

# Supernovae powered by magnetars that transform into black holes

---

Takashi Moriya (National Astronomical Observatory of Japan)

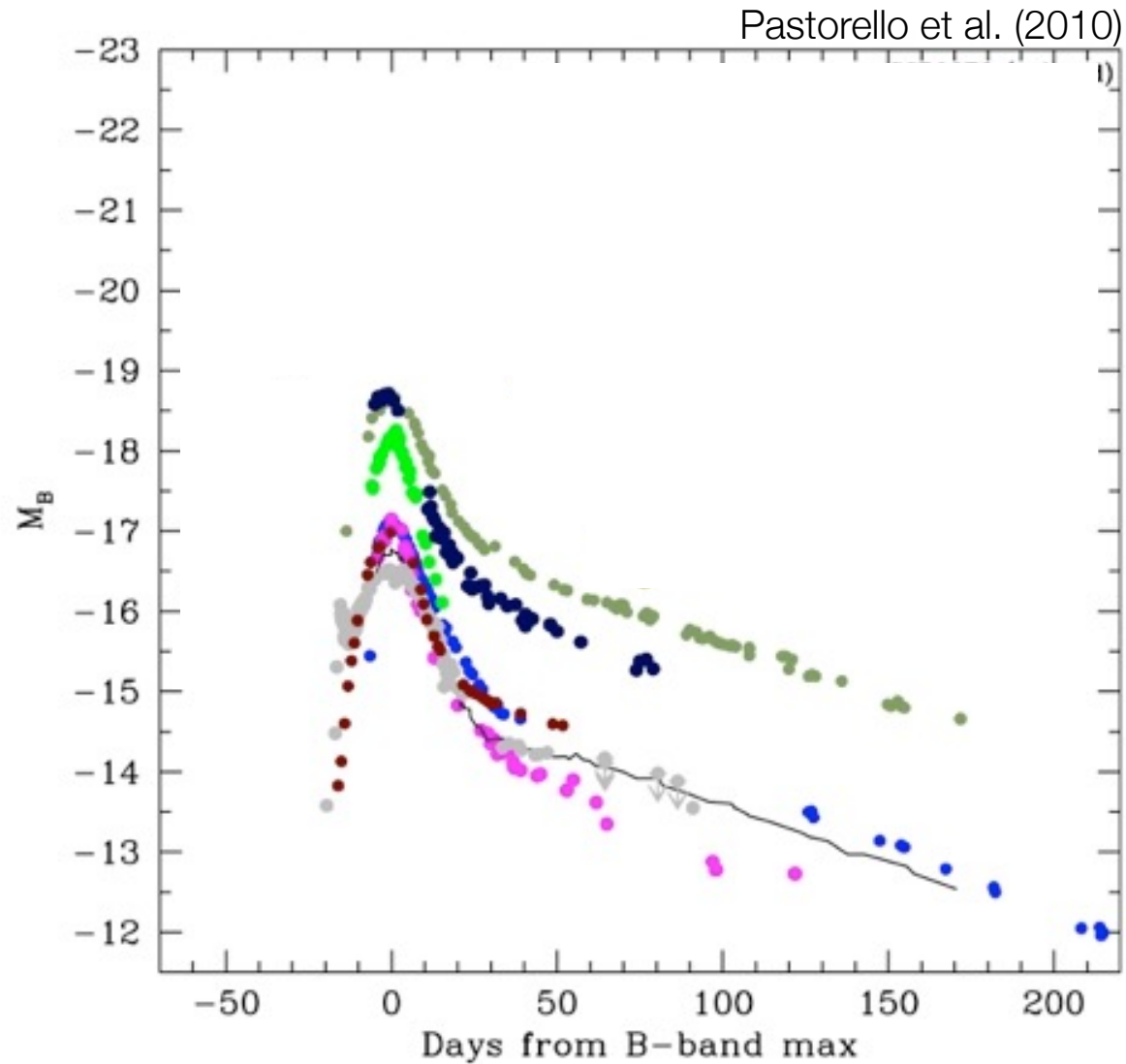
B. Metzger (Columbia), S. Blinnikov (ITEP)

arXiv:1606.09316



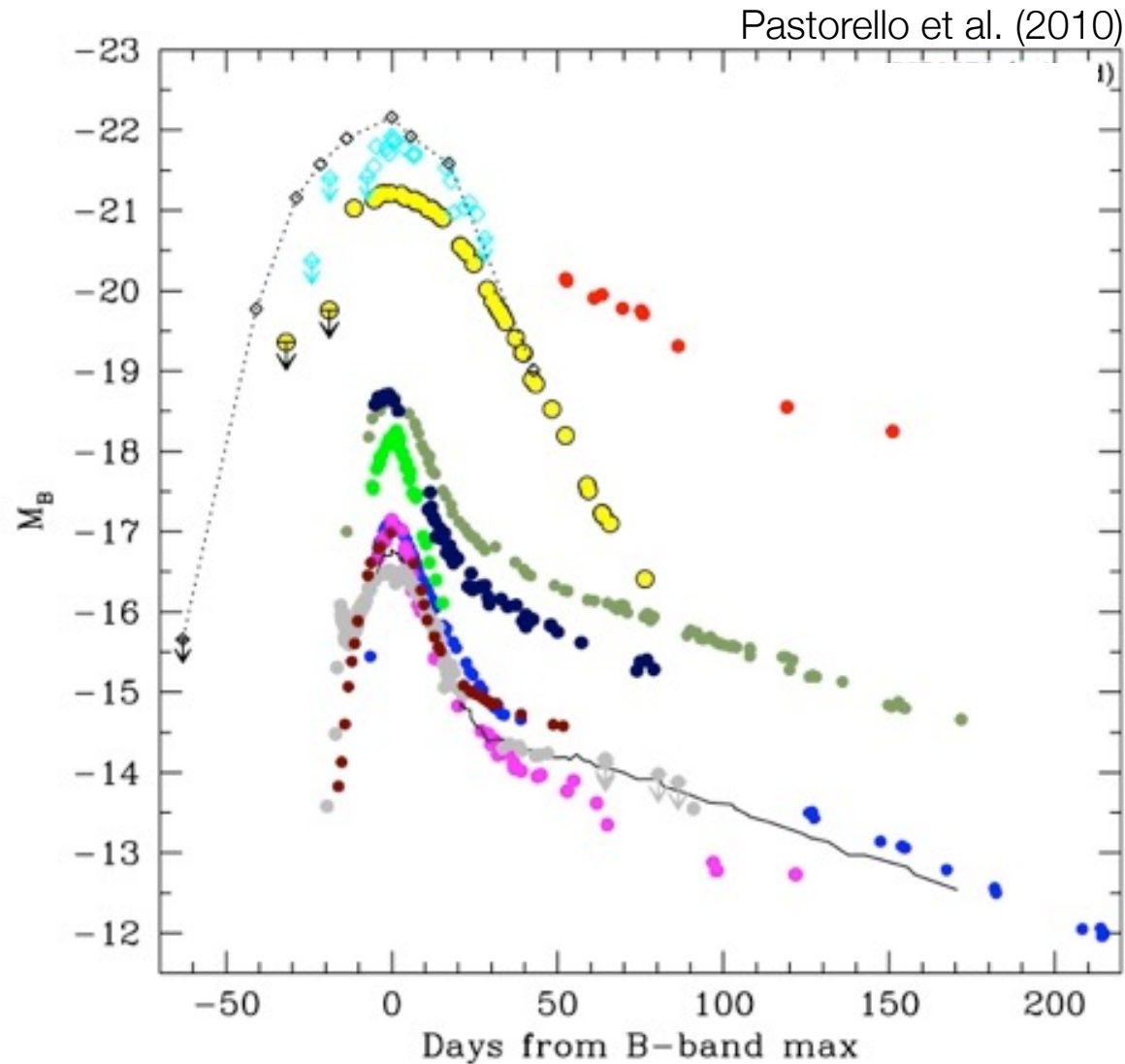
# Classical core-collapse SNe

- until ~ 2006 (but SN 1999as)



# Superluminous SNe (SLSNe)

- SNe brighter than  $\sim 1e44$  erg/s (or -21 mag in optical)



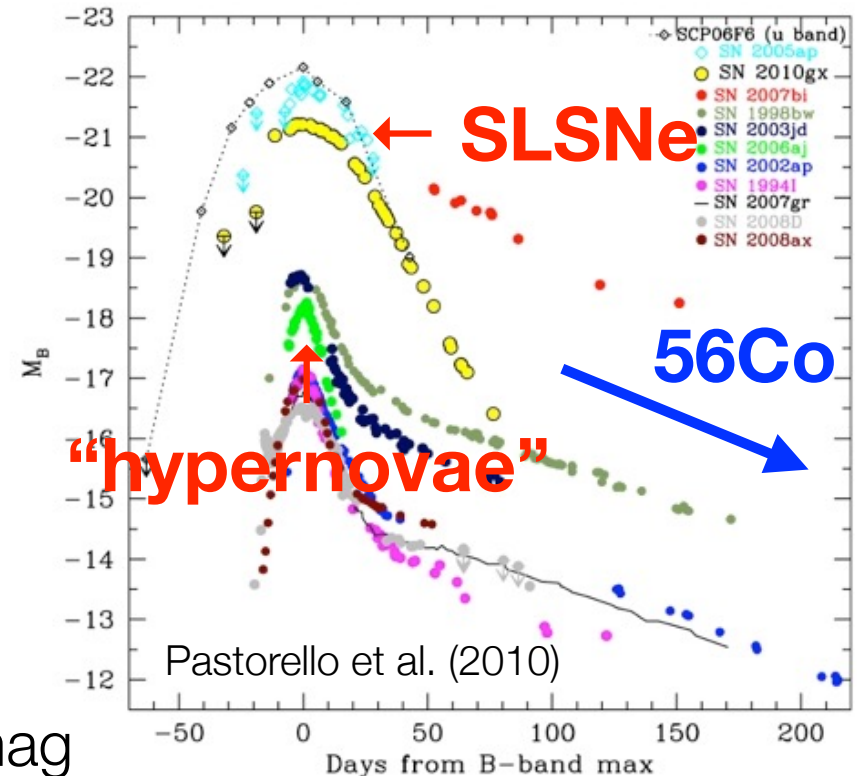
# SLSN properties

Supernova	Redshift	Absolute peak (mag)	Radiated energy (ergs)
SN 2007bi	0.1289	-21.35	1 to $2 \times 10^{51}$
SN 1999as	0.12	-21.4	
CSS100217	0.147	-23.07	$1.3 \times 10^{52}$
SN 2008fz	0.133	-22.34	$1.4 \times 10^{51}$
SN 2008am	0.2338	-22.39	$2 \times 10^{51}$
SN 2008es	0.205	-22.21	$1.1 \times 10^{51}$
SN 2006gy	0.019	-22.0	2.3 to $2.5 \times 10^{51}$
SN 2003ma	0.289	-21.52	$4 \times 10^{51}$
SN 2006tf	0.074	< -20.7	$7 \times 10^{50}$
SN 2005ap	0.2832	-22.73	$1.2 \times 10^{51}$
SCP 06F6	1.189	-22.53	$1.7 \times 10^{51}$
PS1-10ky	0.956	-22.53	0.9 to $1.4 \times 10^{51}$
PS1-10awh	0.908	-22.53	0.9 to $1.4 \times 10^{51}$
PTF09atu	0.501	-22.03	
PTF09cnd	0.258	-22.03	$1.2 \times 10^{51}$
SN 2009jh	0.349	-22.03	
SN 2006oz	0.376	-21.53	
SN 2010gx	0.230	-21.23	$6 \times 10^{50}$

Gal-Yam (2012)

# “Hypernovae” v.s. SLSNe

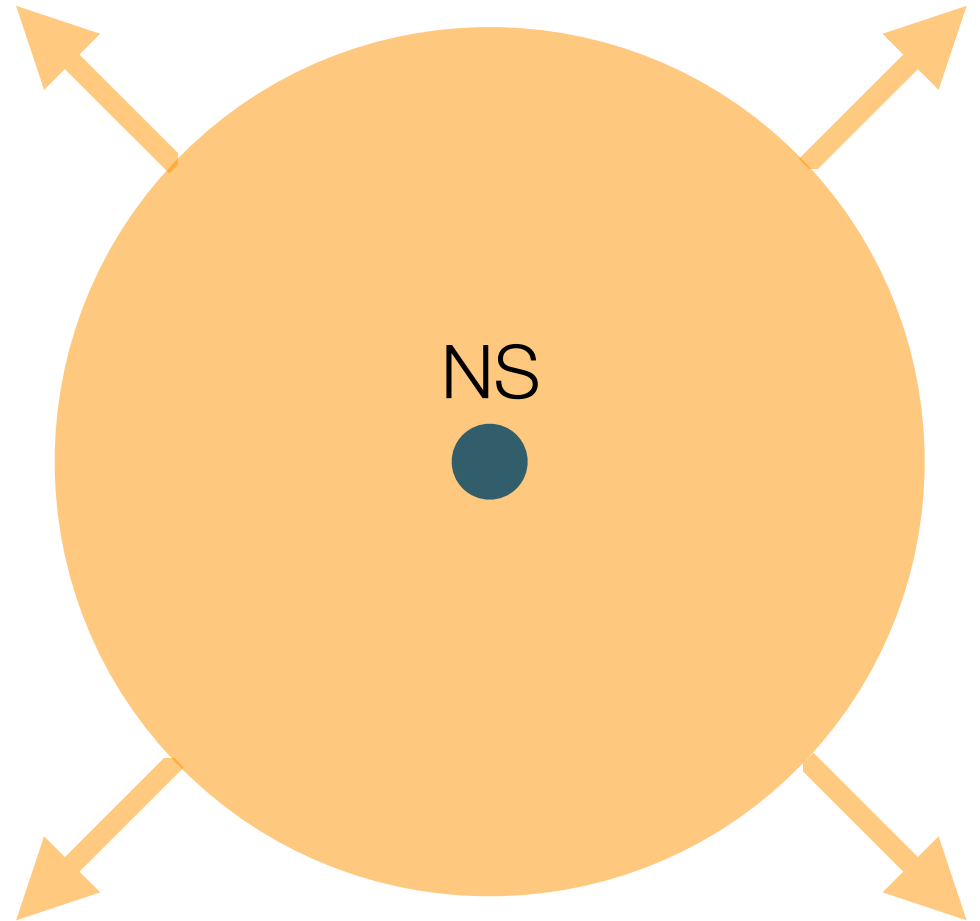
- “hypernovae”
  - broad-line “Type Ic” SNe
  - peak luminosity: -18 - -19 mag
    - $^{56}\text{Ni}$  mass: 0.5 - 1  $M_{\text{sun}}$
  - red spectra
- SLSNe
  - Type II and Ic
  - peak luminosity: more than  $\sim -21$  mag
    - $^{56}\text{Ni}$  mass:  $> \sim 5 M_{\text{sun}}$  (often more than 10  $M_{\text{sun}}$ )
  - blue spectra



