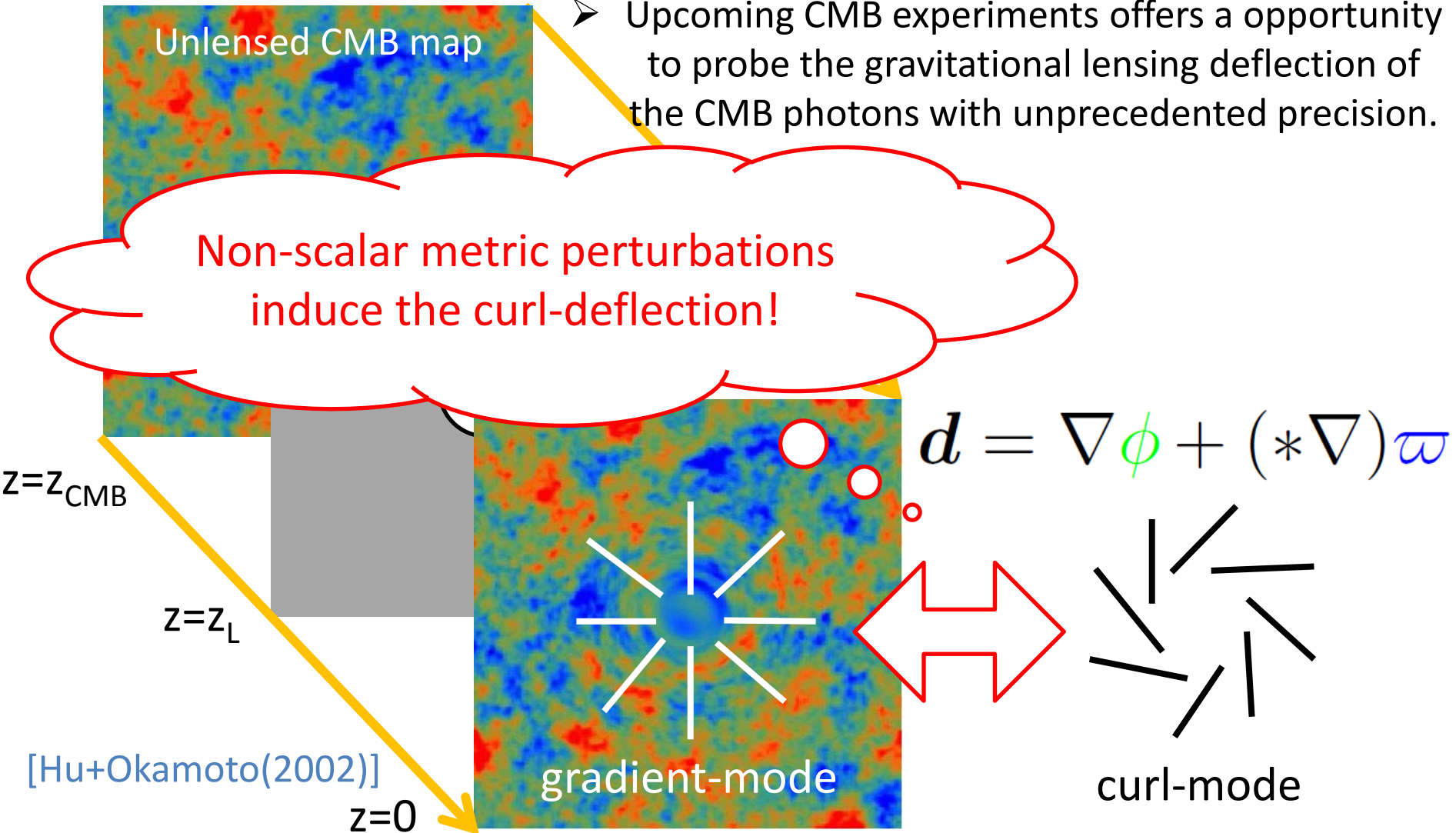
comoving distance

$$C_\ell^{\text{BB}} = \frac{1}{2\pi} \frac{(\ell-1)!(\ell+2)!}{(\ell+1)!(\ell-2)!} \int_0^\infty k^2 dk \int_0^{\chi_S} d\chi \int_0^{\chi_S} d\chi' \times \frac{j_\ell(k\chi)}{\chi} \frac{j_\ell(k\chi')}{\chi'} P_{\sigma_g \sigma_g}(k; \eta_0 - \chi, \eta_0 - \chi')$$

