# Search for anomalous alignments of structures in PLANCK data using Tensor Minkowski Functionals 

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Based on Chingangbam, Yogendran, Joby P. K. et al. 2017, and Joby, Ganesan and Chingangbam, in preparation

## Motivation

- Scalar Minkowski Functionals (MF) have been used to test Gaussianity of random fields.
- The trace of the rank two, translation invariant Tensor Minkowski Functionals (TMFs) yield the usual scalar MFs.
- Thus TMFs contain all the information that exists in the scalar MFs.
- Additionally, TMFs contain shape information.
- TMFs can be used for studying the statistical isotropy of fields.


## Motivation

- Numerical work using stereographic projection has showed disagreement between PLANCK 2015 data and $\Lambda$ CDM model at $4-\sigma$ for CMB E mode.
(Vidhya and Chingangbam, 2017)
- Computation of TMF directly on the sphere is important before probing the causes for the disagreement.


## TMFs in flat 2D space

Consider a closed curve in flat 2D space


$$
\begin{aligned}
& \text { where, } \\
& \hat{n} \text { is the unit normal to the curve. } \\
& \vec{T} \text { is tangent vector to the curve. } \\
& \vec{r} \text { is the position vector. } \\
& \kappa \text { denotes the local curvature of the curve. }
\end{aligned}
$$

Translation invariant second rank Tensor Minkowski Functional:

$$
W_{2}^{1,1} \equiv \frac{1}{4} \oint_{C} \vec{r} \otimes \hat{n} \kappa d s
$$

## Closed curves on the sphere

For a field defined on the sphere:

- Red regions have field value greater than mean.
- Blue regions have field value less than the mean.
- Boundaries of iso colour regions form closed curves on the sphere.


## TMF for closed curves on the sphere



- $\hat{n}$ is the unit normal to the curve at the given point on the curve.
- $\vec{T}$ is the tangent vector.
- Re-express $W_{2}^{1,1}$ in terms of $\vec{T}$ and generalize the definition to curved space as,

$$
\mathcal{W}_{1} \equiv \frac{1}{4} \oint_{C} \hat{T} \otimes \hat{T} d s
$$

## Measuring Statistical Isotropy

- Consider a random field, $u$, defined on the unit sphere.
- Choose a threshold, $\nu_{t}$ to obtain the excursion set.
- Find $\mathcal{W}_{1}$ for the boundaries of individual connected regions and holes.
- Obtain $\overline{\mathcal{W}}_{1}$ by averaging individual elements of $\mathcal{W}_{1}$ over all the structures.
- Let $\Lambda_{1}$ and $\Lambda_{2}$ denote the two eigenvalues of $\overline{\mathcal{W}}_{1}$, such that, $\Lambda_{1} \leq \Lambda_{2}$.
- The alignment parameter is defined as,

$$
\alpha \equiv \frac{\Lambda_{1}}{\Lambda_{2}}
$$

## Interpreting $\alpha$

- $\alpha$ measures the net orientation of the structures.
- $\alpha=1$ means that the field is statistically isotropic.
- For a circle, $\alpha=1$, irrespective of the radius.
- For a pair of identical ellipses with their major axes aligned
in the same direction, $0<\alpha<1$.

- For a pair of identical ellipses with their major axes aligned
perpendicularly to each other, $\alpha=1$.



## Analytic computation of $\alpha$ on the sphere

For a pixelated map of a field $u$ defined on the unit sphere, $\overline{\mathcal{W}}_{1}$ is given by,

$$
\overline{\mathcal{W}}_{1}=\frac{1}{4} \int_{S^{2}} d a \delta\left(u-\nu_{t}\right) \frac{1}{|\nabla u|}\left(\begin{array}{cc}
u_{; 2}^{2} & u_{; 1} u_{; 2} \\
u_{; 1} u_{; 2} & u_{; 1}^{2}
\end{array}\right)
$$

where,

- $u_{; i}$ is the $i^{t h}$ component of the covariant derivative of $u$,
- $|\nabla u|=\left(u_{; 1}, u_{; 2}\right)$
- da is the area element,
- $\delta\left(u-\nu_{t}\right)=\frac{1}{\Delta \nu_{t}}$ when $u$ lies between $\nu_{t}-\Delta \nu_{t} / 2$ and $\nu_{t}-\Delta \nu_{t} / 2$ and zero otherwise,
- $\kappa=\frac{2 u_{i 1} u_{; 2}-u_{i}^{2} u_{i 22}-u_{i 2}^{2} u_{i 11}}{|\nabla u|^{3}}$

Chingangbam, Yogendran, Joby P. K. et al. (2017)

## Method

- In this work we apply the analytic $\alpha$ expression to PLANCK 2015 CMB data.
- Study the behaviour of $\alpha$ upon masking the CMB field.
- We use the common mask UT78 for temperature which has a sky fraction of $77.6 \%$.
- CMB FFP9 simulations provided by PLANCK which include instrumental effects such as the beam effect have been used in this analysis.