# Why are Type Ic SLSNe bright?

---

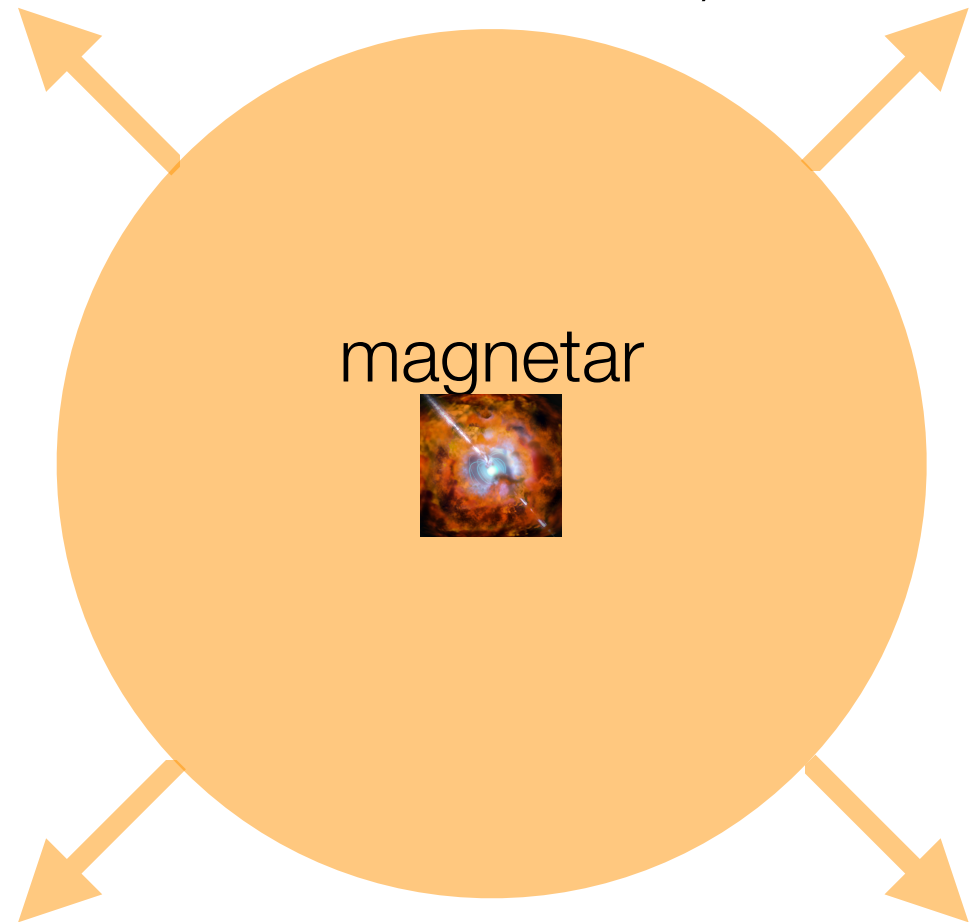
- rotational radiation of magnetars (Kasen & Bildsten 2010, ...)



# Why are Type Ic SLSNe bright?

---

- rotational radiation of magnetars (Kasen & Bildsten 2010, ...)



# Why are Type Ic SLSNe bright?

---

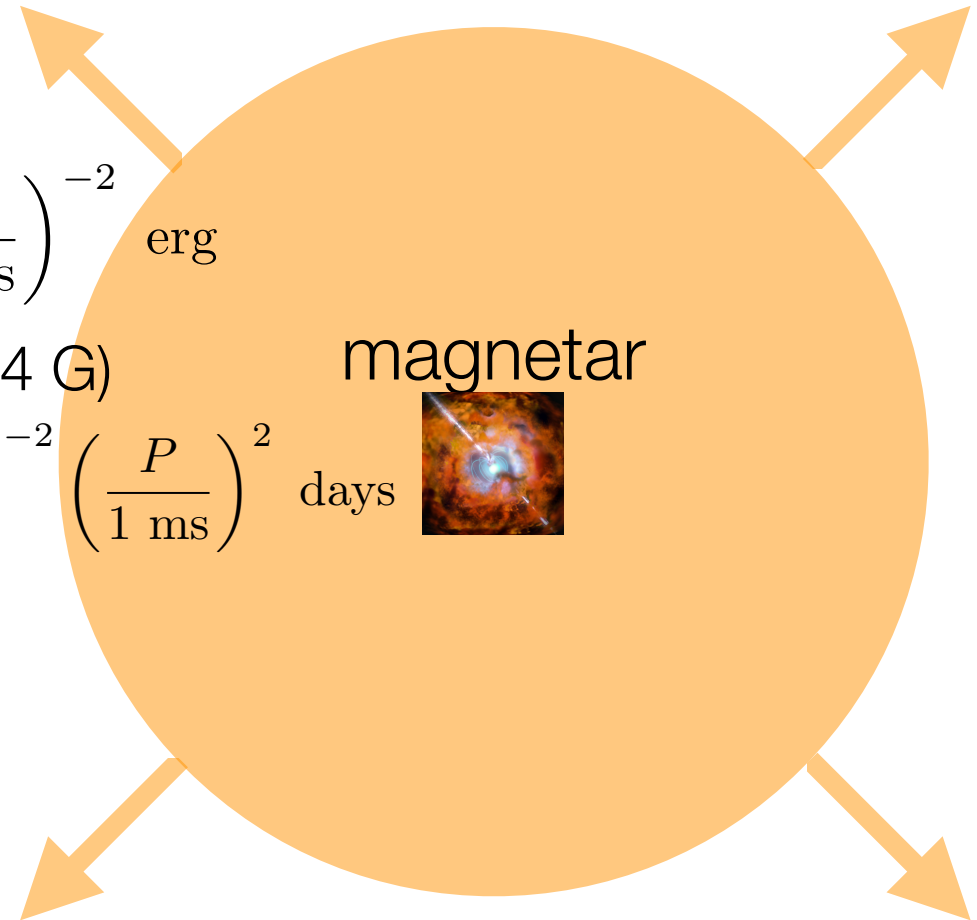
- rotational radiation of magnetars (Kasen & Bildsten 2010, ...)
  - neutron stars with
    - rapid rotation ( $\sim 1$  ms)

$$E_{\text{rot}} = \frac{1}{2} I_{\text{NS}} \Omega^2 \simeq 2 \times 10^{52} \left( \frac{P}{1 \text{ ms}} \right)^{-2} \text{ erg}$$

- strong magnetic field ( $\sim 10^{14}$  G)

$$t_m = \frac{6 I_{\text{NS}} c^3}{B_{\text{dipole}}^2 R_{\text{NS}}^6 \Omega^2} \simeq 5 \left( \frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right)^{-2} \left( \frac{P}{1 \text{ ms}} \right)^2 \text{ days}$$

magnetar





# Why are Type Ic SLSNe bright?

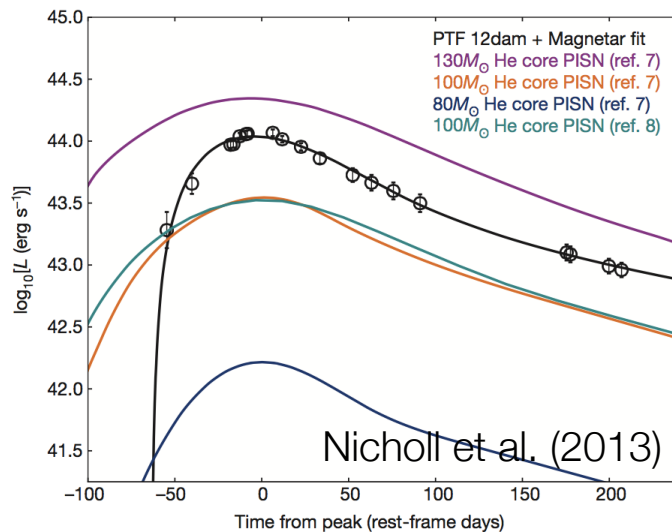
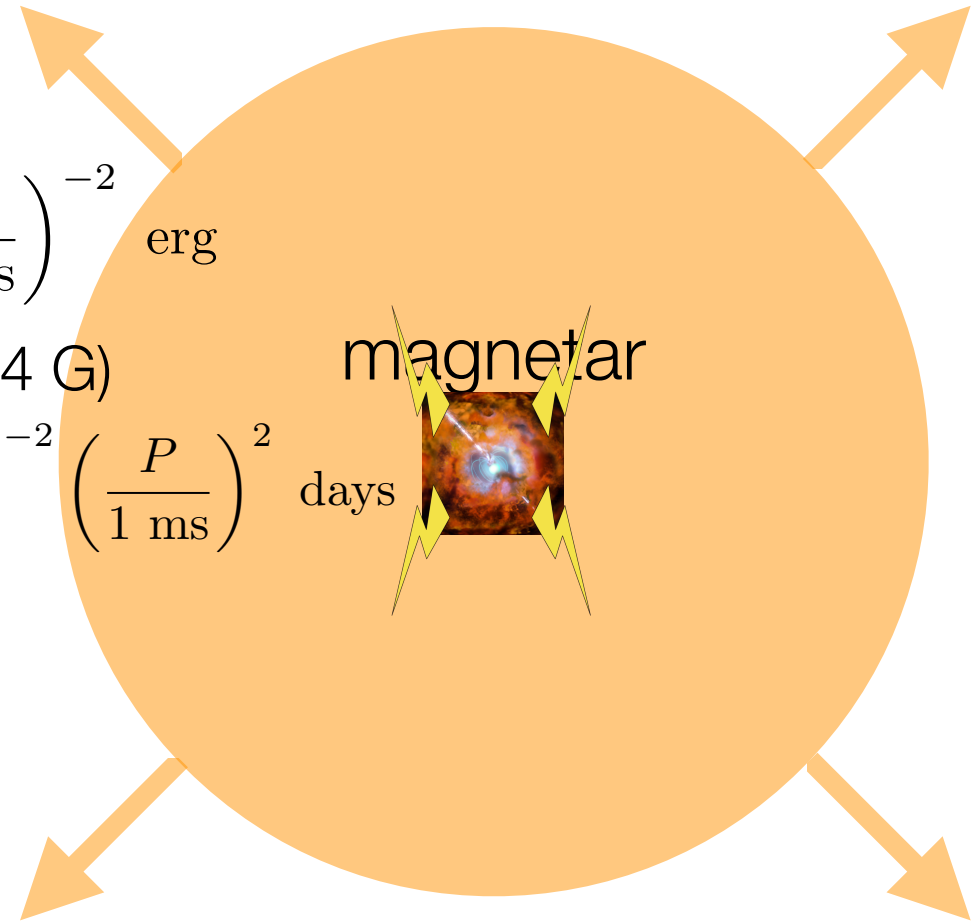
- rotational radiation of magnetars (Kasen & Bildsten 2010, ...)
- neutron stars with

- rapid rotation ( $\sim 1$  ms)

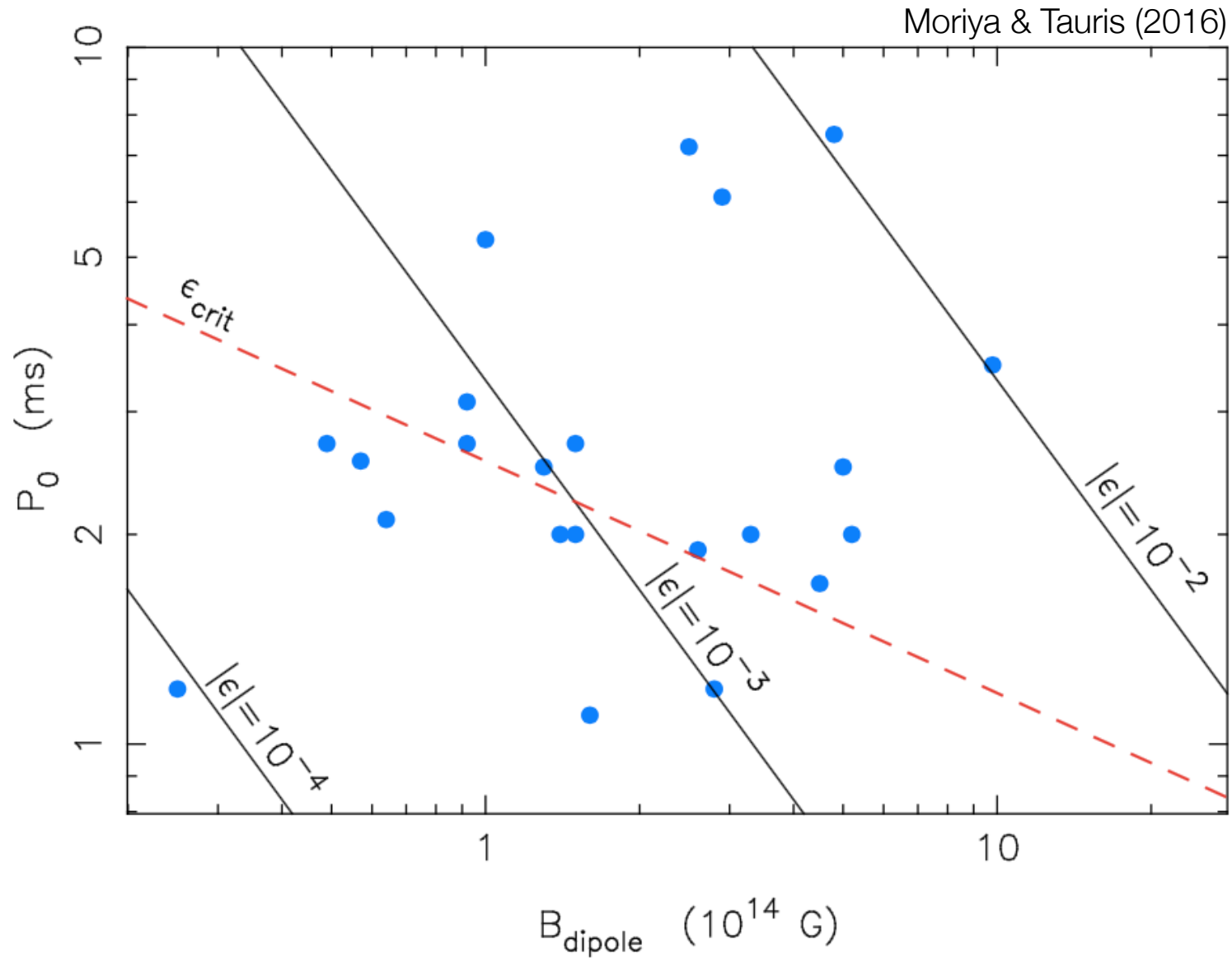
$$E_{\text{rot}} = \frac{1}{2} I_{\text{NS}} \Omega^2 \simeq 2 \times 10^{52} \left( \frac{P}{1 \text{ ms}} \right)^{-2} \text{ erg}$$

- strong magnetic field ( $\sim 1e14$  G)

$$t_m = \frac{6 I_{\text{NS}} c^3}{B_{\text{dipole}}^2 R_{\text{NS}}^6 \Omega^2} \simeq 5 \left( \frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right)^{-2} \left( \frac{P}{1 \text{ ms}} \right)^2 \text{ days}$$

