

Flavored B-L model and B anomalies

*Chengcheng Han
Kavli IPMU*

Based on arXiv:1705.03858(PLB), with Peter Cox and Tsutomu Yanagida

CosPA 2017

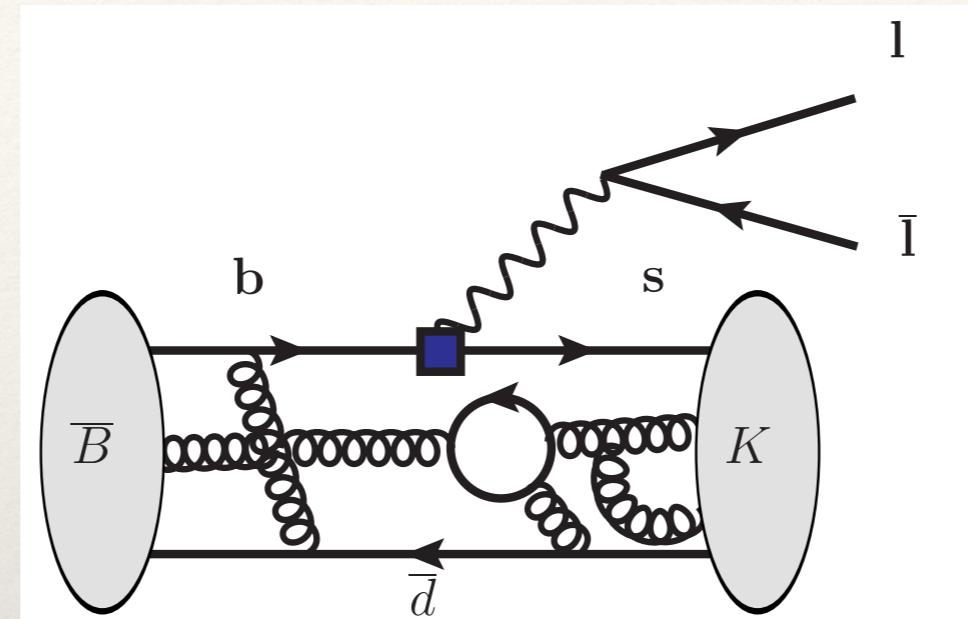
Yukawa Institute for Theoretical Physics, Kyoto University, JAPAN

Outline

- Briefly overview the B anomaly
- The flavored B-L model and phenomenologies
- Summary and conclusions

B decay anomaly

LFU violation



$$R_K = \frac{Br(B \rightarrow K\mu\mu)}{Br(B \rightarrow Kee)_{[1,6]}} = 0.745^{+0.090}_{-0.074}(\text{stat}) \pm 0.036(\text{syst})$$

LHCb arXiv:1406.6482

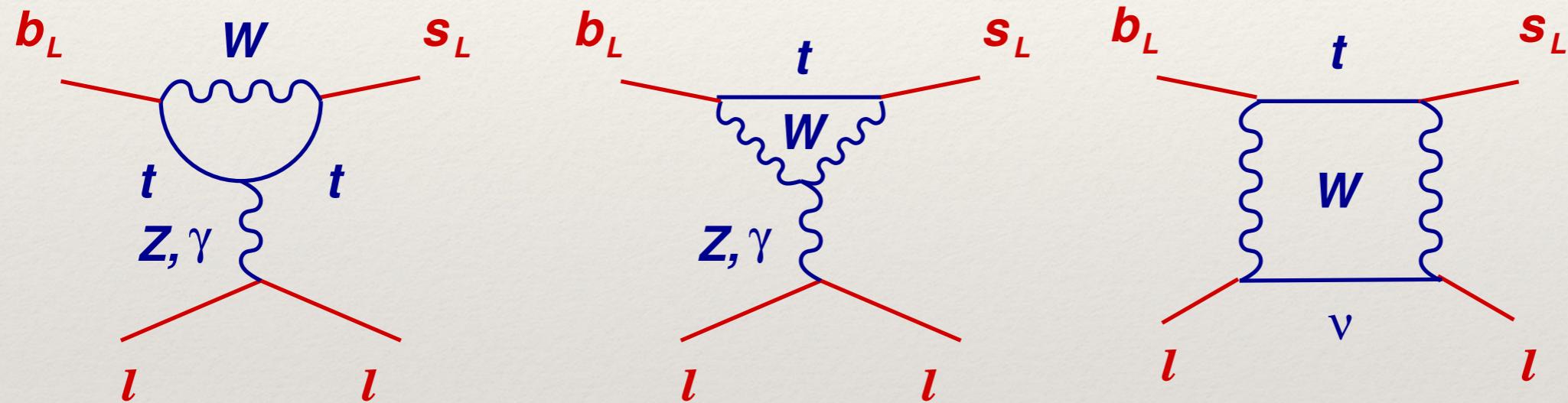
LHCb arXiv:1705.05802

$$R_{K^*} = \frac{Br(B \rightarrow K^*\mu\mu)}{Br(B \rightarrow K^*ee)_{[1,6]}} = 0.69^{+0.11}_{-0.07}(\text{stat}) \pm 0.05(\text{syst})$$
$$[0.05, 1] = 0.66^{+0.11}_{-0.07}(\text{stat}) \pm 0.024(\text{syst})$$

Combined ~ 4 sigma tension with SM($R_{\text{SM}}=1$)

SM contribution

$$\text{EFT } \mathcal{L}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* (C_9^l \mathcal{O}_9^l + C_{10}^l \mathcal{O}_{10}^l + C_\nu \mathcal{O}_\nu) ,$$



$$\mathcal{O}_9^l = \frac{\alpha}{4\pi} (\bar{s} \gamma_\mu b_L) (\bar{l} \gamma_\mu l) , \quad C_9^{SM} \sim -C_{10}^{SM} \sim 4.1$$

$$\mathcal{O}_{10}^l = \frac{\alpha}{4\pi} (\bar{s} \gamma_\mu b_L) (\bar{l} \gamma_\mu \gamma^5 l) ,$$

$$\mathcal{O}_\nu^{ij} = \frac{\alpha}{2\pi} (\bar{s} \gamma_\mu b_L) (\bar{\nu}^i \gamma_\mu \nu_L^j) .$$

New Physics

$$\mathcal{L}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left(C_9^l \mathcal{O}_9^l + C_{10}^l \mathcal{O}_{10}^l + C_T \mathcal{O}_T + C_S \mathcal{O}_S + \dots \right),$$

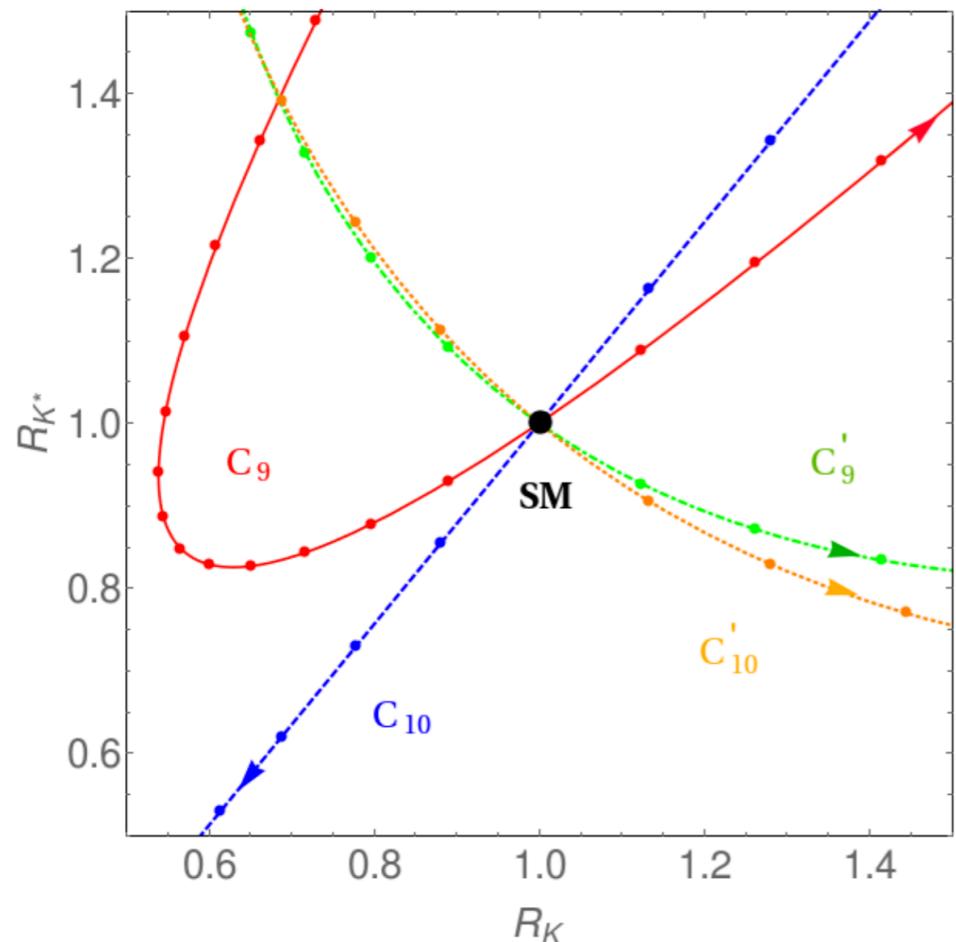
$$\mathcal{O}'_9 = \frac{e^2}{(4\pi)^2} [\bar{s}\gamma_\mu P_R b][\bar{l}\gamma^\mu l], \quad \mathcal{O}'_{10} = \frac{e^2}{(4\pi)^2} [\bar{s}\gamma_\mu P_R b][\bar{l}\gamma^\mu \gamma_5 l],$$

