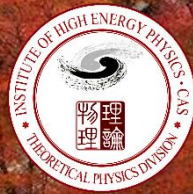
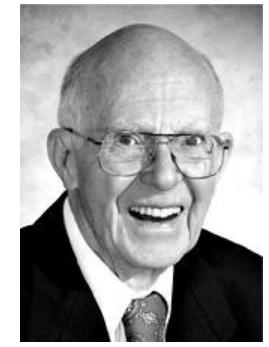
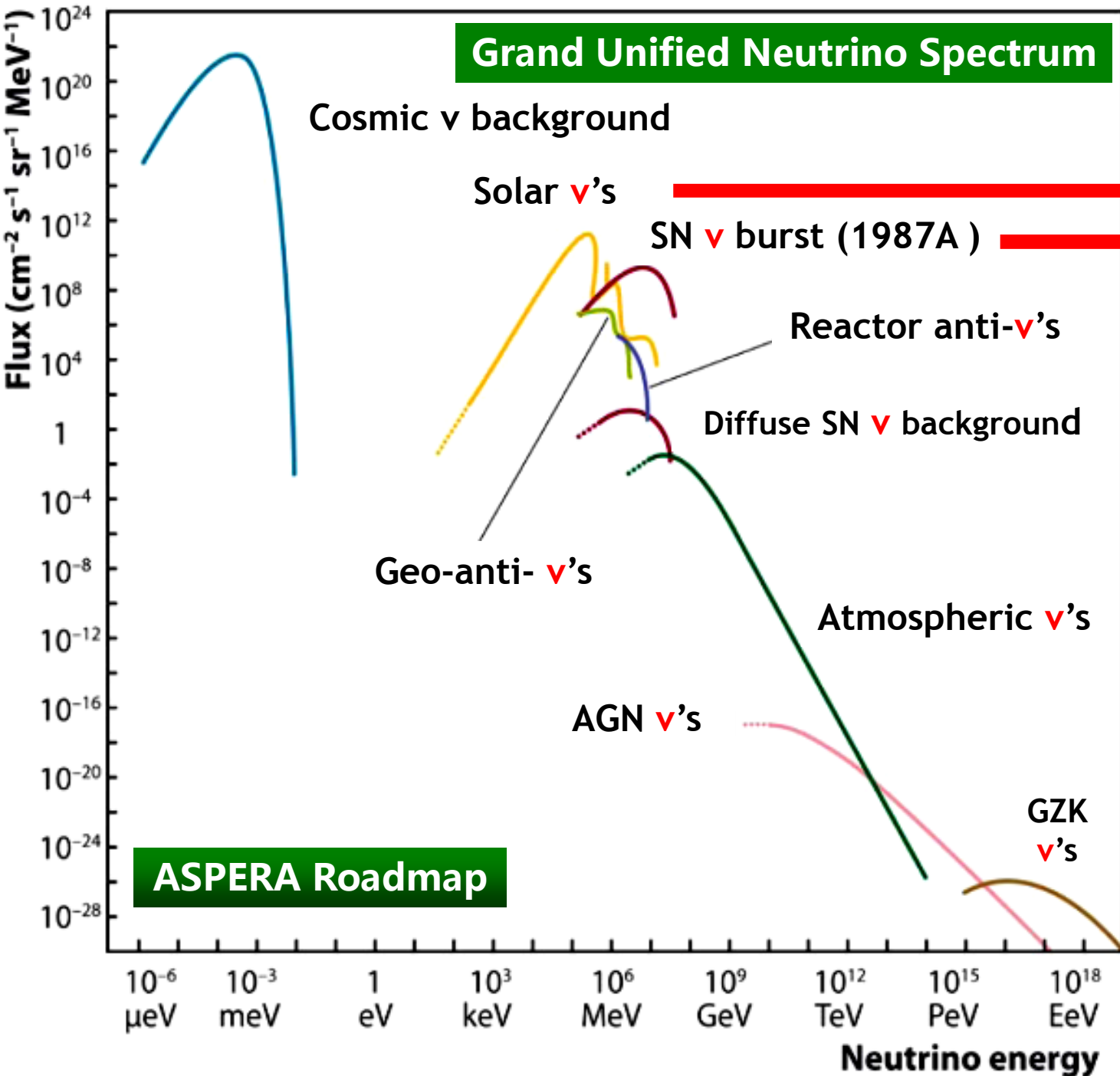


What can we learn from astrophysical neutrinos?

**Shun Zhou
(IHEP, Beijing)**



**CosPA 2017 @ Yukawa Institute for Theoretical Physics, Kyoto University, Kyoto
2017-12-11**

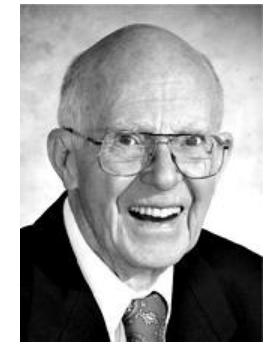
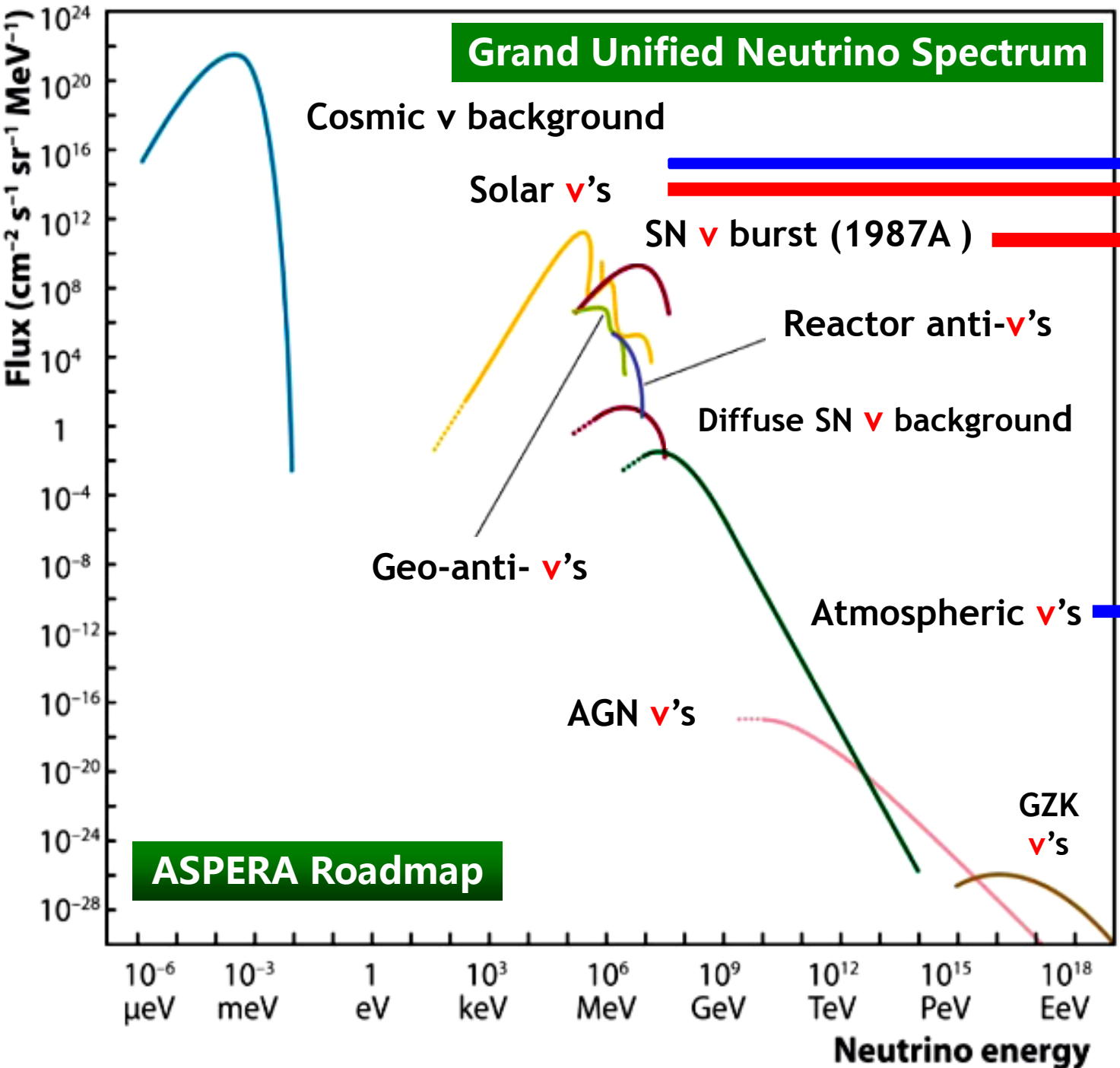


R. Davis Jr.



M. Koshiba

Nobel Prize in 2002 "for the detection of cosmic neutrinos"



R. Davis Jr.



M. Koshiba

Nobel Prize in 2002 "for the detection of **cosmic neutrinos**"

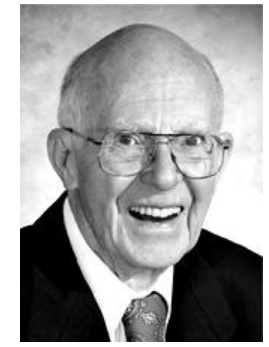
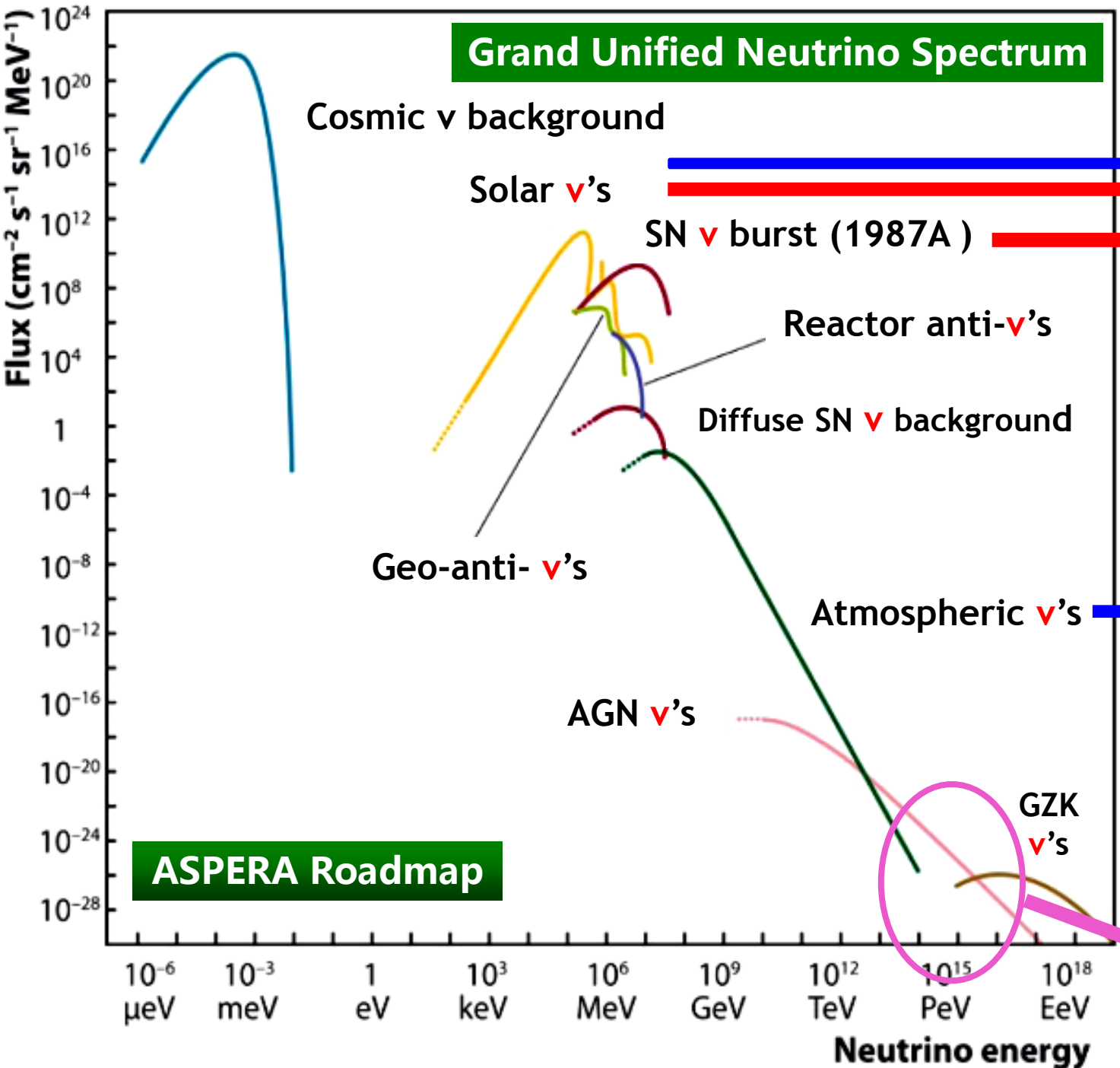


T. Kajita



A.B. McDonald

Nobel Prize in 2015 "for the discovery of neutrino oscillations, which shows that **neutrinos have mass**"



R. Davis Jr.



M. Koshiba

Nobel Prize in 2002 "for the detection of **cosmic neutrinos**"



T. Kajita

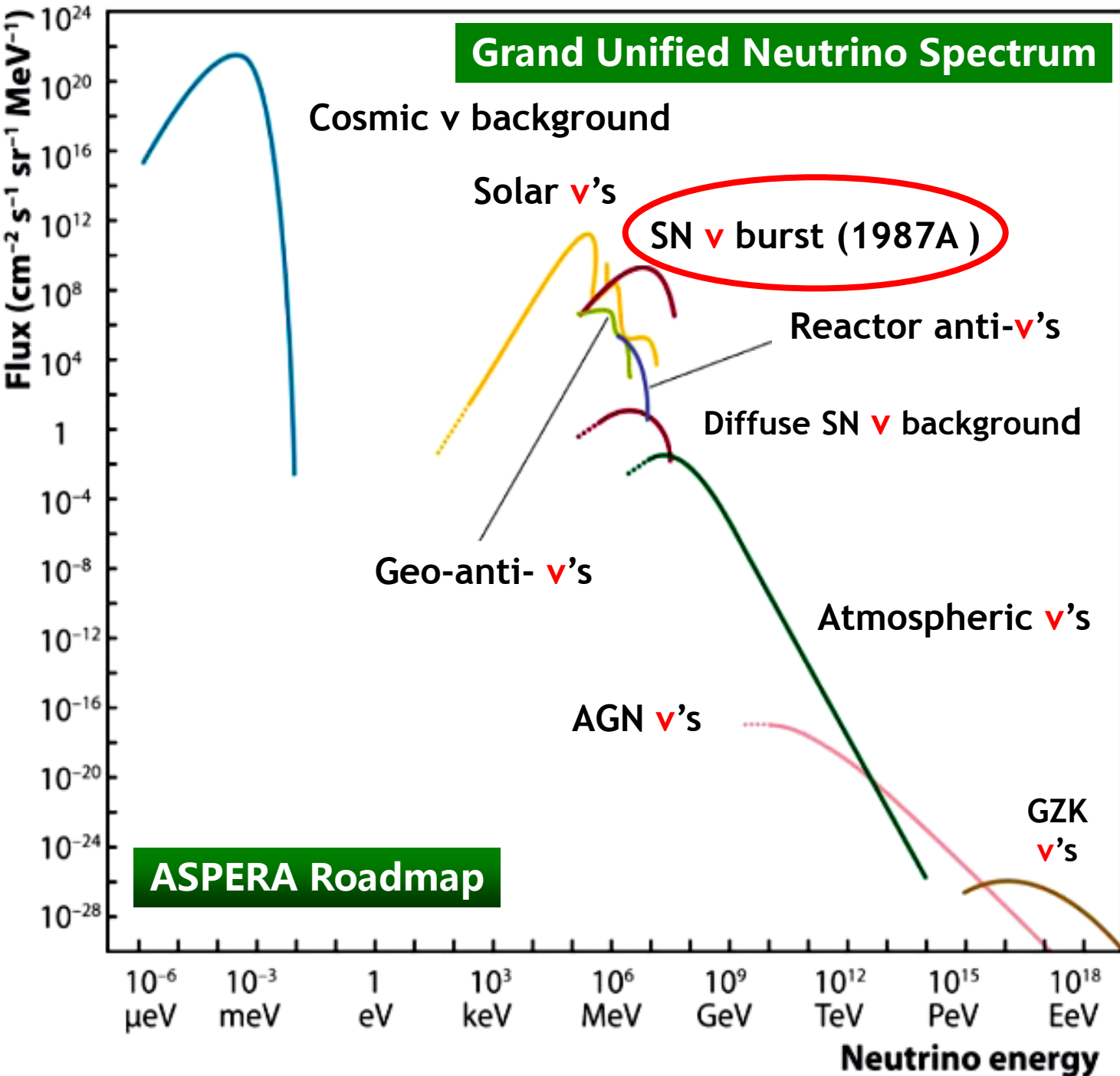


A.B. McDonald

Nobel Prize in 2015 "for the discovery of neutrino oscillations, which shows that **neutrinos have mass**"

Discovery of **PeV** cosmic neutrinos at **IceCube**

a review by **E. Waxman**

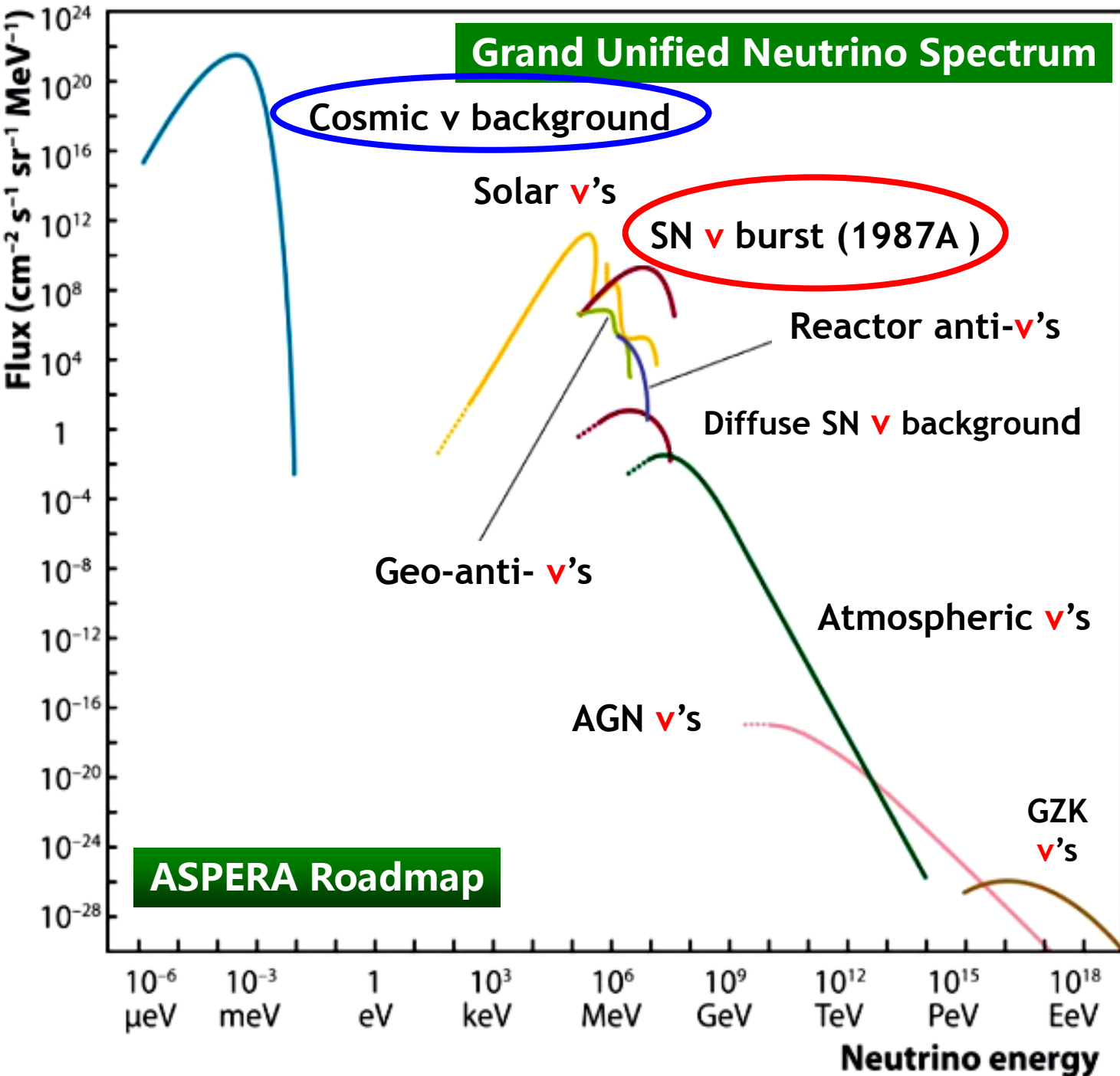


Outline



from the **discovery in 1987** to **precision measurements** of supernova neutrinos