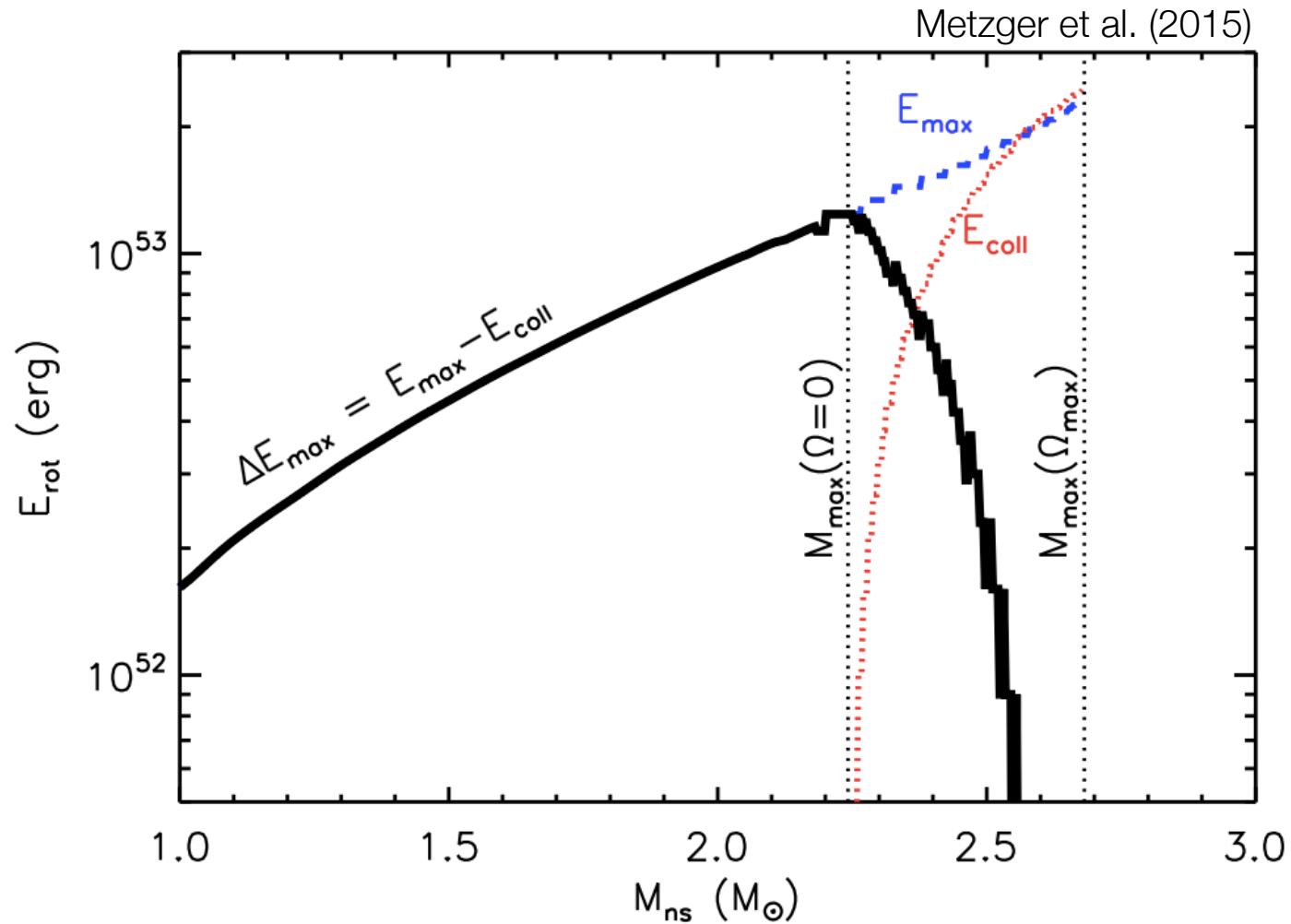


# Estimated initial dipole magnetic field and spin



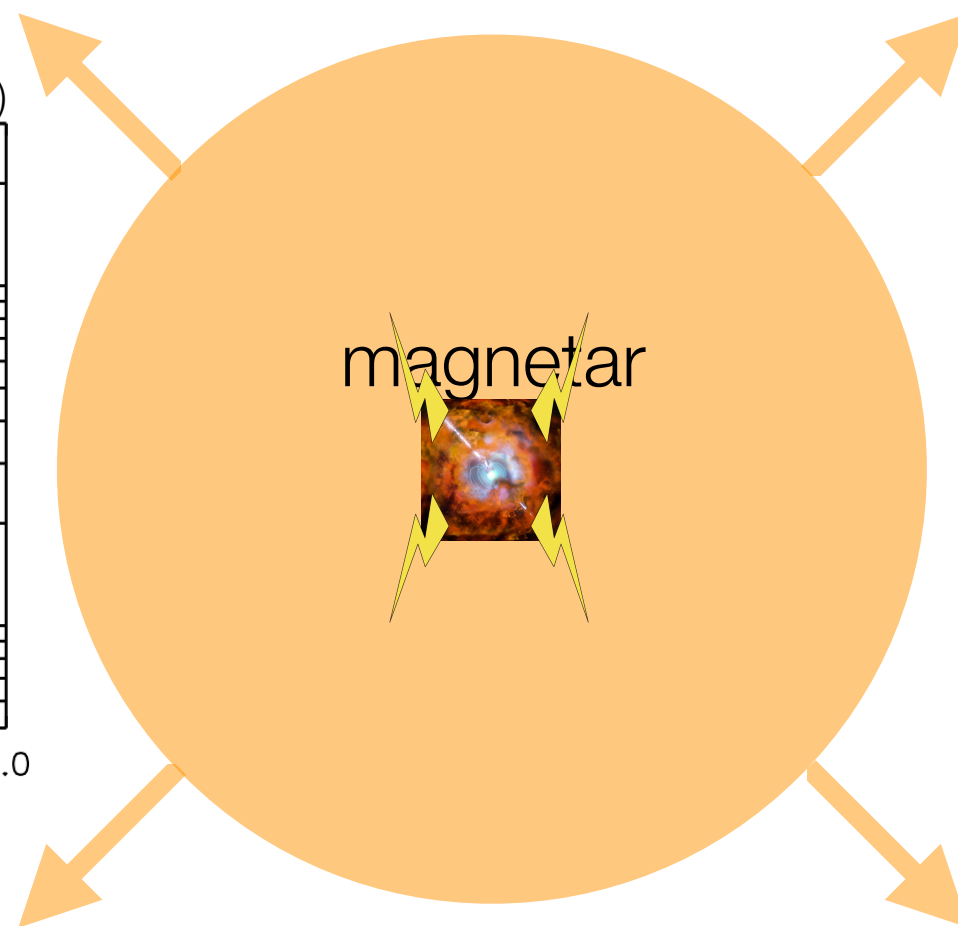
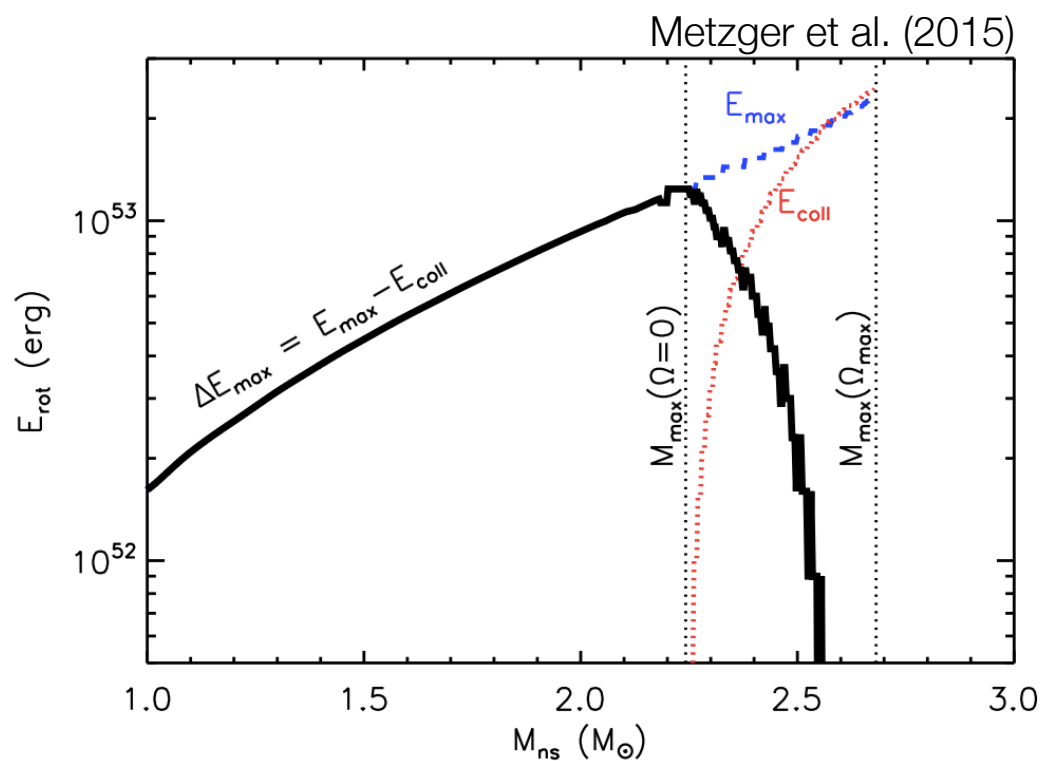
# Maximum rotational energy of NSs

- for an EoS consistent with the current NS constraints



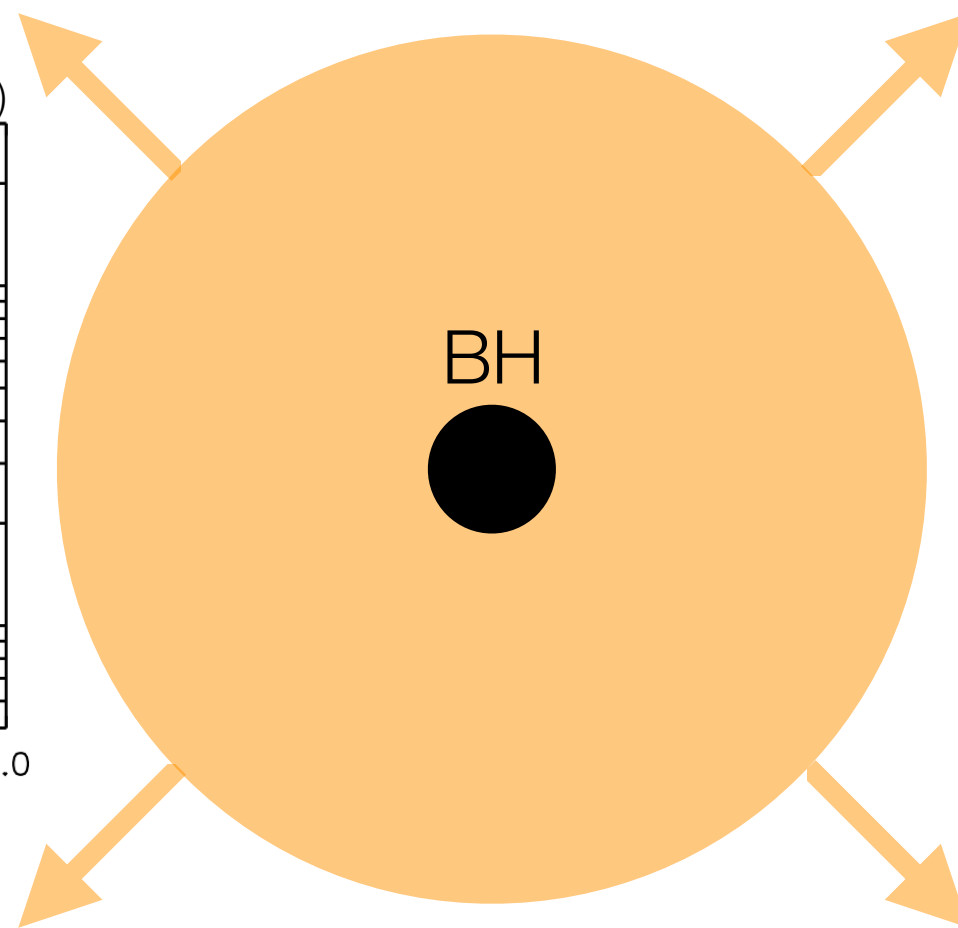
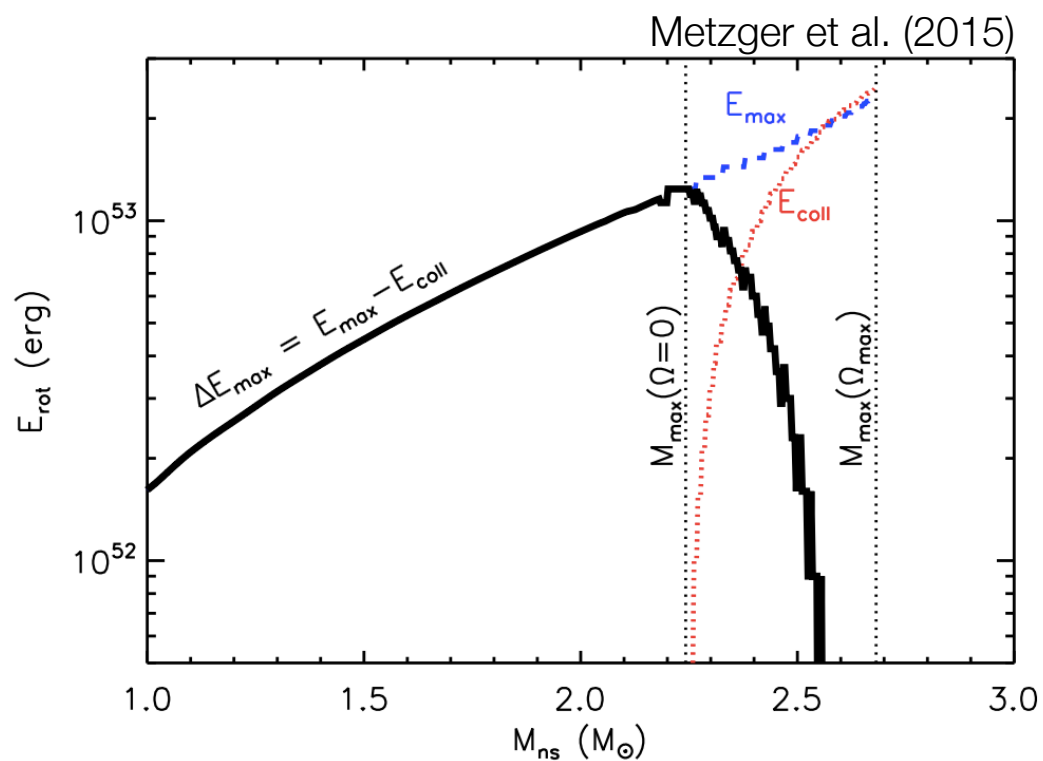
# Magnetars may collapse to BHs