$$\mathcal{O}_S^{(\prime)} = \frac{e^2}{(4\pi)^2} [\bar{s}P_{R(L)} b][\bar{l}l], \quad \mathcal{O}_P^{(\prime)} = \frac{e^2}{(4\pi)^2} [\bar{s}P_{R(L)} b][\bar{l}\gamma_5 l],$$

$$\mathcal{O}_T = \frac{e^2}{(4\pi)^2} [\bar{s}\sigma_{\mu\nu} b][\bar{l}\sigma^{\mu\nu} l], \quad \mathcal{O}_{T5} = \frac{e^2}{(4\pi)^2} [\bar{s}\sigma_{\mu\nu} b][\bar{l}\sigma^{\mu\nu} \gamma_5 l].$$

Bobeth, Hiller, Piranishvili, '07

Fit results



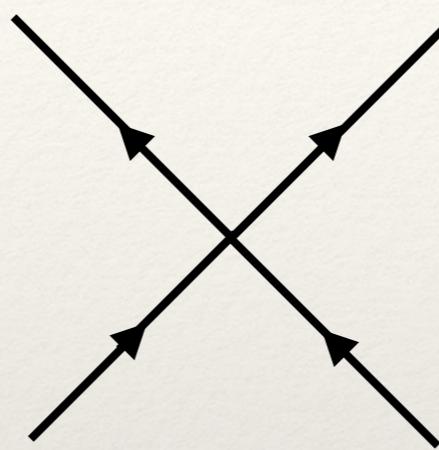
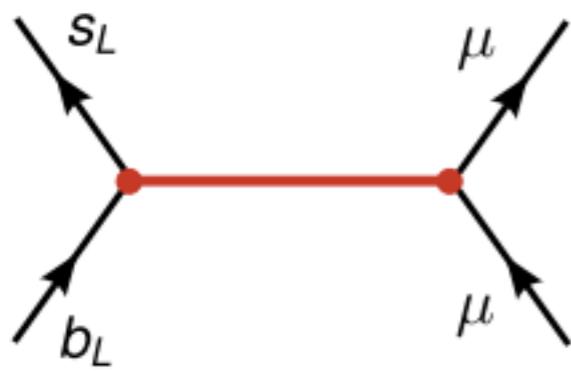
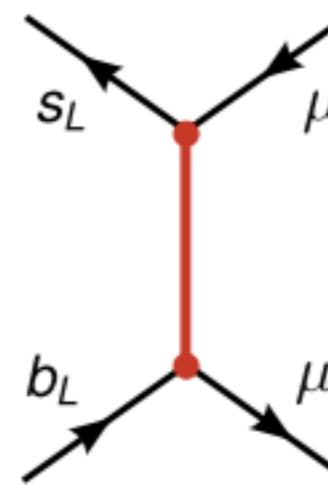
Coeff.	best fit	1σ	2σ	pull
C_9^μ	-1.59	[-2.15, -1.13]	[-2.90, -0.73]	4.2σ
C_{10}^μ	+1.23	[+0.90, +1.60]	[+0.60, +2.04]	4.3σ
C_9^e	+1.58	[+1.17, +2.03]	[+0.79, +2.53]	4.4σ
C_{10}^e	-1.30	[-1.68, -0.95]	[-2.12, -0.64]	4.4σ
$C_9^\mu = -C_{10}^\mu$	-0.64	[-0.81, -0.48]	[-1.00, -0.32]	4.2σ
$C_9^e = -C_{10}^e$	+0.78	[+0.56, +1.02]	[+0.37, +1.31]	4.3σ
$C_9'^\mu$	-0.00	[-0.26, +0.25]	[-0.52, +0.51]	0.0σ
$C_{10}'^\mu$	+0.02	[-0.22, +0.26]	[-0.45, +0.49]	0.1σ
$C_9'^e$	+0.01	[-0.27, +0.31]	[-0.55, +0.62]	0.0σ
$C_{10}'^e$	-0.03	[-0.28, +0.22]	[-0.55, +0.46]	0.1σ

Geng, Grinstein, Jager, Martin Camalich, Ren & Shi,
1704.05446

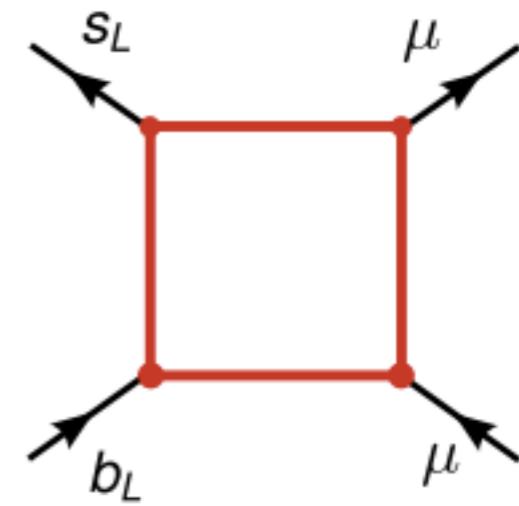
Altmannshofer, Stangl & Straub, 1704.05435

A left hand coupling in the quark sector is preferred

New Physics interpretations

 $\bar{s}\gamma_\rho b_L$  $\bar{\mu}\gamma^\rho \mu$  Z' 

Lepto-quark



Loop

Scale of new physics

W. Almannshofer, Talk at Aspen, Jan. 2016

generic tree

$$\frac{1}{\Lambda_{\text{NP}}^2} (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$$

$$\Lambda_{\text{NP}} \simeq 35 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$$

MFV tree

$$\frac{1}{\Lambda_{\text{NP}}^2} V_{tb} V_{ts}^* (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$$

$$\Lambda_{\text{NP}} \simeq 7 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$$

generic loop

$$\frac{1}{\Lambda_{\text{NP}}^2} \frac{1}{16\pi^2} (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$$

$$\Lambda_{\text{NP}} \simeq 3 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$$

MFV loop

$$\frac{1}{\Lambda_{\text{NP}}^2} \frac{1}{16\pi^2} V_{tb} V_{ts}^* (\bar{s}\gamma_\nu P_L b)(\bar{\mu}\gamma^\nu \mu)$$

$$\Lambda_{\text{NP}} \simeq 0.6 \text{ TeV} \times (C_9^{\text{NP}})^{-1/2}$$

Questions?

- Quark FCNC in the 2 heaviest families
- Lepton Non-Universal interaction in the lightest 2 generations
- If Z' , better same coupling to first 2 families,
- Anomaly free?

Bunches of solutions

Flavored B-L

Only couple to third generation fermions

- Anomaly cancellation within one-generation
- Two Heavy right hand neutrinos still consistent with seesaw and leptogenesis
- No serious flavor changing neutral currents between first two generations
- B anomaly can be explained!

