

CMB Polarization in Einstein-Aether Theory

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Introduction

◆ Two Big Mysteries of Cosmology

→ Dark Energy & Dark Matter



The Mysteries of Gravity

= Modification of Gravity at long-distance scale?

◆ Modification of Gravity

= Adding the extra degree of freedom

- ✓ scalar : $F(R)$ -gravity, Galileon theory...
- ✓ tensor : massive gravity, bi-gravity theory...
- ✓ vector : ???



Einstein-Aether Theory!

Einstein-Aether Theory

- ✓ Action = Einstein-Hilbert + Fixed Norm Vector Field with a Lorentz Violating VEV. [Jacobson and Mattingly (1999)]

$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} [R + \mathcal{L}_A] + \int d^4x \sqrt{-g} \mathcal{L}_m$$
$$\left[\begin{aligned} \mathcal{L}_A = & -[c_1 \nabla_a A^b \nabla^a A_b + c_2 \nabla_a A^a \nabla_b A^b + c_3 \nabla_a A^b \nabla_b A^a \\ & + c_4 A^a A^b \nabla_a A^c \nabla_b A_c] + \lambda (A^a A_a - 1) \end{aligned} \right]$$

** $A_a \rightarrow \partial_a \phi$: correspondence with Horava-Lifshitz gravity.

- ✓ Stability Analysis of Perturbation [Armendariz-Picon et al. (2010)]

scalar sector : $-2 \leq c_1 + c_4 \leq 0, c_1 + c_2 + c_3 < 0$

tensor sector : $c_1 + c_3 > -1$

vector sector : $2c_1 \leq (c_1 + c_3)^2 (1 + c_1 + c_3)$

What is the observational consequence of this new transverse vector degree?

 **CMB B-mode Polarization!**

CMB B-mode Polarization

- ✓ B-mode is the curl-like component of CMB polarization.
 - ✓ Now, so many projects are attempting to detect the CMB B-mode polarization → exciting era for B-mode physics!
 - ✓ Possible origin of B-mode Polarization
 - ✓ Tensor : Primordial Gravitational Waves (PGW)
 - ✓ Vector : Source models such as cosmic strings or other defects
 - ** In GR without sources, vector mode only decays...
 - ✓ Vector modes are more efficient to produce B-mode polarization than tensor modes. [Hu and White (1997)]
 - ↔ difficult for usual vector models to make a definite prediction.
- Aether Field Perturbation can be generated in a same way as PGW during inflation and the subsequent evolution is clearly understood!

Our Works

- ✓ Formulate the linear perturbation of the vector mode in Einstein Aether theory using the **covariant formalism**.
- ✓ Derive the **initial condition** in the early Radiation Dominated stage considering the ordinary matter and Aether.
- ✓ Calculate the **CMB B-mode Power spectrum** due to the vector mode numerically using modified CAMB code
- ✓ Understand the shape of the spectrum in an **analytic way**.

Basic Quantities for Aether Field

✓ Energy-Momentum Tensor : $T_{ab}^{(A)} = -\frac{1}{\sqrt{-g}} \frac{\delta(\sqrt{-g}\mathcal{L}_A)}{\delta g^{ab}}$

$$T_{ab}^{(A)} = \nabla_c \left(J_{(a}^c A_{b)} - J^c_{(a} A_{b)} - J_{(ab)} A^c \right) + Y_{ab} + \frac{1}{2} g_{ab} \mathcal{L}_A + \lambda A_a A_b + c_4 A^c A^d (\nabla_c A_a)(\nabla_d A_b)$$

$$\left(\begin{array}{l} J^a_b = -[c_1 \nabla^a A_b + c_2 \delta_b^a \nabla_c A^c + c_3 \nabla_b A^a + c_4 A^a A^c \nabla_c A_b] \\ Y_{ab} = -c_1 [(\nabla_c A_a)(\nabla^c A_b) - (\nabla_a A_c)(\nabla_b A^c)] \end{array} \right)$$

✓ Equation of motion for Aether Field :

$$c_1 \nabla_c \nabla^c A_a + c_2 \nabla_a \nabla_b A^b + c_3 \nabla_b \nabla_a A^b + c_4 [\nabla_b (A^b A^c \nabla_c A_a) - A^b (\nabla_b A_c)(\nabla_a A^c)] = -\lambda A_a$$

✓ Fixed Norm Constraint : $A_a A^a = 1$

✓ Einstein equation : $G_{ab} = T_{ab}^{(A)} + 8\pi G T_{ab}^{(m)}$

✓ Useful Parameter Definition

$$c_{13} = c_1 + c_3, \quad c_{14} = c_1 + c_4, \quad \alpha = c_1 + 3c_2 + c_3$$

