A Light Curve Model for Interaction-Powered Supernovae (in progress)

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Interaction-powered Supernovae (Type IIn)





http://astronomy.swin.edu.au/cosmos/T/Type+IIn+Supernova

~10% of SNe that have 'n'-arrow hydrogen line feature(s) in the spectrum

Huge mass loss (probably) needed

 $10^{-3} - 1M_{\odot}/yr$ Wind Velocities, Mass-loss Rates, and References for SNe IIn in the Literature Supernova Unshocked Wind Shocked Wind Mass-loss Rate References Velocity Velocity $(M_{\odot} \text{ yr}^{-1})$ $(km s^{-1})$ $({\rm km}~{\rm s}^{-1})$ 10^{-2} SN 1987F 1506000 [1] [2] $7 \times 10^{-4} - 1.5 \times 10^{-2}$ SN 1988Z 1200 - 1800< 200[3] [4] SN 1994W 1000 ~ 4000 0.3 [5] [6] 10^{-3} SN 1994aj 1000 ~ 3700 [7] SN 1995G ~ 1000 3000-4000 0.1[8] [9] [10] SN 1995N 2500-5000 2×10^{-4} < 500[11] [12] 10^{-3} SN 1996L 1600 4800 [13] 10^{-2} SN 1997ab 90 6600 [14] 8.3×10^{-3} SN 1997eg [15] 1607000 $10^{-4} - 10^{-3}$ SN 1998S 30 - 100N/A [16]-[20] SN 2005gl 420 0.03 1500 [21] 2.2×10^{-4} SN 2005ip 100 - 200[22] 1100 0.1 - 1SN 2006gy 2004000 [23] [24] SN 2006tf 190 2000 0.1 - 4.1[25] $1-2 \times 10^{-2}$ SN 2008iy 5000 Kiewe+12 100[26]

Table 9



Previous works

• Moriya+ (2013)

Analytical (phenomenological) model of the light curve

Simple model & easy to estimate ejecta & CSM parameters from observation

• Dessart+ (2015)

Non-LTE radiation hydrodynamics calculations

Roughly reproduces observational features of superluminous supernovae

Both did not resolve the structure of the shocked region ➤A better model can be obtained by resolving shocked region? (Our work)

Methods: Overview





Tested typical(?) parameter sets for Type IIn

<u>Ejecta</u>

≻ Mass $10M_{\odot}$, energy 10^{51} erg



 \blacktriangleright Density profile: broken power-law of velocity ($\delta = 1, n = 10$)

<u>CSM</u>

Stable wind mass-loss (rate $1 \times 10^{-3} M_{\odot}$ /yr, wind velocity 10 km/s) Inner edge: 10^{14} cm, Outer edge: 10^{16} cm (-> total mass ~ 0.3 M_{\odot})

Matzner & McKee 99

Methods (1): Hydrodynamics

- Self-similar solution by Chevalier (1982)
 p, *ρ*, *v* can be obtained as function of time
 Adiabatic solution, so can't fully incorporate radiation feedback onto hydrodynamics

 (We *try*, by setting *γ* = 1.2 < 4/3)
- Inner shocked ejecta (HIGH density)
- + Outer shocked CSM (LOW density)
- The density contrast depends on the exact profile of ejecta & CSM, and adiabatic index



Methods (2): Initial Condition

Use Chevalier's solution for pressure, density, and velocity

Give (radiation) temperature from pressure by

$$p_{Ch} = \frac{1}{3}aT^4 + \frac{\rho_{Ch}}{\mu m_p}k_BT$$



 μ :set to \approx 0.62 assuming ionized gas of solar abundance

Methods (3): Radiation transfer

Flux-limited radiation transfer

• Optically thick case:

$$F = -\frac{c}{3\kappa\rho}\frac{\partial(aT^4)}{\partial r}$$

(Opacity κ comes from OPAL table assuming solar abundance)

• Optically thin case:

$$F = -acT^4 \operatorname{sign}(\partial T/\partial r)$$

• Obtain flux by interpolation of two cases (Levermore & Pomraning 81)

Methods (4): Radiation production at shock

Radiation from Shock heating:

➢ First the gas temperature increases

$$T_{gas} = \frac{\mu m_p p_{Ch}}{k_B \rho_{Ch}}$$

> Then photons are made by free-free emission within diffusion timescale

$$E_{rad} \sim aT_{rad}^{4} \sim \epsilon_{ff} \cdot \frac{\min(\tau, 1) r}{c} \propto \rho^{2} T_{g}^{\frac{1}{2}} r \cdot \min(\tau, 1)$$
(Free-free emissivity)
$$\star Forw \text{ shock (low density) should be}$$
much less efficient at photon production than rev shock (high density)

Gas -> radiation efficiency

At reverse shock, free-free emission is strong enough that radiation & gas reach equilibrium

$$p_{Ch} \approx \frac{1}{3} a T_{rad}^4$$
 gives temperature T_{rad}

@forward shock, radiation energy
is limited by free-free emission
within diffusion timescale

$$aT_{rad}^4 = \epsilon_{ff} \cdot \frac{\min(\tau, 1) r}{c}$$

(Rad. density that can be supplied by free-free within diff time)

(Rad. density assuming equilibrium $e_{rad} = aT_{rad}^4 \approx 3p_{Ch}$) 1000 forward shock reverse shock 100 Rev shock (gas-rad equilibrium) 10 efficiency $aT_{rad}^4 \approx 3p_{Ch}$ Forward shock (eff. goes down as $\propto t^{-3/4}$) 0.1 0.01 10 20 30 40 50 60 70 80 days

Result (Light curve)

Early phase

Photons diffusing out CSM makes a sharp rise in the light curve

- ▶ Peak luminosity $\approx 10^{43}$ erg/s
- Timescale ~ 10 days

Timescale roughly consistent w/ diffusion timescale in CSM



 $(R_{Sh}: \text{ radius of shock at peak})$



Result (Light curve)

Late phase

Light curve monotonically decays (With time dependence given by Energy release rate @ fwd shock & radiation efficiency @ fwd shock)

- Energy release rate $\propto t^{-3/8}$
- Radiation efficiency $\propto t^{-3/4}$ $\implies L \propto t^{-9/8}$



Conclusion

• We calculated the light curve of interaction-powered supernovae, resolving the radial profile inside the shocked region.

• At the early phase we find the luminosity peak that comes from photons created in the shocked region, diffusing out the CSM.

• At the late phase the luminosity comes from photons generated at the forward (and reverse) shock front, reflecting the decreasing efficiency at the forward shock.

Future work

- We have only tested a small number of parameter set, and results should vastly depend on parameters
- > Parameter survey & comparison with observations important

Our better treatment of the forward shock may have some applications, e.g.
 Calculating non-thermal emission from shock-heated electrons
 Collision-less shock acceleration -> hadronic process (-> high-energy neutrino)
 May give multi-wavelength (messenger?) observational predictions of Type IIn SNe