- sudden loss of the central heating sources



# Magnetars may collapse to BHs

- sudden loss of the central heating sources



# When do magnetars collapse to BHs?

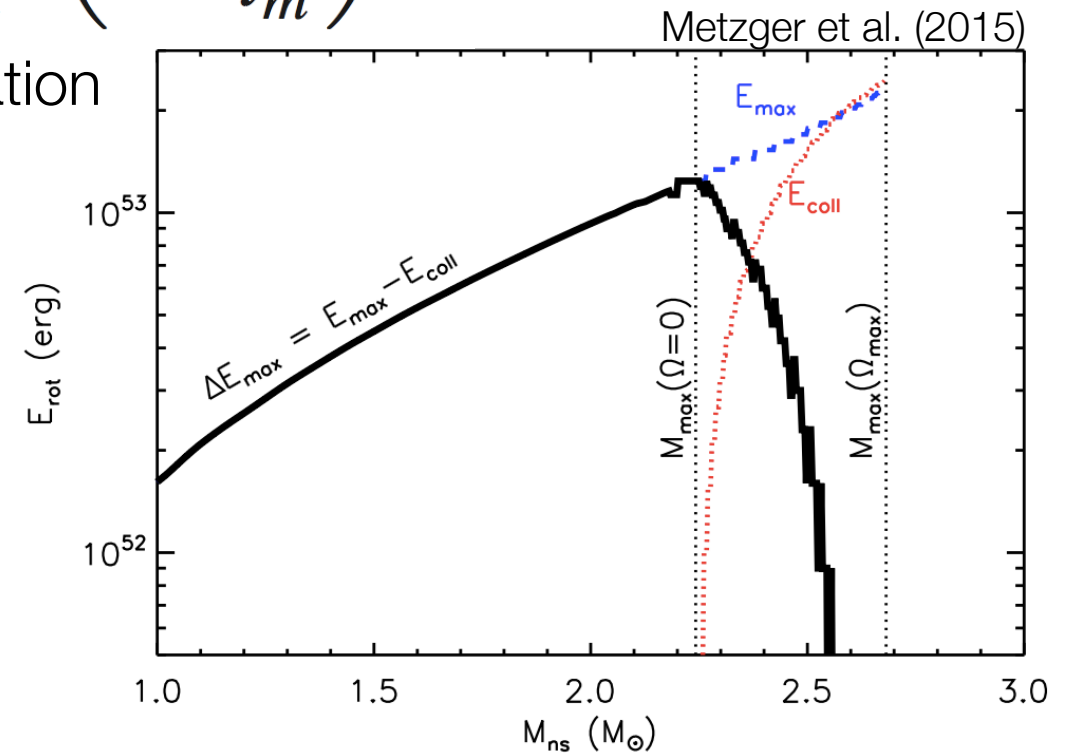
- rotational energy loss by dipole radiation
  - $t_m$ : spin-down timescale
  - $E_m$ : initial rotational energy of magnetars

$$L_{\text{mag}}(t) = \frac{E_m}{t_m} \left( 1 + \frac{t}{t_m} \right)^{-2}$$

- $t_{\text{BH}}$ : time of the BH transformation

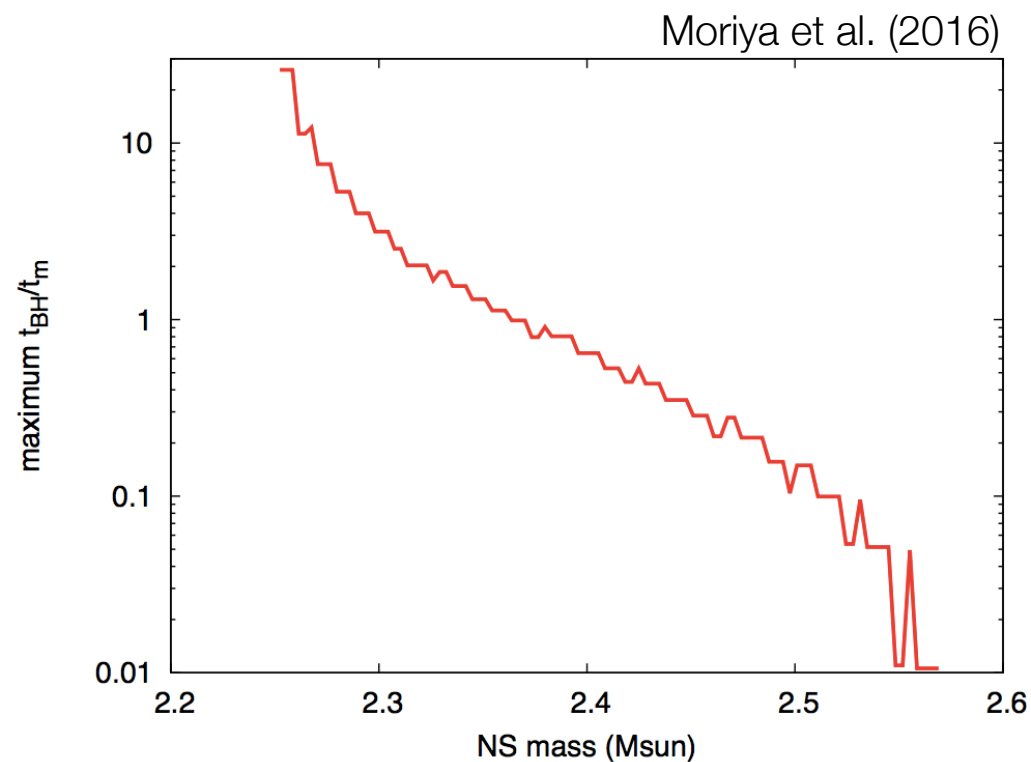
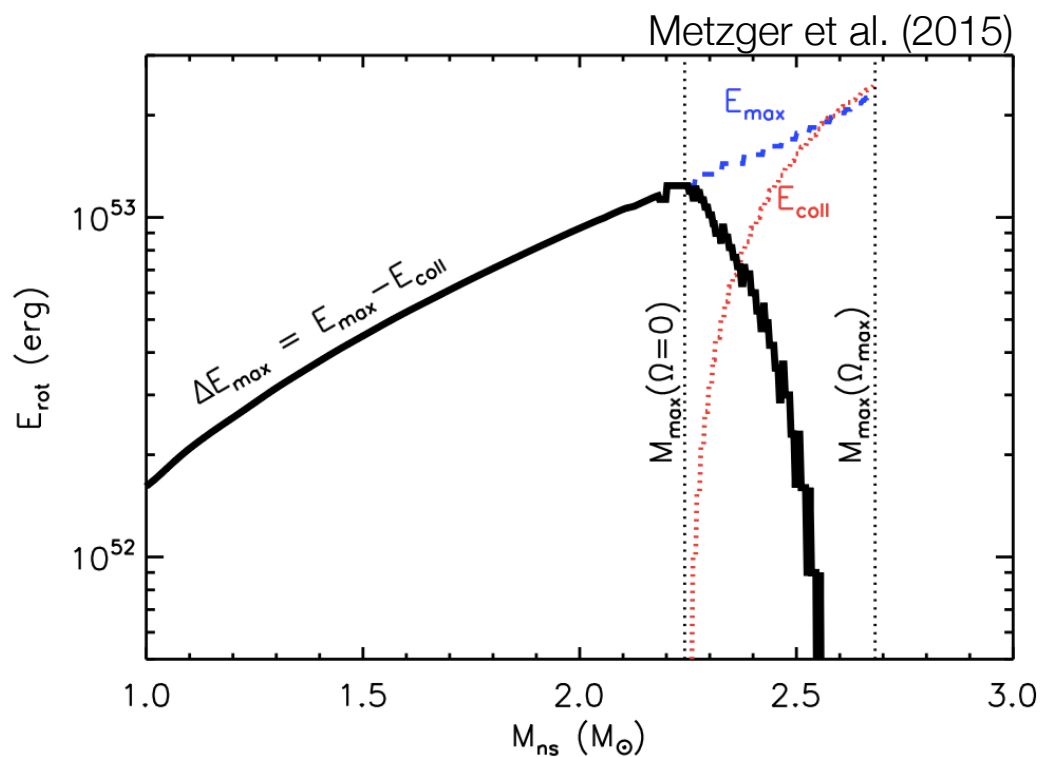
$$t_{\text{BH}} = \frac{\Delta E}{E_{\text{coll}}} t_m$$

$$\Delta E \equiv E_m - E_{\text{coll}}$$



# Time of BH transformation

$$t_{\text{BH}} = \frac{\Delta E}{E_{\text{coll}}} t_m \quad \Delta E \equiv E_m - E_{\text{coll}}$$



- 1/2 of the initial rotational energy is emitted in  $t_m$

# Effect of BH transformation on light curves

---

- semi-analytic method (Arnett 1982)

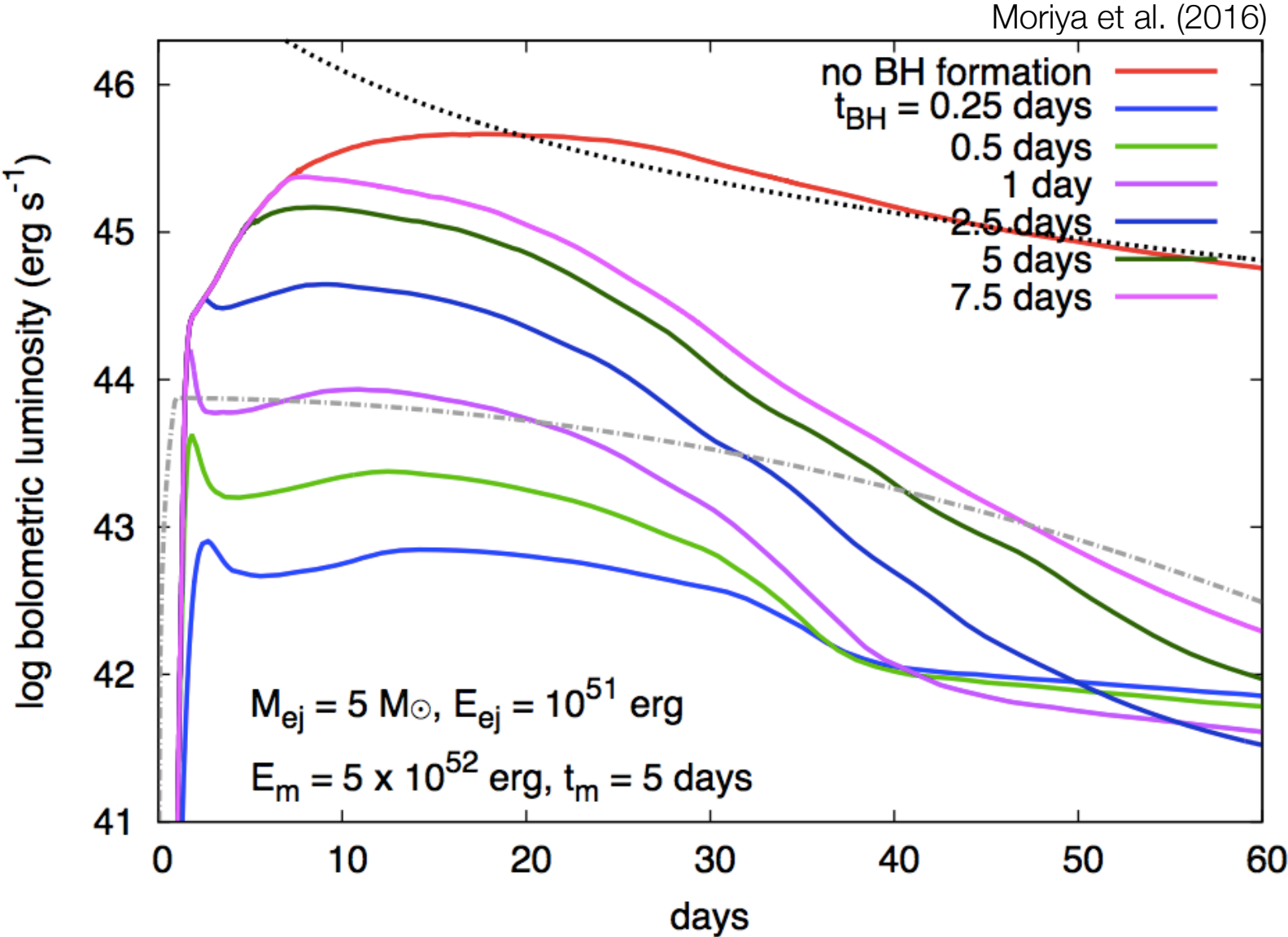
$$L(t) = \int_0^{\frac{t}{\tau_m}} 2\tau_m^{-2} L_{\text{mag}}(t') t' e^{\left(\frac{t'-t}{\tau_m}\right)^2} dt'$$

$$\tau_m = 1.05 \left( \frac{\kappa_e}{\beta c} \right)^{0.5} M_{\text{ej}}^{0.75} E_{\text{ej}}^{-0.25}$$

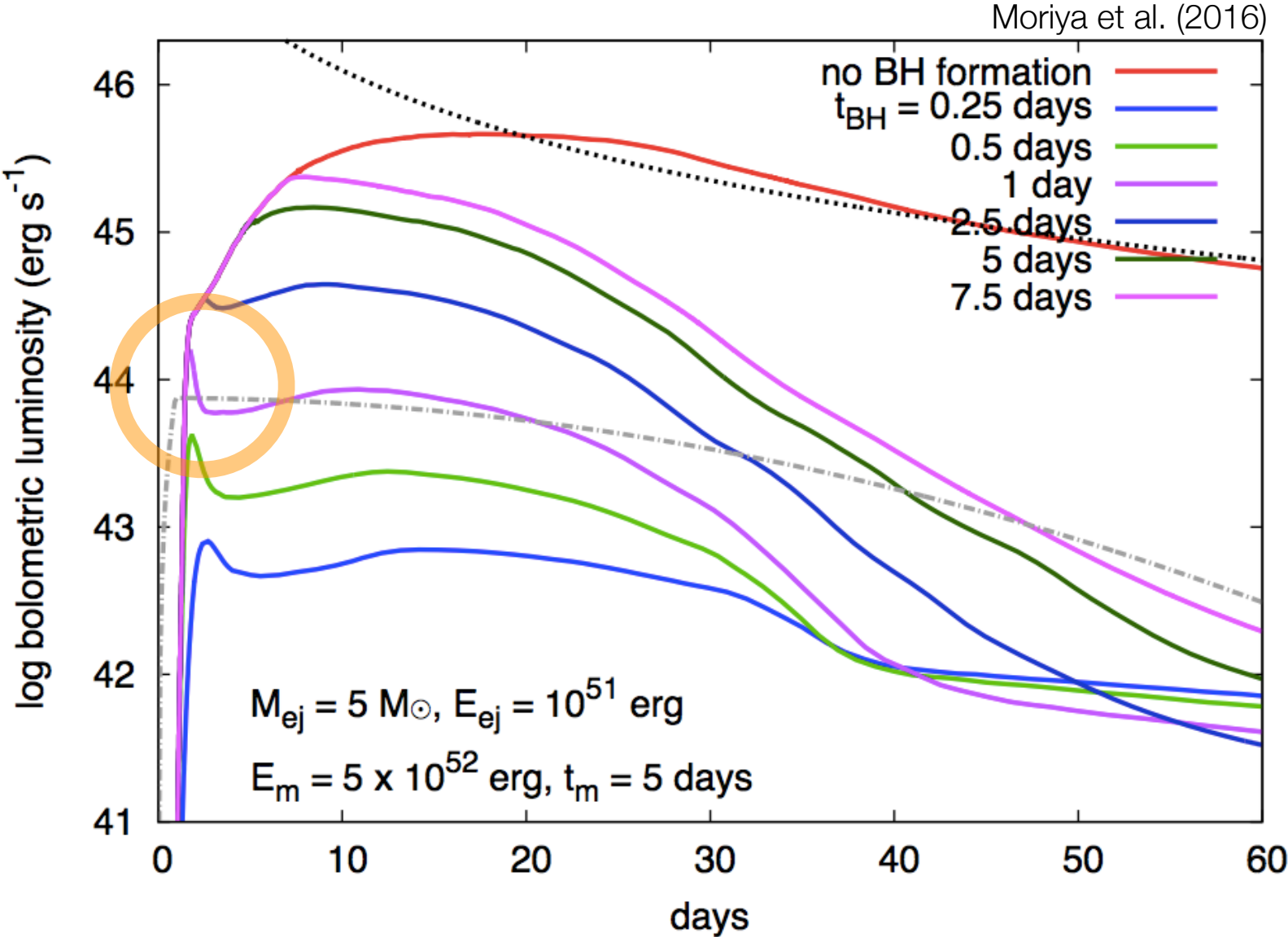
- numerical method
  - 1D radiation hydro code STELLA (Blinnikov et al.)
  - spin-down energy is put as thermal energy
- 100% thermalization of dipole spin-down energy is assumed