- Supernova (SN) Neutrinos
- Detection of Galactic SN ν 's
- Learn more from SN ν 's

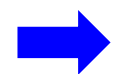


Outline



from the **discovery in 1987** to **precision measurements** of supernova neutrinos

- Supernova (SN) Neutrinos
- Detection of Galactic SN ν 's
- Learn more from SN ν 's



from the **UHE cosmic neutrinos** to **ULE cosmic neutrinos** from Big Bang

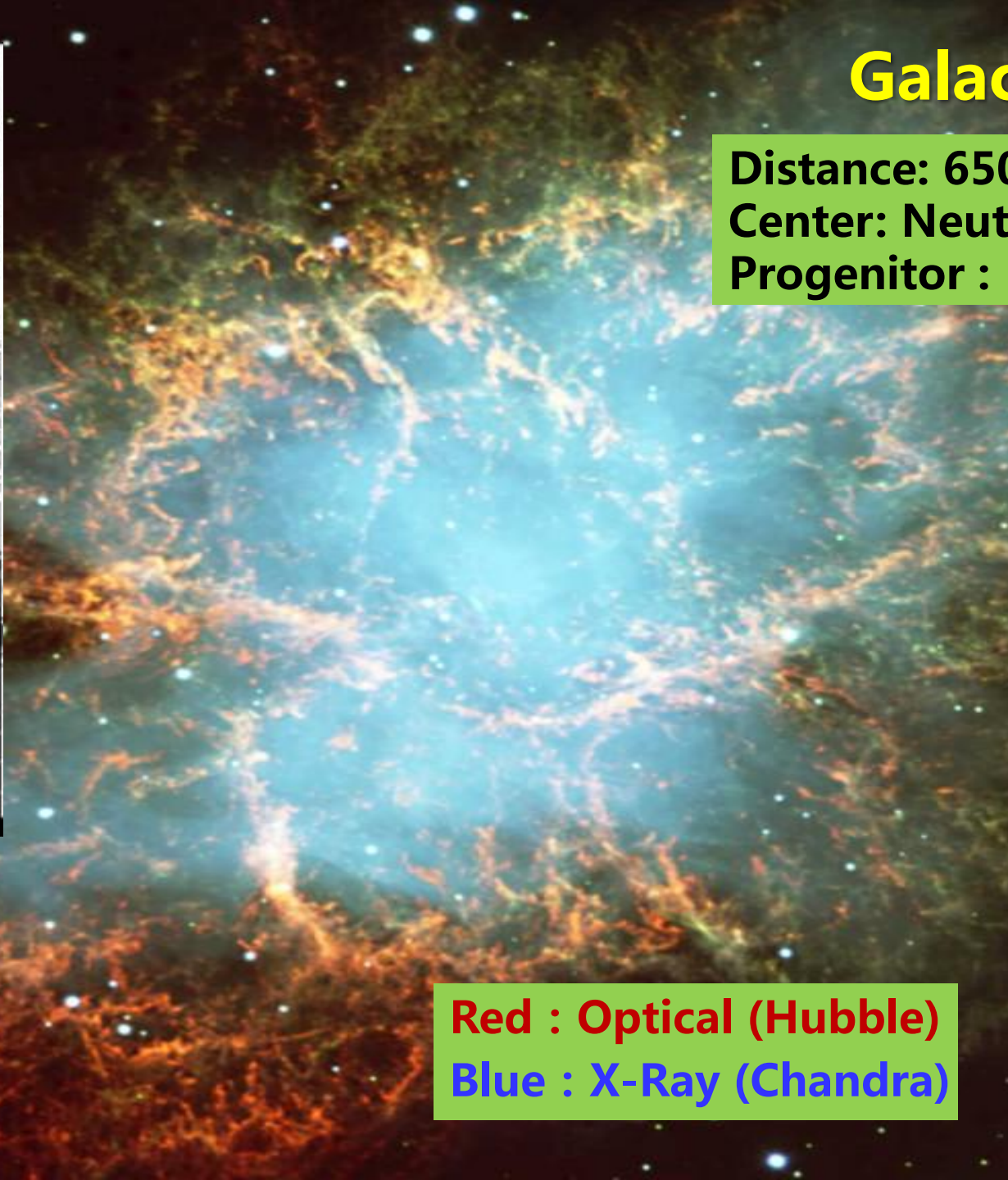
- Cosmic ν Background (CvB)
- Detection of CvB@PTOLEMY
- Physics/Cosmology with CvB

Galactic SN 1054

Distance: 6500 light years (2 kpc)
Center: Neutron Star (R~30 km)
Progenitor: M ~ 10 solar masses

凡十一日没三年三月乙巳出東南方大中祥符四年正月丁丑見南斗魁前天禧五年四月丙辰出軒轅前星西北大如桃速行經軒轅太星入太微垣掩右執法犯次將歷屏星西北凡七十五日入濁没明道元年六月乙巳出東北方近濁有芒彗至丁巳凡十三日没至和元年五月己丑出天關東南可數寸歲餘稍没熙寧二年六月丙辰出箕度中至七月丁卯犯箕乃散三年十一月丁未出天因元祐六年十一月辛亥出參度中犯掩側星壬子犯九游星十二月癸酉入奎至七年三月辛亥乃散紹興八年五月守婁

宋史志卷九



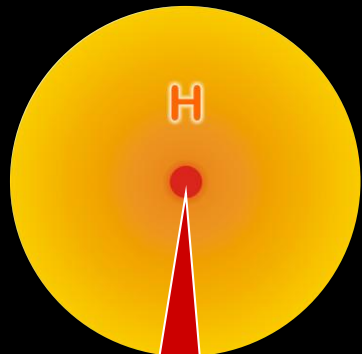
Red : Optical (Hubble)
Blue : X-Ray (Chandra)



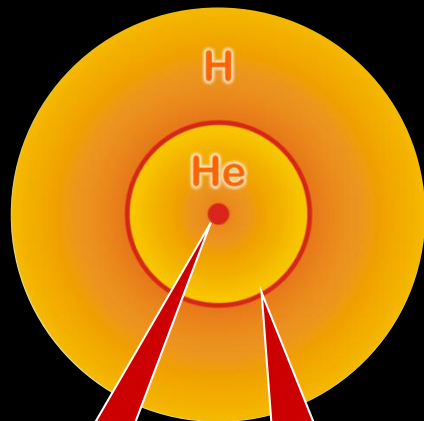
Stellar Collapse and SN Explosion

Main-sequence star Helium-burning star

© G. Raffelt

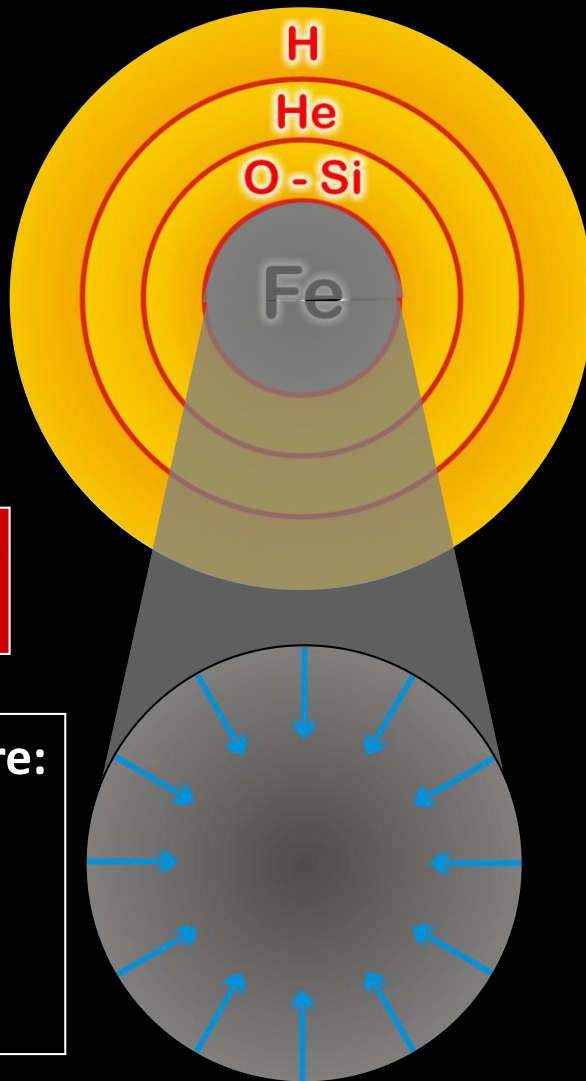


Hydrogen
Burning

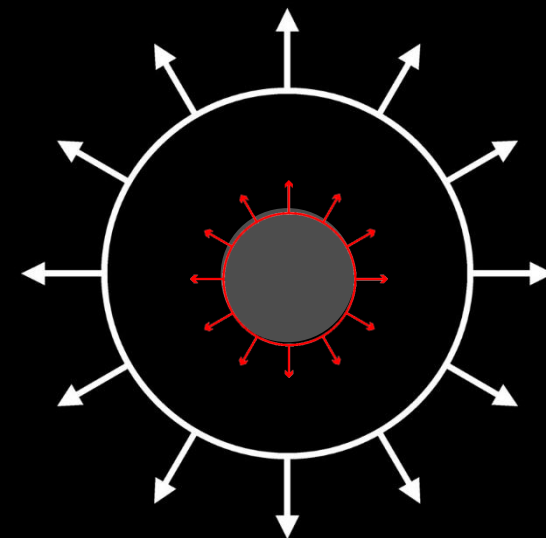


Helium
Burning

Hydrogen
Burning



Grav. binding energy $E_b \approx 3 \times 10^{53}$ erg
99% Neutrinos
1% Kinetic energy of explosion
(1% of this into cosmic rays)
0.01% Photons, outshine host galaxy



1. > 8...10 Solar Mass
2. Collapse → Bounce
3. Shock wave halted
4. ν energy deposited
5. Final SN explosion

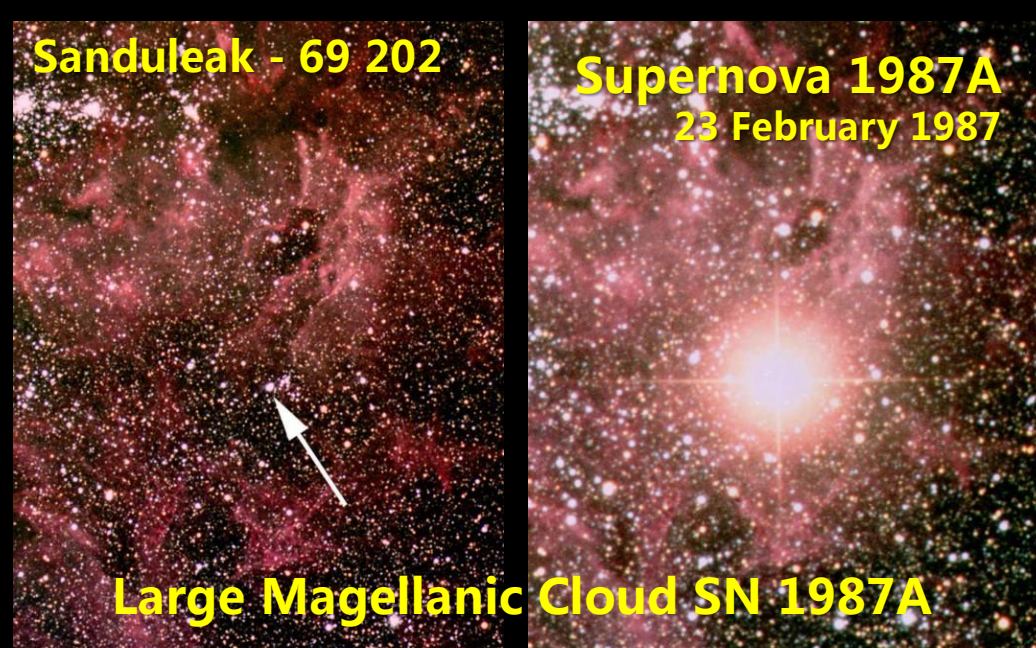
Degenerate iron core:
 $\rho \approx 10^9 \text{ g cm}^{-3}$
 $T \approx 10^{10} \text{ K}$
 $M_{\text{Fe}} \approx 1.5 M_{\text{sun}}$
 $R_{\text{Fe}} \approx 8000 \text{ km}$

Proto-Neutron star:
 $\rho \sim \rho_{\text{nuc}} = 3 \times 10^{14} \text{ g cm}^{-3}$
 $T \sim 30 \text{ MeV}$

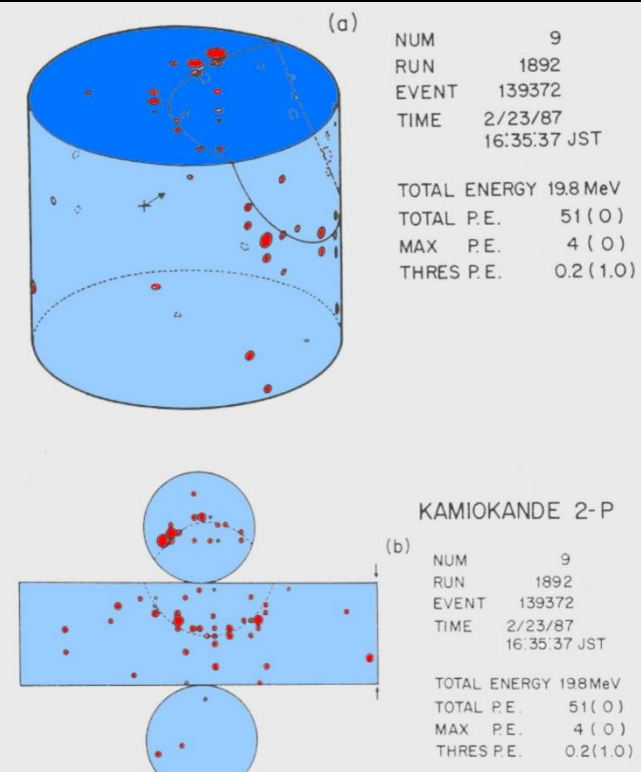
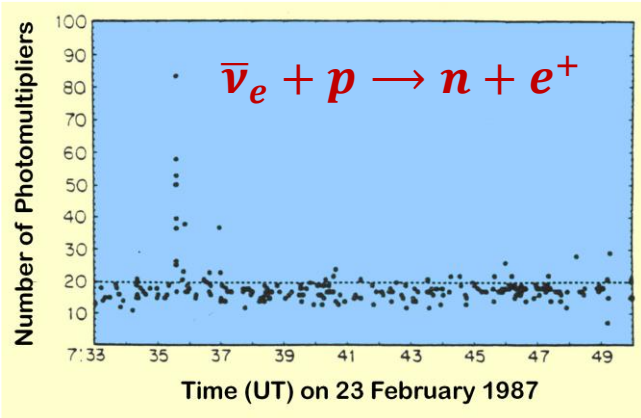
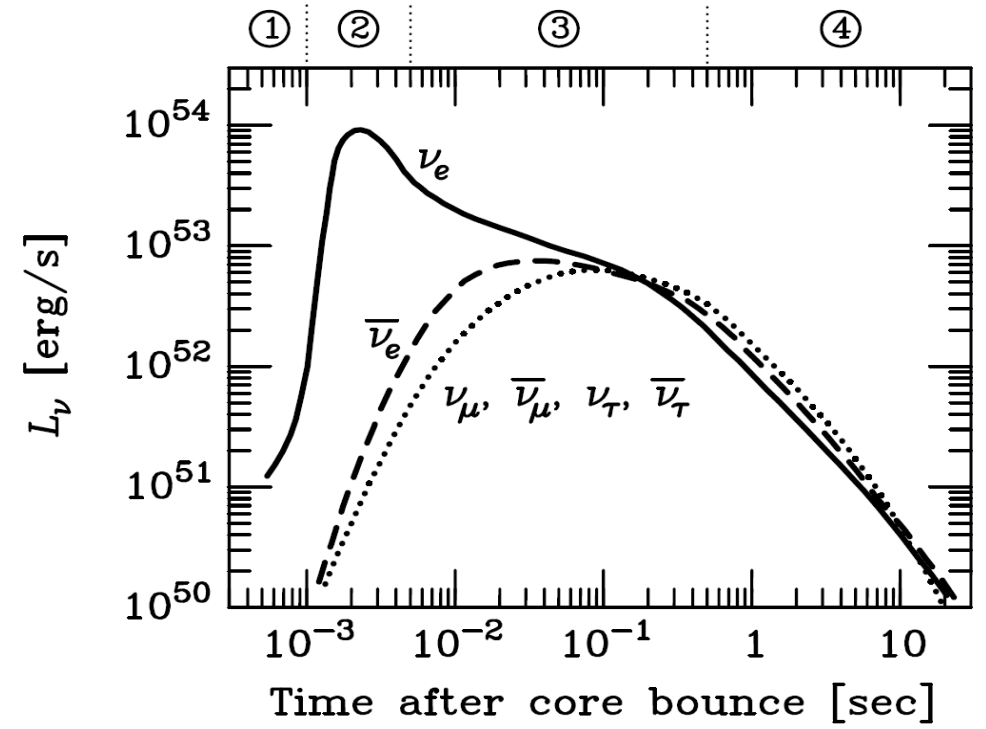
Neutrinos from SN 1987A