Background Cosmology

[Carroll and Lim (2004)]

Flat FRW Background : $ds^2 = a^2(\eta)[d\eta^2 - d\vec{x}^2]$, $A^a = \left(\frac{1}{a}, 0, 0, 0\right)$

$$\text{Friedman equations : } \mathcal{H} = \frac{8\pi G_{\text{cos}}}{3} a^2 \rho, \quad \mathcal{H}' = -\frac{4\pi G_{\text{cos}}}{3} a^2 (\rho + 3p)$$

$$G_{\text{cos}} = \frac{G}{1 - (c_1 + 3c_2 + c_3)/2}$$

Effect of Aether can be incorporated into the cosmological Gravitational Constant (G.C.).

Note that the fundamental G.C. , which appears in the analysis of perturbations, is different from the G.C. of the background cosmology.

** The G.C. that the local experiment measures is $G_N = \frac{G}{1 + (c_1 + c_4)/2}$

➔ Constraint from BBN : $\left| \frac{G_{\text{cos}} - G_N}{G_N} \right| < 0.1$

Primordial Spectrum of Aether Vector Perturbation

[Armendariz-Picon et al. (2010)]

Second order action without matter

$$S^{(2)} = \int d\eta d^3x \frac{(-c_{14})}{16\pi G} \left[\xi'^2 - c_v^2 \partial_i \xi \partial^i \xi - \frac{\alpha}{c_{14}} (\mathcal{H}^2 - \mathcal{H}') \xi^2 \right]$$

$$\left(\xi = a \boxed{V} \quad \begin{array}{l} \text{Vector Perturbation} \\ \text{for Aether Field} \end{array} \quad c_v^2 = \frac{c_1}{c_{14}} \left[1 - \frac{c_{13}^2}{2c_1(1+c_{13})} \right] \right)$$

- ✓ During inflation $a \propto (-\eta)^q$, we have the positive frequency solution,

$$V = \frac{\sqrt{8\pi G}}{2} \sqrt{\frac{\pi}{-c_{14}}} \frac{\sqrt{-\eta}}{a} H_{\nu/2}^{(1)}(-c_v k \eta) \quad \left(\nu = \sqrt{1 - 4q(q+1) \frac{\alpha}{c_{14}}} \right)$$

- ✓ Defining the dimensionless power spectrum as $\langle |V(k, \eta)|^2 \rangle \equiv \frac{2\pi^2}{k^3} \mathcal{P}_V(k, \eta)$, we can derive

$$\mathcal{P}_V(k, \eta) = \mathcal{A}_V(\eta) \times \left(\frac{k}{k_0} \right)^{3-\nu}$$

** The amplitude and tilt depends on the details of the inflation ($\nu = 1$ for dS inflation).

Covariant Formalism

[Hawking (1966), Olsen (1976), ...]

✓ 1+3 Decomposition using

- ✓ Fundamental observer : u_a ($u_a u^a = 1$)
- ✓ Projection tensor : $h_{ab} = g_{ab} - u_a u_b$

Merits

- Gauge invariant
- Non-linear effect
- CAMB !

✓ 10 Fundamental Variables

✓ Energy-Momentum Tensor

$$\rho = T_{ab} u^a u^b, \quad p = -\frac{1}{3} h^{ab} T_{ab}, \quad q_a = T_{bc} h_a^b u_b, \quad \pi_{ab} = \left[h_{(a}^c h_{b)}^d - \frac{1}{3} h^{cd} h_{ab} \right] T_{cd}$$

✓ Covariant Derivative of 4-velocity

$$K_a = \dot{u}_a = u_b \nabla^b u_a, \quad \Theta = D^a u_a, \quad \sigma_{ab} = D_{\langle a} u_{b \rangle}, \quad \omega_{ab} = D_{[a} u_{b]}$$

$$q_a^{(i)} = (\rho^{(i)} + p^{(i)}) v_a^{(i)}$$

✓ Decomposition of Weyl Tensor

$$E_{ab} = C_{acbd} u^c u^d, \quad H_{ab} = \frac{1}{2} \epsilon_{acd} C^cd_{be} u^e$$

** In Background FRW universe, only Θ, ρ, p have non-vanishing values.

✓ Projective Covariant Derivative

$$\dot{Q}^{a\dots}_{b\dots} = u^c \nabla_c Q^{a\dots}_{b\dots}$$

$$D_c Q^{a\dots}_{b\dots} = h^f_c h^a_d \dots h^e_b \dots \nabla_f Q^{d\dots}_{e\dots}$$

Basic Equations

□ From Ricci identities and contracted Bianchi identities using Einstein equation, we obtain the following equations.