# Effect of BH transformation on light curves



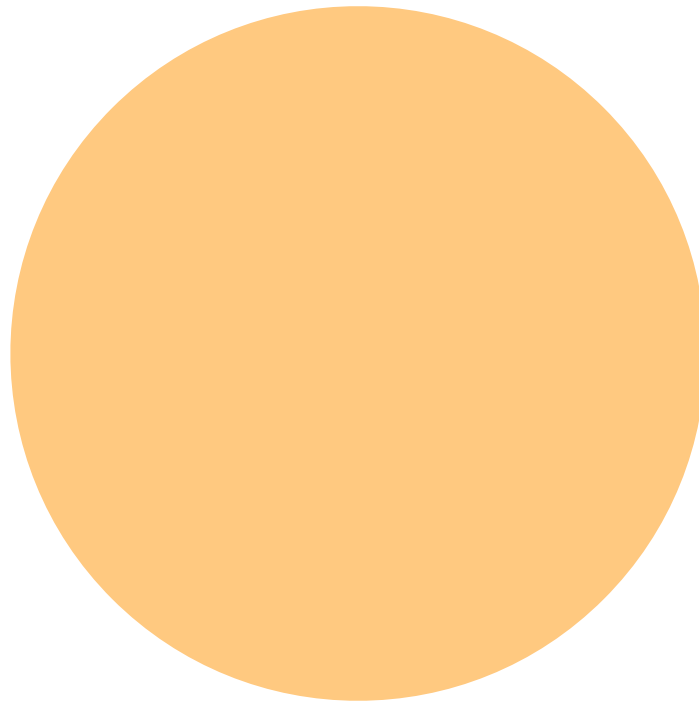
# Effect of BH transformation on light curves



# Magnetar-driven shock breakout

---

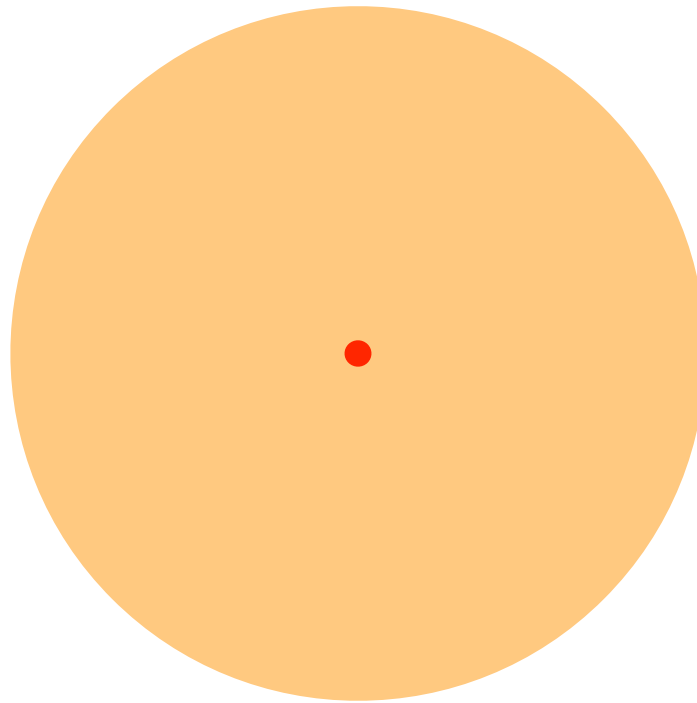
- magnetar-driven shock breakout (Kasen et al. 2016)



# Magnetar-driven shock breakout

---

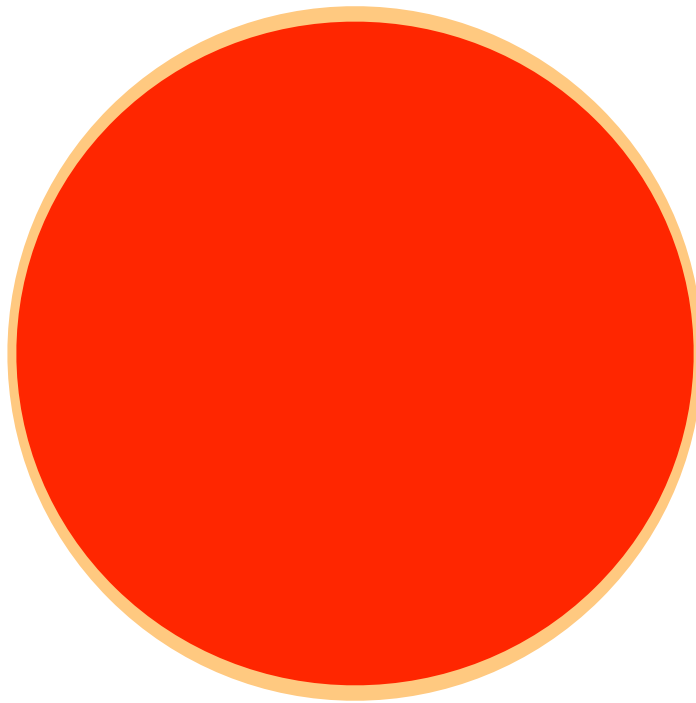
- magnetar-driven shock breakout (Kasen et al. 2016)



# Magnetar-driven shock breakout

---

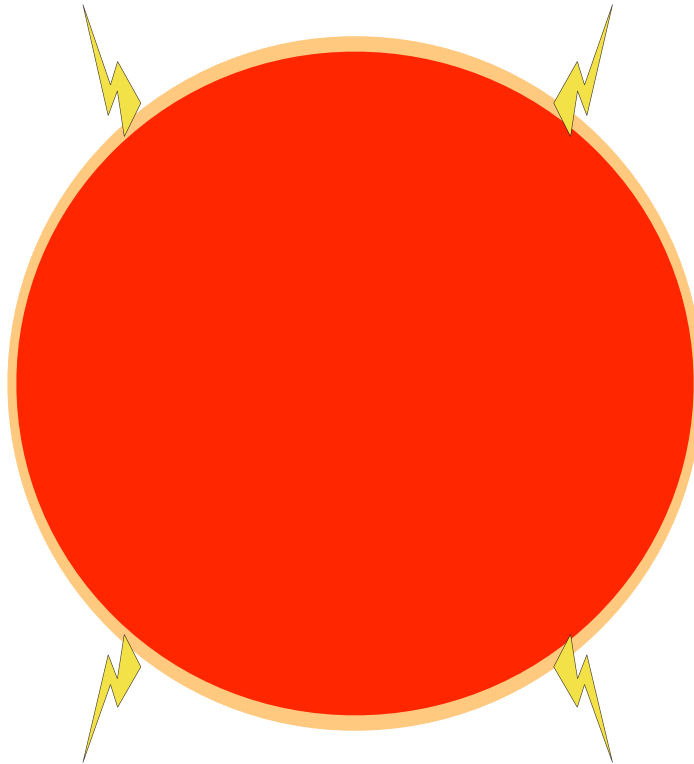
- magnetar-driven shock breakout (Kasen et al. 2016)



# Magnetar-driven shock breakout

---

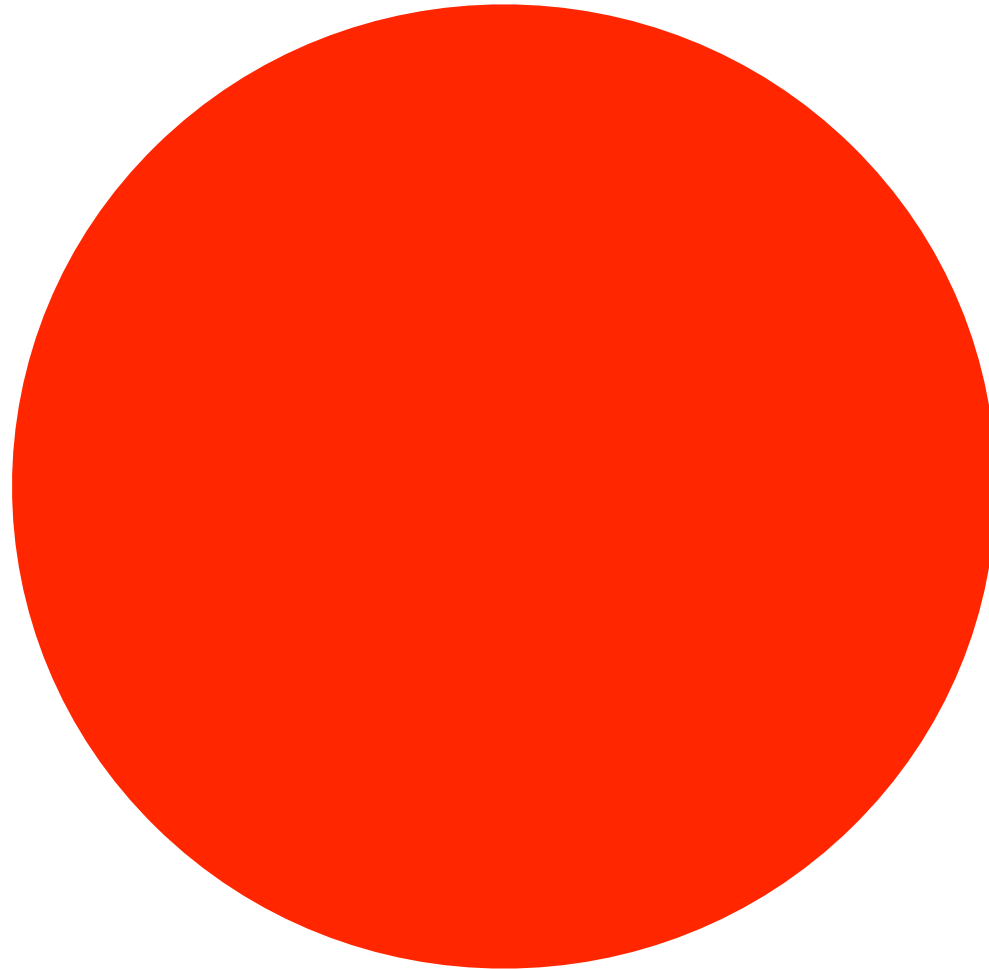
- magnetar-driven shock breakout (Kasen et al. 2016)



# Magnetar-driven shock breakout

---

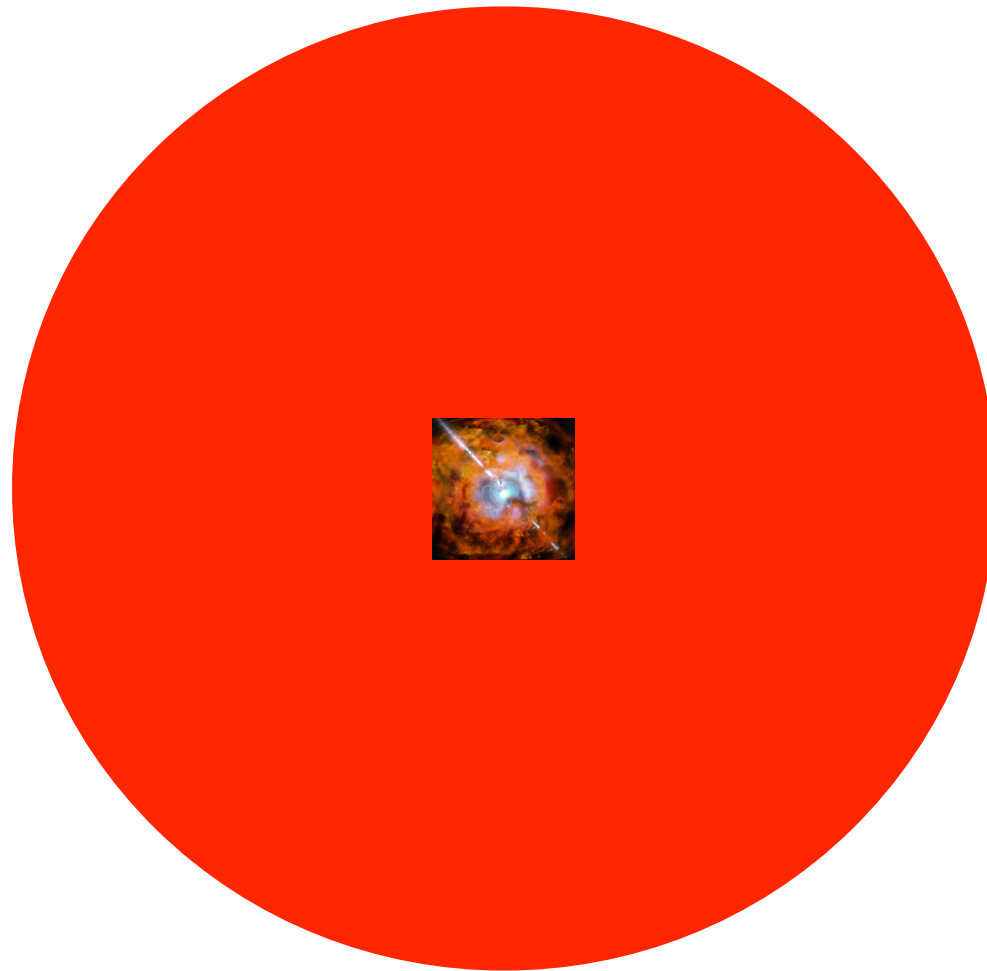
- magnetar-driven shock breakout (Kasen et al. 2016)



# Magnetar-driven shock breakout

---

- magnetar-driven shock breakout (Kasen et al. 2016)