Sanduleak - 69 202

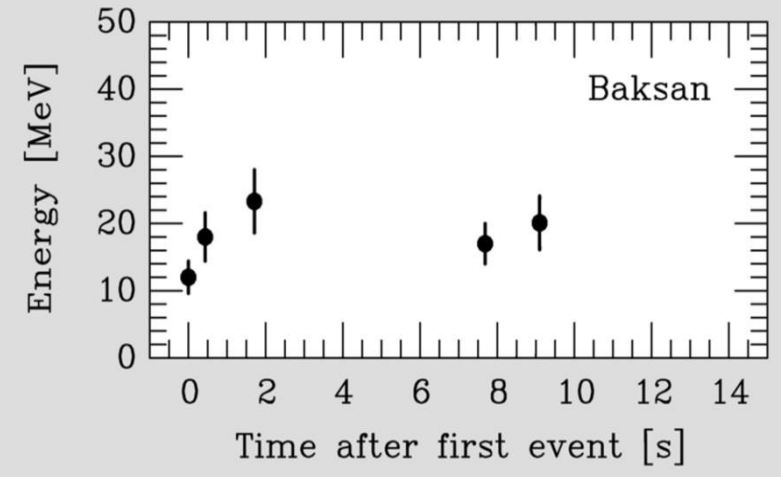
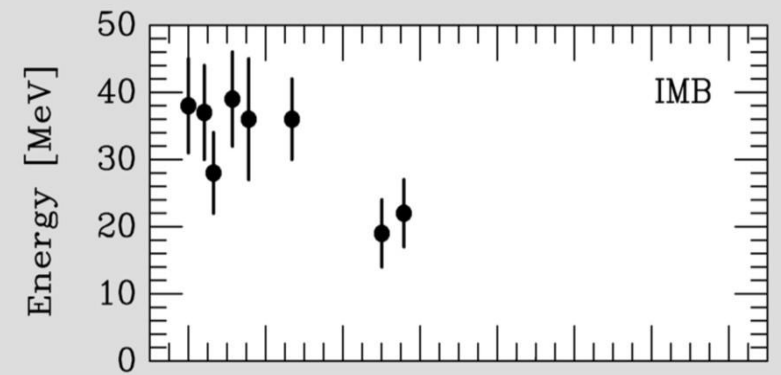
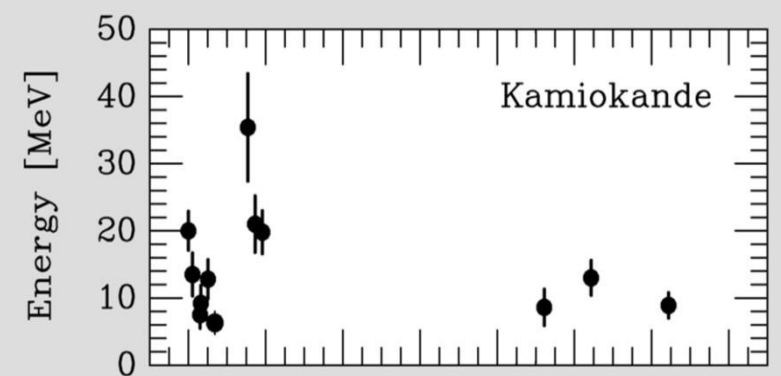
Supernova 1987A
23 February 1987



Large Magellanic Cloud SN 1987A



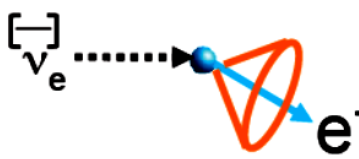
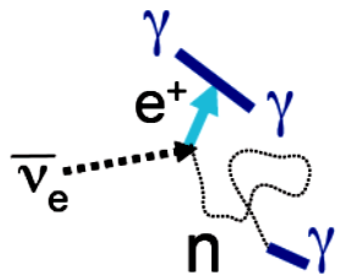
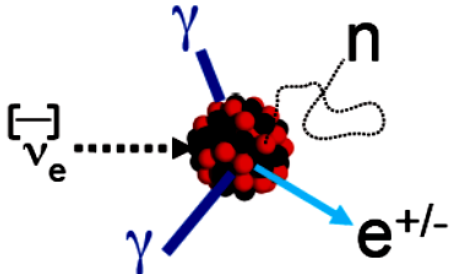
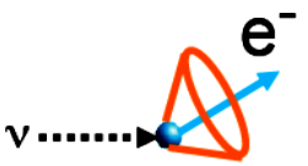
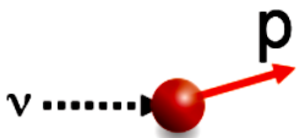
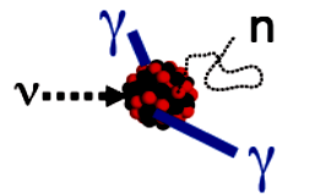
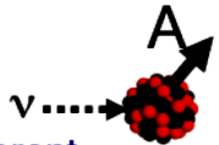
Hirata et al., PRD 38 (1988) 448



Detection of SN Neutrinos

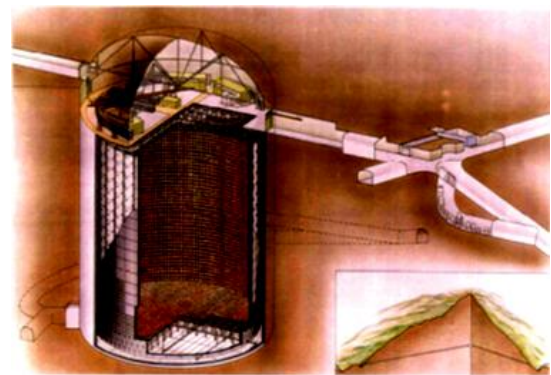
Supernova-relevant neutrino interactions

Scholberg,
@SNOBS 2017

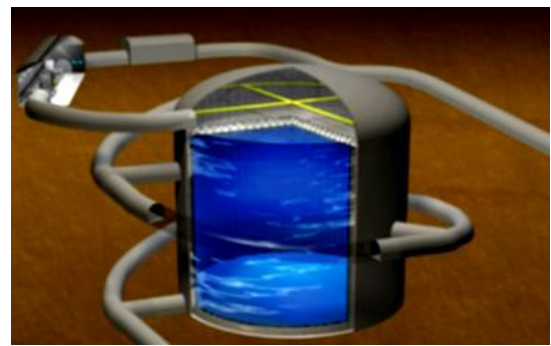
	Electrons	Protons	Nuclei
Charged current	<p>Elastic scattering</p> $\nu + e^- \rightarrow \nu + e^-$ 	<p>Inverse beta decay</p> $\bar{\nu}_e + p \rightarrow e^+ + n$ 	$\nu_e + (N, Z) \rightarrow e^- + (N - 1, Z + 1)$ $\bar{\nu}_e + (N, Z) \rightarrow e^+ + (N + 1, Z - 1)$ 
Neutral current	 <p>Useful for pointing</p>	<p>Elastic scattering</p>  <p>very low energy recoils</p>	$\nu + A \rightarrow \nu + A^*$  <p>Coherent elastic (CEvNS)</p> 

Various possible ejecta and deexcitation products

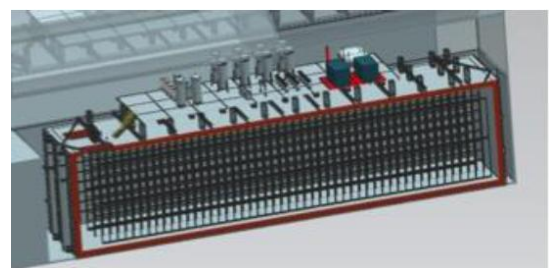
Super-Kamiokande



Hyper-Kamiokande



LAr-TPC DUNE



Summary of SN Neutrino Detectors

Detector	Type	Location	Mass (kton)	Events @ 10 kpc	Status
Super-K	Water	Japan	32	8000	Running
LVD	Scintillator	Italy	1	300	Running
KamLAND	Scintillator	Japan	1	300	Running
Borexino	Scintillator	Italy	0.3	100	Running
IceCube	Long string	South Pole	(600)	(10 ⁶)	Running
Baksan	Scintillator	Russia	0.33	50	Running
HALO	Lead	Canada	0.079	20	Running
Daya Bay	Scintillator	China	0.33	100	Running
NOvA	Scintillator	USA	15	3000	Running
MicroBooNE	Liquid argon	USA	0.17	17	Running
SNO+	Scintillator	Canada	1	300	Under construction
DUNE	Liquid argon	USA	40	3000	Future
Hyper-K	Water	Japan	540	110,000	Future
JUNO	Scintillator	China	20	6000	Under construction
PINGU/GEN-2	Long string	South pole	(600)	(10 ⁶)	Future

Scholberg @SNOBS 2017

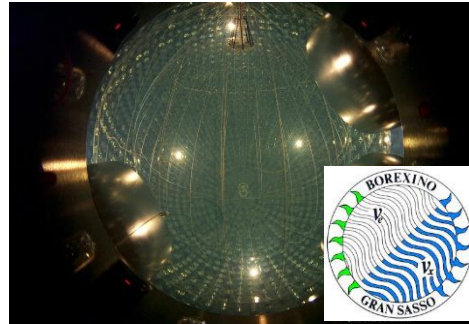
Current & Future Scintillator-based Detectors



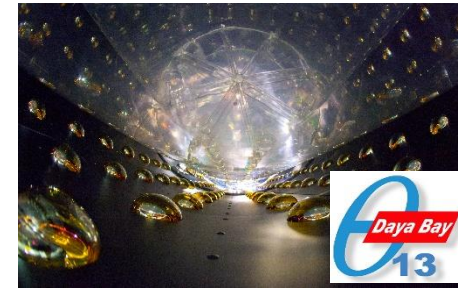
LVD, 1 kt



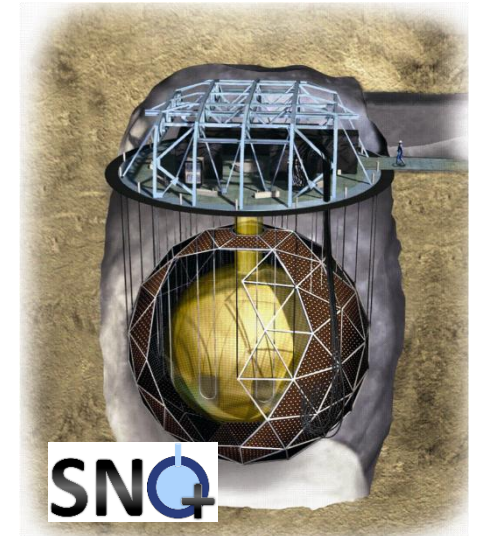
KamLAND, 1 kt



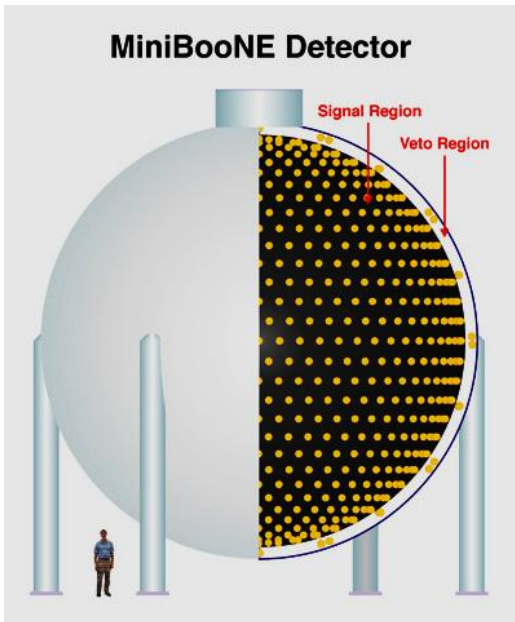
Borexino, 0.3 kt



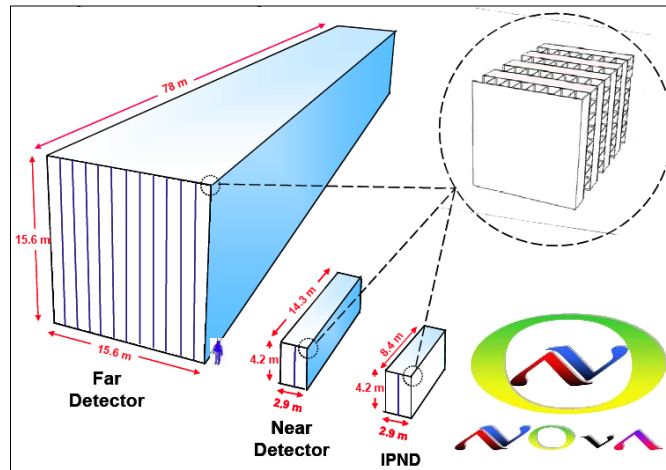
Daya Bay, 0.16 kt



SNO+, 1 kt



MiniBooNE, 0.7 kt



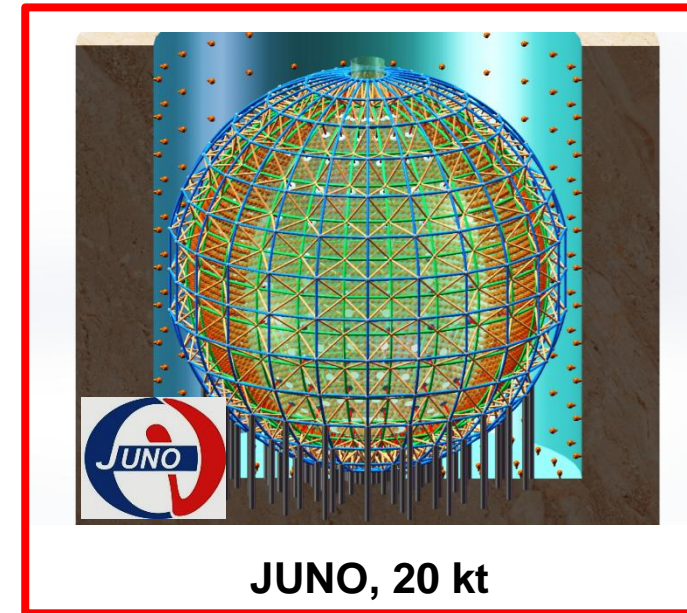
NOvA, 15 kt



Baksan, 0.33 kt



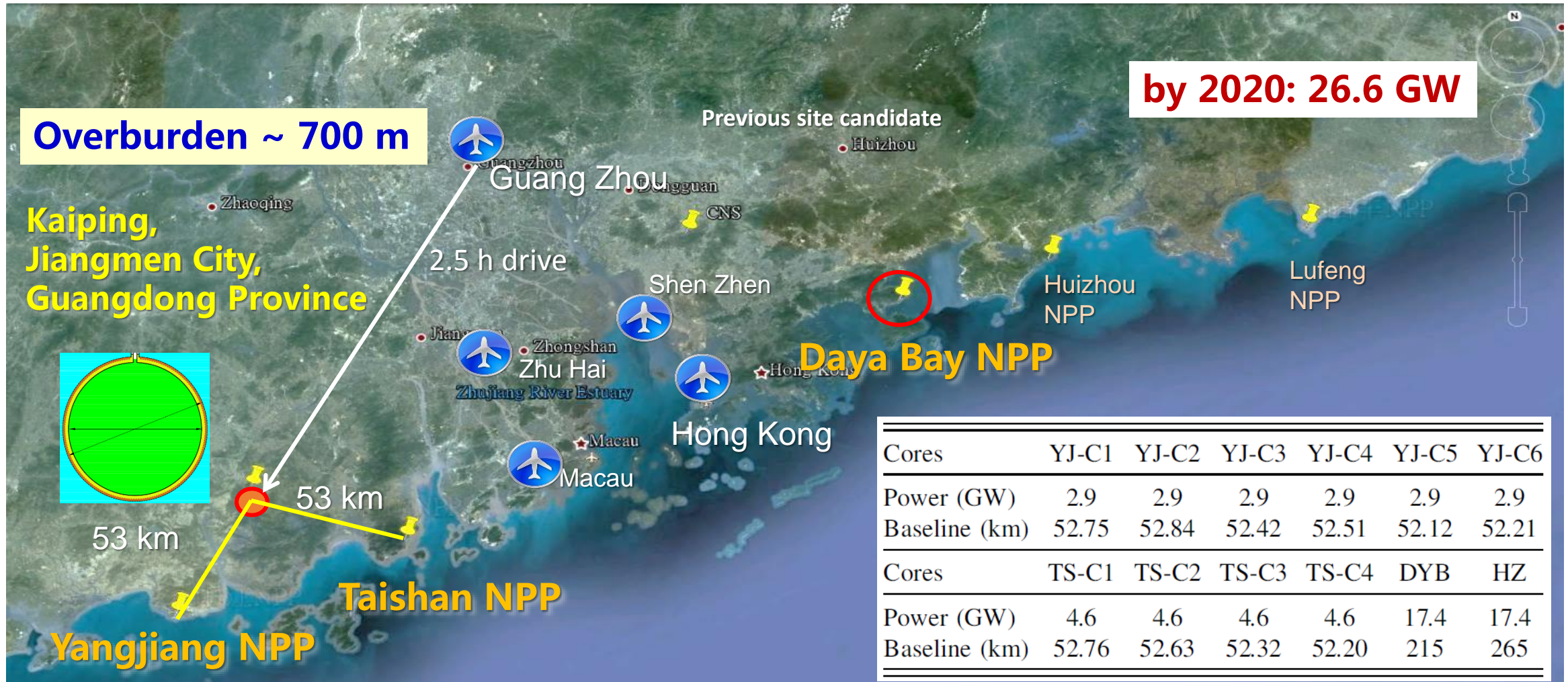
LENA, 50 kt



JUNO, 20 kt

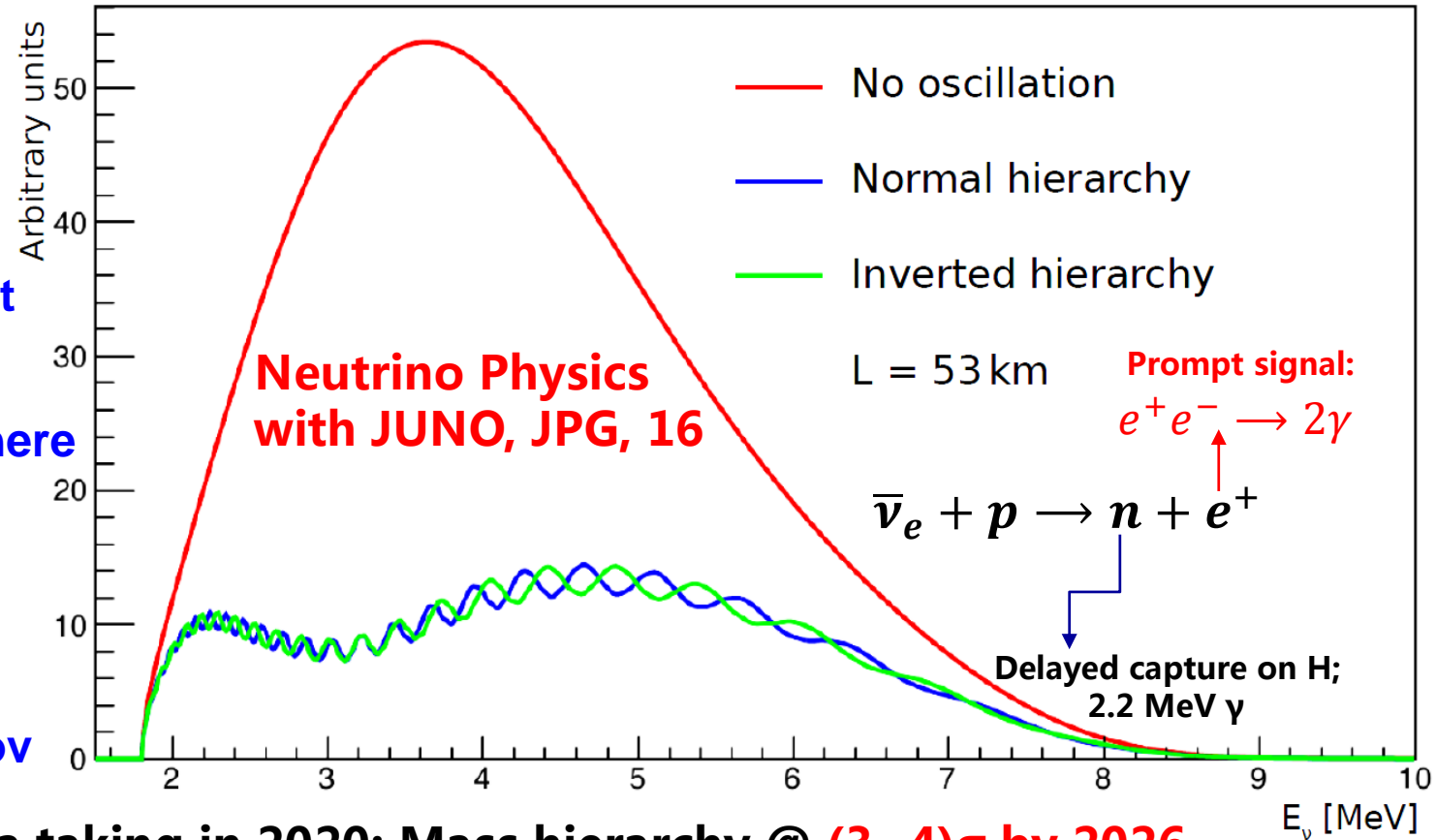
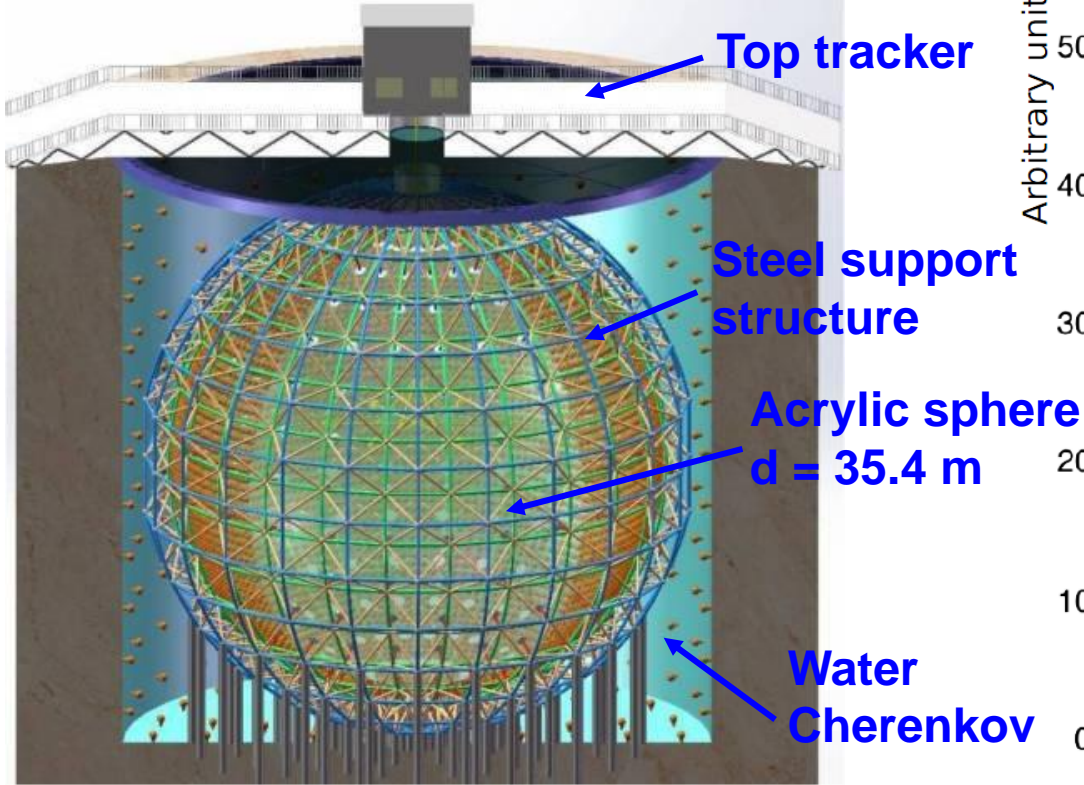
The JUNO Experiment