Constraint eqs.

$$D^a \sigma_{ab} - \frac{1}{2} \text{curl} \Omega_a - \frac{2}{3} D_b \Theta - \kappa q_b = 0$$

$$D^a E_{ab} - \kappa \left(\frac{\Theta}{3} q_b + \frac{1}{3} D_b \rho + \frac{1}{2} D^a \pi_{ab} \right) = 0$$

$$D^a H_{ab} - \frac{1}{2} \kappa [(\rho + p) \Omega_b + \text{curl} q_b] = 0$$

$$H_{ab} - \text{curl} \sigma_{ab} + \frac{1}{2} D_{\langle a} \Omega_{b \rangle} = 0$$

Evolution eqs.

$$\dot{\Omega}_a + \frac{2}{3} \Theta \Omega_a = \text{curl} K_a$$

$$\dot{\sigma}_{ab} + \frac{2}{3} \Theta \sigma_{ab} = -E_{ab} - \frac{1}{2} \kappa \pi_{ab}$$

$$\dot{E}_{ab} + \Theta E_{ab} = \text{curl} H_{ab} + \frac{\kappa}{2} \left[\dot{\pi}_{ab} - (\rho + p) \sigma_{ab} + \frac{\Theta}{3} \pi_{ab} \right]$$

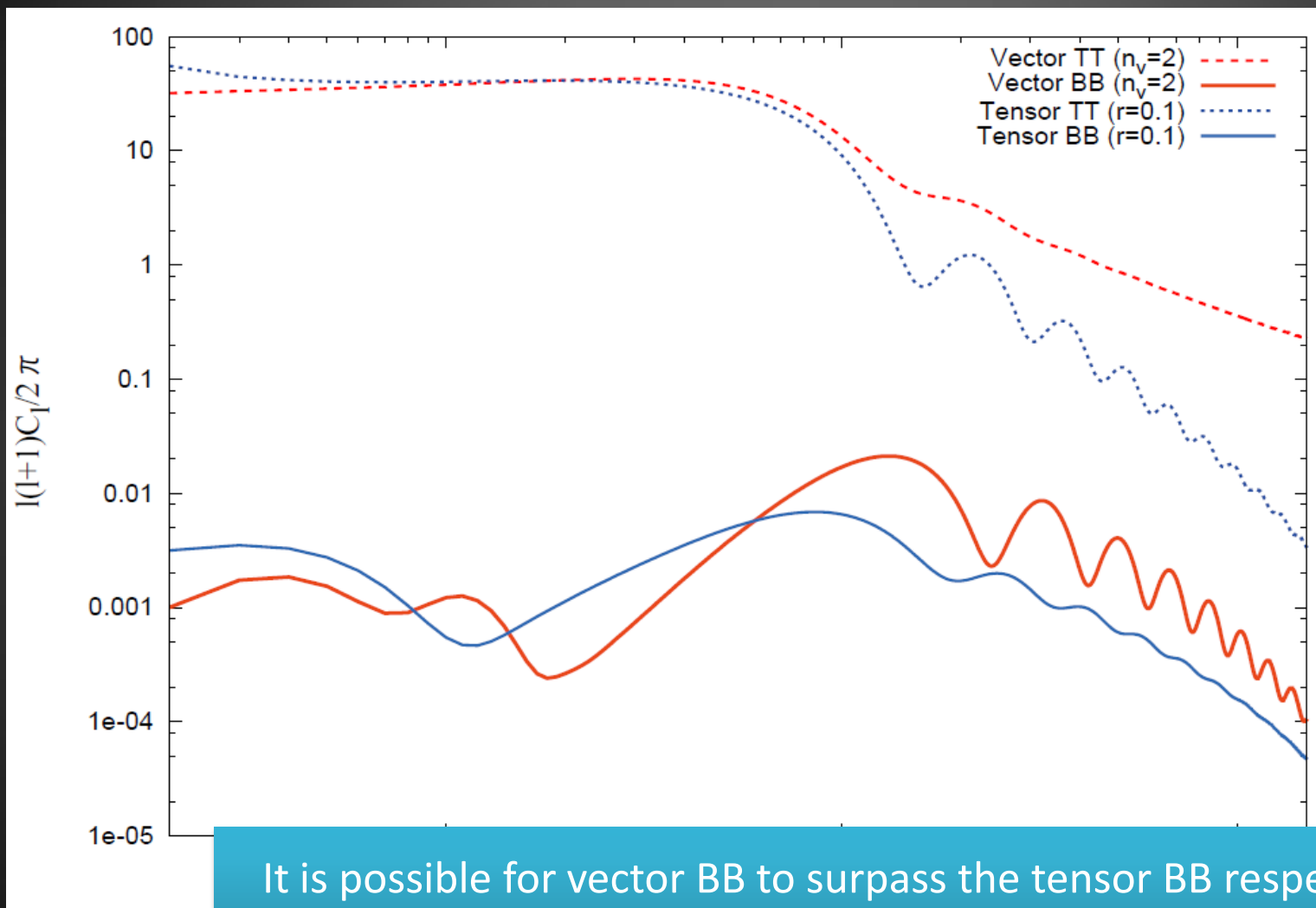
$$\dot{H}_{ab} + \Theta H_{ab} = -\text{curl} E_{ab} - \frac{\kappa}{2} \text{curl} \pi_{ab}$$