Couplings

Flavor eigenstate

$$T^q = \frac{1}{3} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad T^l = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

Into mass eigenstate (we only assume left-hand mixing here)

$$T_{d_L} = U_{d_L}^\dagger T^q U_{d_L} = \frac{1}{3} \begin{pmatrix} 0 & 0 & 0 \\ 0 & s_{\theta_q}^2 & -s_{\theta_q} c_{\theta_q} \\ 0 & -s_{\theta_q} c_{\theta_q} & c_{\theta_q}^2 \end{pmatrix}$$

$$T_{e_L} = U_{e_L}^\dagger T^l U_{e_L} = -1 \begin{pmatrix} 0 & 0 & 0 \\ 0 & s_{\theta_l}^2 & -s_{\theta_l} c_{\theta_l} \\ 0 & -s_{\theta_l} c_{\theta_l} & c_{\theta_l}^2 \end{pmatrix}$$

- To explain the B anomaly, we need opposite sign coupling with b-s and mu-mu, $\theta_q < 0$
- If assuming $U_{d_L} = V_{CKM}$, it will give wrong sign!

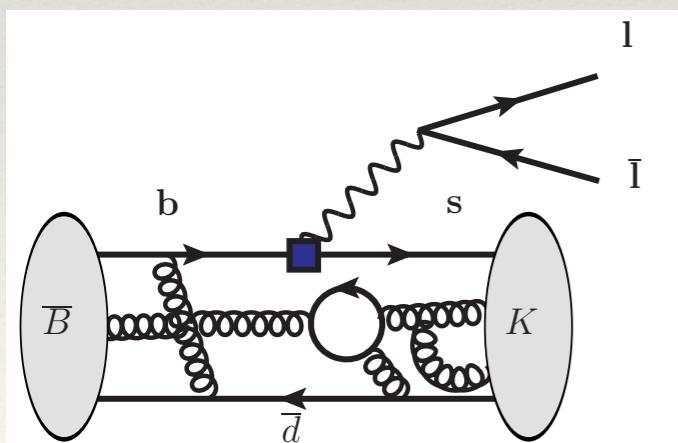
LOW ENERGY Phenomenology

$$S = \int d^4x \left\{ \frac{1}{2} Z_h^\mu (\partial^2 + M^2) Z_{h,\mu} - g_h Z_h^\mu J_\mu^h \right\}$$

$$\text{On-shell } Z_h \int d^4x \left(-\frac{1}{2} \frac{g_h^2 J_h^2}{M^2} + \mathcal{O}(\partial^2/M^2) \right)$$

LFU violation process

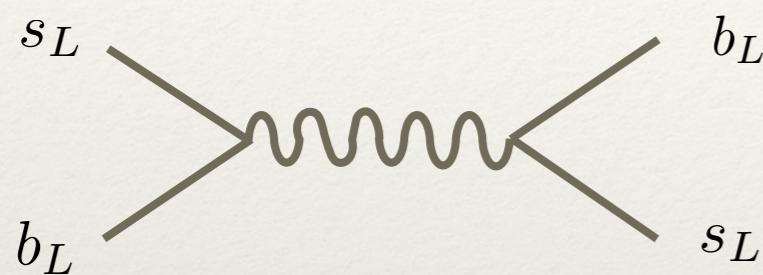
$$\mathcal{L} = -\frac{g^2 s_{\theta_q} c_{\theta_q} s_{\theta_l}^2}{3M^2} (\bar{s} \gamma^\rho b_L)(\bar{\mu} \gamma_\rho \mu_L),$$



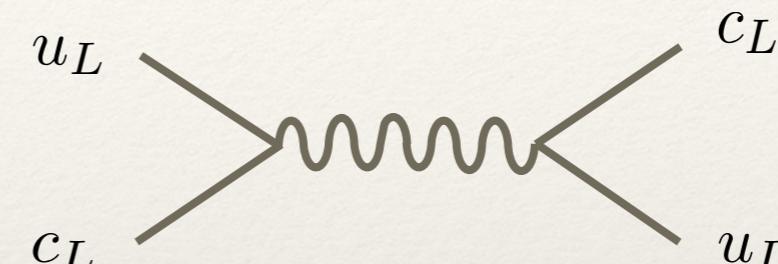
$$\delta C_9^\mu = -\delta C_{10}^\mu = -\frac{\pi}{\alpha \sqrt{2} G_F V_{tb} V_{ts}^*} \frac{g^2 s_{\theta_q} c_{\theta_q} s_{\theta_l}^2}{3M^2}$$

Constraints

$B_s - \bar{B}_s$ mixing



$D^0 - \bar{D}^0$ mixing



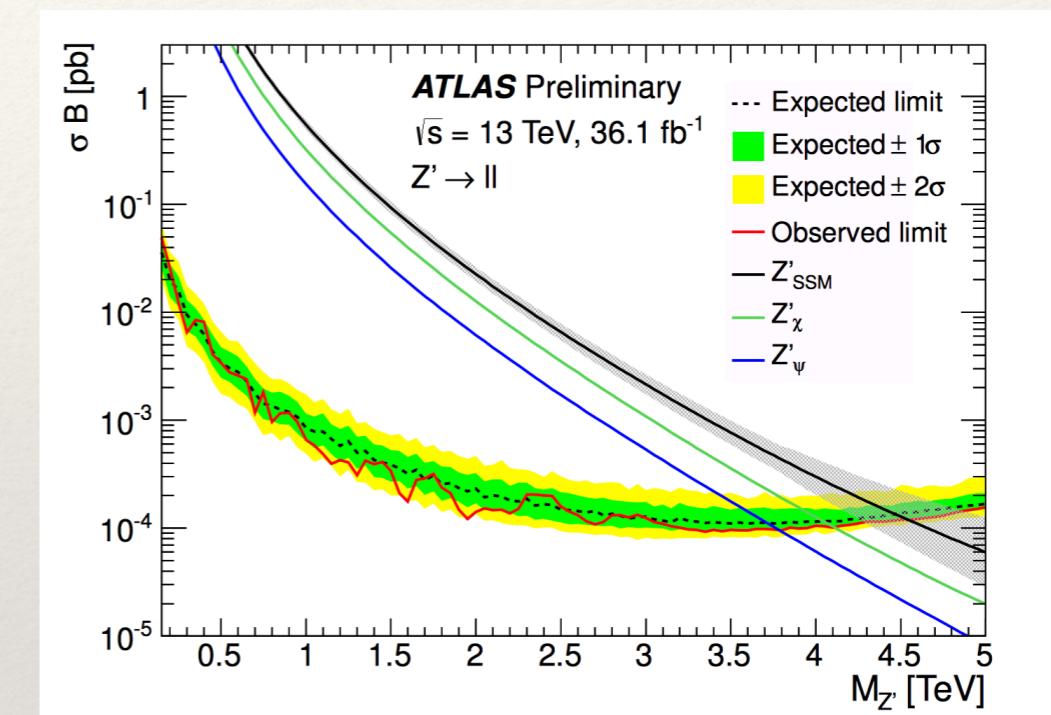
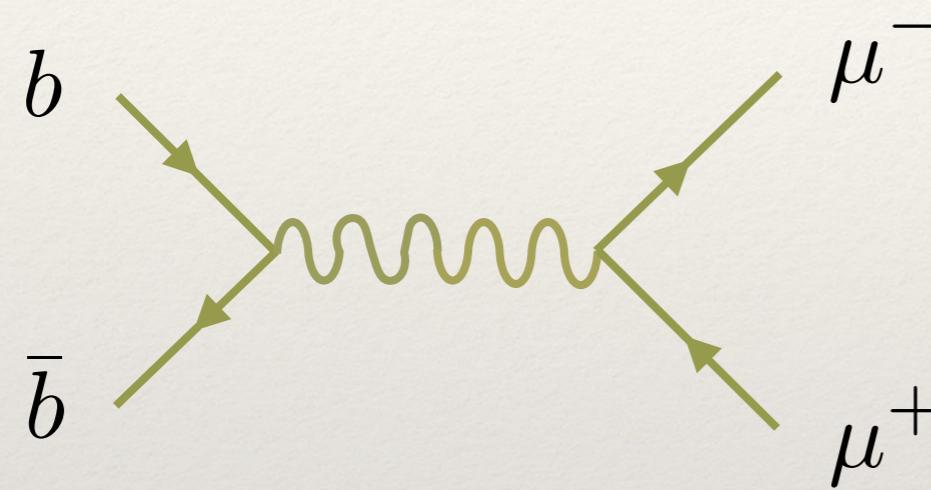
$$\mathcal{L} = -\frac{g^2 s_{\theta_q}^2 c_{\theta_q}^2}{18M^2} (\bar{s} \gamma^\mu b_L)^2 - \frac{g^2 c_D^2}{18M^2} (\bar{u} \gamma^\mu c_L)^2 \quad c_D \equiv (V_{ub} c_{\theta_q} - V_{us} s_{\theta_q}) (V_{cb}^* c_{\theta_q} - V_{cs}^* s_{\theta_q})$$