NPP	Daya Bay	Huizhou	Lufeng	Yangjiang	Taishan
Status	Operational	Planned	Planned	Under construction	Under construction
Power	17.4 GW	17.4 GW	17.4 GW	17.4 GW	18.4 GW



The JUNO Experiment

Jiangmen **U**nderground **N**eutrino **O**bservatory

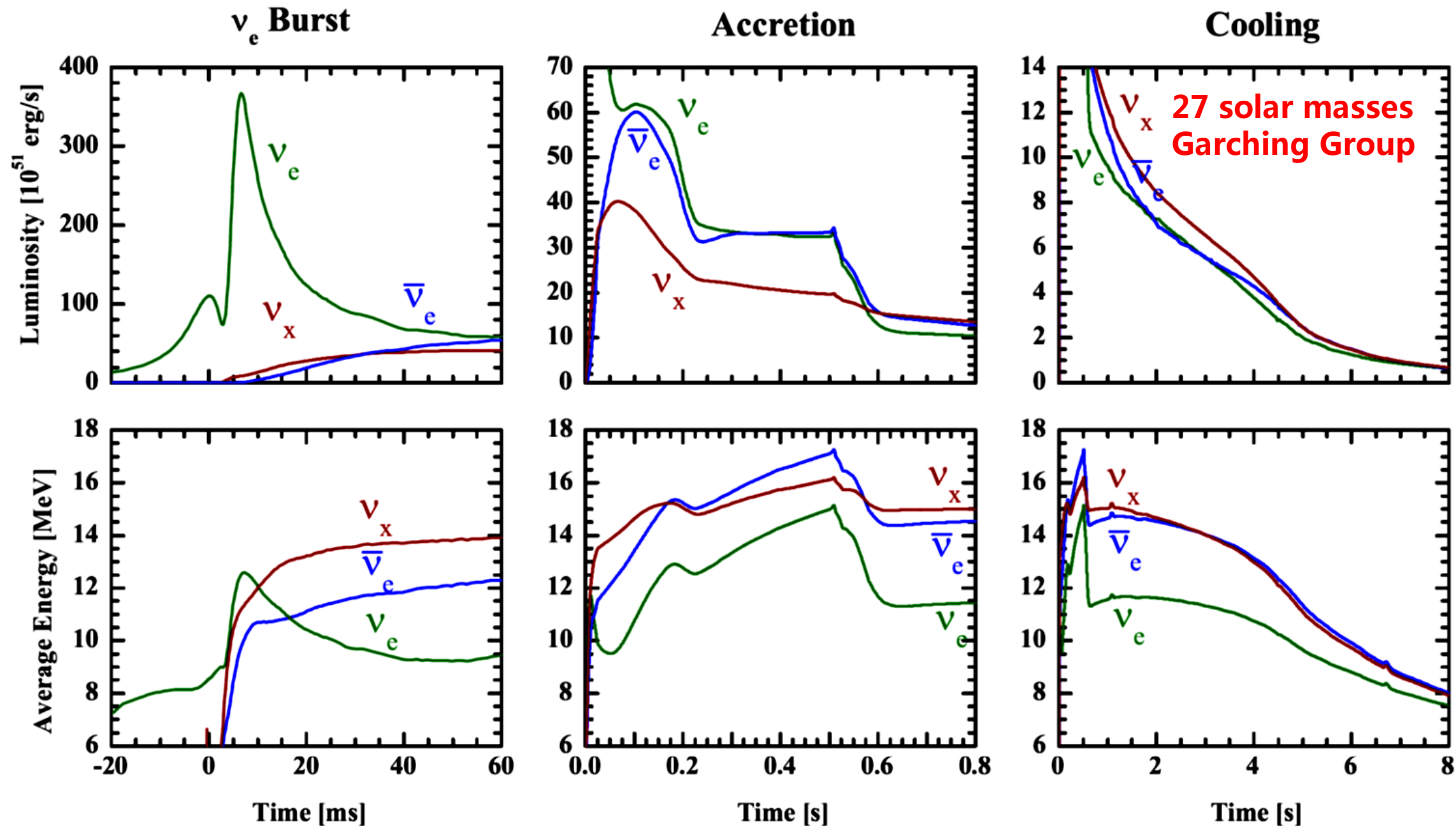


- 20 kiloton LS detector
- 3% energy resolution@ 1 MeV
- 700 m underground
- 18,000 20" + 25,000 3" PMTs
- 53 km to the NPPs

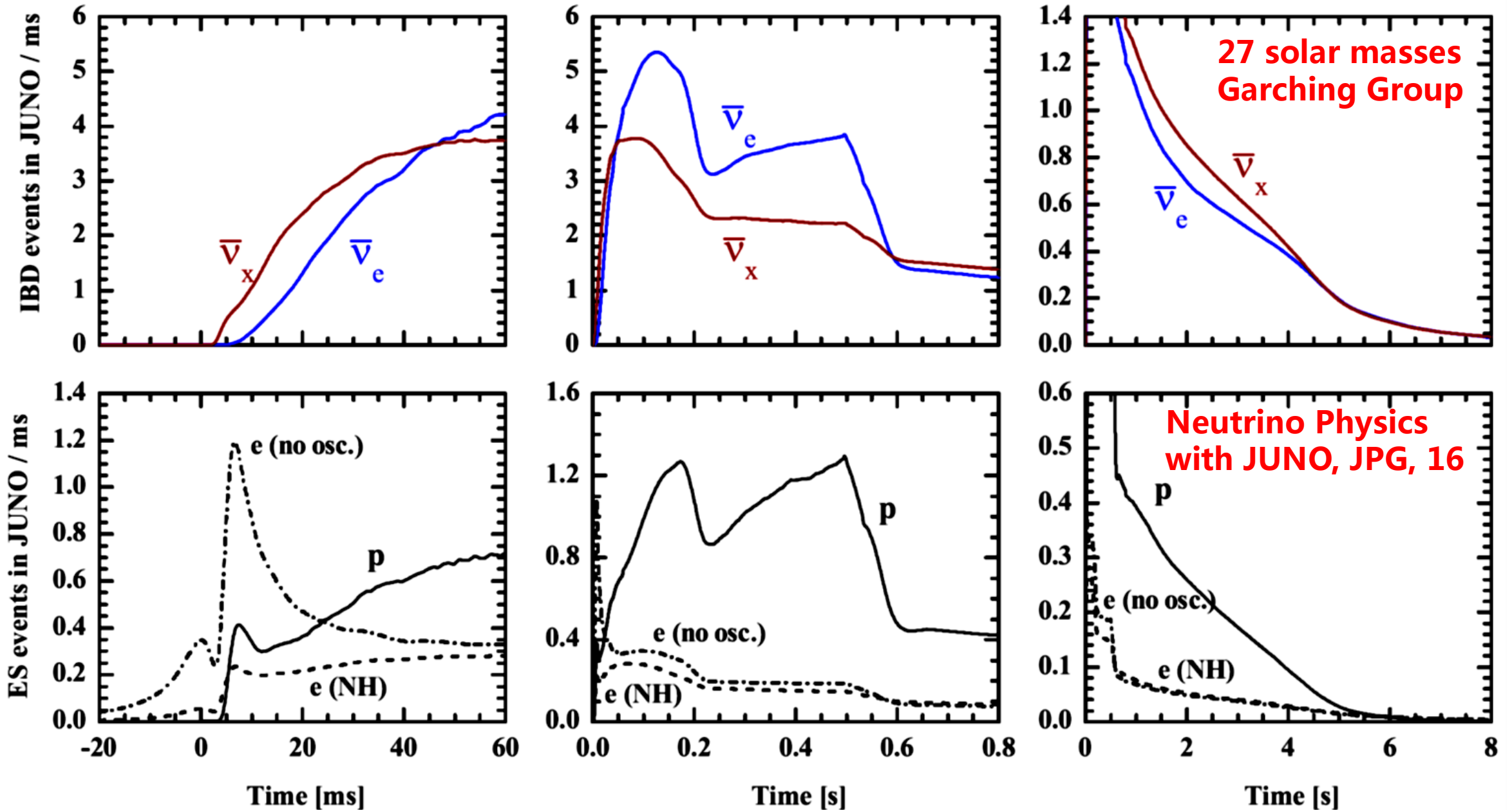
Data taking in 2020; Mass hierarchy @ $(3\sim 4)\sigma$ by 2026

	KamLAND	Borexino	JUNO
LS mass	1 kt	0.3 kt	20 kt
Energy Resolution	$6\%/\sqrt{E}$	$5\%/\sqrt{E}$	$3\%/\sqrt{E}$
Light yield	250 p.e./MeV	511 p.e./MeV	1200 p.e./MeV

SN Models & Neutrino Spectra



SN Neutrino Event Rates

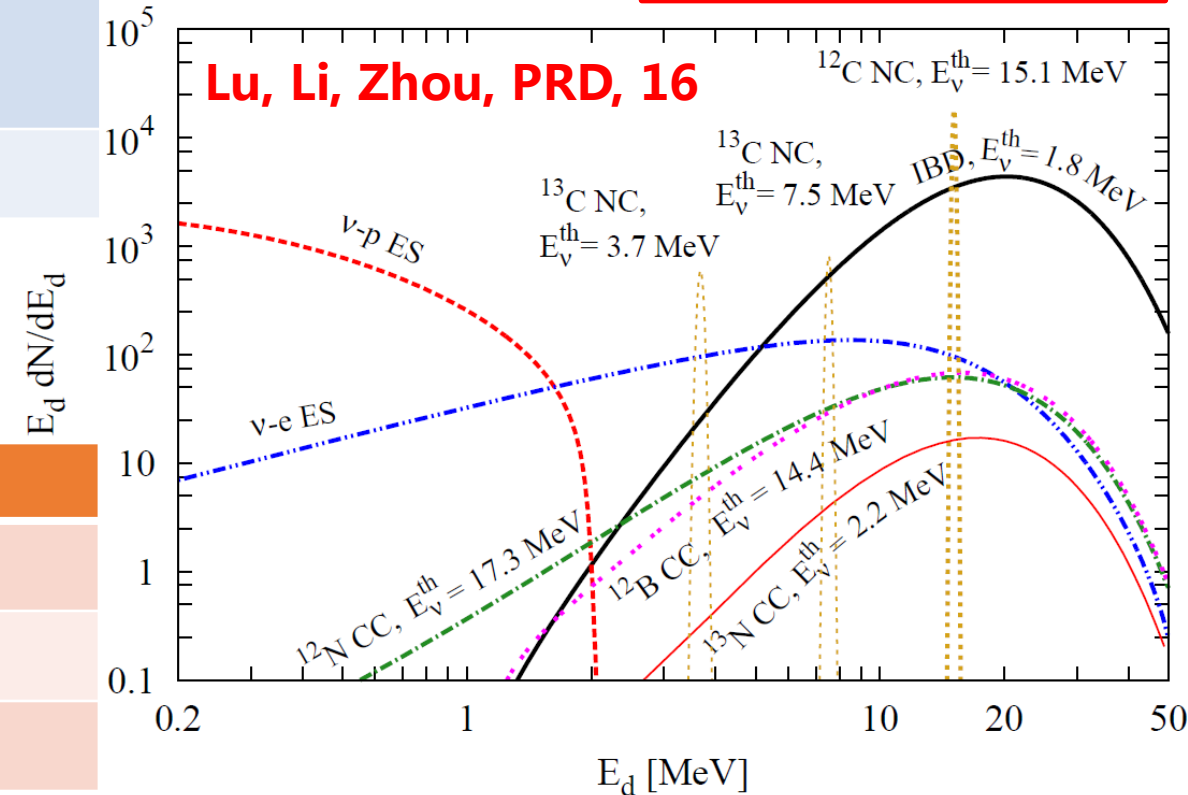


Reaction channel	Interaction type	Sensitive to
$\bar{\nu}_e + p \rightarrow e^+ + n$	CC	$\bar{\nu}_e$
$\nu + p \rightarrow \nu + p$	NC	ν_x
$\nu + e^- \rightarrow \nu + e^-$	CC+NC	ν_e
$\bar{\nu}_e + {}^{12}\text{C} \rightarrow e^+ + {}^{12}\text{B}$ (14.39 MeV, 20 ms)	CC	$\bar{\nu}_e$
$\nu_e + {}^{12}\text{C} \rightarrow e^- + {}^{12}\text{N}$ (17.34 MeV, 11 ms)	CC	ν_e
$\nu + {}^{12}\text{C} \rightarrow \nu + {}^{12}\text{C}^*$	NC	ν_x

- Elastic ν -p scattering important
- Advantage of LS: low threshold
Beacom, Farr, Vogel, PRD, 02;
Dasgupta, Beacom, PRD, 11

Event Spectra @ JUNO

Average E ($\nu_e, \bar{\nu}_e, \nu_x$)
(12, 14, 16) MeV



Natural abundance of ${}^{13}\text{C}$ is about 1.1%

Fukugita *et al.*, PLB, 90; Suzuki *et al.*, PRD, 12

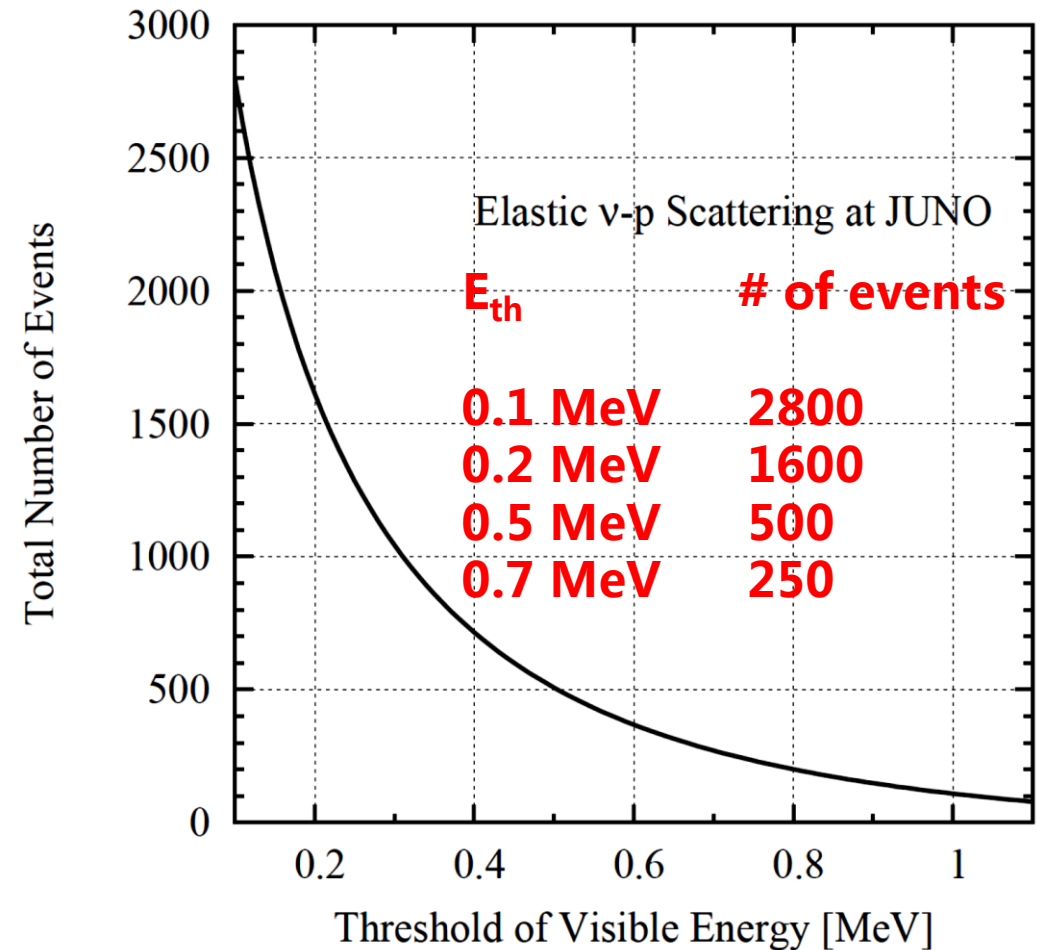
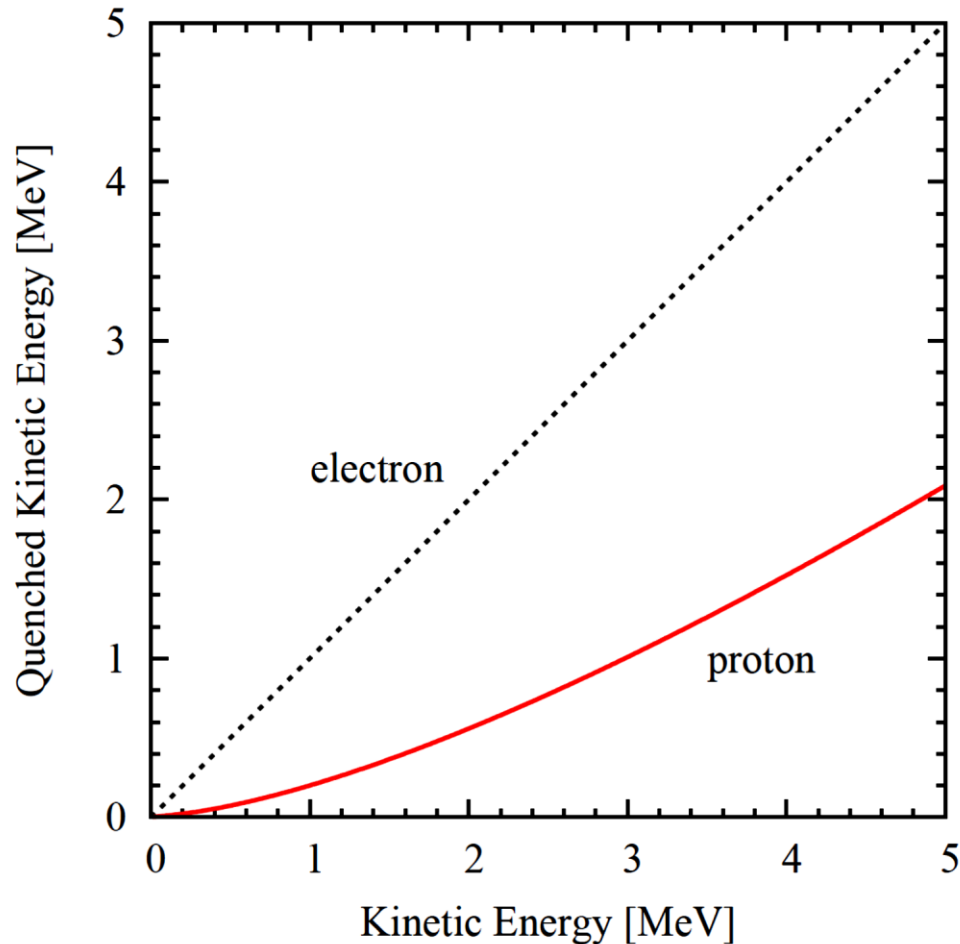
Reaction channel	Interaction type	Sensitive to
$\bar{\nu}_e + {}^{13}\text{C} \rightarrow e^+ + {}^{13}\text{B}$	CC	$\bar{\nu}_e$
$\nu_e + {}^{13}\text{C} \rightarrow e^- + {}^{13}\text{N}$	CC	ν_e
$\nu + {}^{13}\text{C} \rightarrow \nu + {}^{13}\text{C}^*$	NC	ν_x

