Thermal Gravitational Contribution to Dark Matter Production

Yong TANG University of Tokyo

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based on Wu & Tang, 1708.05138, 1604.04701

Evidence of Dark Matter



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New Interaction

Weakly Interacting Massive Particle

- Mass around ~100GeV
- Coupling ~ 0.5
- Correct relic abundance $\Omega \sim 0.3$
- Searches for CDM
 - Collider qq > XXj
 - Direct Xq > Xq
 - Indirect XX > qq
- Theoretically interesting Yong TANG(U.Tokyo)





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What if only Gravity?



- Gravitational interaction is very weak.
- One may wonder whether DM can be produced.
- We shall show gravity can be strong enough to play...



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Gravitational Contributions

• Non-Thermal (*well-studied*)

e.g. Ema, Jinno, Mukaida&Nakayama, 1502.02475,1604.08898 and refs. therein

- Expansion of cosmic background
- QFT in curved spacetime
- Vacuum Fluctuation $n_X \propto H^3$
 - Bogoliubov transformation
- Thermal scattering
 - EFT for *E*<<*M*_p

$$\mathcal{L}_{\rm int} = \frac{\kappa}{2} h_{\mu\nu} T^{\mu\nu},$$

 $\kappa = \sqrt{32\pi G}$

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Wu&Tang, 1604.04701, 1708.05138

Gary, Sandora, Sloth & Palessandro, 1511.03278, 1709.09688

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EFT in Einstein's Gravity

Einstein-Hilbert action

$$S = \int \mathcal{L} d^4 x , \quad \mathcal{L} = \sqrt{-g} \left[\frac{1}{16\pi G} R + L_m \right]$$

$$\mathcal{L}_{\rm int} = \frac{\kappa}{2} h_{\mu\nu} T^{\mu\nu},$$

Energy-Momentum Tensor

$$\begin{split} T_{S}^{\mu\nu} &= -\eta^{\mu\nu}\partial^{\alpha}S^{\dagger}\partial_{\alpha}S + \eta^{\mu\nu}m_{S}^{2}S^{\dagger}S + \partial^{\mu}S^{\dagger}\partial^{\nu}S + \partial^{\nu}S^{\dagger}\partial^{\mu}S, \\ T_{F}^{\mu\nu} &= -\eta^{\mu\nu}\left(\overline{F}i\partial\!\!\!/ F - m_{F}\overline{F}F\right) + \frac{1}{2}\overline{F}i\gamma^{\mu}\partial^{\nu}F + \frac{1}{2}\overline{F}i\gamma^{\nu}\partial^{\mu}F \\ &\quad + \frac{1}{2}\eta^{\mu\nu}\partial^{\alpha}\left(\overline{F}i\gamma_{\alpha}F\right) - \frac{1}{4}\partial^{\mu}\left(\overline{F}i\gamma^{\nu}F\right) - \frac{1}{4}\partial^{\nu}\left(\overline{F}i\gamma^{\mu}F\right), \\ T_{V}^{\mu\nu} &= \eta^{\mu\nu}\left(\frac{1}{4}F^{\alpha\beta}F_{\alpha\beta} - \frac{1}{2}m_{V}^{2}V^{\alpha}V_{\alpha}\right) - \left(F^{\mu\alpha}F^{\nu}{}_{\alpha} - m_{V}^{2}V^{\mu}V^{\nu}\right), \\ T_{\gamma}^{\mu\nu} &= \frac{1}{4}\eta^{\mu\nu}F^{\alpha\beta}F_{\alpha\beta} - F^{\mu\alpha}F^{\nu}{}_{\alpha}. \end{split}$$

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Annihilation Processes

- Boltzmann Equation $\dot{n} + 3Hn \equiv \frac{d(a^3n)}{a^3dt} = C_{col}$
- Reduced to

$$\frac{d(a^3n)}{a^3dt} = \frac{g^2T}{32\pi^4} \int ds \,\sigma \sqrt{s(s-4m^2)} K_1\left(\frac{\sqrt{s}}{T}\right),$$

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• The core

$$\sigma = \frac{\kappa^4}{32\pi s \left(Sg_i^2\right)} \frac{\left|\vec{p}_f\right|}{\left|\vec{p}_i\right|} \mathcal{A}$$