CMB lensing

➤ Upcoming CMB experiments offers a opportunity to probe the gravitational lensing deflection of the CMB photons with unprecedented precision.

Non-scalar metric perturbations induce the curl-deflection!



Formula for vector-induced curl-mode

[DY+Namikawa+Taruya, 1205.2139]

$$ds^2 = a^2(\eta) \left[-d\eta^2 + 2B_i d\eta dx^i + (\delta_{ij} + 2H_{(i|j)}) dx^i dx^j \right]$$

$$\sigma_{g,i} \equiv \dot{H}_i - B_i : \text{vector perturbations}$$

$$C_\ell^{\overline{\omega}\overline{\omega}} = \frac{1}{2\pi} \frac{(\ell-1)!}{(\ell+1)!} \int_0^\infty k^2 dk \int_0^{\chi_{\text{CMB}}} d\chi \int_0^{\chi_{\text{CMB}}} d\chi' \\ \times \frac{j_\ell(k\chi)}{\chi} \frac{j_\ell(k\chi')}{\chi'} P_{\sigma_g \sigma_g}(k; \eta_0 - \chi, \eta_0 - \chi')$$

Note : Considering a single light source, we can extract the vector-induced gradient-/curl-modes from the vector-induced shear fields.

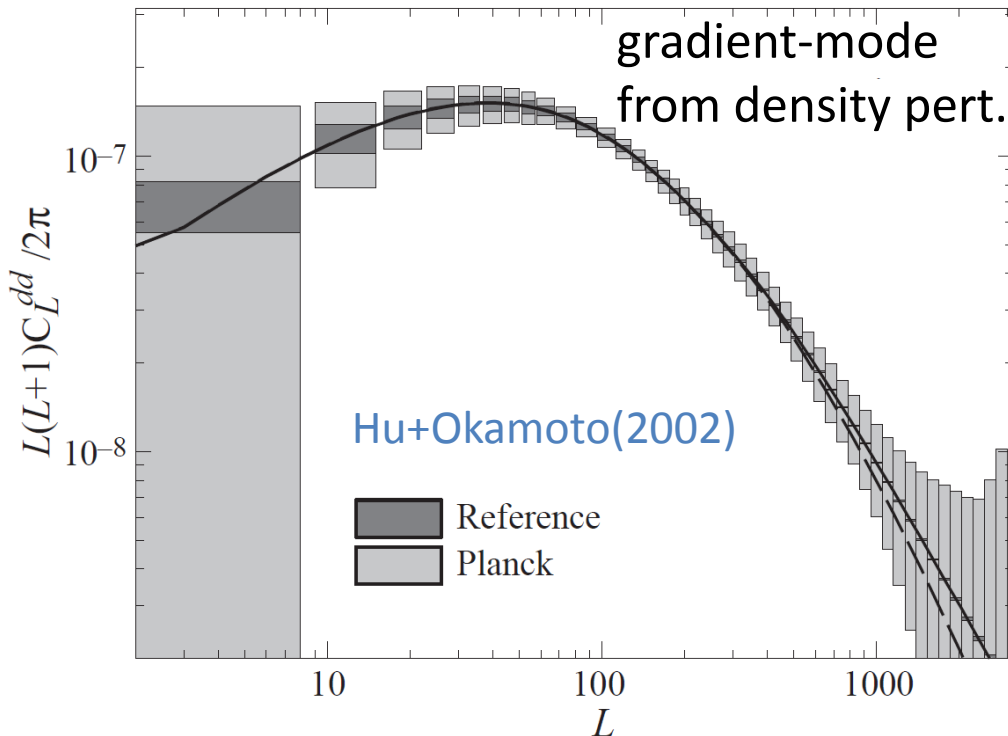
$$C_\ell^{\phi\phi} = 4 \frac{(\ell-2)!}{(\ell+2)!} C_\ell^{\text{EE}}, \quad C_\ell^{\overline{\omega}\overline{\omega}} = 4 \frac{(\ell-2)!}{(\ell+2)!} C_\ell^{\text{BB}}$$