Elastic ν -p Scattering in LS Detectors

Beacom, Farr, Vogel, PRD, 02; Dasgupta, Beacom, PRD, 11; Lu, Li, Zhou, PRD, 16

Quenching effects on the proton recoil energy $T_p \leq 2 E^2/m_p$

- Elastic ν -p scattering important
- Advantage of LS: low threshold



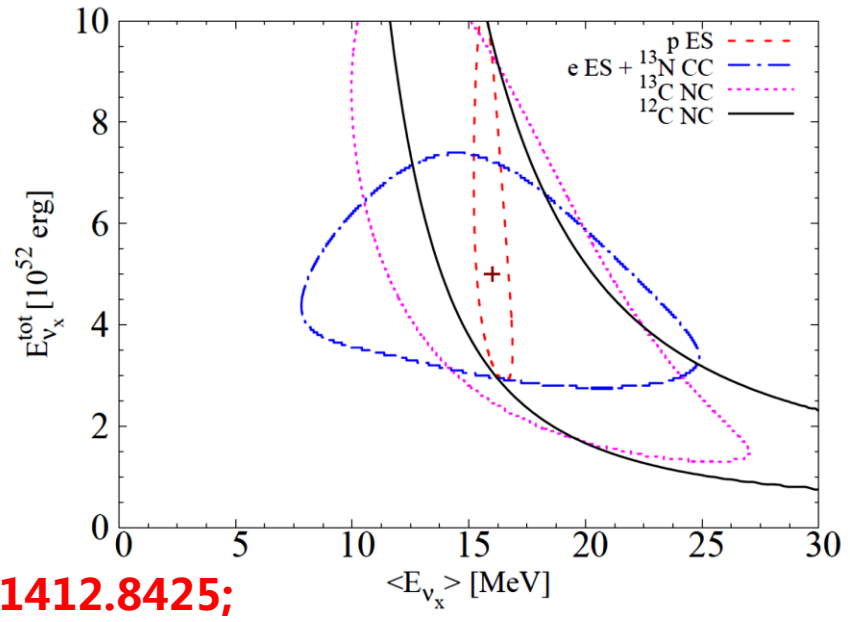
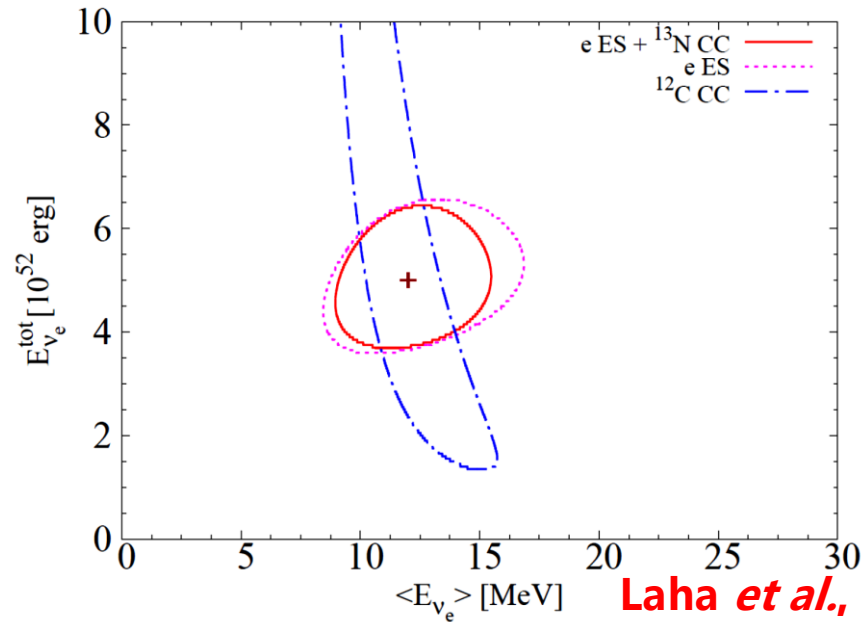
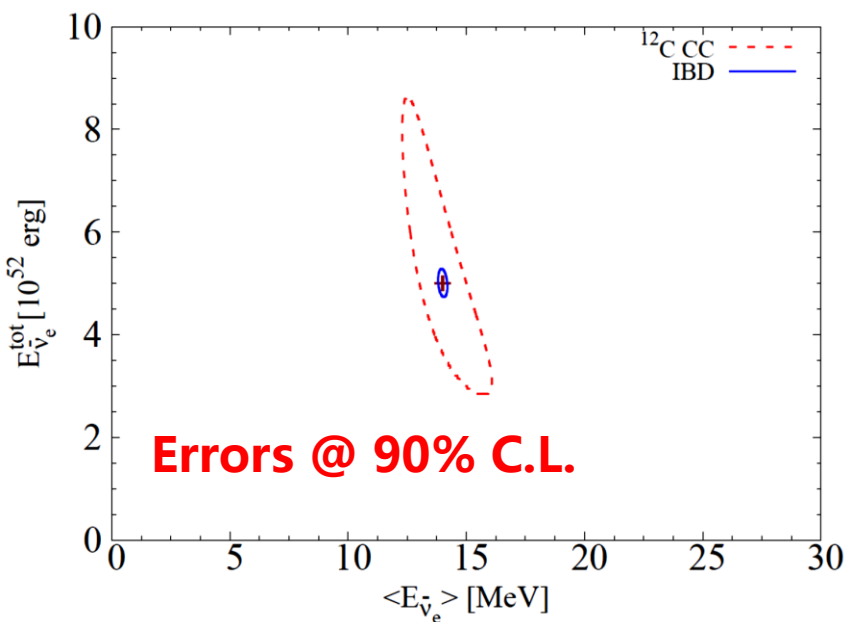
SN Neutrinos @ LS Detectors

Channel	Type	Number of SN Neutrino Events at JUNO			
		No Oscillations	Normal Ordering	Inverted Ordering	
$\bar{\nu}_e + p \rightarrow e^+ + n$	CC	4573	4775	5185	
		1578	1578	1578	
$\nu + p \rightarrow \nu + p$	ES	ν_e	107	354	278
		$\bar{\nu}_e$	179	214	292
		ν_x	1292	1010	1008
		314	316	316	
$\nu_e + e \rightarrow \nu_e + e$	ES	ν_e	157	159	158
		$\bar{\nu}_e$	61	61	62
		ν_x	96	96	96
$\nu_e + {}^{12}\text{C} \rightarrow e^- + {}^{12}\text{N}$	CC	43	134	106	
$\bar{\nu}_e + {}^{12}\text{C} \rightarrow e^+ + {}^{12}\text{B}$	CC	86	98	126	
		352	352	352	
$\nu + {}^{12}\text{C} \rightarrow \nu + {}^{12}\text{C}^*$	NC	ν_e	27	76	61
		$\bar{\nu}_e$	43	50	65
		ν_x	282	226	226
$\nu_e + {}^{13}\text{C} \rightarrow e^- + {}^{13}\text{N}$	CC	19	29	26	
	$3/2^- (5/2^-)$	23(15)	23(15)	23(15)	
$\nu + {}^{13}\text{C} \rightarrow \nu + {}^{13}\text{C}^*$	NC	ν_e	3(1)	4(3)	4(2)
		$\bar{\nu}_e$	3(2)	4(2)	4(3)
		ν_x	17(12)	15(10)	15(10)

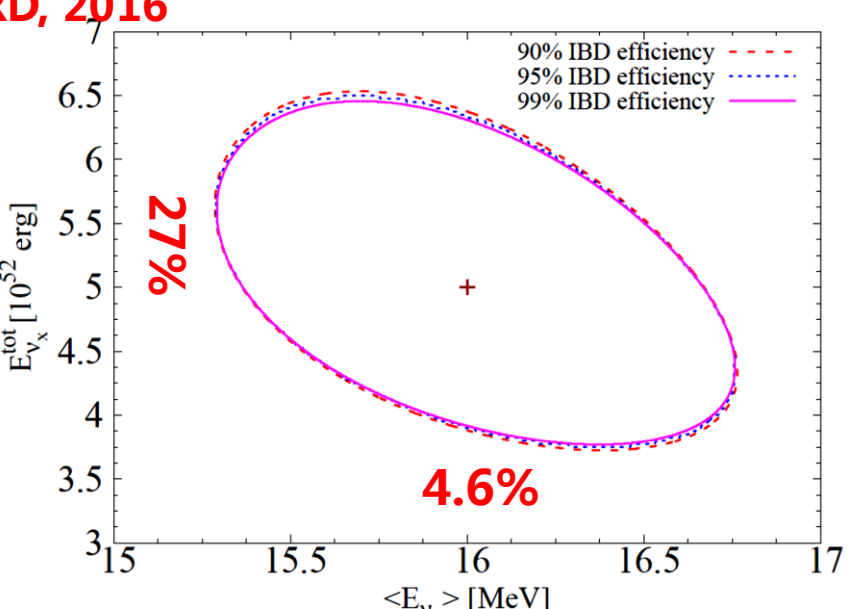
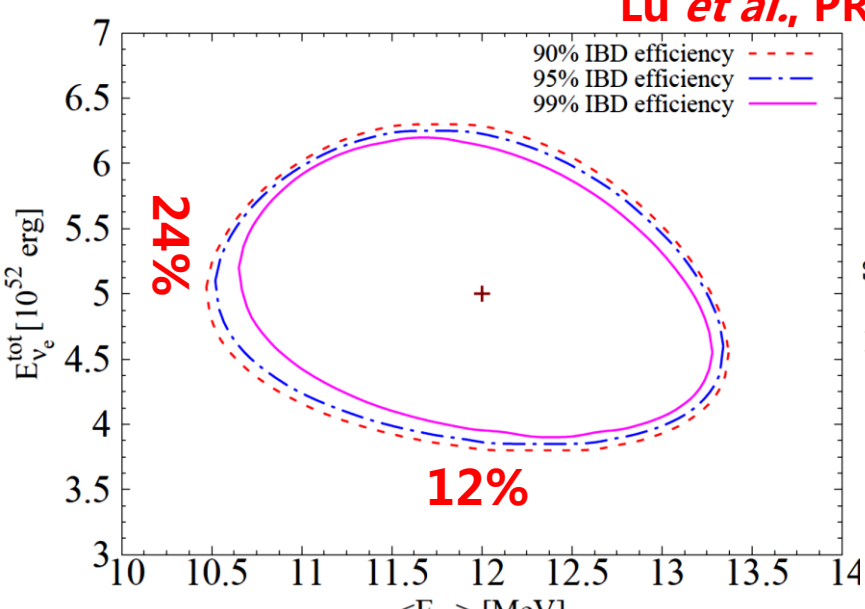
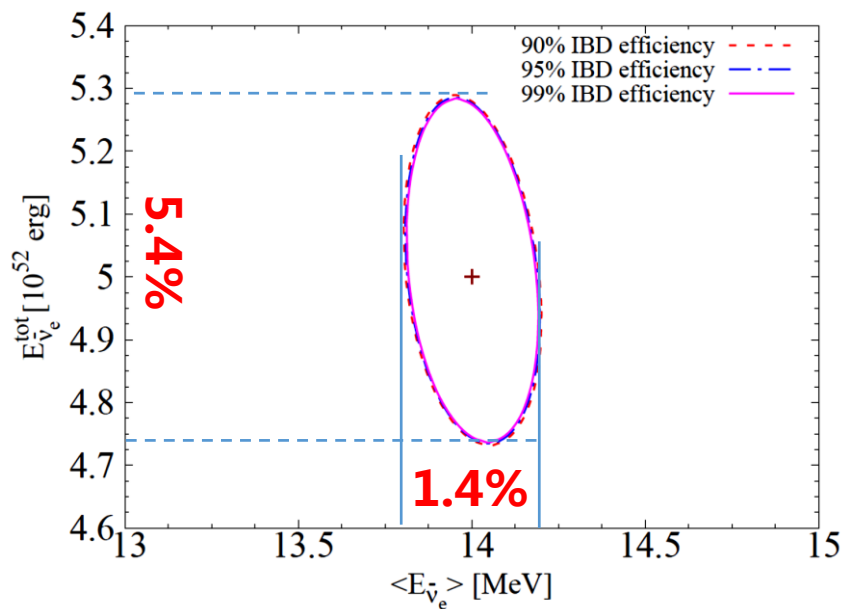
Detection channels	ν Flavors	Efficiency	Backgrounds	Systematics
IBD	$\bar{\nu}_e$	95%	None	Detection 2%
${}^{12}\text{C}$ -CC	$\bar{\nu}_e$ and ν_e	90%	None	Detection 2%
p ES	$\bar{\nu}_e, \nu_e$ and ν_x	99%	e ES	Detection 2%
				Cross section 20%
			k_B	3%
e ES	$\bar{\nu}_e, \nu_e$ and ν_x	99%	${}^{13}\text{N}$ -CC+IBD+ p ES	Detection 2%
${}^{13}\text{N}$ -CC	ν_e	100%	e ES+IBD	Detection 2%
				Cross section 20%
${}^{12}\text{C}$ -NC	$\bar{\nu}_e, \nu_e$ and ν_x	100%	e ES+IBD	Detection 2%
				Cross section 20%
${}^{13}\text{C}$ -NC	$\bar{\nu}_e, \nu_e$ and ν_x	100%	e ES+IBD	Detection 2%
				Cross section 20%

- IBD for $\bar{\nu}_e$ + sub-leading effects from ${}^{12}\text{C}$ CC
- Elastic ν -e scattering for $\nu_e + {}^{12}\text{C}$ CC
- Elastic ν -p scattering for $\nu_x + e$ ES
- A global analysis of all reaction channels?
Laha *et al.*, 1412.8425; Lu *et al.*, PRD, 2016

SN Neutrinos @ LS Detectors



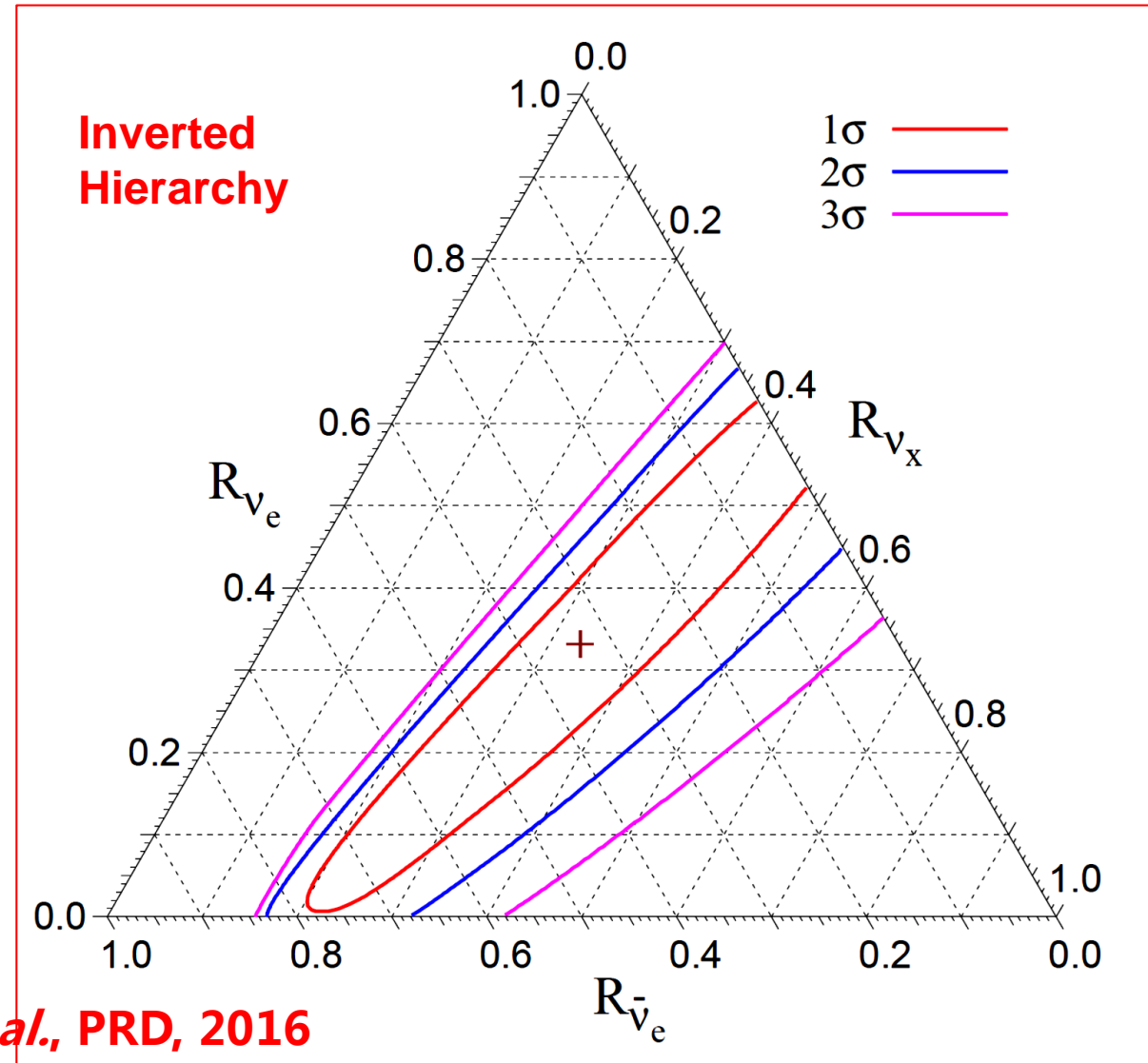
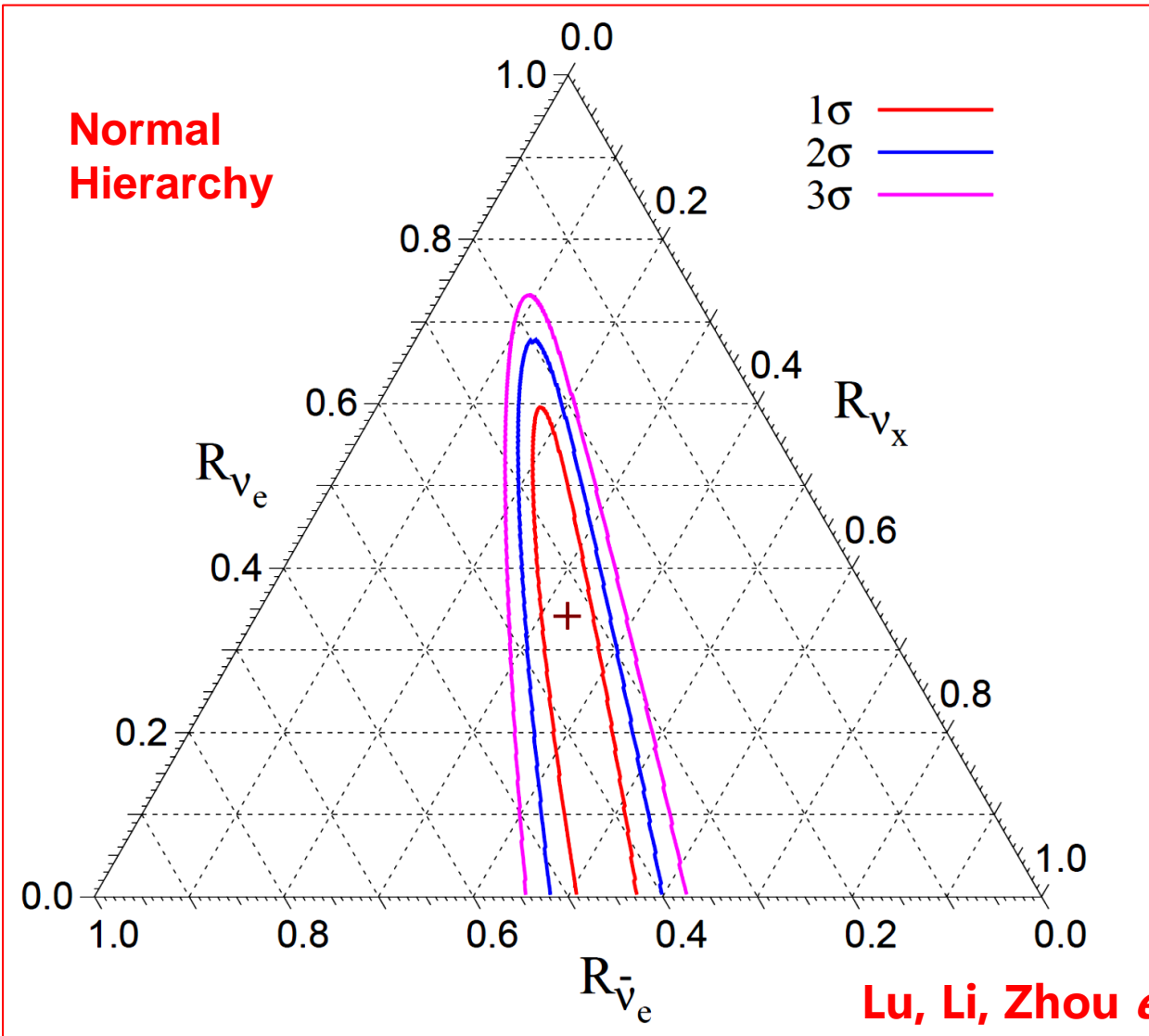
Laha *et al.*, 1412.8425;
Lu *et al.*, PRD, 2016