$$\mathcal{L}_{\text{LFV}} = \frac{g^2}{M^2} s_{\theta_l}^3 c_{\theta_l} \bar{\tau} \gamma^\rho \mu_L \bar{\mu} \gamma_\rho \mu_L$$

LHC dimuon resonance search

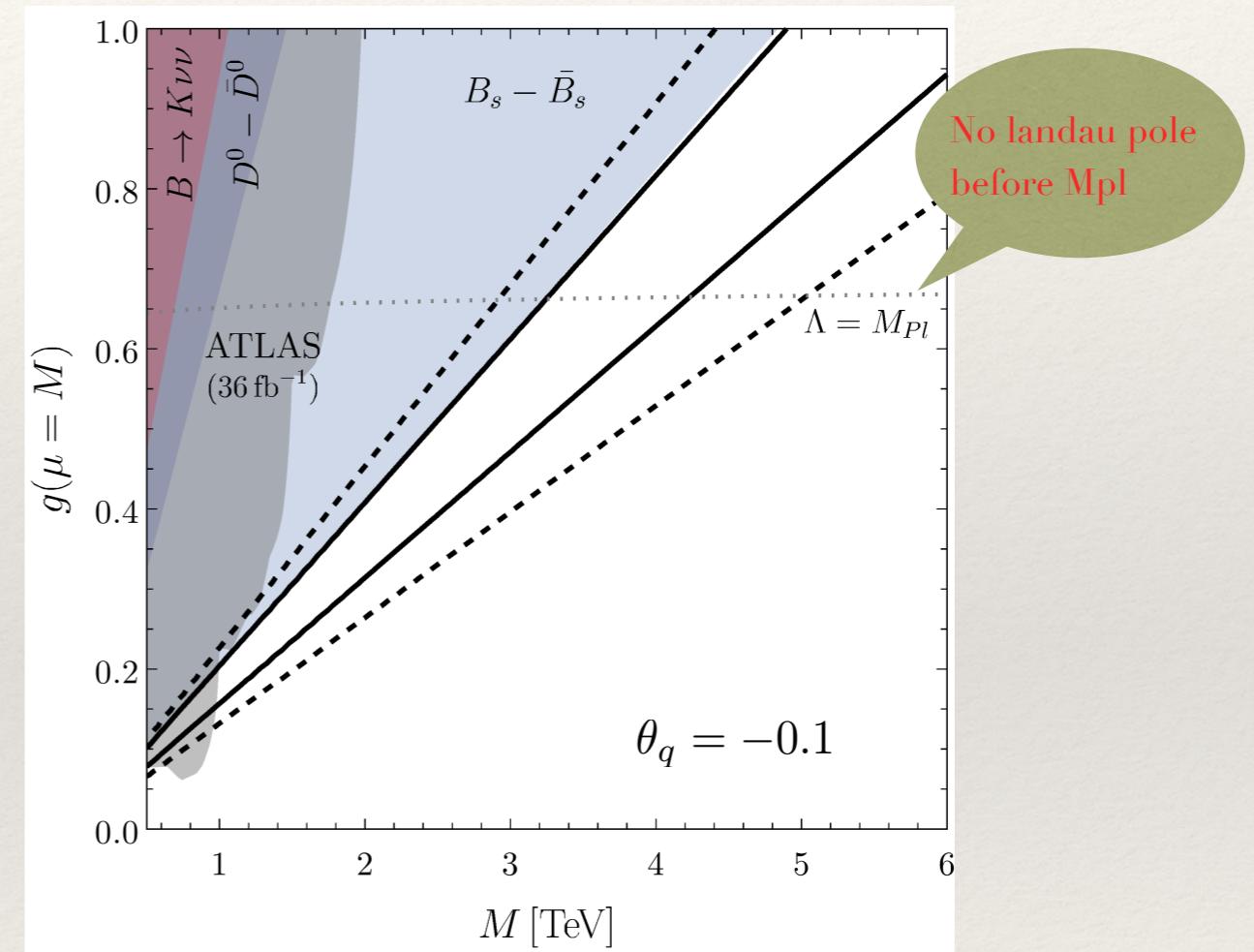
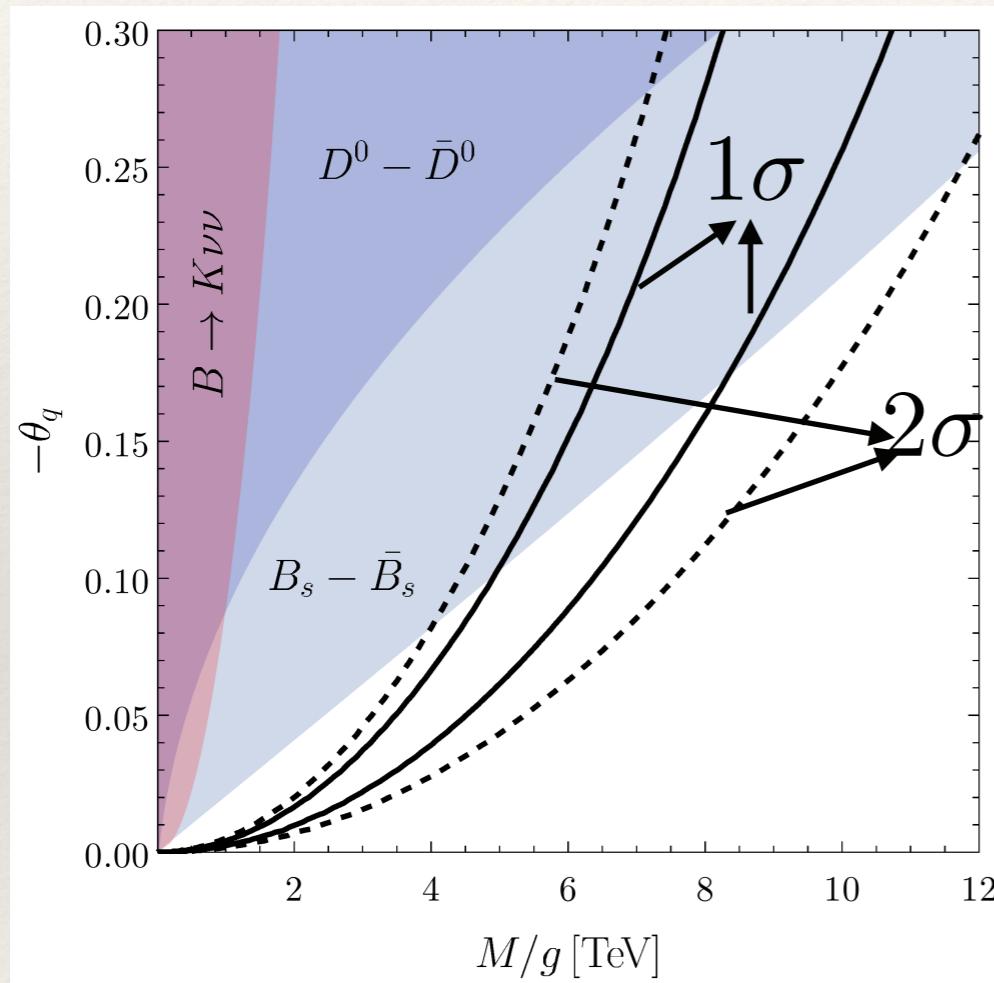
ATLAS-CONF-2017-027