CMB-lensing reconstruction

Lensing induces higher-order term, i.e., **non-Gaussianity**, in measured CMB anisotropies.



- NG can be used to reconstruct the lensing angular power spectrum.

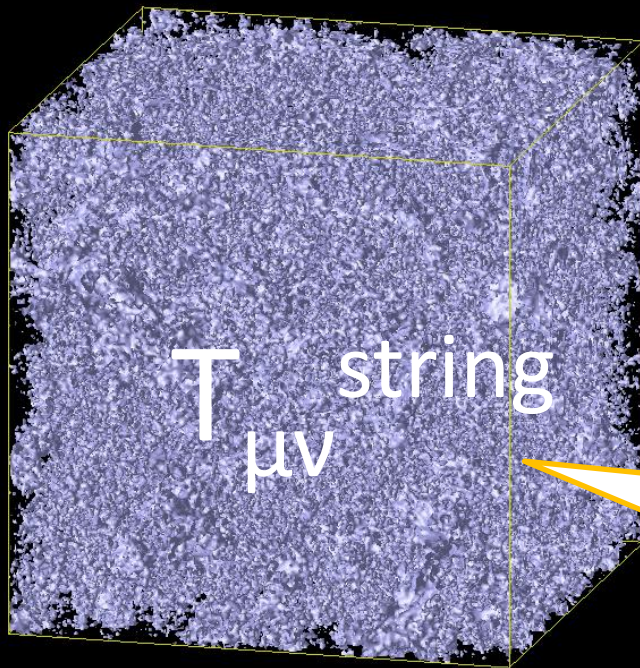


$$\mathbf{d} = \nabla \phi + (*\nabla)\varpi$$

✓ Gradient-mode ϕ
[Hu+Okamoto (2002),...]

✓ Curl-mode ω
[Cooray+ (2005),
Namikawa+DY+Taruya (2011)]

A cosmic string network as a source of vector perturbations



[Hiramatsu, DY +, in prep.]