Test of Energy Equipartition Hypothesis

Including only the MSW matter effects in the SN

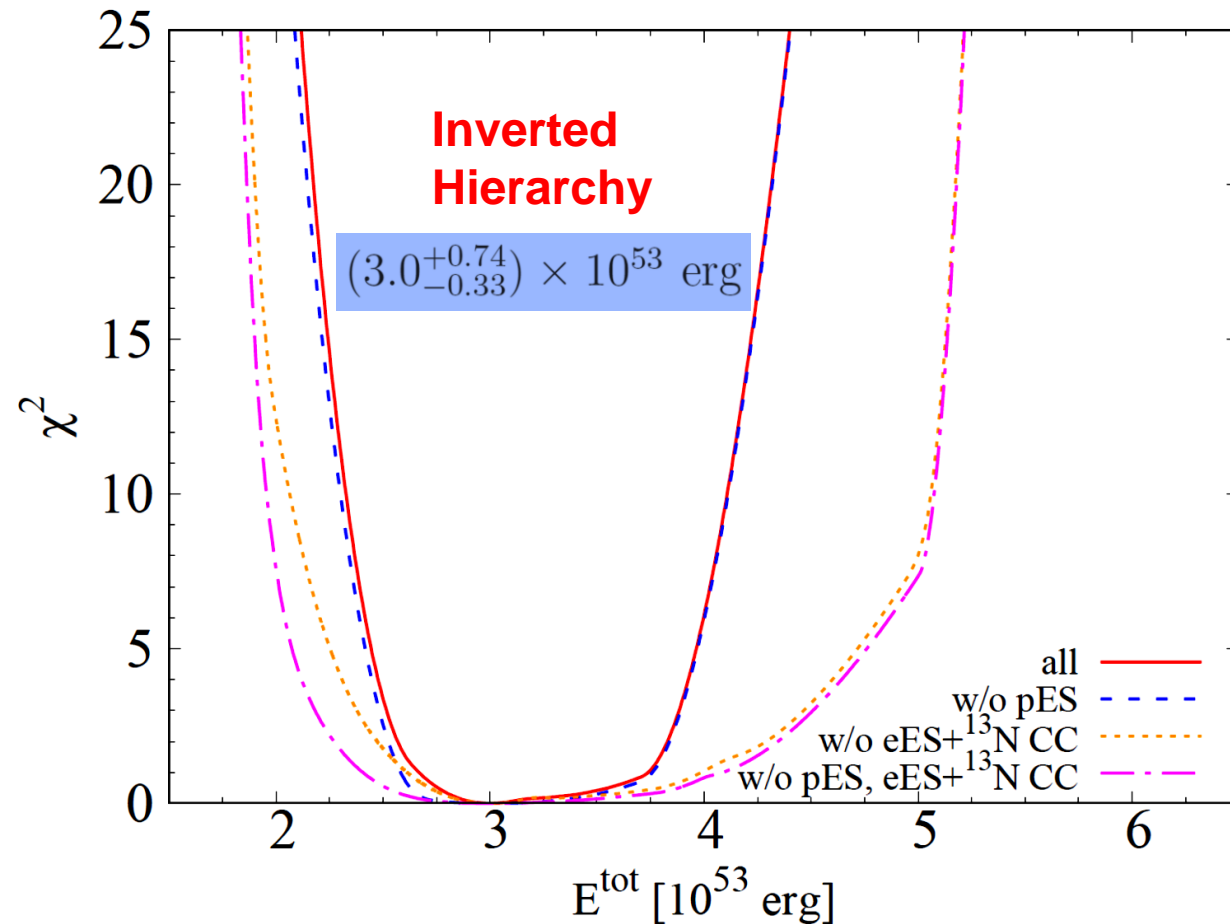
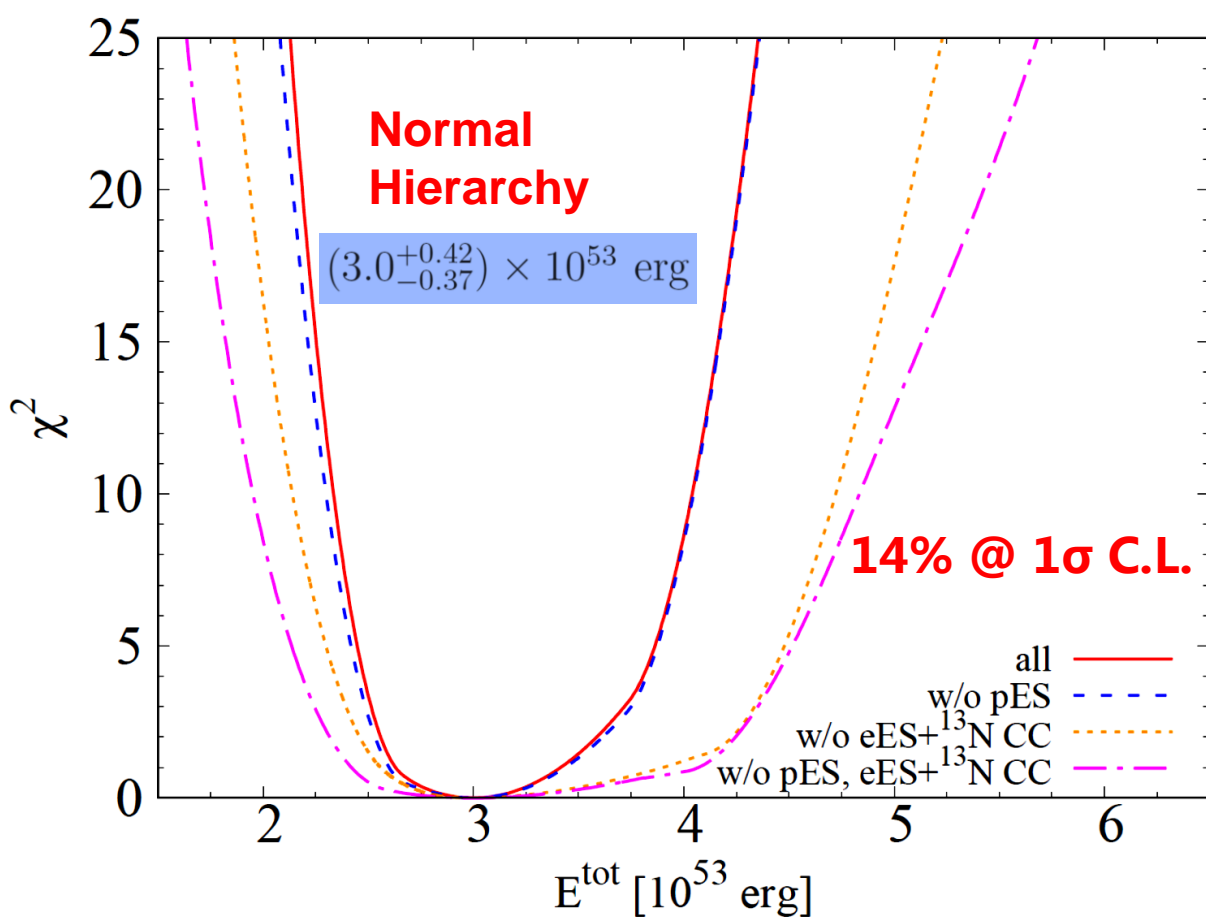
For collective flavor conversions, talk by M.R. Wu



Total Gravitational Binding Energy

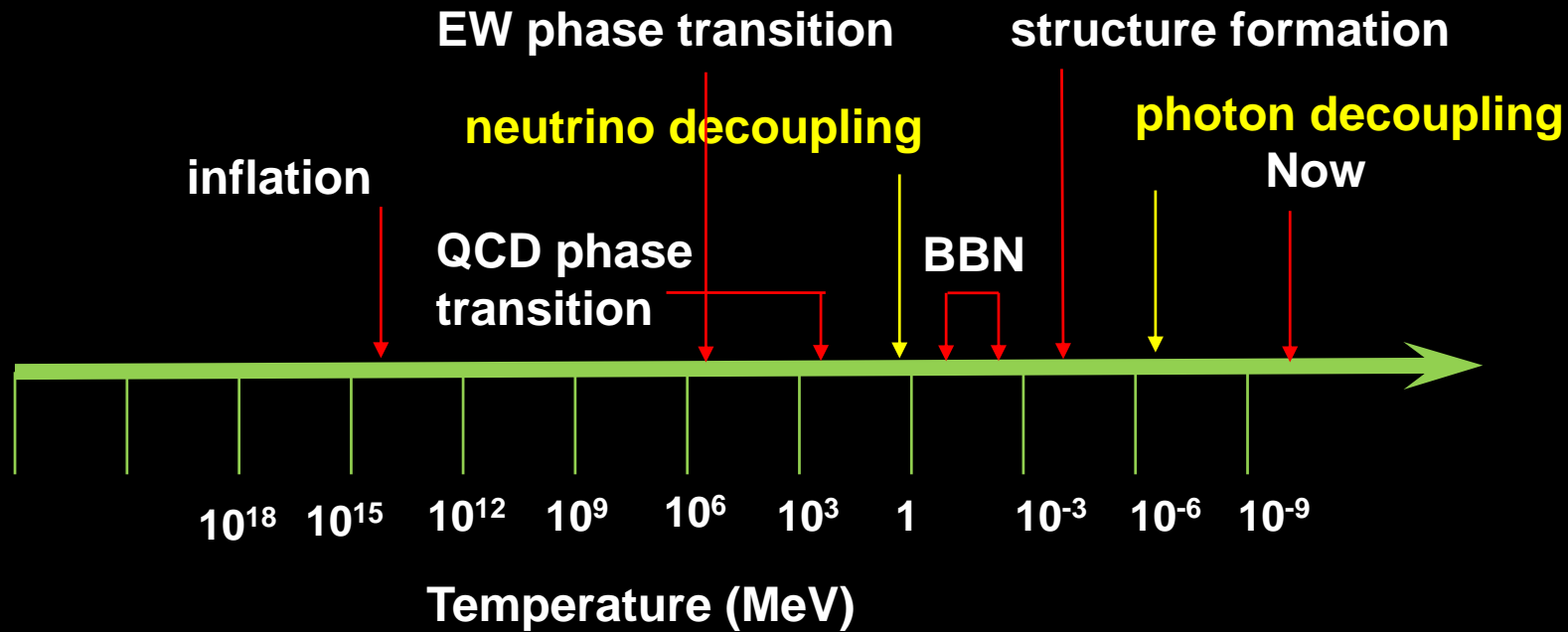
Including only the MSW effects in the SN, and fixing the spectral indices at $\gamma = 3$

Lu, Li, Zhou *et al.*, PRD, 2016



- With a high-statistics measurement of a galactic SN: **time & energy spectra, all flavors, ...**
- More works need to be done: **collective flavor conversions, explosion mechanisms, ...**

Formation of Cosmic ν Background



□ ν in thermal equilibrium
@ high temperature

$$\nu_\alpha \nu_\beta \leftrightarrow \nu_\alpha \nu_\beta$$

$$\nu_\alpha \bar{\nu}_\beta \leftrightarrow \nu_\alpha \bar{\nu}_\beta$$

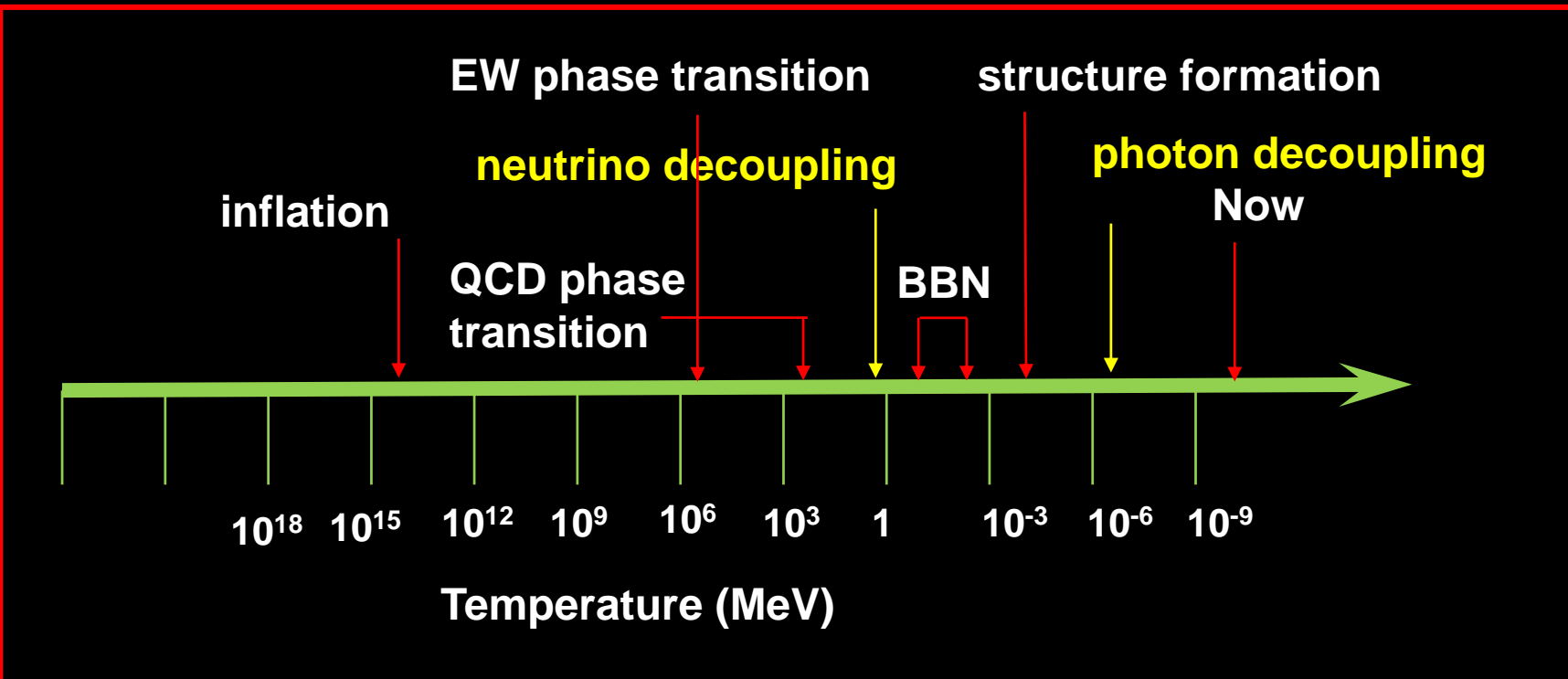
$$\nu_\alpha e^- \leftrightarrow \nu_\alpha e^-$$

$$\nu_\alpha \bar{\nu}_\alpha \leftrightarrow e^+ e^-$$

$$e^+ e^- \leftrightarrow \gamma \gamma$$

$$T_\nu = T_e = T_\gamma$$

Formation of Cosmic ν Background



□ ν in thermal equilibrium @ high temperature

$$\begin{aligned} \nu_\alpha \nu_\beta &\leftrightarrow \nu_\alpha \nu_\beta \\ \nu_\alpha \bar{\nu}_\beta &\leftrightarrow \nu_\alpha \bar{\nu}_\beta \\ \nu_\alpha e^- &\leftrightarrow \nu_\alpha e^- \\ \nu_\alpha \bar{\nu}_\alpha &\leftrightarrow e^+ e^- \\ e^+ e^- &\leftrightarrow \gamma\gamma \end{aligned}$$