No light quark(u,d) coupling,
limit could be much weaker

Favored parameter space

$\theta_l = \frac{\pi}{2}$ is assumed



- ❖ Z prime mass around few TeV
- ❖ LHC limit less than 2 TeV, could extend in the future to probe the z prime

Summary

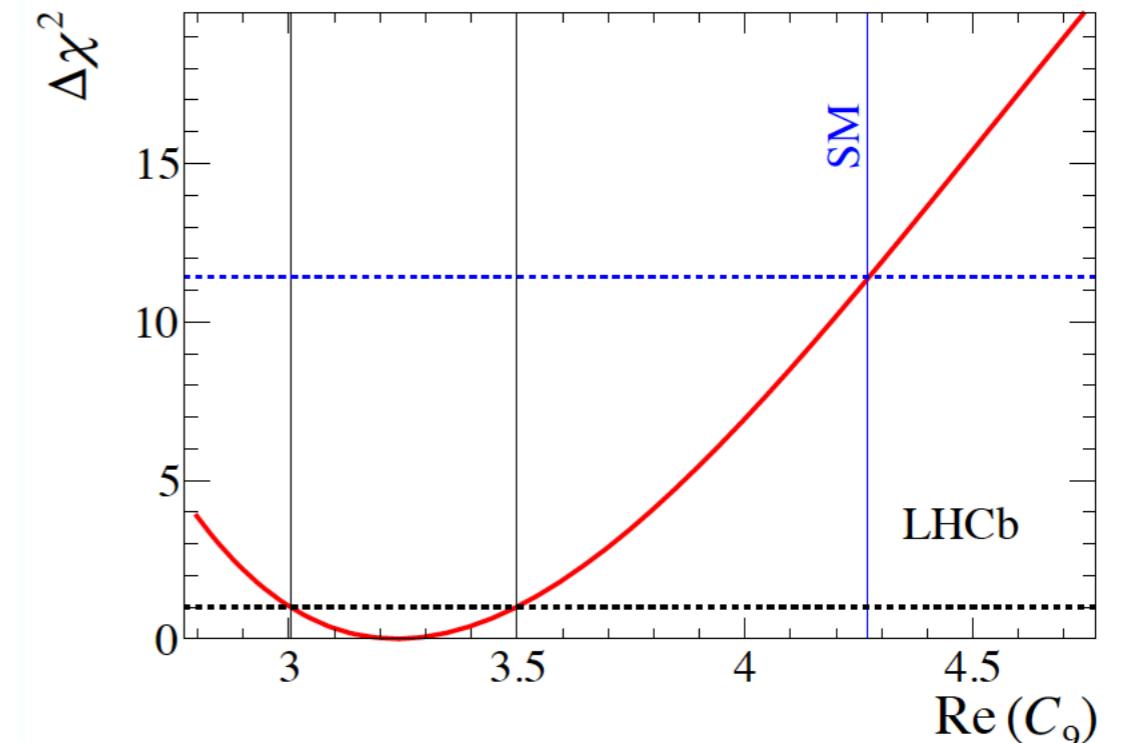
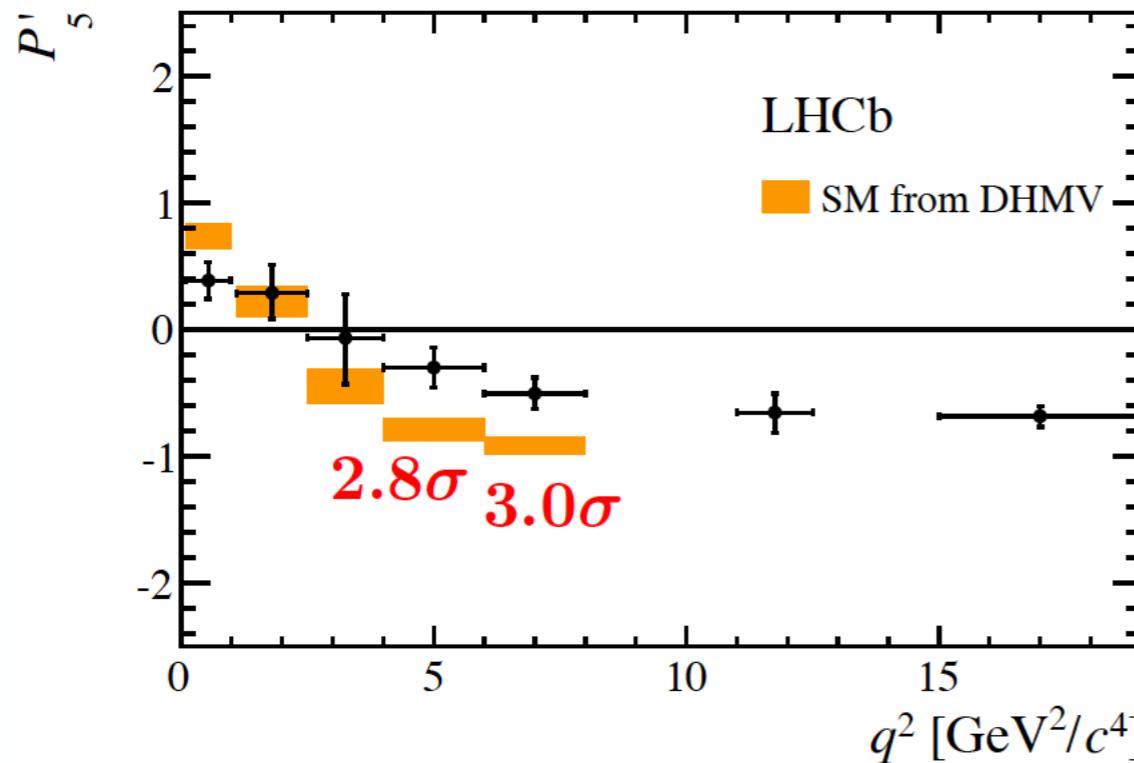
- ❖ B anomaly may indicate a new gauge sector beyond SM
- ❖ This new gauge boson may also play the role of mediator of dark matter(see Peter Cox's talk)
- ❖ Let us wait and see!

Back up

- In 2013, LHCb [1 fb⁻¹] observed a 3.7 sigma discrepancy between the data and the SM in one bin for P5'.
- In 2015, LHCb [3 fb⁻¹] confirmed it with a 3 sigma deviation in each of two bins.
- LHCb also observed a systematic deficit with respect to the SM predictions for the BRs of several decays, such as $B_s \rightarrow \phi\mu\mu$.
- In 2016, Belle confirmed the P5' anomaly.
- Recent ATLAS and CMS data show a good overall agreement with the LHCb results.

Anomaly?