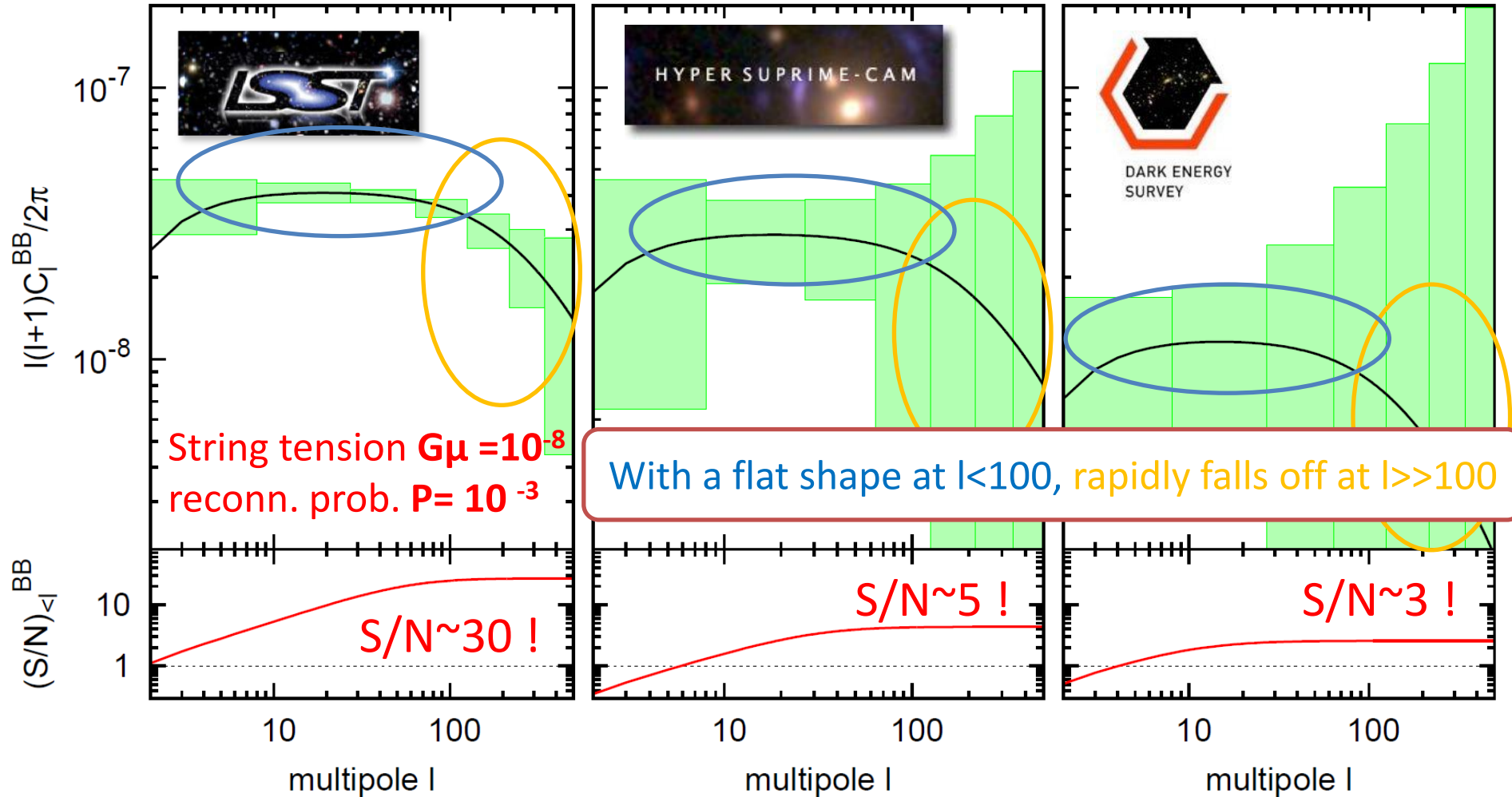
“COSMIC STRINGS” : relics from the early universe and exist in a wide variety of particle-physics models.

The vector metric perturbations, “ $\sigma_{g,i}$ ”, are sourced by the non-vanishing vector mode of string stress-energy tensor.