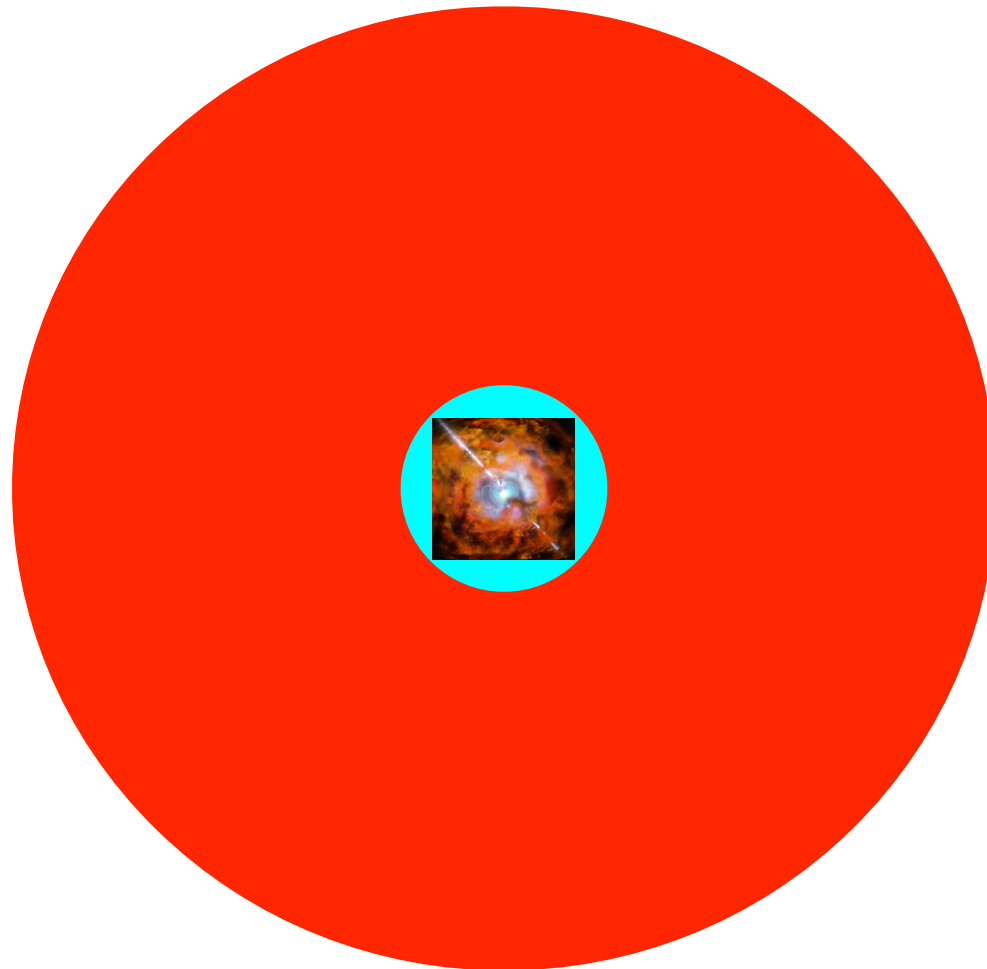




# Magnetar-driven shock breakout

---

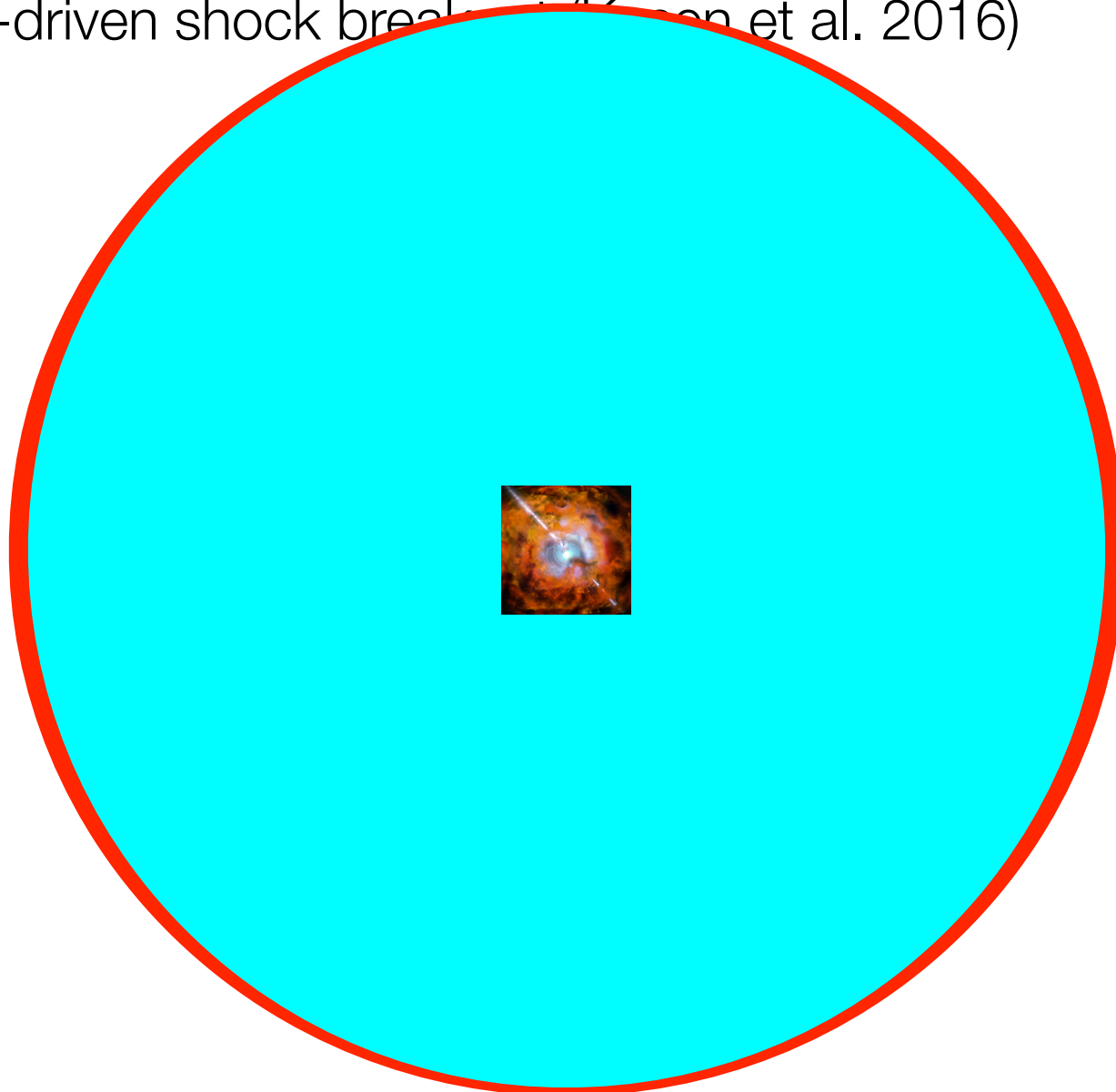
- magnetar-driven shock breakout (Kasen et al. 2016)



# Magnetar-driven shock breakout

---

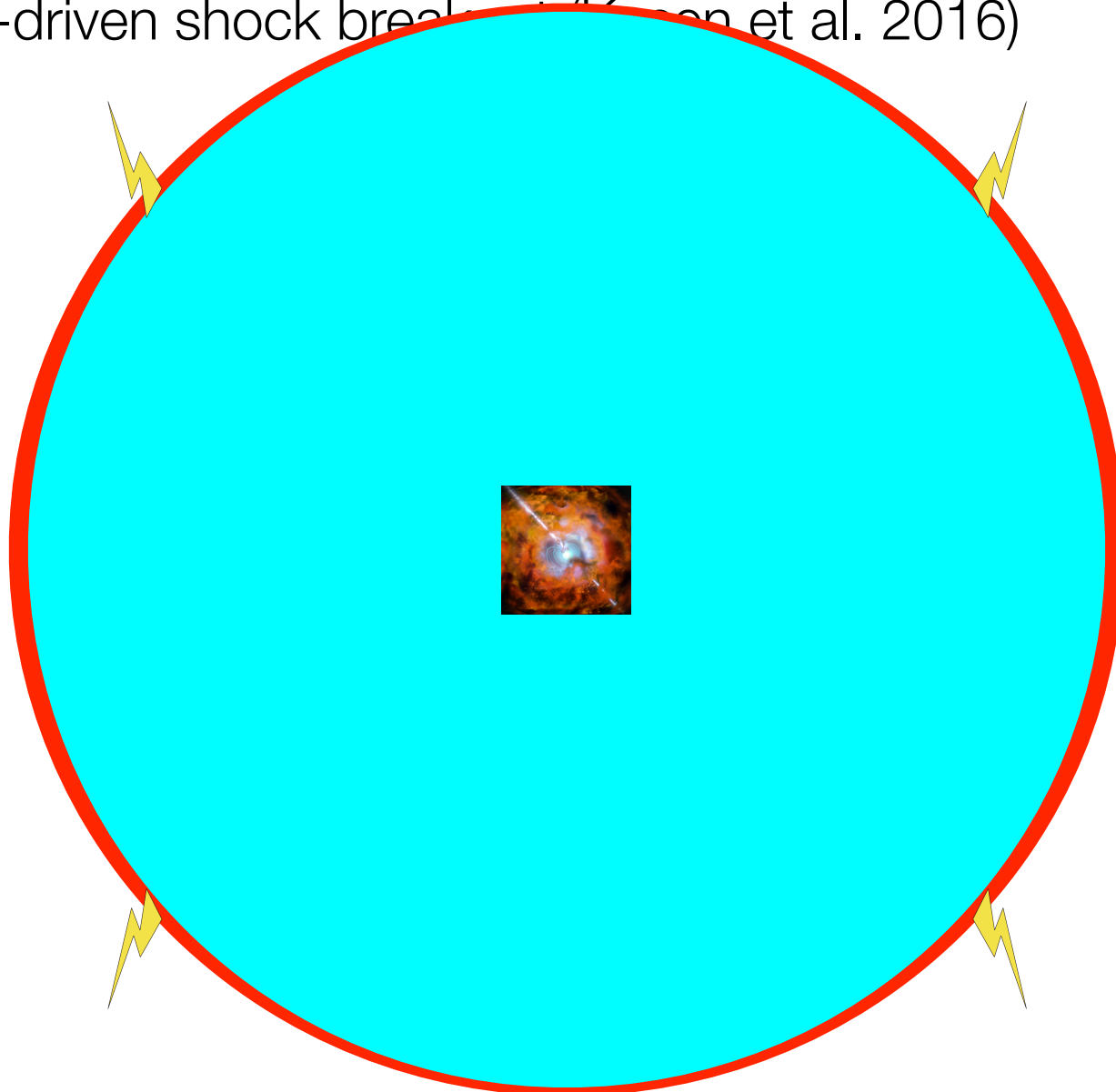
- magnetar-driven shock breakout (Kluźniak et al. 2016)



# Magnetar-driven shock breakout

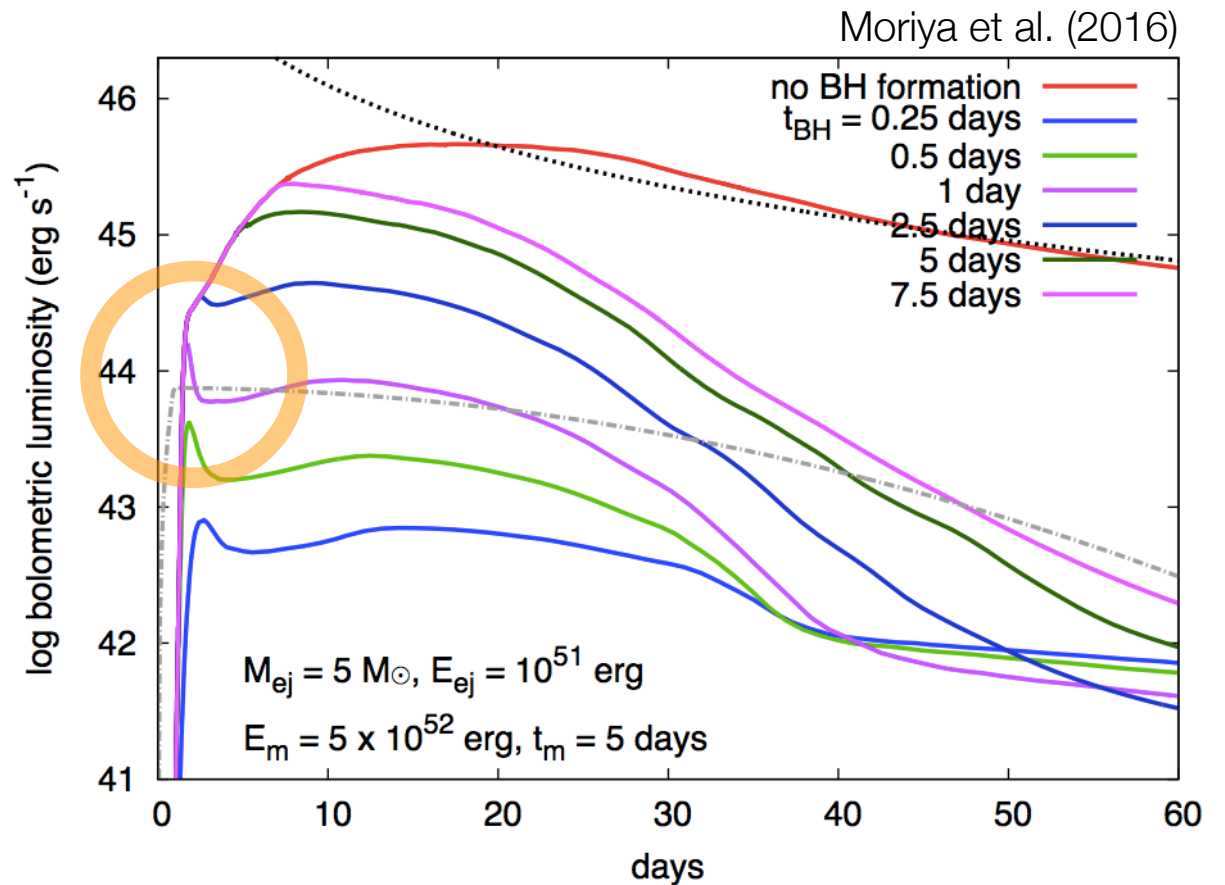
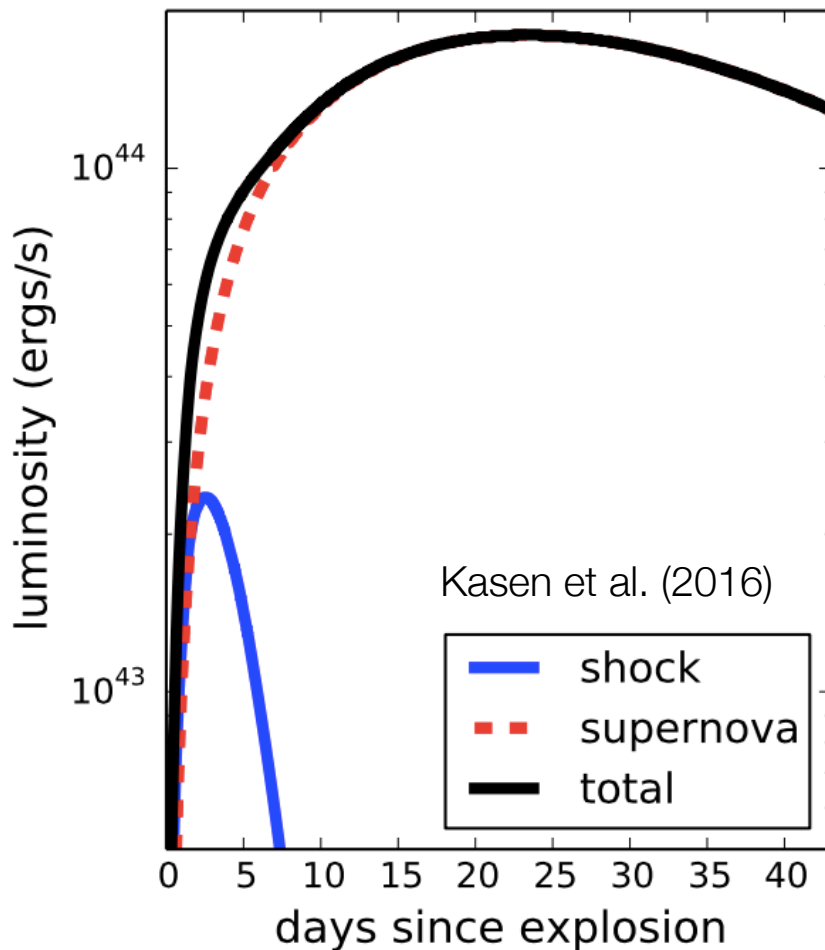
---