LHCb, 1512.04442



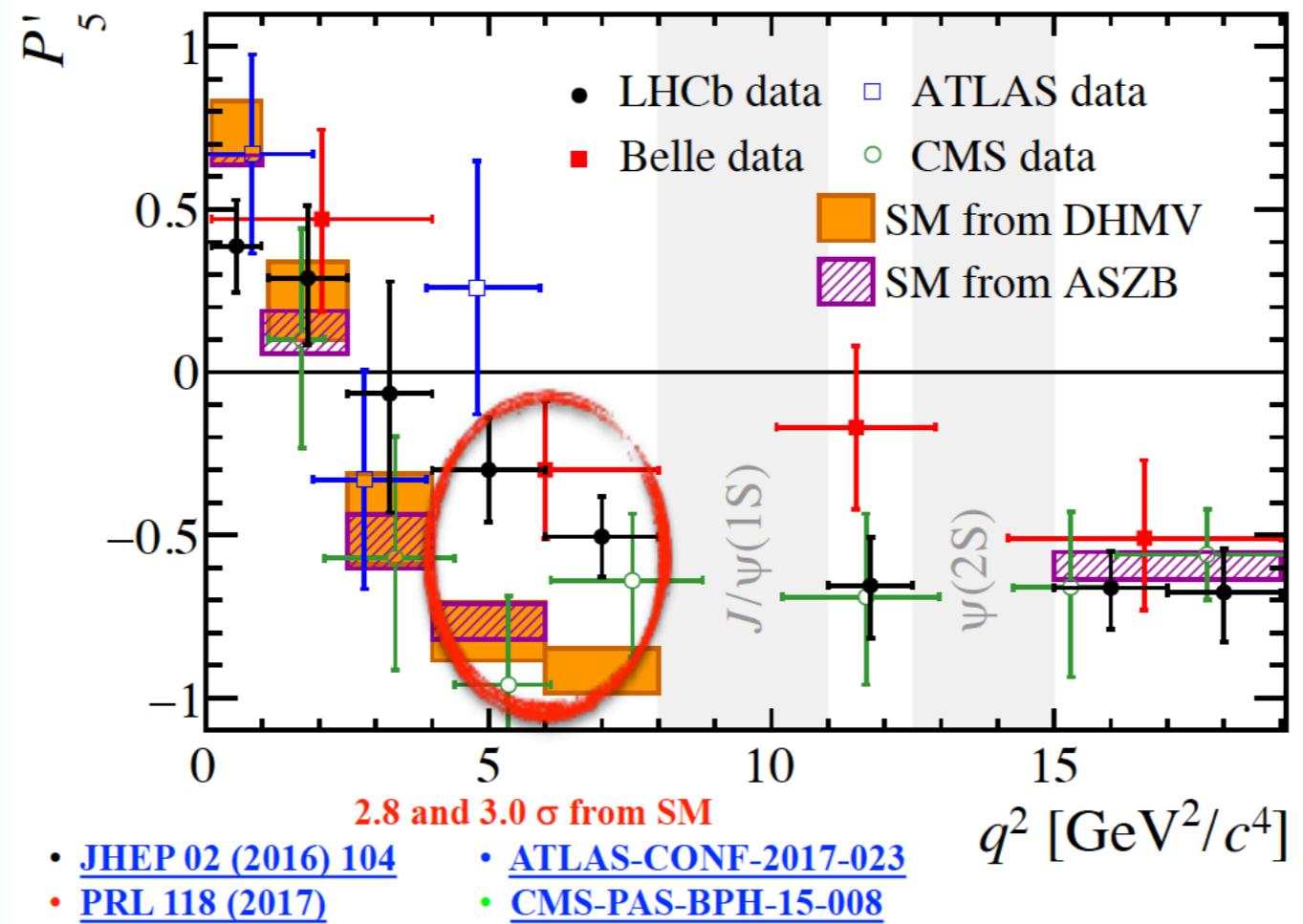
DHMV = Descotes-Genon, Hofer, Matias & Virto (2014)

$$C_9^{\text{NP}} < 0$$

$$O_9 = (\bar{s}\gamma_\mu P_L b)(\bar{\ell}\gamma^\mu \ell)$$

$$|C_9^{\text{NP}}/C_9^{\text{SM}}| \sim 25\%$$

P5' anomaly in $B \rightarrow K^* \mu^+ \mu^-$



Patterns of NP in $R_{K^{(*)}}$

● $B \rightarrow K \ell \ell$

$$\frac{d\Gamma_K}{dq^2} = \mathcal{N}_K |\vec{k}|^3 f_+(q^2)^2 \left(|C_{10}^\ell + C_{10}'^\ell|^2 + \left| C_9^\ell + C_9'{}^\ell + 2 \frac{m_b}{m_B + m_K} C_7 \frac{f_T(q^2)}{f_+(q^2)} - 8\pi^2 h_K \right|^2 \right) \\ + \mathcal{O}\left(\frac{m_\ell^4}{q^4}\right) + \frac{m_\ell^2}{m_B^2} \times O(\alpha_s, \frac{q^2}{m_B^2} \times \frac{\Lambda}{m_b}),$$

f_i, V_i, T_i : form factors

h_i : hadronic contributions

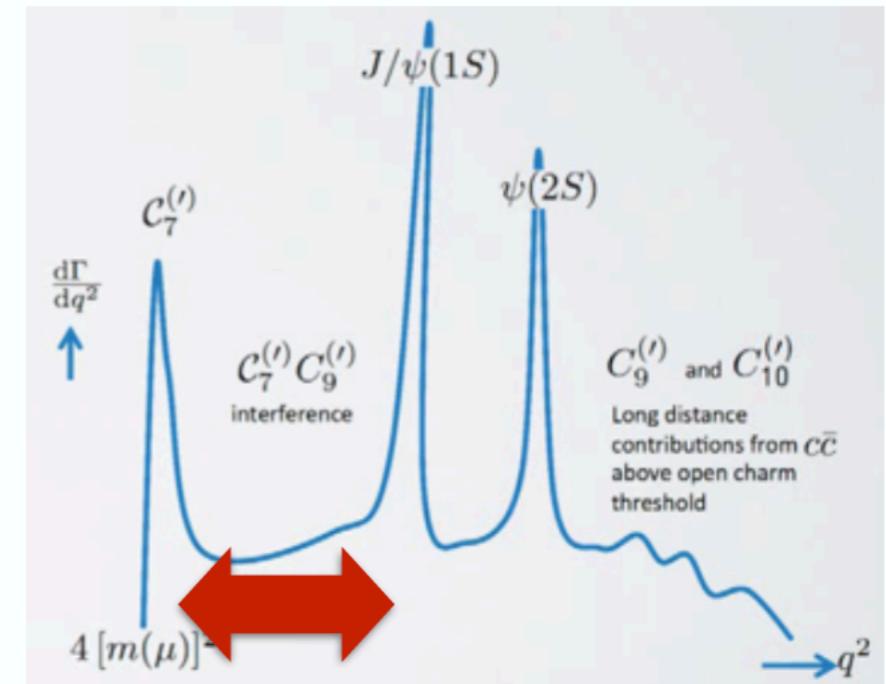
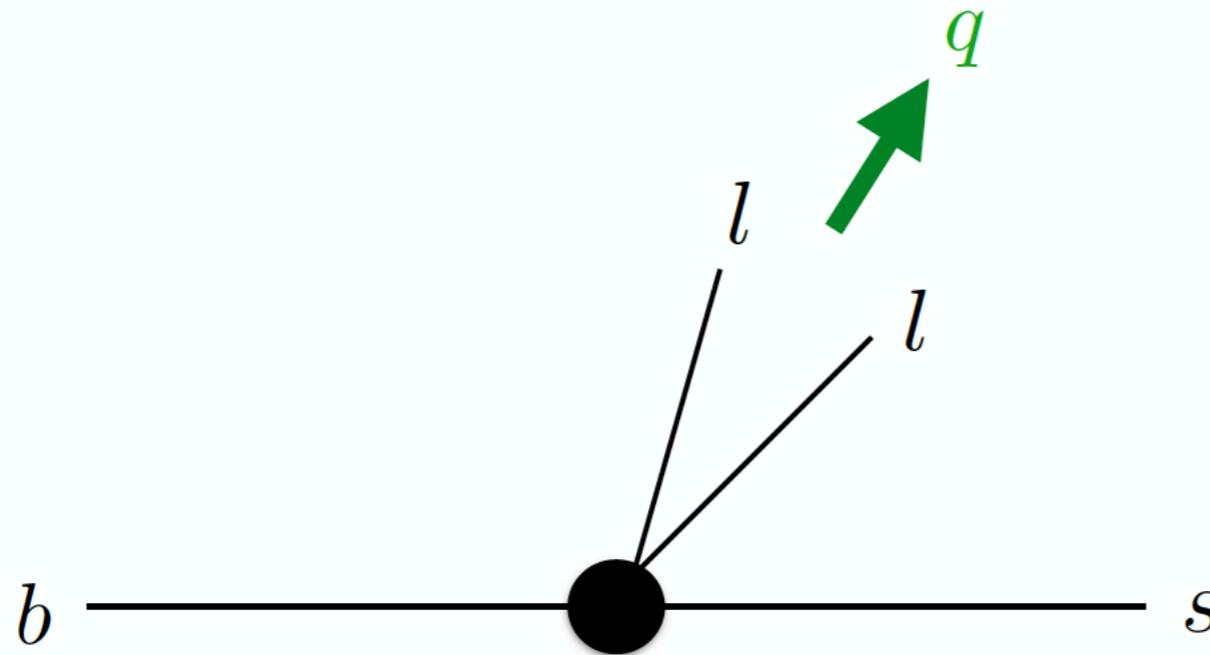
● $B \rightarrow K^* \ell \ell$