$$T_\nu = T_e = T_\gamma$$

□ neutrino decoupling

$$\Gamma < H @ T \sim 1 \text{ MeV}$$

Weak interactions

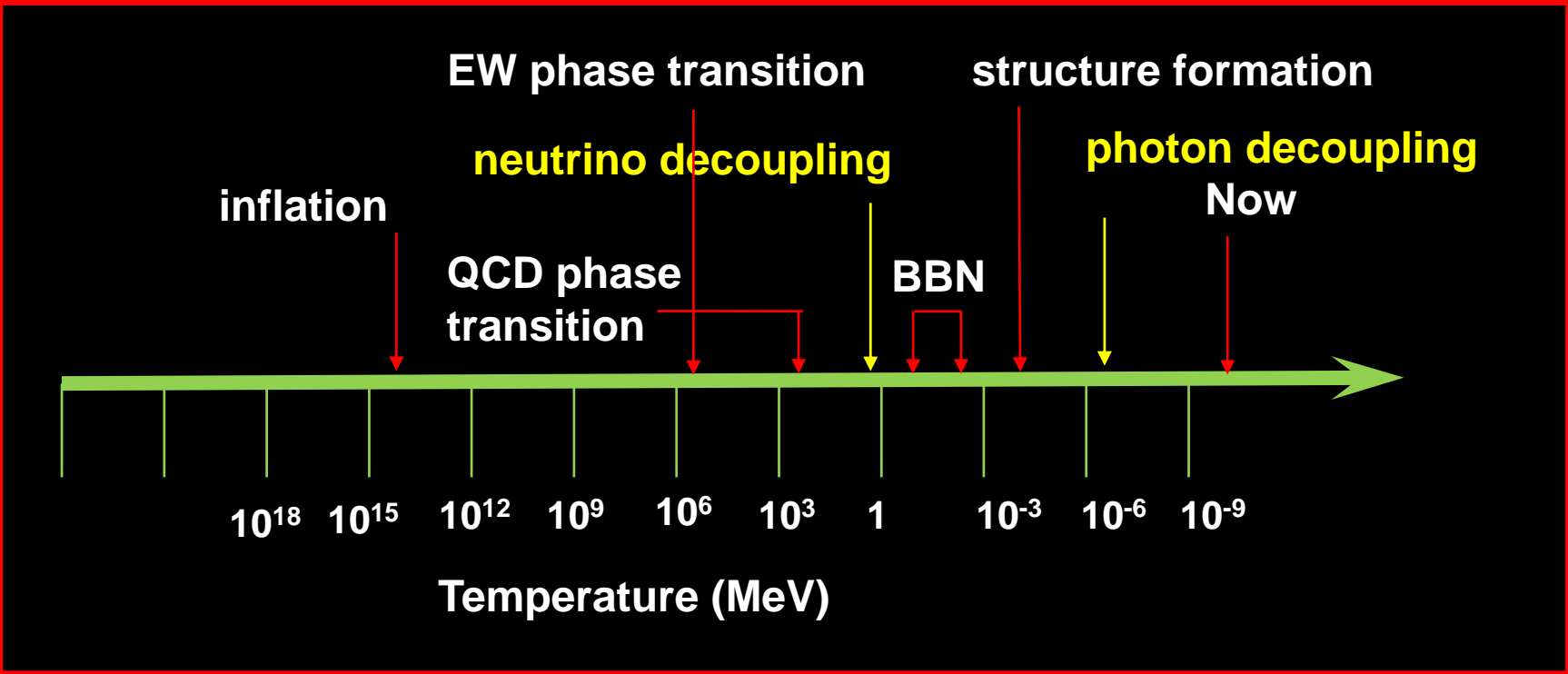
$$\Gamma \approx G_F^2 T^5$$

Hubble expansion

$$H = \frac{\sqrt{g_*} T^2}{M_{\text{pl}}}$$

Fermi-Dirac spectrum

Formation of Cosmic ν Background



□ ν in thermal equilibrium @ high temperature

$$\begin{aligned} \nu_\alpha \nu_\beta &\leftrightarrow \nu_\alpha \nu_\beta \\ \nu_\alpha \bar{\nu}_\beta &\leftrightarrow \nu_\alpha \bar{\nu}_\beta \\ \nu_\alpha e^- &\leftrightarrow \nu_\alpha e^- \\ \nu_\alpha \bar{\nu}_\alpha &\leftrightarrow e^+ e^- \\ e^+ e^- &\leftrightarrow \gamma\gamma \end{aligned}$$

$$T_\nu = T_e = T_\gamma$$

□ neutrino decoupling

$$\Gamma < H @ T \sim 1 \text{ MeV}$$

Weak interactions

$$\Gamma \approx G_F^2 T^5$$

Hubble expansion

$$H = \frac{\sqrt{g_*} T^2}{M_{\text{pl}}}$$

Fermi-Dirac spectrum

□ photon reheating

$$e^+ e^- \leftrightarrow \gamma\gamma @ T < m_e$$

$$T_\nu = \left(\frac{4}{11}\right)^{1/3} T_\gamma$$

□ Basic properties of **CvB**

- $T_0 = 1.95 \text{ K}$ and $\langle p \rangle = 3T_0 = 5 \times 10^{-4} \text{ eV}$
- number density $n = 56 \text{ cm}^{-3}$ per species

Acceleration in the neutrino wind

Shvartsman et al., 82; Smith, Lewin, 83; Duda et al., 01

$$2 \times 10^{-28} \frac{n_\nu}{\bar{n}_\nu} \frac{10^{-3} c}{v_{\text{rel}}} \frac{\rho_t}{\text{g/cm}^3} \frac{r_t^3}{\lambda} \frac{\text{cm}}{\text{s}^2}$$

Target mass

Current Sensitivity
 $10^{-13} \text{ cm s}^{-2}$

de Broglie
wavelength

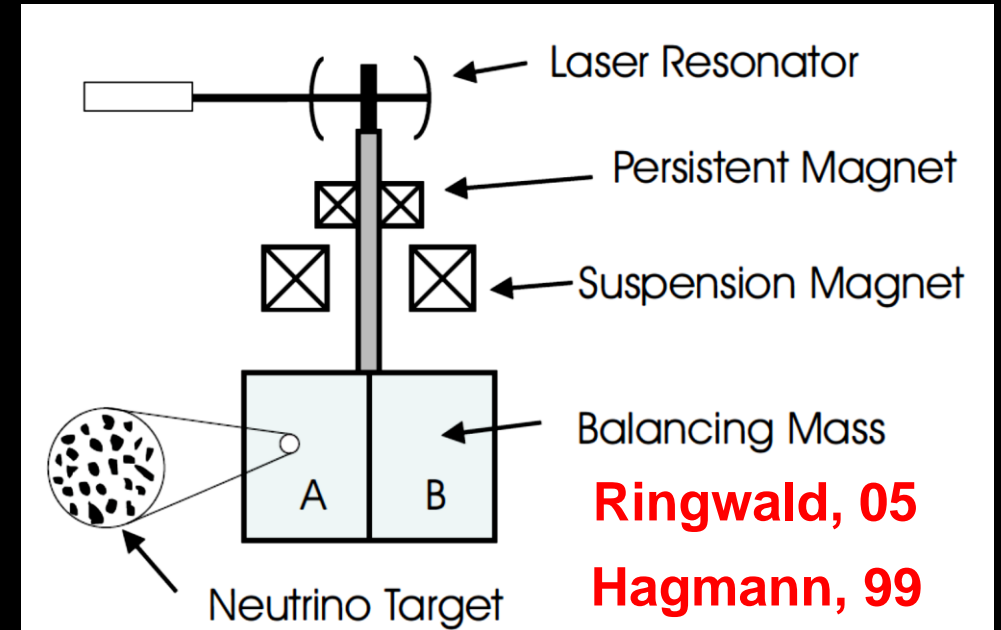
Resonant absorption of EHE neutrinos

Weiler, 82; Eberle et al., 04; Ringwald, 09

$$E_{0,i}^{\text{res}} = \frac{m_Z^2}{2m_{\nu_{0,i}}} = 4.2 \times 10^{12} \left(\frac{\text{eV}}{m_{\nu_i}} \right) \text{ GeV}$$

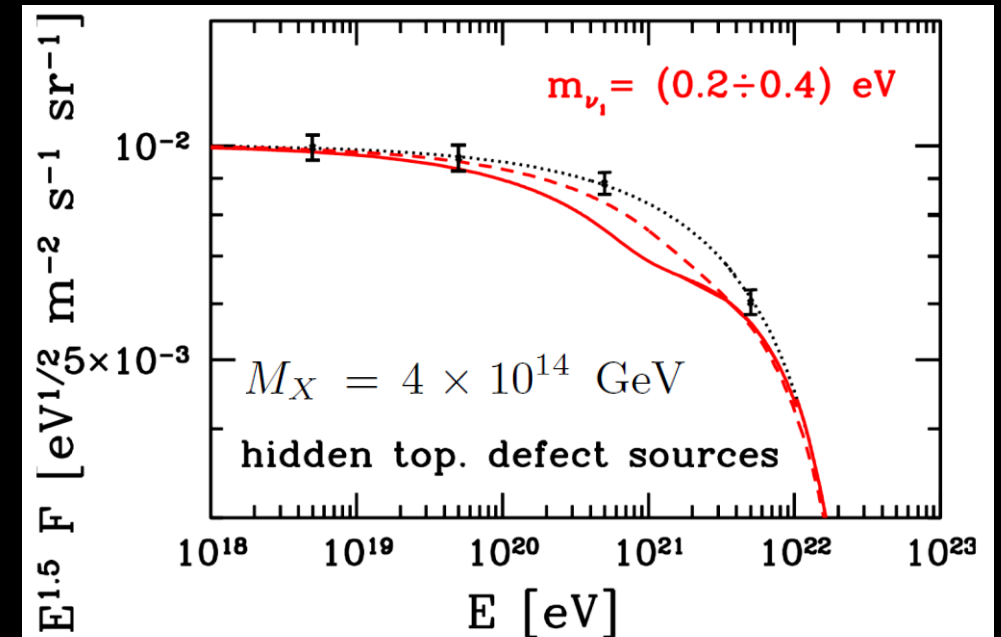
- Sources of EHE neutrinos
- Nearly-degenerate masses
- Z-burst for EHE CR events

Eberle et al., 04

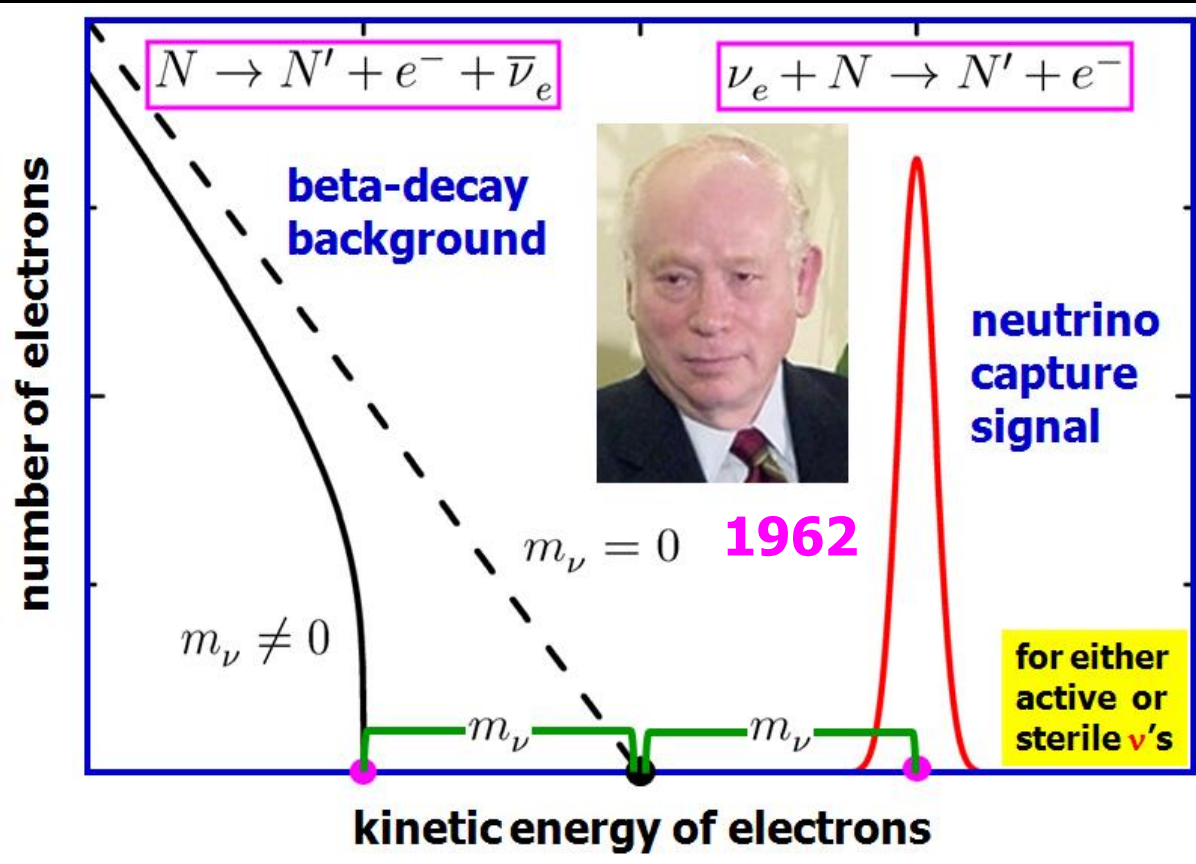


Ringwald, 05

Hagmann, 99



Relic neutrino capture on β -decaying nuclei



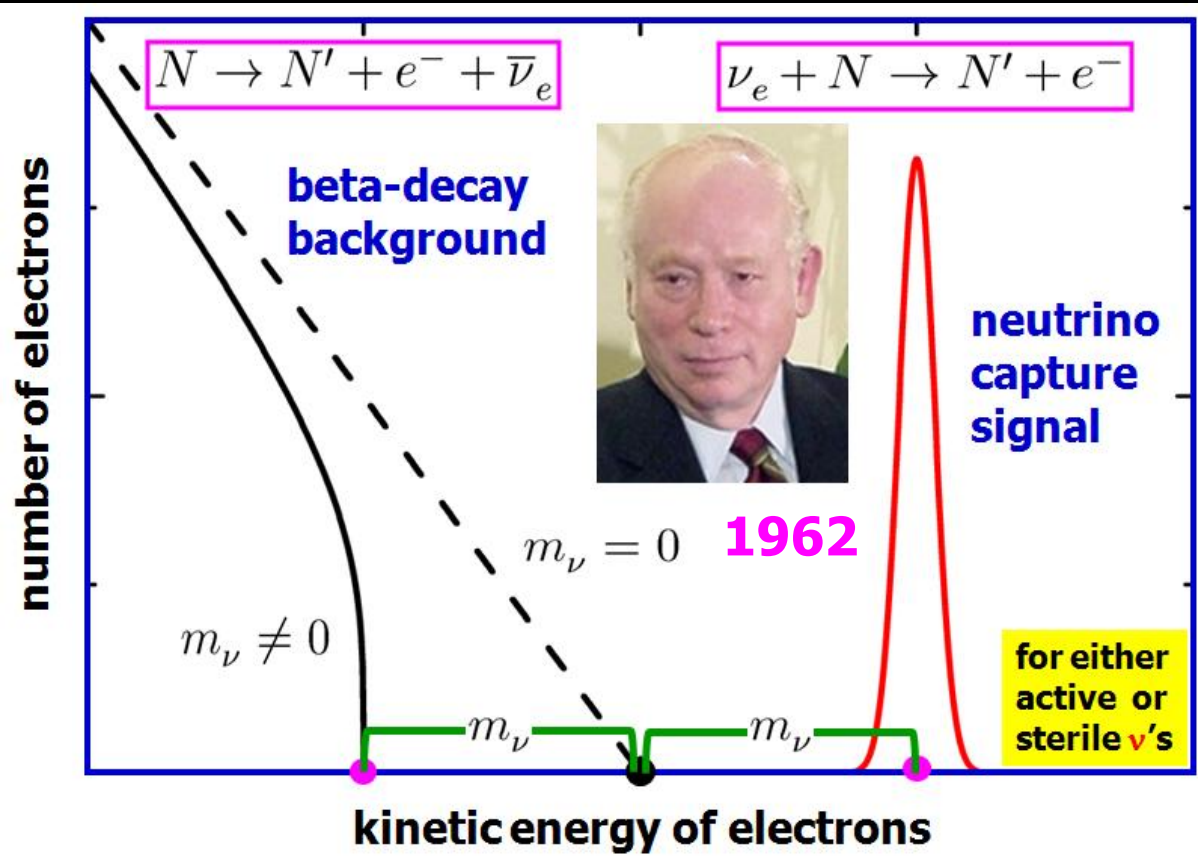
no energy threshold on incident ν 's **mono-**
energetic outgoing electrons

(Irvine & Humphreys, 83; Cucco et al., 07)

At least 2 ν 's cold today **NON-relativistic ν 's!**

Prospects for CvB Detection

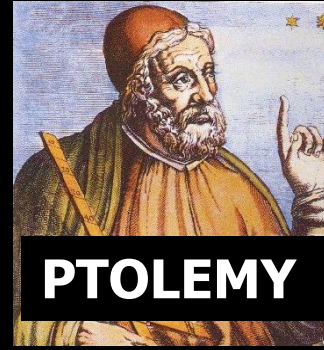
Relic neutrino capture on β -decaying nuclei



no energy threshold on incident ν 's mono-energetic outgoing electrons

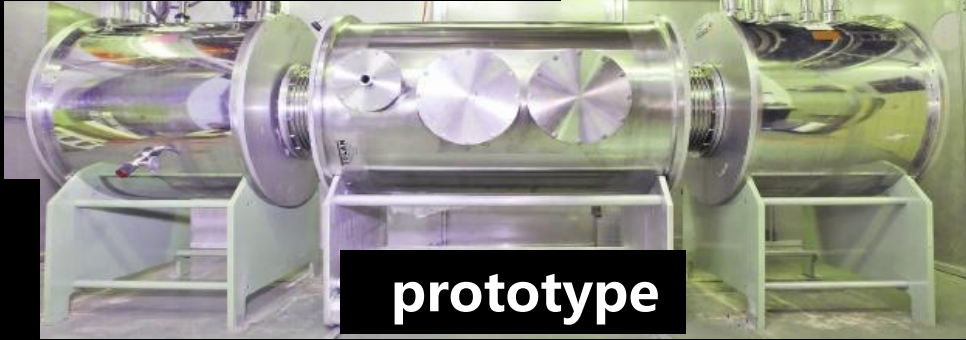
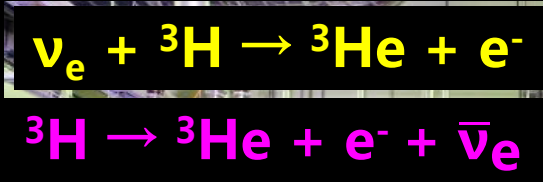
(Irvine & Humphreys, 83; Cucco et al., 07)

At least 2 ν 's cold today **NON-relativistic ν 's!**



PTOLEMY

Princeton Tritium Observ. for Light, Early-Univ., Massive- ν Yield



prototype

Betts et al, arXiv:1307.4738

★ CvB capture rate

$\Gamma_{C\nu B}^D \sim 4 \text{ yr}^{-1}$

$\Gamma_{C\nu B}^M \sim 8 \text{ yr}^{-1}$

D = Dirac

M = Majorana

- ★ first experiment
- ★ 100 g of tritium
- ★ graphene target
- ★ planned energy resolution 0.15 eV

Capture rate of a polarized neutrino state $\nu_j(s_\nu)$ on a free neutron

$$\sigma_j(s_\nu) v_{\nu_j} = \frac{G_F^2}{2\pi} |V_{ud}|^2 |U_{ej}|^2 F(Z, E_e) \frac{m_p}{m_n} E_e p_e A(s_\nu) (f^2 + 3g^2)$$

$$A(s_\nu) \equiv 1 - 2s_\nu v_{\nu_j} = \begin{cases} 1 - v_{\nu_j}, & s_\nu = +1/2 \quad \text{RH Helicity} \\ 1 + v_{\nu_j}, & s_\nu = -1/2 \quad \text{LH Helicity} \end{cases}$$

In the limit $v_{\nu_j} \rightarrow 1$, the state of $s_\nu = +1/2$ cannot be captured

In the limit $v_{\nu_j} \rightarrow 0$, both RH and LH helical states do contribute