## Effect of masking



Figure: Left panel: $\alpha$ for $T$ of CMB for unmasked and masked simulations. Average over 100 realizations. The error bars are the sample variance for the 100 realizations. Right panel: Same as left panel for $E$ mode.

- The total number of structures (hotspots + coldspots) is smaller for threshold values further away from 0 .
- Our ordering of $\Lambda$, along with fewer number of structures at larger thresholds, leads to smaller $\alpha$ values at these thresholds.
- Masking reduces the number of hotspots and coldspots and hence $\alpha$ for masked maps are found to be smaller than full sky maps.


## Results: PLANCK T Low Frequency Bands





Figure: Top: Left panel shows a comparison of $\alpha$ for $T$ of CMB obtained from simulations and from PLANCK 30 GHz cleaned maps. Right panel displays the same for PLANCK 44 GHz cleaned maps. Bottom: Same as top for PLANCK 70 GHz cleaned maps.

## Results: PLANCK T High Frequency Bands



Figure: Top panel shows a comparison of $\alpha$ for $T$ of CMB obtained from simulations and from PLANCK $100 \mathrm{GHz}, 143 \mathrm{GHz}$ and 217 GHz cleaned maps. Bottom panel is the same as top for PLANCK $353 \mathrm{GHz}, 545 \mathrm{GHz}$ and 857 GHz cleaned maps.

## Quantifying the difference between Data and Simulations

- Compute $\alpha$ from PLANCK cleaned CMB maps.
- Compute the mean $\alpha, \bar{\alpha}$, and the variance, $\sigma_{\alpha}$, from 100 simulations.
- For each frequency band, compute the quantity, $\mathcal{D}_{\alpha}^{X}$, which is defined as follows,

$$
\mathcal{D}_{\alpha}^{X}=\left|\frac{\alpha_{o b s}^{X}-\bar{\alpha}^{X}}{\sigma_{\alpha}^{X}}\right|
$$

where, $X$ can be T or E mode of CMB polarization.

- $\mathcal{D}_{\alpha}^{X}$ gives a measure of the deviation in statistical isotropy of the observed data from the theory.


## $\mathcal{D}_{\alpha}$ for PLANCK frequency bands

| $\mathcal{D}_{\alpha}^{T}\left(\nu_{t}=0\right)$ | 30 GHz | 44 GHz | 70 GHz | 100 GHz | 143 GHz | 217 GHz | 353 GHz | 545 GHz | 857 GHz |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| COM | 3.5 | 1.4 | 1.4 | 1.0 | 0.8 | 0.9 | 1.0 | 6.5 | 13.5 |
| NILC | 3.3 | 1.2 | 1.8 | 1.4 | 1.3 | 1.3 | 1.4 | 6.9 | 14 |
| SEVEM | 3.5 | 1.5 | 1.4 | 1.0 | 0.9 | 1.0 | 1.0 | 6.6 | 13.6 |
| SMICA | 3.4 | 1.7 | 1.5 | 1.2 | 1.0 | 1.0 | 1.1 | 6.6 | 13.6 |

Table: $\mathcal{D}_{\alpha}\left(\nu_{t}=0\right)$ for CMB Temperature for various PLANCK datasets.

## Conclusions

- No significant alignment found in the PLANCK CMB Temperature data in the mid frequency range.
- Disagreement between $\Lambda$ CDM model and PLANCK observations is seen at frequencies below 44 GHz and those above 545 GHz where foreground removal is expected to be less effective.
- $\mathcal{D}_{\alpha}\left(\nu_{t}=0\right)$ values for CMB Temperature obtained analytically here are comparable to those from Vidhya and Chingangbam (2017).

Ongoing work:

- Comparison of analytic $\alpha$ computed from data and simulations, for E mode of CMB polarization.