$$\frac{d\Gamma_{\bar{K}^*}}{dq^2} = \frac{d\Gamma_\perp}{dq^2} + \frac{d\Gamma_0}{dq^2}$$

$$\frac{d\Gamma_0}{dq^2} = \mathcal{N}_{K^*0} |\vec{k}|^3 V_0(q^2)^2 \left(|C_{10}^\ell - C_{10}'^\ell|^2 + \left| C_9^\ell - C_9'{}^\ell + \frac{2m_b}{m_B} C_7 \frac{T_0(q^2)}{V_0(q^2)} - 8\pi^2 h_{K^*0} \right|^2 \right) + \mathcal{O}\left(\frac{m_\ell^2}{q^2}\right),$$

$$\frac{d\Gamma_\perp}{dq^2} = \mathcal{N}_{K^*\perp} |\vec{k}| q^2 V_-(q^2)^2 \left(|C_{10}^\ell|^2 + |C_9'{}^\ell|^2 + |C_{10}'{}^\ell|^2 + \left| C_9^\ell + \frac{2m_b m_B}{q^2} C_7 \frac{T_-(q^2)}{V_-(q^2)} - 8\pi^2 h_{K^*\perp} \right|^2 \right) + \mathcal{O}\left(\frac{m_\ell^2}{q^2}\right) + \mathcal{O}\left(\frac{\Lambda}{m_b}\right)$$

Recently many papers on b to s II processes appear on arXiv.



Lepton universality in $B^+ \rightarrow K^+ l^+ l^-$ ($l = e, \mu$) ([1406.6482](#))

Angular analysis of $B_0 \rightarrow K_0^ \mu^+ \mu^-$* ([1308.1707](#); [1512.04442](#))

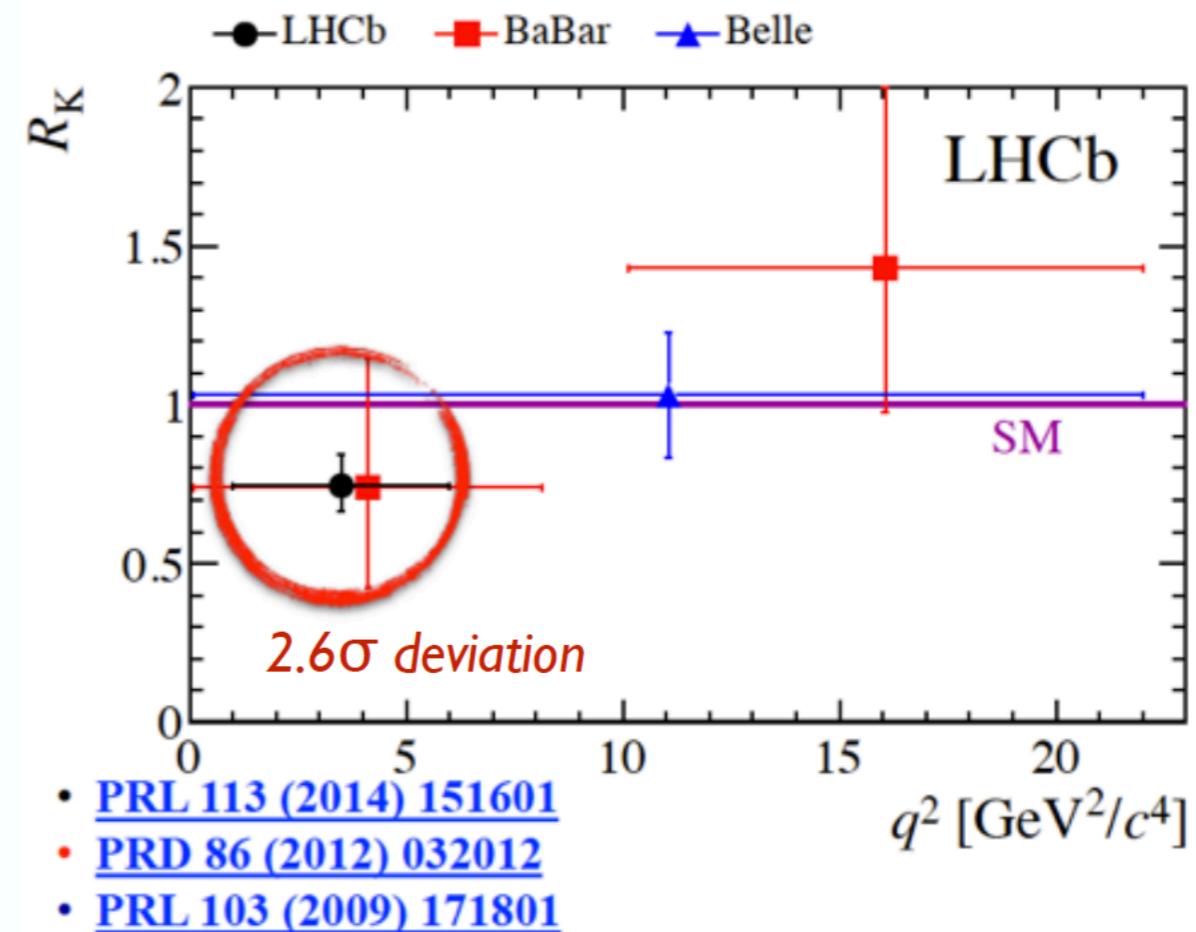
Lepton universality in $B_0 \rightarrow K_0^ l^+ l^-$ ($l = e, \mu$)*

(Talk by Bifani@CERN, 2017.4.18)

Lepton universality in $B^+ \rightarrow K^+ l^+ l^-$ ($l = e, \mu$) (I406.6482)

observable

$$R_K = \frac{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]}{dq^2} dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma[B^+ \rightarrow K^+ e^+ e^-]}{dq^2} dq^2}$$



- $B \rightarrow K ee$ looks consistent (Hiller, Schmaltz I408.1627; LHCb, I406.6482)

$$\mathcal{B}(\bar{B} \rightarrow \bar{K} ee)^{LHCb} = (1.56_{-0.15}^{+0.19+0.06}) \cdot 10^{-7},$$

$$\mathcal{B}(\bar{B} \rightarrow \bar{K} ee)^{SM} = (1.75_{-0.29}^{+0.60}) \cdot 10^{-7},$$

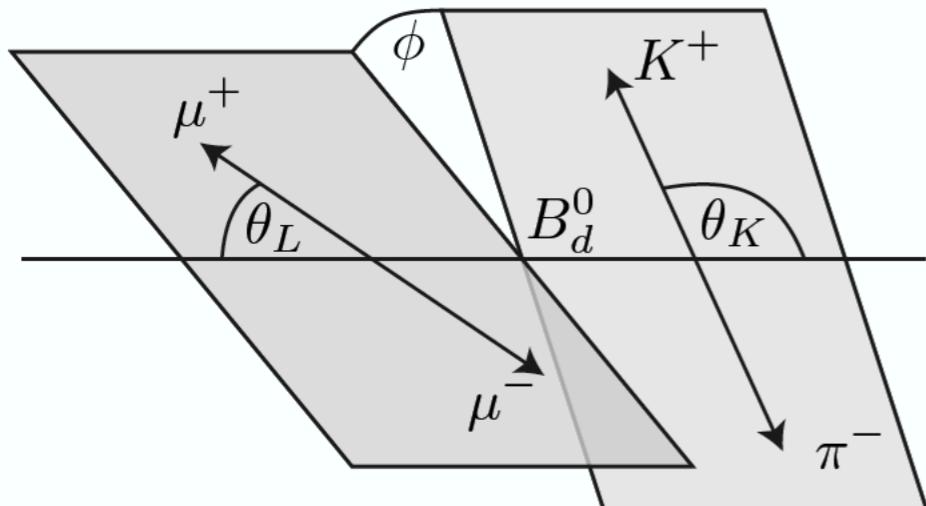
$$\frac{\mathcal{B}(\bar{B} \rightarrow \bar{K} ee)^{LHCb}}{\mathcal{B}(\bar{B} \rightarrow \bar{K} ee)^{SM}} = 0.83 \pm 0.21.$$