Conservation of Energy-Momentum Tensor

$$\dot{\rho} + \Theta(\rho + p) + D^a q_a = 0 \quad \dot{q}_a + \frac{4}{3} \Theta q_a + (\rho + p) K_a - D_a p + D^b \sigma_{ab} = 0$$

** In the above equations, we have dropped the higher order terms
(left only the linear quantity around the flat FRW universe).

Numerical Result from Modified CAMB



$$\mathcal{P}_V = \mathcal{A}_k \left(\frac{k}{k_0} \right)^{n_v}$$

$$\mathcal{P}_T \propto k^0$$

$$c_{13} = -0.3$$

$$c_1 = -0.1$$

$$c_{14} = -0.2$$

$$\alpha = 0.2$$

It is possible for vector BB to surpass the tensor BB respecting the existing observational constraint from the TT amplitude.

Analytic Approach -1-

Photon Multipole Equations in covariant formalism

$$I'_\ell = -k \frac{\ell}{2\ell+1} \left[\frac{\ell+2}{\ell+1} I_{\ell+1} - I_{\ell-1} \right] - \dot{\tau} \left(I_\ell - \frac{4}{3} \delta_{\ell 1} v - \frac{2}{15} \zeta \delta_{\ell 2} \right) + \frac{8}{15} k \sigma \delta_{\ell 2}$$

$$E'_\ell = -\frac{(\ell+3)(\ell+2)\ell(\ell-1)}{(\ell+1)^3(2\ell+1)} k E_{\ell+1} - \frac{\ell}{2\ell+1} k E_{\ell-1} - \frac{1}{\ell(\ell+1)} k B_\ell - \dot{\tau} \left(E_\ell - \frac{2}{15} \zeta \delta_{\ell 2} \right)$$

$$B'_\ell = -\frac{(\ell+3)(\ell+2)\ell(\ell-1)}{(\ell+1)^3(2\ell+1)} k B_{\ell+1} + \frac{\ell}{2\ell+1} k B_{\ell-1} - \frac{2}{\ell(\ell+1)} k E_\ell - \dot{\tau} B_\ell$$

➔ Integral Solution for B-mode Polarization

$$B_\ell(\eta_0) = -\frac{\ell-1}{\ell+1} \int^{\eta_0} d\eta \dot{\tau} e^{-\tau} \frac{\ell j_\ell[k(\eta_0 - \eta)]}{k(\eta_0 - \eta)} \zeta$$

$$\zeta \equiv \frac{3}{4} I_2 - \frac{9}{2} E_2$$

Note that $\dot{\tau} e^{-\tau} \sim \delta(\eta - \eta_{rec})$, it is important to know ζ at the recombination epoch.

