Stable cosmology in chameleon bigravity

CosPA 27, 17.12.14 M. Oliosi (YITP)

Based on

Stable cosmology in chameleon bigravity

arXiv 1711.04655, with A. De Felice, S. Mukohyama, and Yota Watanabe

Outline

- 1. Describe the theory
- 2. Our goal: realistic background cosmology
- 3. Construct the tools
 - i. The action
 - ii. Scaling solutions
 - iii. Stability
- 4. Numerics and results
- 5. Conclusion

Chameleon bigravity

De Felice, Uzan, Mukohyama, 1702.04490

- A theory of 2 gravitons and 1 scalar field
- This extends **massive bigravity** and addresses: Hassan and Rosen, 1109.3515
 - a. Higuchi bound $m_T^2 \gtrsim H^2$
 - b. No fine-tuning needed to pass solar system tests...
 - c. ... and to have an interesting phenomenology

De Felice, Gumrukcuoglu, Mukohyama, Tanahashi, Tanaka, 1404.0008

Goal of the work

Show that the theory can accommodate a "realistic" background cosmology !



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Does everything work as planned?

- Higuchi bound
- Stability
- Modes

The action

Chameleon bigravity side

$$s^{\mu}{}_{\nu} \equiv \left(\sqrt{g^{-1}f}\right)^{\mu}{}_{\nu}$$

$$S_{\rm EH} = \frac{M_g^2}{2} \int R[g] \sqrt{-g} d^4 x + \frac{\kappa M_g^2}{2} \int R[f] \sqrt{-f} d^4 x ,$$

$$S_m = M_g^2 m^2 \int \sum_{i=0}^4 \beta_i(\phi) e_i[s] \sqrt{-g} d^4 x ,$$

$$S_\phi = -\frac{1}{2} \int g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi \sqrt{-g} d^4 x$$

Matter side

$$S_{\rm mat} = \int \mathcal{L}_{\rm mat}(\psi, \tilde{g}_{\mu\nu}) d^4x$$

$$\tilde{g}_{\mu\nu} \equiv A(\phi)^2 g_{\mu\nu}$$

Scaling solutions

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Exponential couplings $A(\phi) = e^{\beta \phi/M_g}, \quad \beta_i(\phi) = -c_i e^{-\lambda \phi/M_g}$

Exact radiation dominated and Λ-dominated solutions

$$ds_g^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j, \quad ds_f^2 = \xi^2(t) \left[-c^2(t)dt^2 + a^2(t)\delta_{ij}dx^i dx^j \right]$$

$$\frac{d\varphi}{dN_e} = \frac{n}{\lambda}, \quad \frac{1}{h}\frac{dh}{dN_e} = -\frac{2}{n}, \qquad \varphi \equiv \phi/M_g, \quad h \equiv \frac{H}{m}, \qquad \xi \equiv const, \quad c = const \qquad N_e = \ln\left(a(t)/a_i\right)$$

Dust-dominated, under condition

$$\beta \left(\lambda^2 - \frac{3c}{c + \kappa \xi^2} \right) = 0$$

• When $\beta \ll 1$ yields an **approximate scaling solution**.

Scaling solutions

Scaling solutions under homogeneous perturbations

$$\begin{cases} \ln h = \ln h_0 - \frac{n}{2}N_e + \epsilon h^{(1)}, \\ \varphi = \frac{nN_e}{\lambda} (1 + \epsilon \varphi^{(1)}), \\ \xi = \bar{\xi} + \epsilon \xi^{(1)}, \\ c = c^{(0)} + \epsilon c^{(1)}, \end{cases}$$

yields

$$\varphi^{(1)''} + \left(1 + \frac{2}{N_e}\right)\varphi^{(1)'} + \mathcal{A}_r\varphi^{(1)} = 0, \qquad ' \equiv \frac{d}{dI}$$
$$\varphi^{(1)''} + \left(\frac{3}{2} + \frac{2}{N_e}\right)\varphi^{(1)'} + \mathcal{A}_m\varphi^{(1)} = 0 \qquad \qquad \mathcal{A}_i > 0$$

 $\overline{\mathbb{V}_e}$

Inhomogeneous perturbations

ADM splitting

 $ds_g^2 = -\mathcal{N}^2 dt^2 + \gamma_{ij} (\mathcal{N}^i dt + dx^i) (\mathcal{N}^j dt + dx^j), \quad ds_f^2 = -\tilde{\mathcal{N}}^2 dt^2 + \tilde{\gamma}_{ij} (\tilde{\mathcal{N}}^i dt + dx^i) (\tilde{\mathcal{N}}^j dt + dx^j)$