- magnetar-driven shock breakout (Kluźniak et al. 2016)



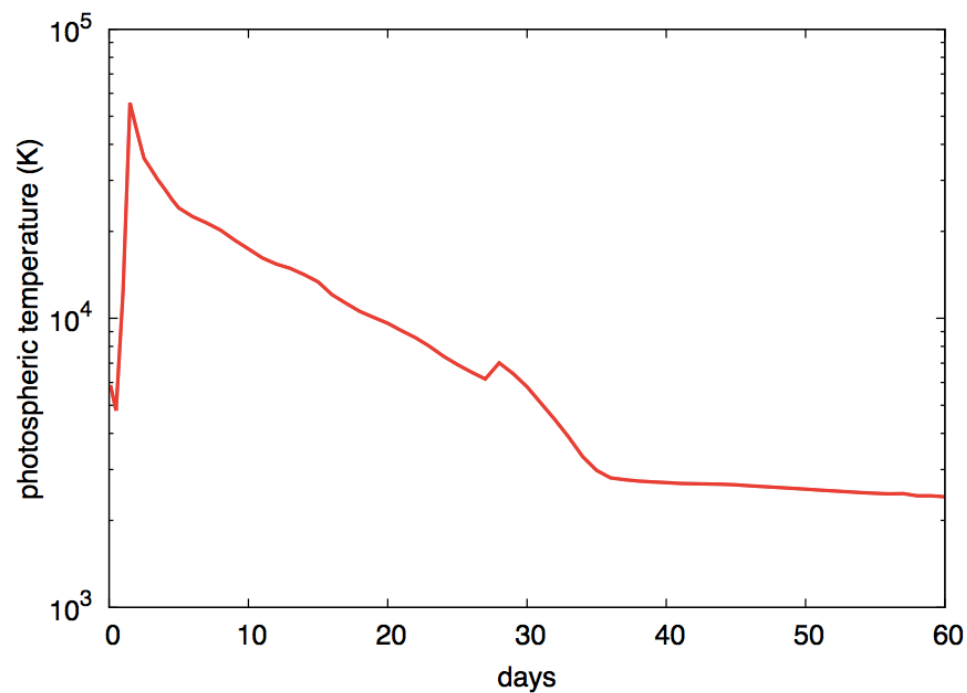
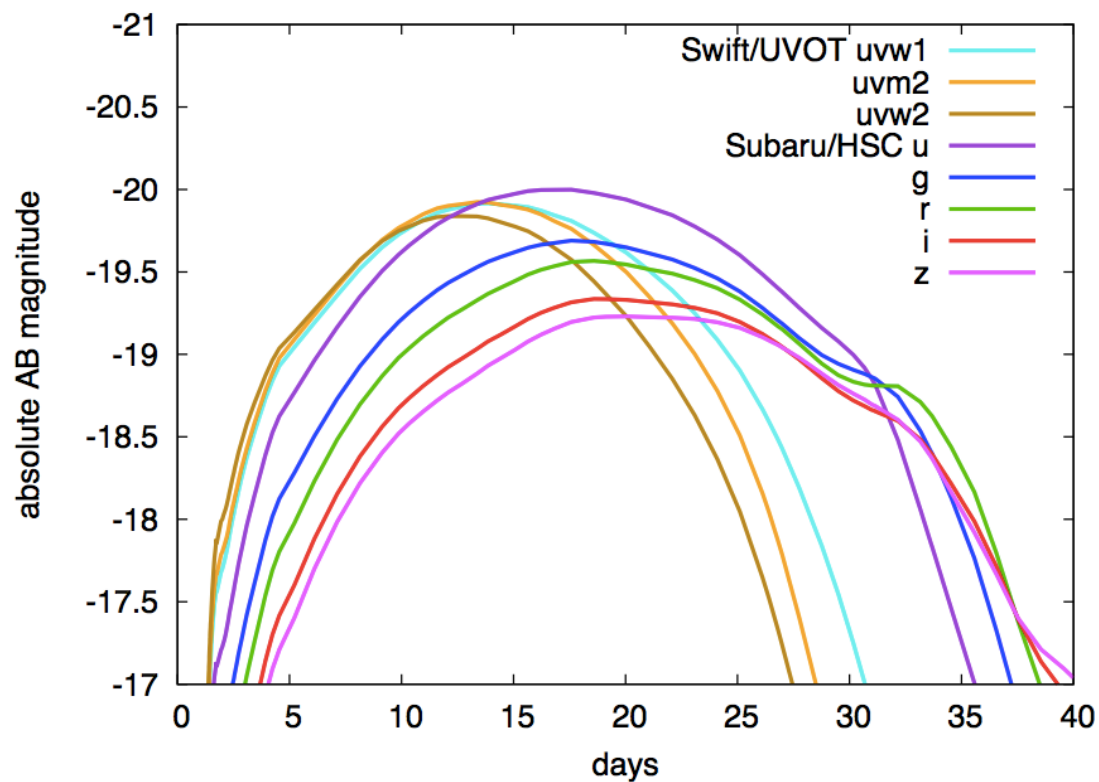
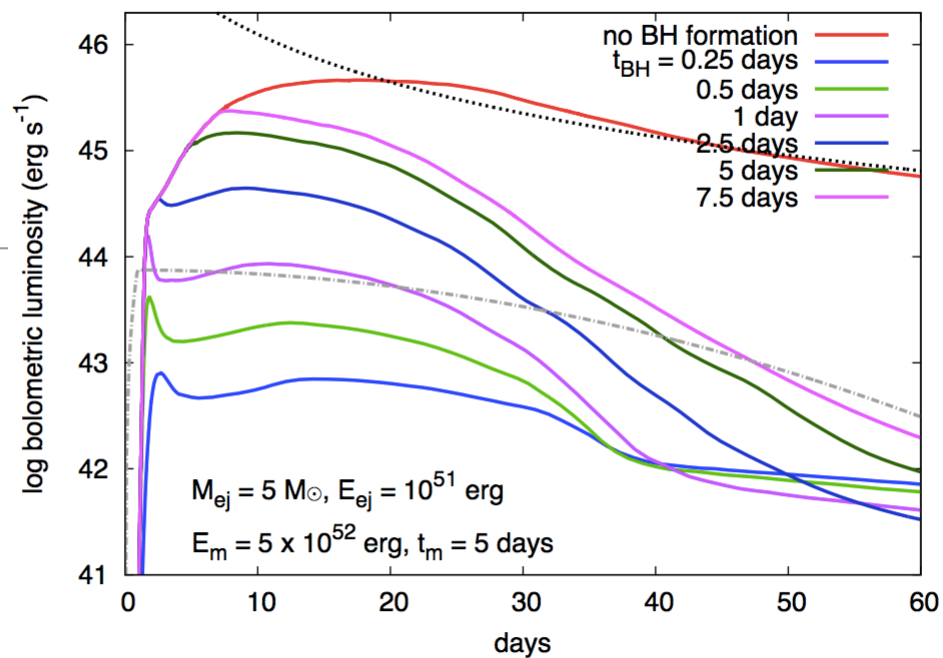
# Magnetar-driven shock breakout

- It can occur in magnetar-powered SNe without the BH transformation
  - but it is not significant



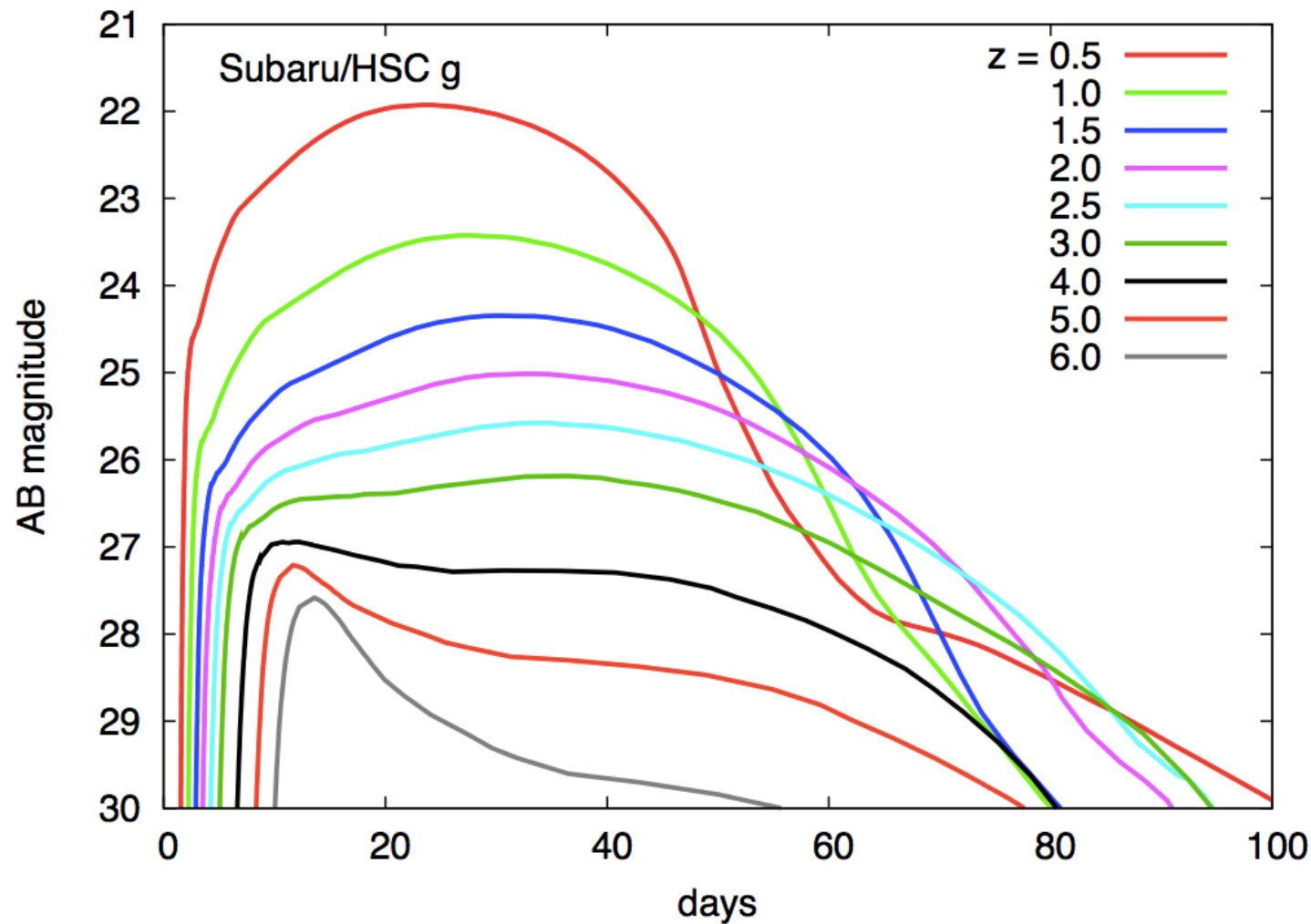
# Optical and NUV light curves

- faint in optical and NUV
- high photospheric temperature



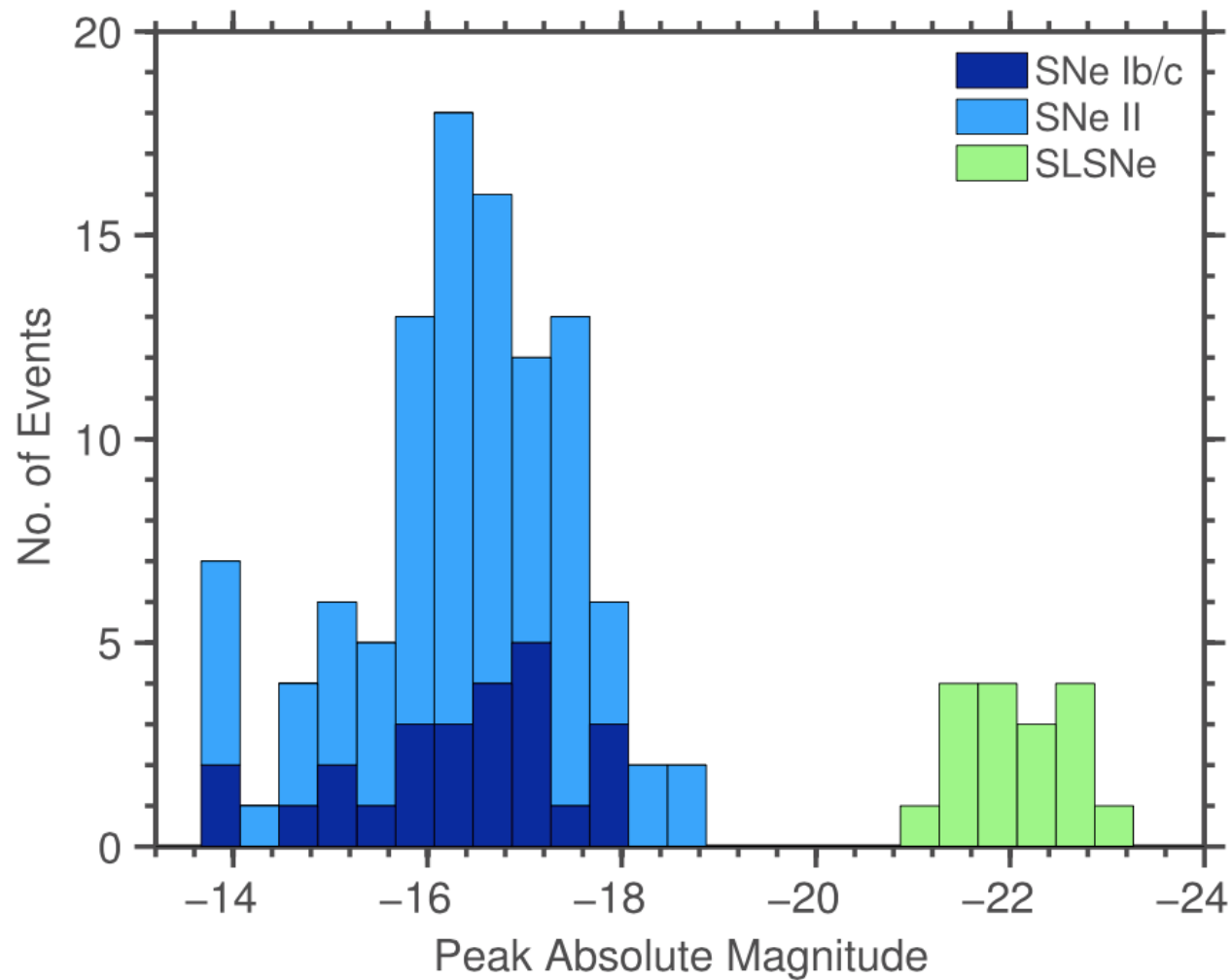
# Detecting magnetar-driven shock breakout