Long, Lunardini, Sabancilar, 14; Lisanti, Safdi, Tully, 14

Total Rate

$$\Gamma_{\text{CvB}} = \sum_j \left[\sigma_j \left(+\frac{1}{2} \right) v_{\nu_j} \mathbf{n}_j(\mathbf{v}_{\text{hR}}) + \sigma_j \left(-\frac{1}{2} \right) v_{\nu_j} \mathbf{n}_j(\mathbf{v}_{\text{hL}}) \right] N_T$$

Capture rate of a polarized neutrino state $\nu_j(s_\nu)$ on a free neutron

$$\sigma_j(s_\nu) v_{\nu_j} = \frac{G_F^2}{2\pi} |V_{ud}|^2 |U_{ej}|^2 F(Z, E_e) \frac{m_p}{m_n} E_e p_e A(s_\nu) (f^2 + 3g^2)$$

$$A(s_\nu) \equiv 1 - 2s_\nu v_{\nu_j} = \begin{cases} 1 - v_{\nu_j}, & s_\nu = +1/2 \quad \text{RH Helicity} \\ 1 + v_{\nu_j}, & s_\nu = -1/2 \quad \text{LH Helicity} \end{cases}$$

In the limit $v_{\nu_j} \rightarrow 1$, the state of $s_\nu = +1/2$ cannot be captured

In the limit $v_{\nu_j} \rightarrow 0$, both RH and LH helical states do contribute

Long, Lunardini, Sabancilar, 14; Lisanti, Safdi, Tully, 14

Total Rate

$$\Gamma_{\text{CvB}} = \sum_j \left[\sigma_j \left(+\frac{1}{2} \right) v_{\nu_j} n_j(\nu_{\text{hR}}) + \sigma_j \left(-\frac{1}{2} \right) v_{\nu_j} n_j(\nu_{\text{hL}}) \right] N_T$$

Dirac

Majorana

Decoupling

$$\begin{array}{ll} n(\nu_L) = n(z) & n(\nu_L) = n(z) \\ n(\bar{\nu}_R) = n(z) & n(\nu_R) = n(z) \end{array}$$



Nowadays

$$\begin{array}{ll} n(\nu_{\text{hL}}) = n_0 & n(\nu_{\text{hL}}) = n_0 \\ n(\bar{\nu}_{\text{hR}}) = n_0 & n(\nu_{\text{hR}}) = n_0 \end{array}$$

$$\bar{\sigma} \approx 3.8 \times 10^{-45} \text{ cm}^2$$

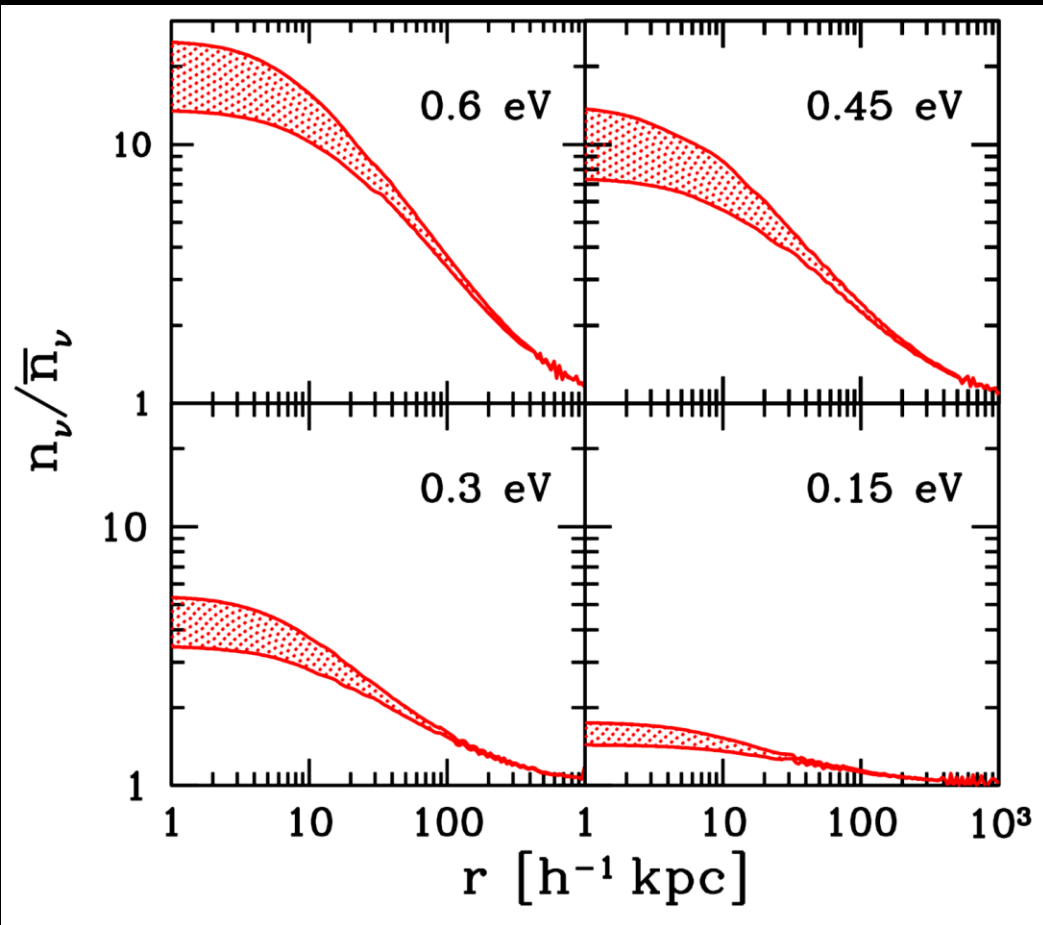
$$\Gamma_{\text{CvB}}^{\text{D}} = \bar{\sigma} n_0 N_T$$

$$\Gamma_{\text{CvB}}^{\text{M}} = 2\bar{\sigma} n_0 N_T$$

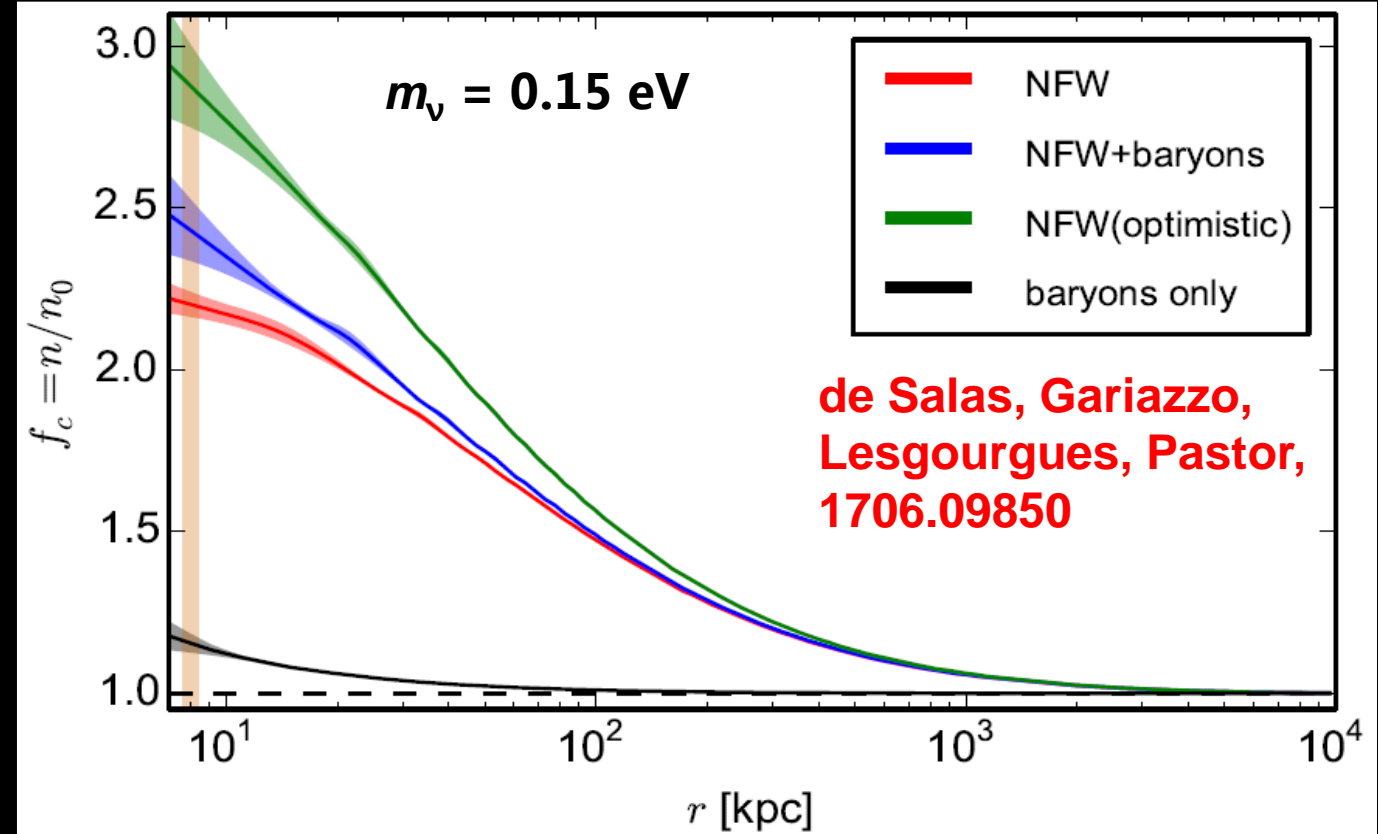
Neutrino Clustering

Ringwald, Wong, 04

Clustering in the Milky Way

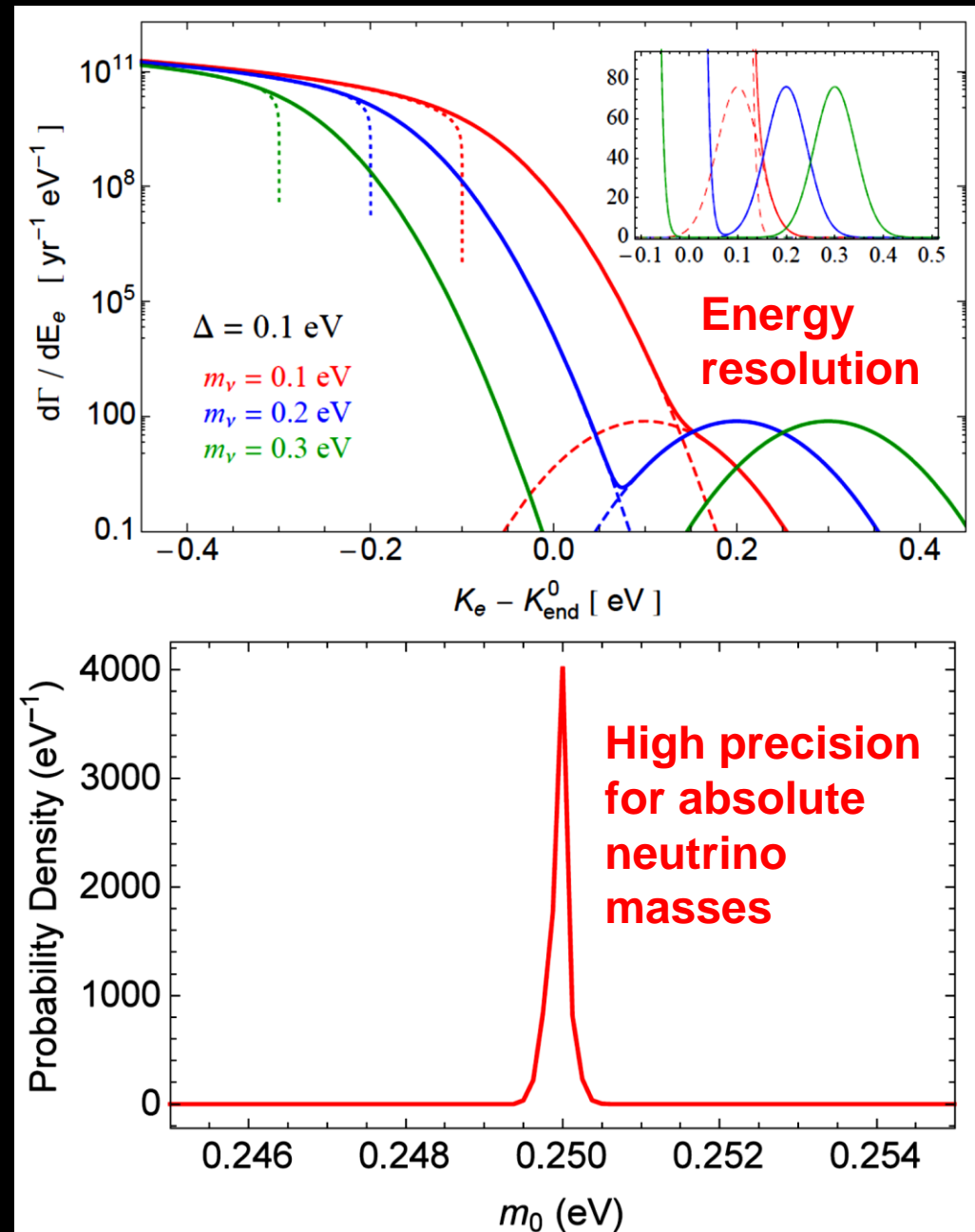


At the Earth, larger by a factor of 1 to 20

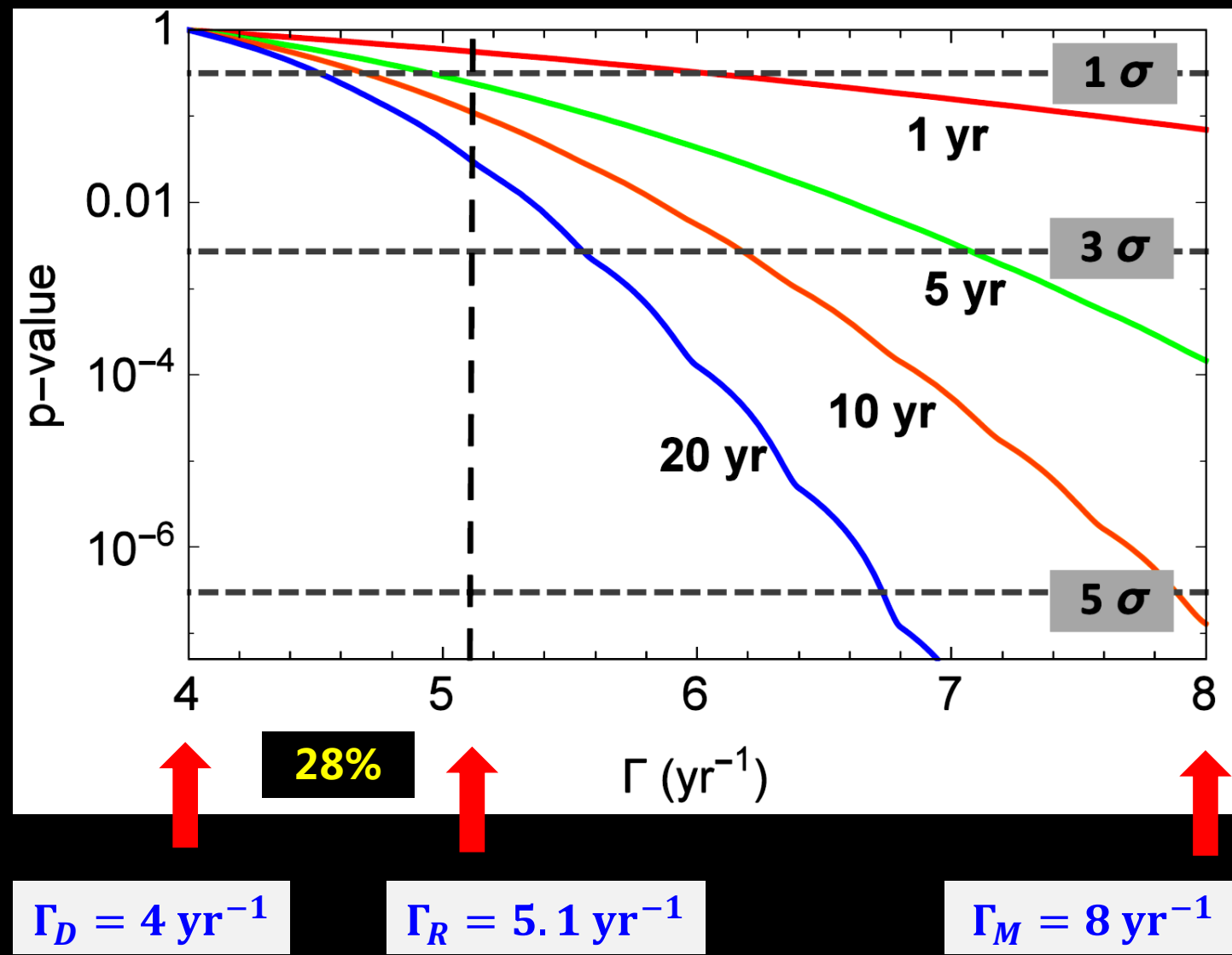


de Salas, Gariazzo,
Lesgourgues, Pastor,
1706.09850

matter halo	overdensity f_c $f_1 \simeq f_2 \simeq f_3$	Γ_{tot}^D (yr ⁻¹)	Γ_{tot}^M (yr ⁻¹)
any	no clustering	4.06	8.12
NFW(+bar)	2.18 (2.44)	8.8 (9.9)	17.7 (19.8)
NFW optimistic	2.88	11.7	23.4
EIN(+bar)	1.68 (1.87)	6.8 (7.6)	13.6 (15.1)
EIN optimistic	2.43	9.9	19.7



- Nominal setup for PTOLEMY: $100 \text{ g } ^3\text{H}$, $\Delta = 0.15 \text{ eV}$
- Absolute masses, Dirac vs. Majorana, Extra ν 's, ...



- The next galactic core-collapse SN will be measured by a number of neutrino detectors, and a high-statistics real-time measurement is helpful in understanding explosion mechanisms and probing the intrinsic properties of massive neutrinos
- Give the priority to Diffuse SN Neutrino Background, a guaranteed source of SN neutrinos. We have SK with Gd doping, and JUNO (available within 3 years) also has a good chance.

Syst. uncertainty BG	5%		20%	
	rate only	spectral fit	rate only	spectral fit
$\langle E_{\bar{\nu}_e} \rangle$				
12 MeV	2.3σ	2.5σ	2.0σ	2.3σ
15 MeV	3.5σ	3.7σ	3.2σ	3.3σ
18 MeV	4.6σ	4.8σ	4.1σ	4.3σ
21 MeV	5.5σ	5.8σ	4.9σ	5.1σ

**Neutrino Physics
with JUNO, JGD, 16**

- Very promising to detect the relic neutrinos from the Big Bang in the PTOLEMY experiment. It is time to have a serious look at theoretical predictions for local number densities, the detection rates and physics potentials for elementary particles and cosmology

- The next galactic core-collapse SN will be measured by a number of neutrino detectors, and a high-statistics real-time measurement is helpful in understanding explosion mechanisms and probing the intrinsic properties of massive neutrinos
- Give the priority to Diffuse SN Neutrino Background, a guaranteed source of SN neutrinos. We have SK with Gd doping, and JUNO (available within 3 years) also has a good chance.

Syst. uncertainty BG $\langle E_{\bar{\nu}_e} \rangle$	5%		20%	
	rate only	spectral fit	rate only	spectral fit
12 MeV	2.3σ	2.5σ	2.0σ	2.3σ
15 MeV	3.5σ	3.7σ	3.2σ	3.3σ
18 MeV	4.6σ	4.8σ	4.1σ	4.3σ
21 MeV	5.5σ	5.8σ	4.9σ	5.1σ

**Neutrino Physics
with JUNO, JPG, 16**

- Very promising to detect the relic neutrinos from the Big Bang in the PTOLEMY experiment. It is time to have a serious look at theoretical predictions for local number densities, the detection rates and physics potentials for elementary particles and cosmology

Thanks a lot for your attention!