SUPERSYMMETRY BREAKING AND SINGULARITY IN DYNAMICAL M-BRANE BACKGROUND

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[]] Introduction

Dynamical branes in string theory

brane collision

(Gibbons & Lu & Pope, Phys.Rev.Lett. 94 (2005) 131602) (Maeda & Minamitsuji & Ohta & Uzawa, Phys. Rev. D82 (2010)046007) (Uzawa, Phys.Rev. D90 (2014) 025024)

 COSMIC Big-Bang of our universe (Chen, et al., Nucl.Phys. B732 (2006) 118-135) (Minamitsuji & Ohta & Uzawa, Phys. Rev. D82 (2010)086002))

• black hole in expanding universe (Maeda & Ohta & Uzawa, JHEP 0906 (2009) 051) (Maeda & Nozawa, Phys.Rev. D81 (2010) 044017) The cosmological scenario from the time dependent solution until the present have been much explored.

 However, the study of SUSY breaking in terms of dynamical solution is much less extensive.

 One motivation for the present work is to improve this situation.

The dynamical D3-brane solution preserves 1/4 SUSY in the conifold background. (H. Kodama & K. Uzawa, JHEP 0507 061 (2005))

\Rightarrow Question

Do supersymmetries preserve in the dynamical M-brane background?

[2] Preserved supersymmetry (11d SUGRA)

The 11-dimensional action is invariant under local SUSY transformations :

 e^{A}_{M} : graviton Ψ_{M} : gravitino, A_{MNP} : 3-form gauge potential

 $\delta e^A{}_M = \bar{\varepsilon} \, \Gamma^A \, \Psi_M \,,$

 $\delta A_{MNP} = -3\,\bar{\varepsilon}\,\Gamma_{[MN}\,\Psi_{P]}\,,$

$$\delta \Psi_M = D_M \varepsilon$$
$$= \left[\nabla_M + \frac{1}{12 \cdot 4!} \left(\Gamma_M F_{MNPQ} \Gamma^{MNPQ} - 12 F_{MNPQ} \Gamma^{NPQ} \right) \right] \varepsilon$$

Dynamical M2-brane solution :

$$ds^{2} = \left(c_{\mu}x^{\mu} + c + \frac{M}{r^{6}}\right)^{-2/3} \eta_{\mu\nu}(\mathbf{X})dx^{\mu}dx^{\nu}$$

$$+ \left(c_{\mu}x^{\mu} + c + \frac{M}{r^{6}}\right)^{1/3} \left(dr^{2} + r^{2}d\Omega_{(7)}^{2}\right)$$

$$\mathbf{8} - \mathbf{dim \ transverse \ space \ to \ brane}$$

$$F_{r\mu\nu\rho} = -\frac{6M}{r^{7}} \left(c_{\mu}x^{\mu} + c + \frac{M}{r^{6}}\right)^{-2} \varepsilon_{\mu\nu\rho}, \quad \Psi_{M} = 0$$

• Solution for dynamical background $h^{-1/3} c_{\mu} \Gamma^{\mu} \varepsilon = 0$, $(1 \pm h \Gamma_0 \Gamma_1 \Gamma_2) \varepsilon = 0$, $h = c_{\mu} x^{\mu} + c + \frac{M}{r^6}$.

•Integrability condition $[\nabla_M, \nabla_N] \varepsilon = 0$ gives

$$c_{\mu} c^{\mu} = 0$$

Dynamical spacetime
(1)
$$M \neq 0, c_{\mu}c^{\mu} = 0, c_{0} \neq 0, c_{1} \neq 0, c_{2}^{2} = c_{0}^{2} - c_{1}^{2}$$
:
 $\frac{1}{4}$ SUSY
(2) $M = 0$ (or $r \rightarrow \infty$), $c_{\mu}c^{\mu} = 0, c_{0} \neq 0, c_{1} \neq 0, c_{2}^{2} = c_{0}^{2} - c_{1}^{2}$: $\frac{1}{2}$ SUSY, plane wave
(3) $M \neq 0, c_{0} \neq 0, c_{1} = 0, c_{2} = 0$: Non SUSY
(4) $c_{\mu} = 0, c = 0$: Static, Maximal SUSY

Dynamical M2-brane background ($c_{\mu} c^{\mu} = 0$)



Dynamical M2-brane background $(c_1 = c_2 = 0)$



[3] SUSY breaking and enhancement of SUSY

(1) SUSY solution: $h=h(\tau, x^i, r), \tau/\tau_0=(ct)^{2/3}$

$$-\left(c\,t+c_{i}x^{i}+\frac{M}{r^{6}}\right)^{-\frac{2}{3}}dt^{2}+\cdots$$
$$=-\left[1+\left(\frac{\tau}{\tau_{0}}\right)^{-\frac{3}{2}}\left(c_{i}x^{i}+\frac{M}{r^{6}}\right)\right]^{-\frac{2}{3}}d\tau^{2}+\cdots$$

(2) As time increases (for $c_i x^i \ll M/r^6$),

$$1 + \left(\frac{\tau}{\tau_0}\right)^{-\frac{3}{2}} \left(c_i x^i + \frac{M}{r^6}\right) \rightarrow 1 + \left(\frac{\tau}{\tau_0}\right)^{-\frac{3}{2}} \frac{M}{r^6}$$

(3) $h(\tau, x^i, r)$ (SUSY) $\rightarrow h(\tau, r)$ (Non SUSY)

Time evolution (But toy model !!)







[4] Summary and comments

(1) The dynamical M2-brane background preserves the $\frac{1}{4}$ supersymmetry. For vanishing M2-brane charge, we also find $\frac{1}{2}$ SUSY solution.

(2) The solutions of field equations cannot give a homogeneous expansion at constant r unless supersymmetries are completely broken.

(3) Although the solution itself is by no means realistic, its interesting behavior suggests a possibility that the Universe preserved originally SUSY and begin to evolve toward a Universe without SUSY.