Massless limit $\sigma \propto \kappa^4 s$ Wu&Tang **1604.04701**

$$|\vec{p}_i| = \sqrt{s^2/4 - m^2}, \ |\vec{p}_f| = \sqrt{s^2/4 - M^2}$$

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Various Contributions

Scalar

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$$\begin{split} \mathcal{A}\left(S \to S\right) &= \frac{7m^4 M^4}{30s^2} - \frac{m^2 M^2}{30s} \left(m^2 + M^2\right), \\ &\quad + \frac{1}{40} \left(m^4 + 4m^2 M^2 + M^4\right) + \frac{s}{120} \left(m^2 + M^2\right) + \frac{s^2}{240}, \\ \mathcal{A}\left(F \to S\right) &= -\frac{7m^4 M^4}{15s^2} - \frac{m^2 M^2}{60s} (M^2 - 4m^2) \\ &\quad + \frac{1}{60} \left(2M^4 + 3m^2 M^2 - 3m^4\right) - \frac{s}{240} (4M^2 - m^2) + \frac{s^2}{480}, \\ \mathcal{A}\left(V \to S\right) &= \frac{101m^4 M^4}{30s^2} - \frac{m^2 M^2}{10s} \left(11M^2 + m^2\right) \\ &\quad + \frac{1}{120} \left(19M^4 + 76m^2 M^2 + 49m^4\right) - \frac{7s}{120} \left(m^2 + M^2\right) + \frac{s^2}{80}, \\ \mathcal{A}\left(\gamma \to S\right) &= \frac{1}{120} \left(s - 4M^2\right)^2, \end{split}$$

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Various Contributions

• Fermion $\mathcal{A}(F \to F) = \frac{14m^4M^4}{15s^2} + \frac{m^2M^2}{30s} (m^2 + M^2),$ $-\frac{1}{120} (8m^4 - 3m^2M^2 + 8M^4) - \frac{s}{120} (m^2 + M^2) + \frac{s^2}{160},$ $\mathcal{A}(V \to F) = -\frac{101m^4M^4}{15s^2} + \frac{m^2M^2}{20s} (44M^2 - m^2)$ $-\frac{1}{60} (19M^4 - 19m^2M^2 - 26m^4) - \frac{s}{240} (7M^2 + 52m^2) + \frac{13s^2}{480},$ $\mathcal{A}(\gamma \to F) = \frac{1}{120} (s - 4M^2) (3s + 8M^2),$ • Vector

$$\begin{split} \mathcal{A}\left(V \to V\right) &= \frac{2983m^4M^4}{30s^2} - \frac{293m^2M^2}{10s} \left(m^2 + M^2\right), \\ &+ \frac{1}{120} \left(257m^4 + 1188m^2M^2 + 257M^4\right) - \frac{37s}{40} \left(m^2 + M^2\right) + \frac{29s^2}{240}, \\ \mathcal{A}\left(\gamma \to V\right) &= \frac{13}{120} \left(s - 4M^2\right)^2, \\ \mathcal{A}\left(\gamma \to \gamma\right) &= \frac{s^2}{10}. \end{split}$$

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Parameter Space



Effects of Inflation

 The temperature after inflation is determined by the reheating process, usually the decay of the inflaton.

$$\ddot{\phi} + 3H\dot{\phi} + \Gamma_{\phi}\dot{\phi} + V'(\phi) = 0, \qquad V(\phi) = \frac{1}{2}m_{\phi}^{2}\phi^{2}$$

$$T_{R} = \sqrt{\Gamma_{\phi}M_{P}}$$
Another important effect is from *inflaton annihilation*

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reheating

Annihilation from Inflaton

 The energy density during inflation is much lower than Planck scale

$$\sigma = \frac{\kappa^4}{32\pi s \left(Sg_i^2\right)} \frac{\left|\vec{p}_f\right|}{\left|\vec{p}_i\right|} \mathcal{A}$$

$$\phi$$
 X $h_{\mu\nu}$ X

 $m = m_{\phi}$

 $M = M_X$

- Scalar $\mathcal{A} = \frac{1}{32} \left[2(1 6\zeta)m^2 + M^2 \right]^2$
- Fermion $\frac{1}{16}M^2(m^2 M^2)$
- $\frac{1}{32} \left(4m^4 4m^2 M^2 + 3M^4 \right)$ • Vector
- helicity suppression

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Massless vector 0

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Parameter Space

For massive scalar and vector



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Possible Signatures

• If unstable, decay products can be shown as anomalies in astrophysical observables



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Summary

- Gravitational contributions to dark matter production can be important for non-WIMP case
- We consider the contribution due to thermal **SM** particles' gravitational annihilation
- Inflation plays two important roles
 - Reheating temperature
 - Inflaton's gravitational annihilation
- Possible astrophysical signatures if DM decay.

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Thanks for your attention.

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