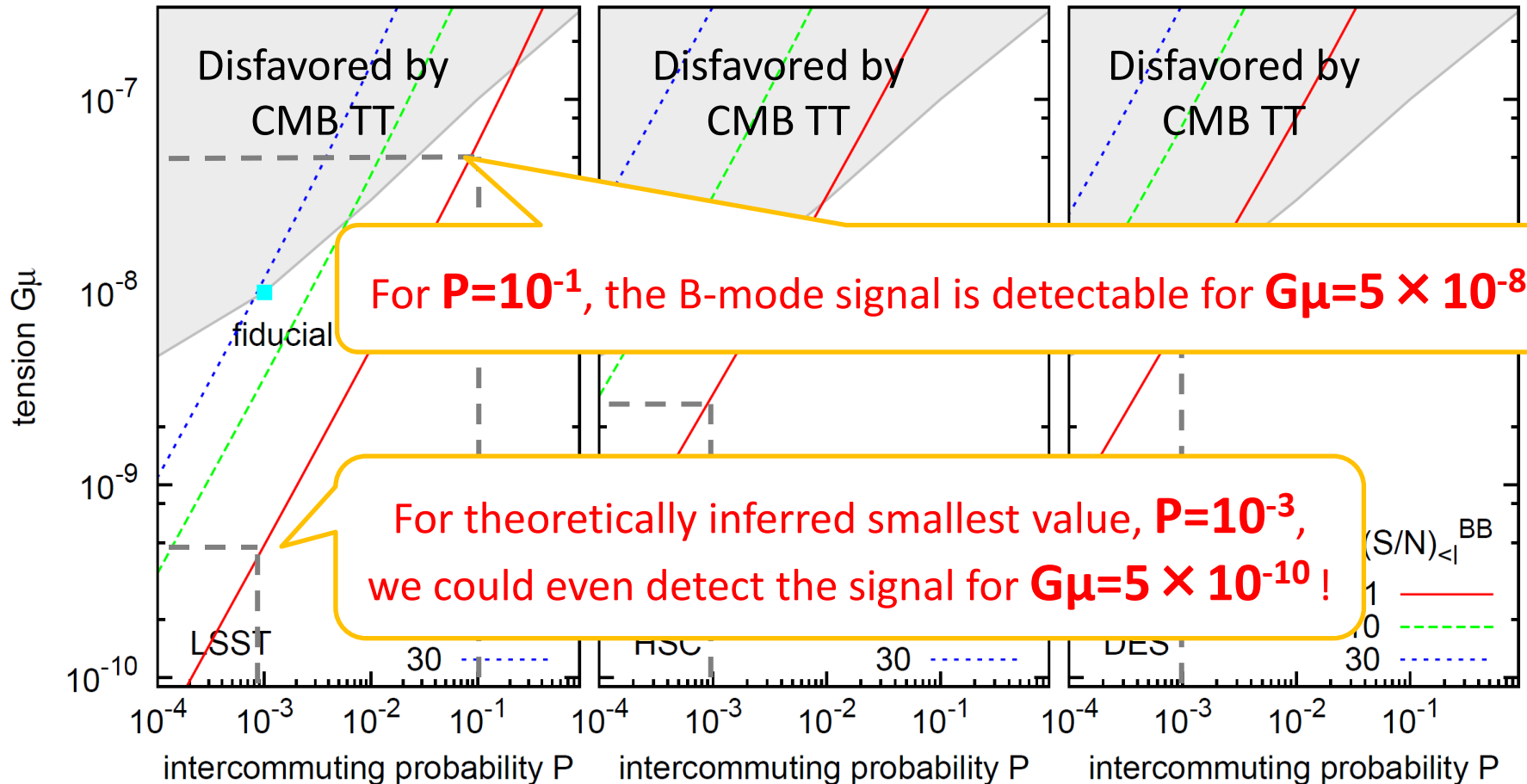
(For convenience, we use the simple analytic model for the network evolution and the correlations.)

B-mode shear from cosmic strings

[DY+Namikawa+Taruya, 1205.2139]