Perturbations

$$\begin{split} \phi &= \bar{\phi} + \delta \phi \quad \psi_{\alpha} = \bar{\psi}_{\alpha} + \delta \psi_{\alpha} \\ \mathcal{N} &= N(1 + \Phi), \quad \mathcal{N}_{i} = N_{i} + \delta N_{i}, \quad \gamma_{ij} = a^{2} \delta_{ij} + \delta \gamma_{ij}, \\ \tilde{\mathcal{N}} &= \tilde{N}(1 + \Phi), \quad \tilde{\mathcal{N}}_{i} = \tilde{N}_{i} + \delta \tilde{N}_{i}, \quad \tilde{\gamma}_{ij} = \tilde{a}^{2} \delta_{ij} + \delta \tilde{\gamma}_{ij} \end{split}$$

Decomposition in SO(3) representations

$$\delta N_{i} = Na(\partial_{i}\hat{B} + B_{i}), \quad \delta \gamma_{ij} = a^{2} \left[2\delta_{i}\hat{\Psi} + \left(\partial_{i}\partial_{j} - \frac{\delta_{ij}}{3}\Delta\right)\hat{E} + \partial_{i}\hat{E}_{j} + \hat{h}_{ij} \right]$$

$$\delta \tilde{N}_{i} = \tilde{N}\tilde{a}(\partial_{i}\hat{B} + B_{i}), \quad \delta \tilde{\gamma}_{ij} = \tilde{a}^{2} \left[2\delta_{ij}\hat{\Psi} + \left(\partial_{i}\partial_{j} - \frac{\delta_{ij}}{3}\Delta\right)\hat{E} + \partial_{i}\hat{E}_{j} + \hat{h}_{ij} \right]$$

vector

Inhomogeneous perturbations

2x2 tensor modes

 $c_{T1} = 1, c_{T2} = c$ 2 massive modes $m_T^2 = m^2 \Gamma \frac{c + \kappa \xi^2}{\kappa \xi}$ & 2 massless modes

Non trivial no-ghost condition: c > 0

Non trivial no-ghost condition: _____<u>I > 0</u>

1x2 vector modes

2 massive modes

 $m_V^2 = m_T^2$

 $c_V = m^2 \Gamma \frac{c+1}{2\xi I}$

Non trivial nogradient instability condition: $\Gamma > 0$

2 scalar modes

- massive modes
- non trivial sound speeds

+ matter modes

Non trivial no-ghost condition (large expression)

Non trivial nogradient instability condition (large expression)

Equations for numerics

Set of equations to integrate

$$\left\{egin{aligned} h' &= h'(h,\xi,arphi,arphi')\,, \ arphi'' &= arphi''(h,\xi,arphi,arphi')\,, \ \xi' &= \xi'(h,\xi,arphi,arphi')\,, \end{aligned}
ight.$$

Start near a radiation scaling solution

$$h'_i \approx -2h_i, \quad \varphi'_i \approx \varphi'_{\rm sc} = \frac{4}{\lambda}, \quad \varphi''_i \approx 0, \quad \xi'_i \approx 0$$

Parameters for numerics

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We chose some parameters so that J > 0 is always satisfied

$$c_3c_1 - c_2^2 = \mathfrak{A}, \quad c_1 + 2c_2 + c_3 = \mathfrak{B}$$

Finally we chose the example parameters

$$\begin{split} c_{\rm in} &= \frac{101}{100} \,, \quad c_{\rm V,in}^2 = 1 \,, \quad \mathfrak{A} = 1 \,, \quad \mathfrak{B} = 1 \,, \\ \Omega_{\Lambda i}^{\rm EF} &= 1 \times 10^{-30} \,, \quad \Omega_{di}^{\rm EF} = 1 \times 10^{-5} \,, \quad \Omega_{ki}^{\rm EF} = \frac{3}{200} \,, \quad \Omega_{Vi}^{\rm EF} = \frac{1}{200} \,, \\ \beta &= 1 \times 10^{-2} \,, \quad \lambda = \frac{40}{3} \end{split}$$

Numerical results

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Evolution as planned !

- Radiation dust Λ domination
- Scaling solutions stable
- Small numerical errors





What about the Higuchi bound and the sound-speeds?...

Numerical results

Again just as planned !

- $m_T^2 \gg H^2$ at all times
- Positive sound-speeds, close to 1
- No-ghost conditions satisfied





Promising!

Proof of existence for a stable cosmology in chameleon bigravity !

Summary

- i. Chameleon bigravity solves the fine-tuning problems of bigravity and extends its reach
- ii. Scaling solutions were described
- iii. Stability conditions under homogeneous and inhomogeneous perturbations were found
- iv. The model propagates 2x2 tensor, 1x2 vector, 2 scalar + matter modes
- v. Numerical integration and example background cosmology achieved

Future outlook

A promising model, with avenues for further study !

E.g. constraints from:

- i. More precise background cosmology
- ii. Evolution of perturbations
- iii. Solar-system tests
- iv. Modification of wave-forms of GW due to graviton oscillation



Thank you for your attention !

light



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