- we need those at high redshifts



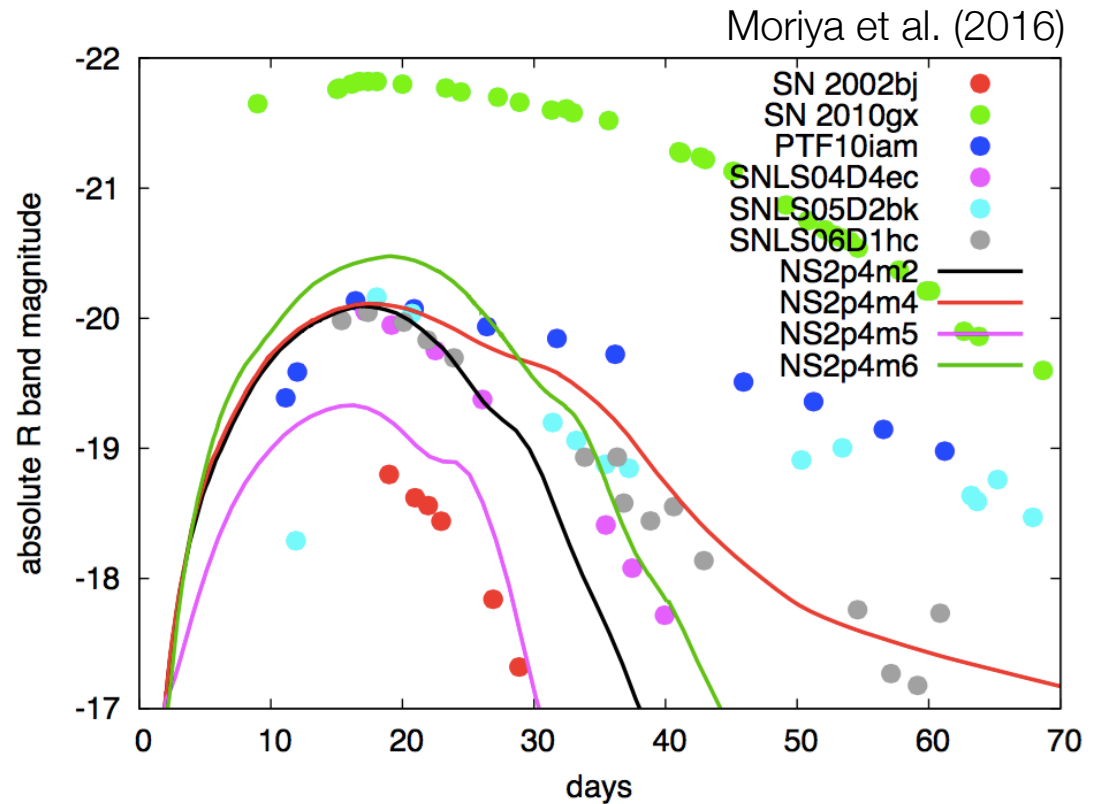
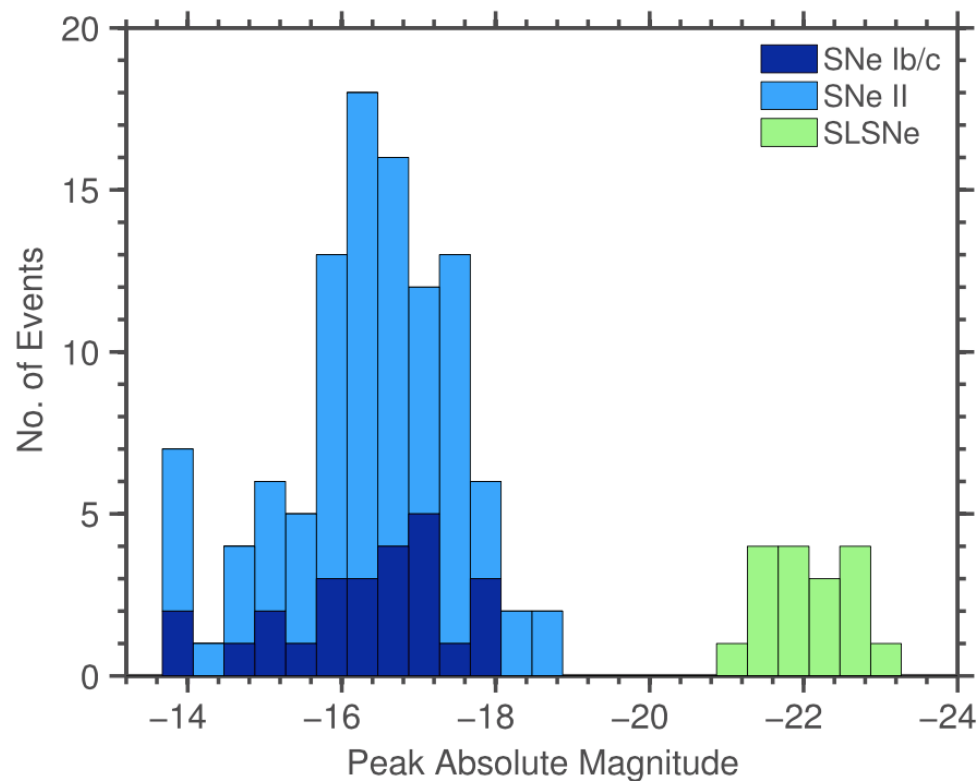
# SN - SLSN gap

- Arcavi et al. (2016)
  - possible luminosity jump?



# Some transients in the gap are observed

- Arcavi et al. (2016)
  - reported several transients in the gap
  - they could not fit them by the standard magnetar-powered SN models



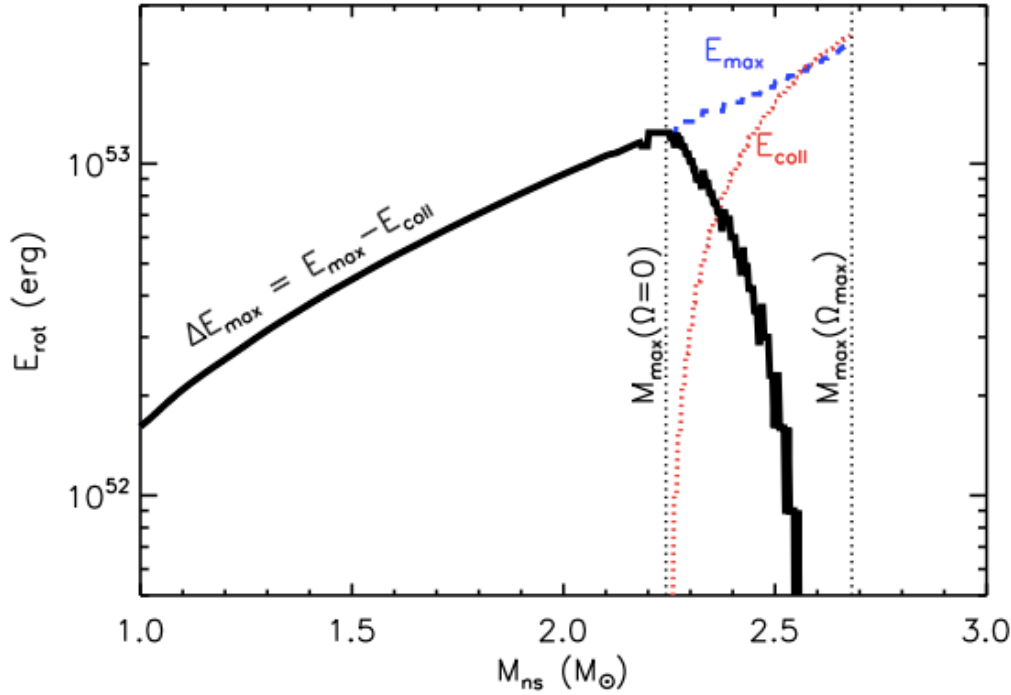


# Summary

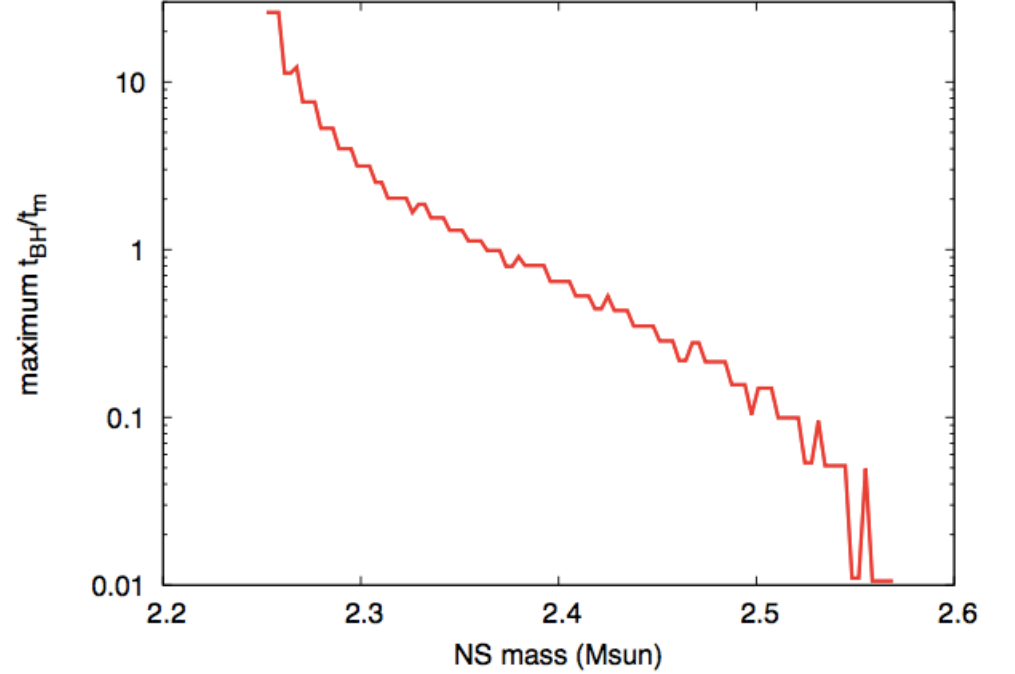
---

- some SNe, like SLSNe, may be powered by magnetar spin-down
- supramassive magnetars have minimum rotational energy to support themselves
  - they transform into BHs when they lose too much rotational energy
  - SNe powered by supramassive magnetars may suddenly lose their energy input because of the BH transformation
- the BH transformation can have significant effect on the observational properties of SNe powered by magnetars
  - magnetar-driven shock breakout signals can be more significant
  - they can be observed as the SN - SLSN gap transients
- if confirmed...
  - SLSNe to the gap transients ratio may indicate the mass fraction of normal mass NSs to supramassive NSs

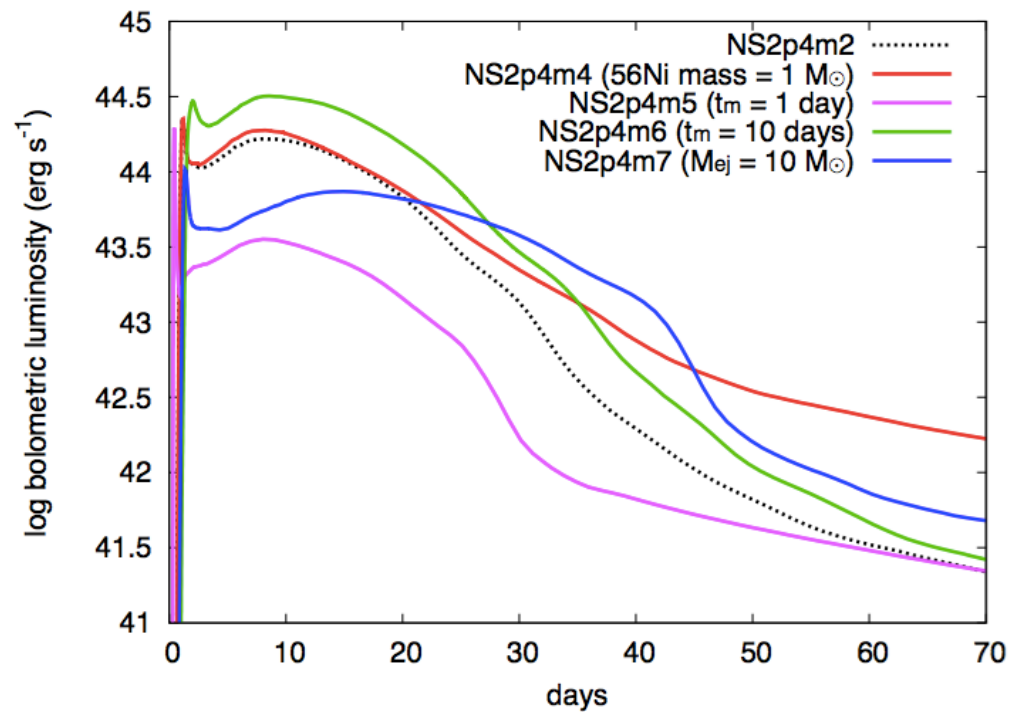
---



**Figure 4.** Maximum extractable rotational energy from an NS,  $\Delta E_{\max} \equiv E_{\max} - E_{\text{coll}}$  (black solid line), as a function of the NS mass  $M$ , where  $E_{\max}$  (blue dashed line) is the maximum rotational energy at the mass-shedding limit and  $E_{\text{coll}}$  (red dotted line) is the minimum rotational energy required for support of a supramassive NS against collapse to a BH. The structure of the solid-body rotating NS is calculated using the `rns` code assuming a parametrized piecewise polytropic EOS with an adiabatic index  $\Gamma = 3$  above the break density of  $\rho_1 = 10^{14.7} \text{ g cm}^{-3}$  at a pressure of  $P_1 = 3.2 \times 10^{34} \text{ dyn cm}^{-2}$ . We find that the maximum value of  $\Delta E_{\max}$  lies within the relatively narrow range of  $0.9\text{--}1.65 \times 10^{53} \text{ erg}$  across a wide range of  $\Gamma - P_1$  consistent with constraints on the maximum mass of a non-rotating NS.



**Figure 1.** Maximum allowed ratio of BH formation time  $t_{\text{BH}}$  to magnetar spin-down time  $t_m$ , as a function of the NS gravitational mass, based on Figure 4 in Metzger et al. (2015). The structure of the solid-body rotating NS is calculated using the `rns` code (Stergioulas & Friedman 1995) assuming a parametrized piecewise polytropic EOS with an adiabatic index  $\Gamma = 3$  above the break density of  $\rho_1 = 10^{14.7} \text{ g cm}^{-3}$  at a pressure of  $P_1 = 3.2 \times 10^{34} \text{ dyn cm}^{-2}$  (Margalit et al. 2015). The chosen EOS results in a  $1.4 M_{\odot}$  NS radius of 10.6 km and maximum non-rotating mass of  $\approx 2.24 M_{\odot}$ , consistent with observational constraints.



**Figure 7.** Numerical LC models with different SN ejecta and magnetar properties. The magnetar initial rotational energy ( $E_m = 1.1 \times 10^{53}$  erg) and the NS mass ( $2.4 M_\odot$ , i.e.,  $E_{\text{coll}} = 9.3 \times 10^{52}$  erg) are the same as those in NS2p4m2. The difference between NS2p4m2 ( $M_{\text{ej}} = 5 M_\odot$ ,  $E_{\text{ej}} = 10^{51}$  erg,  $^{56}\text{Ni}$  mass of  $0.1 M_\odot$ , and  $t_m = 5$  days) and the other models are indicated in the figure.