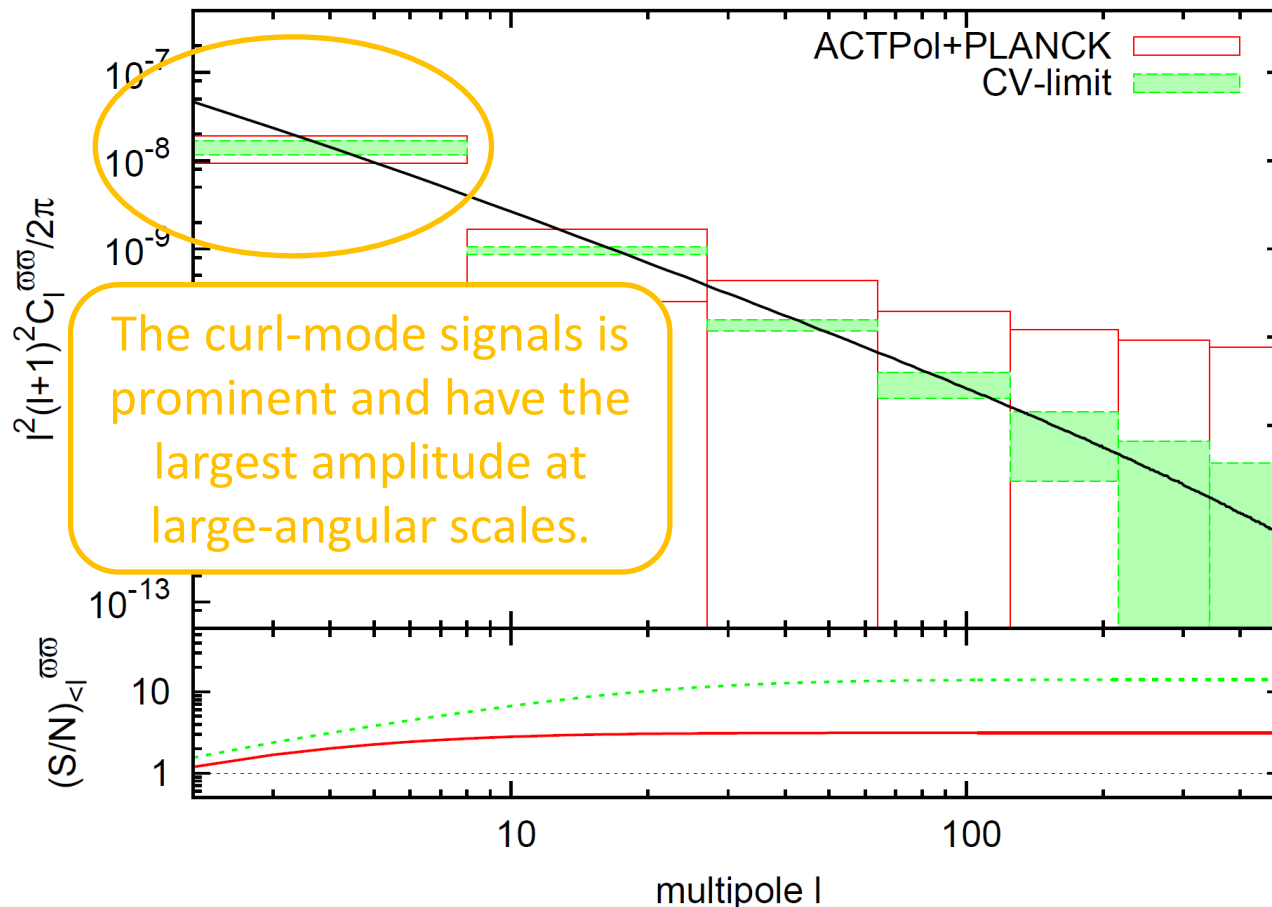


S/N for string-induced B-mode shear



Curl-deflection from cosmic strings

[DY+Namikawa+Taruya,
1205.2139]