with $1 \text{ GeV}^2 \leq q^2 < 6 \text{ GeV}^2$

Angular analysis of $B_0 \rightarrow K_0^* \mu^+ \mu^-$

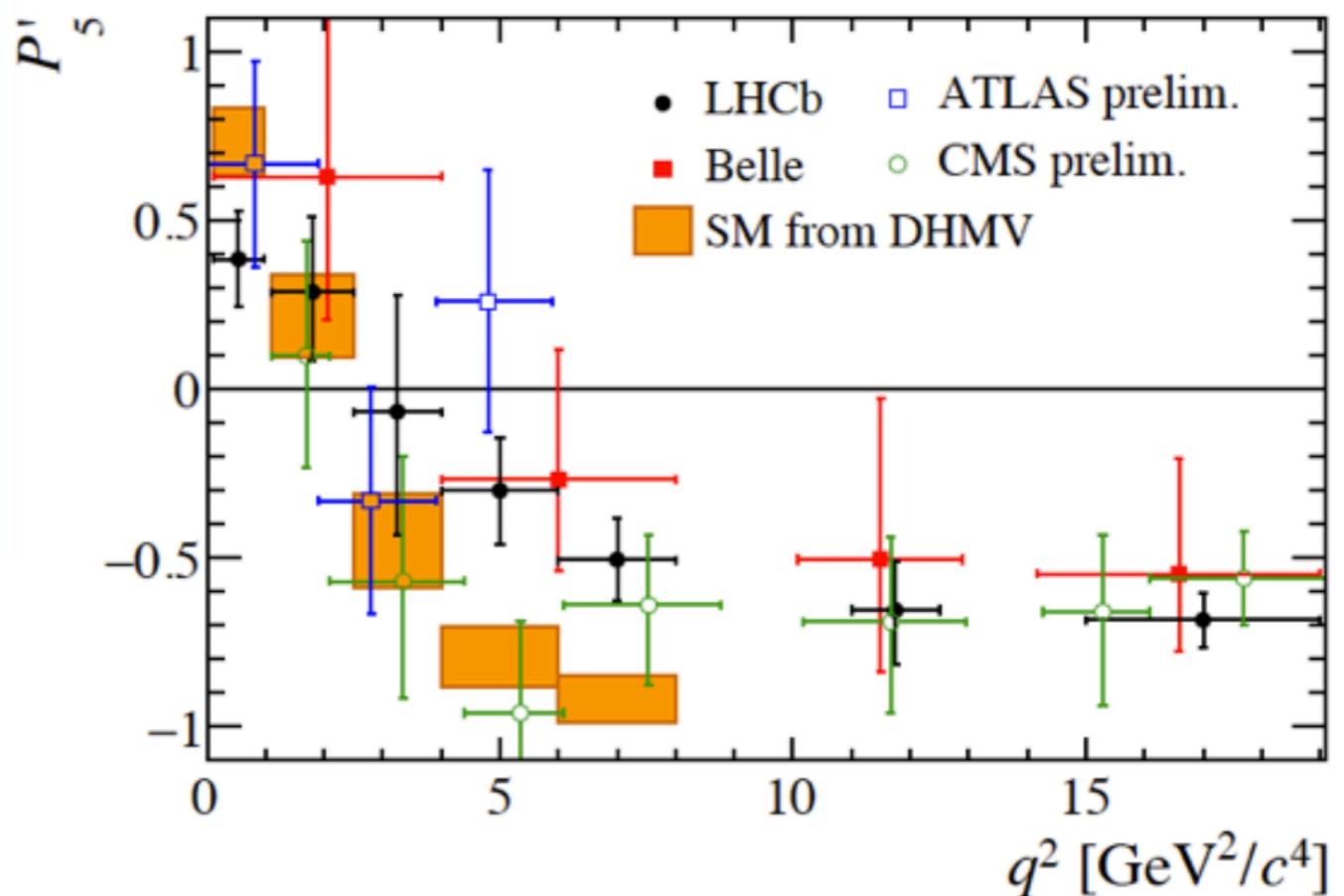
(1308.1707; 1512.04442)

;ATLAS-CONF-2017-023;CMS-PAS-BPH-15-008)



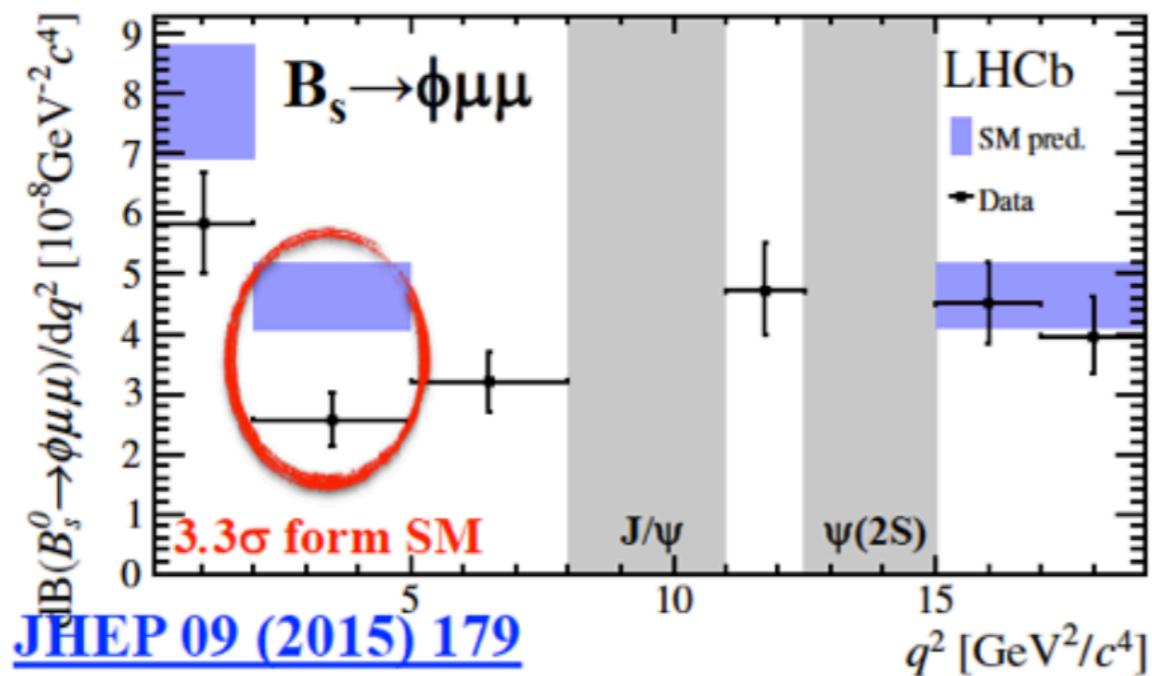
$$\begin{aligned} & \frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d\cos\theta_\ell d\cos\theta_K d\phi dq^2} \\ &= \frac{9}{32\pi} \left[\frac{3}{4}(1 - F_L) \sin^2\theta_K + F_L \cos^2\theta_K + \frac{1}{4}(1 - F_L) \sin^2\theta_K \cos 2\theta_\ell \right. \\ &\quad - F_L \cos^2\theta_K \cos 2\theta_\ell + S_3 \sin^2\theta_K \sin^2\theta_\ell \cos 2\phi \\ &\quad + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos\phi + S_5 \sin 2\theta_K \sin\theta_\ell \cos\phi \\ &\quad + S_6 \sin^2\theta_K \cos\theta_\ell + S_7 \sin 2\theta_K \sin\theta_\ell \sin\phi \\ &\quad \left. + S_8 \sin 2\theta_K \sin 2\theta_\ell \sin\phi + S_9 \sin^2\theta_K \sin^2\theta_\ell \sin 2\phi \right], \end{aligned}$$

(Blake,Gersabeck, et al., 1703.10005)



Branching ratio of $B_s \rightarrow \phi\mu^+\mu^-$

(1506.08777)



[JHEP 09 \(2015\) 179](#)

Branching ratio of $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$

(1503.07138)