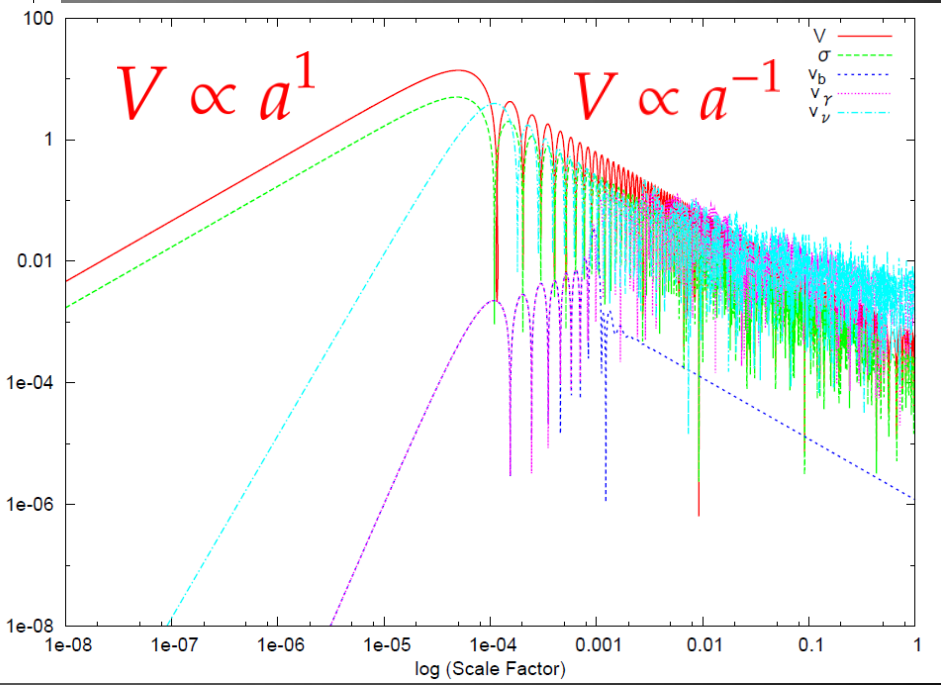
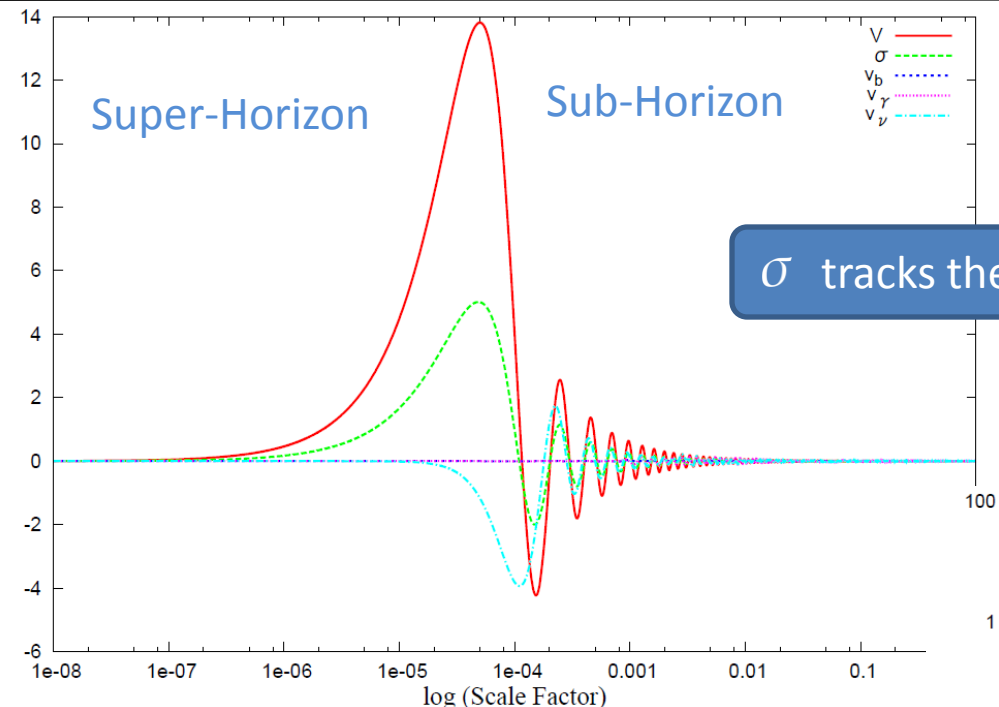
➔ Tight-Coupling Expansion

Evolution of Variables

$$\alpha = -c_{14}$$

$$k = 0.1 [\text{Mpc}^{-1}]$$

σ tracks the evolution of V



Aether vs Scale-factor
in normal plot

Aether vs Scale-factor
in log plot

Summary and Discussion

- ✓ We have computed the power spectrum of CMB B-mode Polarization in the Einstein-Aether Theory.
- ✓ To do that, we have employed the covariant formalism and derived the regular initial conditions.
- ✓ The shape of the spectrum is clearly understood in an analytic way using tight-coupling expansion.
- ✓ More quantitative argument for whether the spectrum can be observed or not and the search for the consistent parameter range considering other experiments will be future works.
- ✓ Are there any cosmological Aether effects? (e.g. Weak Lensing?)