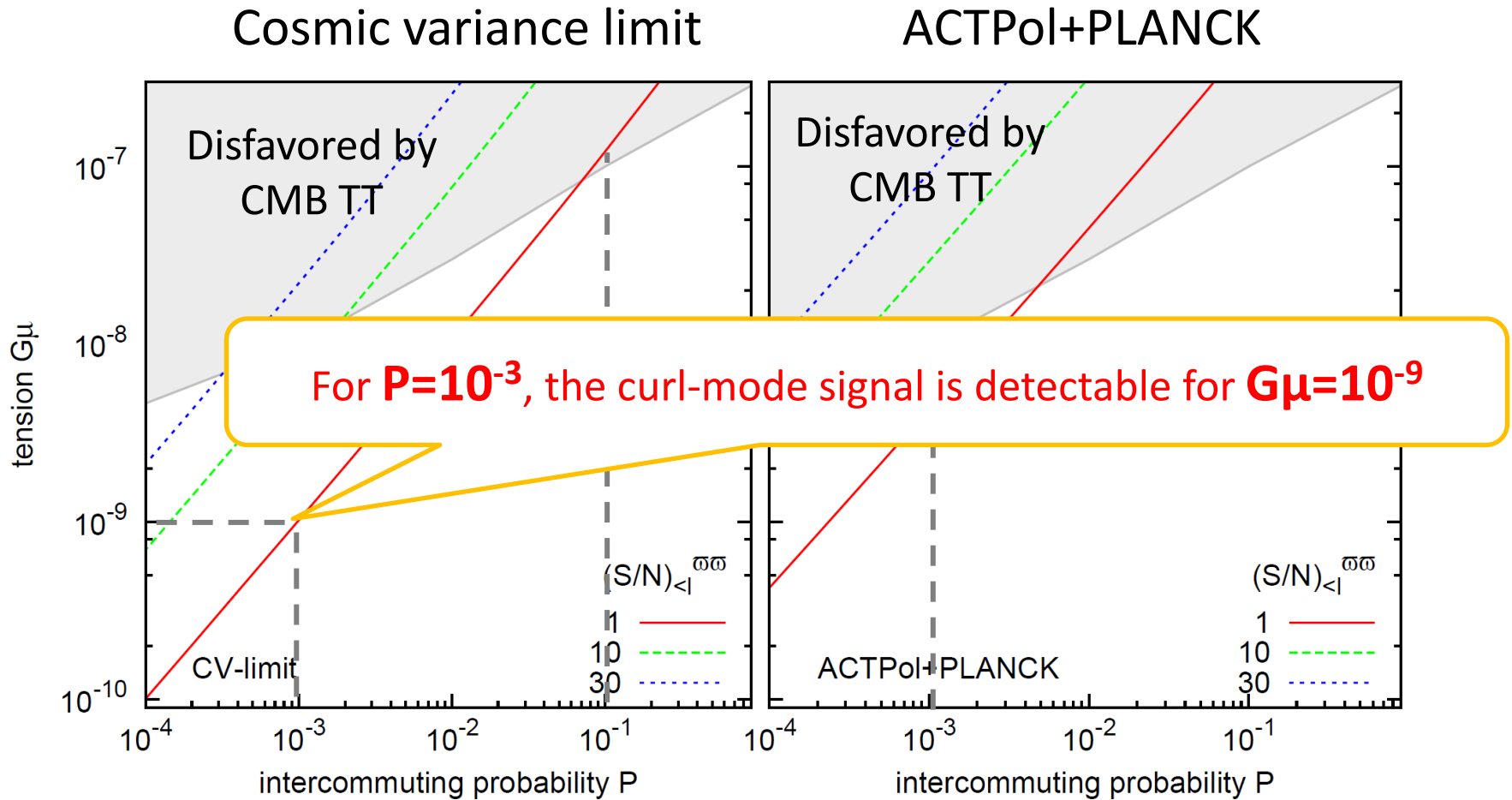
String tension $G\mu = 10^{-8}$
reconn. prob. $P = 10^{-3}$

$S/N \sim 10!$
 $S/N \sim 3!$

The **CURL-MODE** measurement would provide not only a direct probe of cosmic strings, but also a diagnosis helpful to check the systematics in the derived constraints from the CMB TT.

S/N for string-induced curl-deflection

[DY+Namikawa+Taruya, 1205.2139]



Summary

- We study how vector metric perturbations affect weak gravitational lensing and possible signatures.
- As an application, we consider a cosmic string network as a possible source for seeding vector perturbations.
 - ✓ B-mode of shear fields
 - $P < 10^{-1} \rightarrow G\mu > 5 \times 10^{-8}$, $P = 10^{-3} \rightarrow G\mu > 5 \times 10^{-10}$
 - ✓ Curl-mode of deflection angle
 - $P < 10^{-1} \rightarrow G\mu > 4 \times 10^{-9}$, $P = 10^{-3} \rightarrow G\mu > 10^{-9}$
- There are several missing pieces including the contribution of tensor perturbations, the additional correlations, and the other sources